

Nucleon swelling and Compton scattering

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Problem of ^4He point nucleon density distribution

Variable proton size in CMD calculation

Quasifree Compton scattering

Different effective masses for nucleon

non - relativistic $[\frac{1}{2m} p^2 + U_S(r, \varepsilon)]\psi_S = \varepsilon\psi_S$

$$U_S = V_S + iW_S$$

- $\frac{m_S^*(\varepsilon)}{m} = 1 - \frac{d}{d\varepsilon} V_S(\varepsilon)$ (conventional effective mass)

- $\frac{\tilde{m}(\varepsilon)}{m} = [1 + \frac{m}{k} \frac{\partial}{\partial k} V_S]_{k=k(\varepsilon)}^{-1}$ (effective k - mass)

- $\frac{\bar{m}(\varepsilon)}{m} = [1 - \frac{\partial}{\partial \varepsilon} V_S]_{k=k(\varepsilon)}$ (effective E - mass)

$$\frac{m_S^*(\varepsilon)}{m} = \frac{\tilde{m}(\varepsilon)}{m} \cdot \frac{\bar{m}(\varepsilon)}{m}$$

relativistic $[\vec{\alpha} \cdot \vec{p} + \gamma_0(m + U_\sigma + \gamma_0 U_0)]\phi = E\phi$

- $\frac{M^*}{m} = 1 + \frac{V_\sigma}{m}$ (Dirac mass)

- $\frac{m_e^*(\varepsilon)}{m} = 1 - \frac{d}{d\varepsilon} V_e(\varepsilon) = 1 - \frac{V_0}{m}$ (Lorentz mass)

(Schroedinger equivalent)

QCD base

- NJL

$$f_\pi^* / f_\pi \sim m_\sigma^* / m_\sigma$$

$$\sim m_V^* / m_V \sim m_N^* / m_N$$

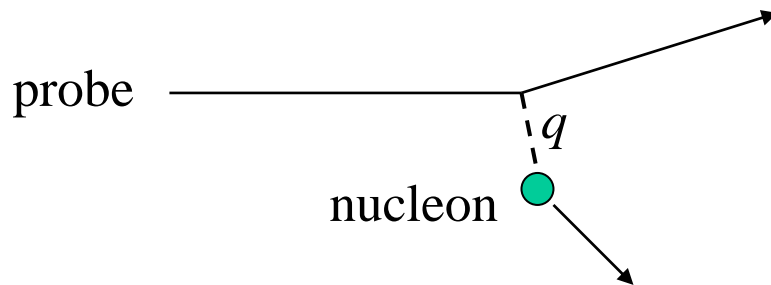
- QCD sum rule

Nucleon form factor (or size) in nuclear medium

Is it changed from free nucleon?

(medium effect other than mass shift)

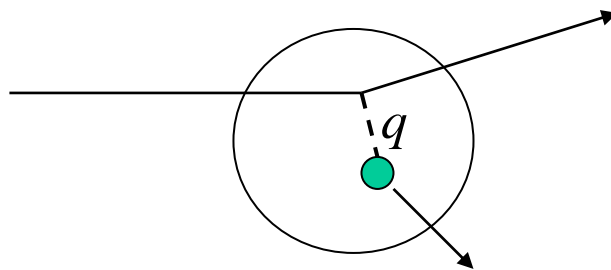
- elastic



[interaction \times form factor]

$$F(q^2) = 1 - \frac{1}{6} q^2 \langle r^2 \rangle + \dots$$

- quasi-elastic



If the exchange particle is a hadron

\rightarrow medium effect of itself

photon \rightarrow MEC

gluon ?

FSI (final state interaction) effect

Quark substructure approach to ${}^4\text{He}$ charge distribution

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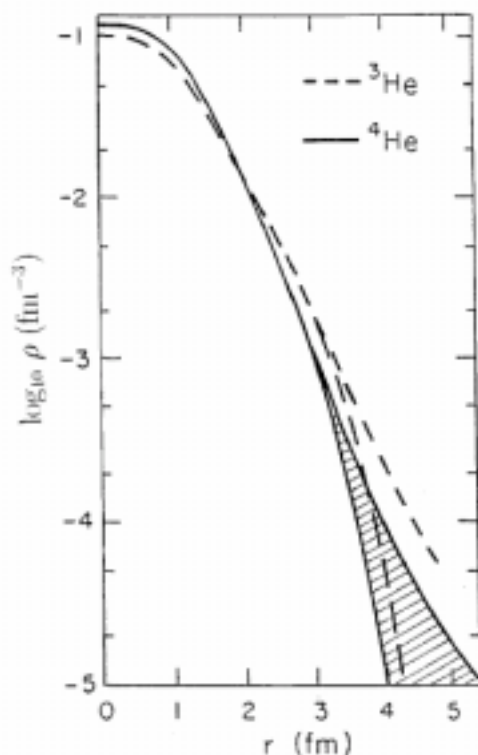


FIG. 1. Model-independent charge distributions for ${}^3\text{He}$ and ${}^4\text{He}$ extracted from experiment. Reproduced from McCarthy *et al.* [1], who state that "the extreme limits of $\rho(r)$ cover the statistical, systematic as well as the completeness error of the data."

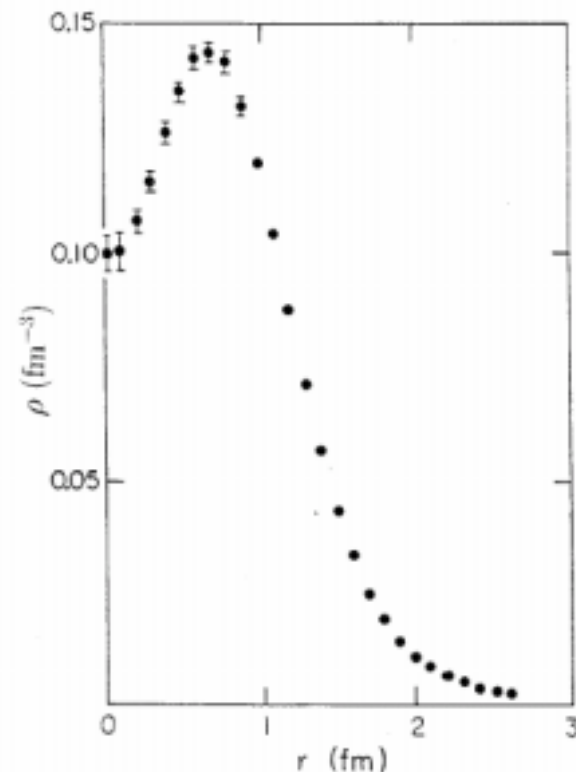


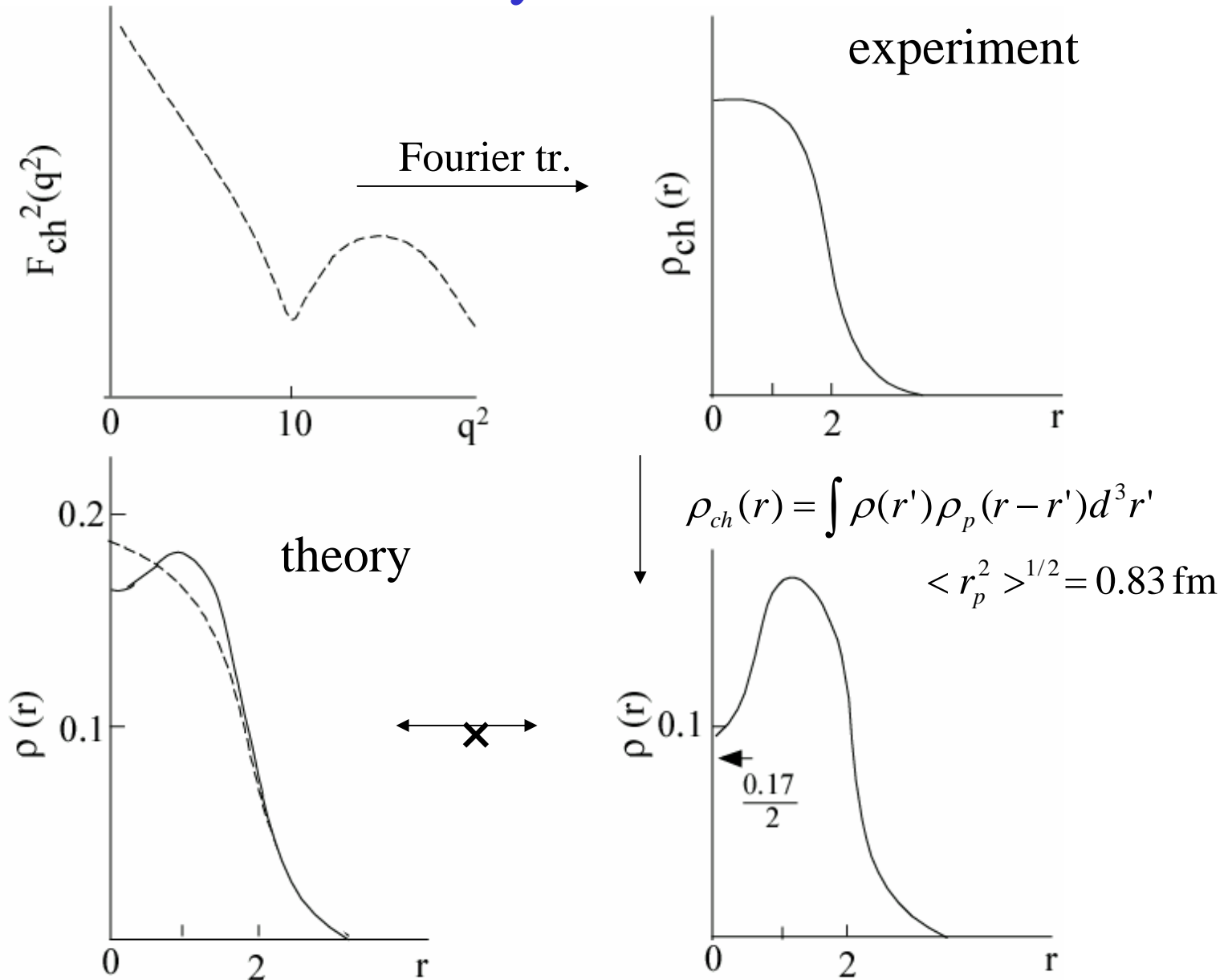
FIG. 2. Point-proton density distribution for ${}^4\text{He}$ obtained by unfolding the free proton form factor, allowing for meson exchange corrections and relativistic effects. Reproduced from Sick [2].

Nuclear density distribution

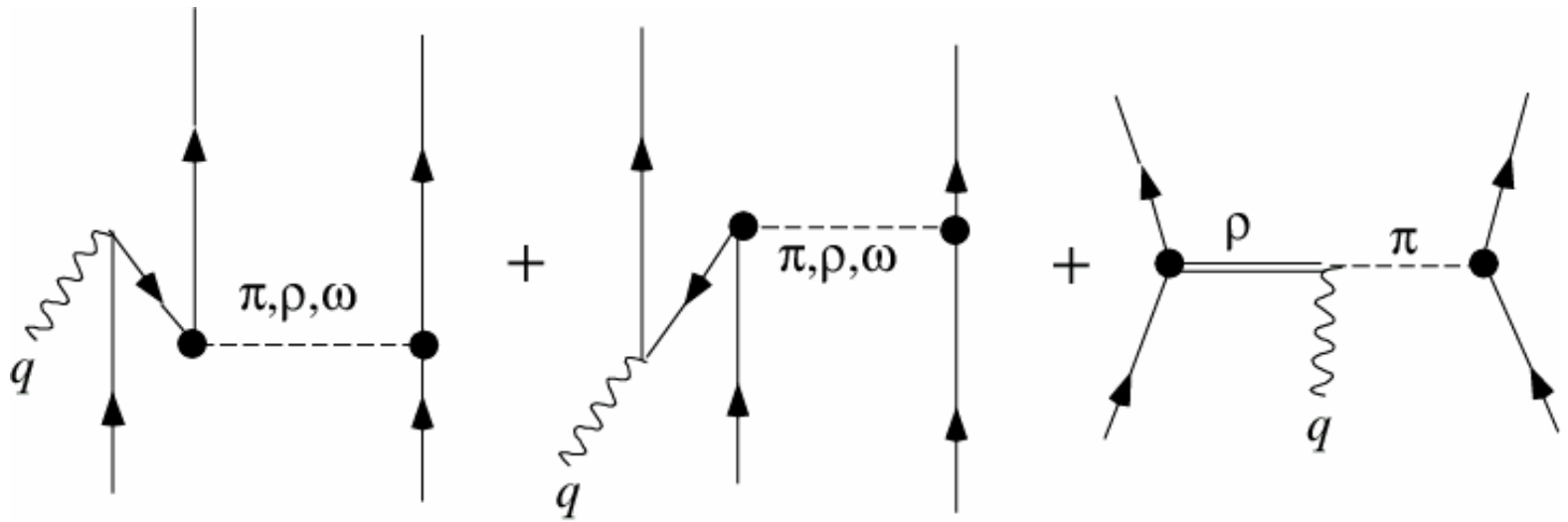
[Experiment]	(e, e') or μ -atom ← model-independent analysis charge density distribution fold $\rho_p + \text{MEC}$ ← unfold proton charge distribution ρ_p point proton density distribution ground state wave function ← Green function MC, etc. (many body) Schroedinger eq. (Dirac eq.)
[Theory]	NN interaction (2-body, 3-body)

Usually unfold (fold) the charge distribution of free proton
→ If proton is “**swelling**” in the nucleus,?

Point nucleon density distribution in ^4He



Usually it is due to MEC (meson exchange current)



Chromodielectric Soliton Model

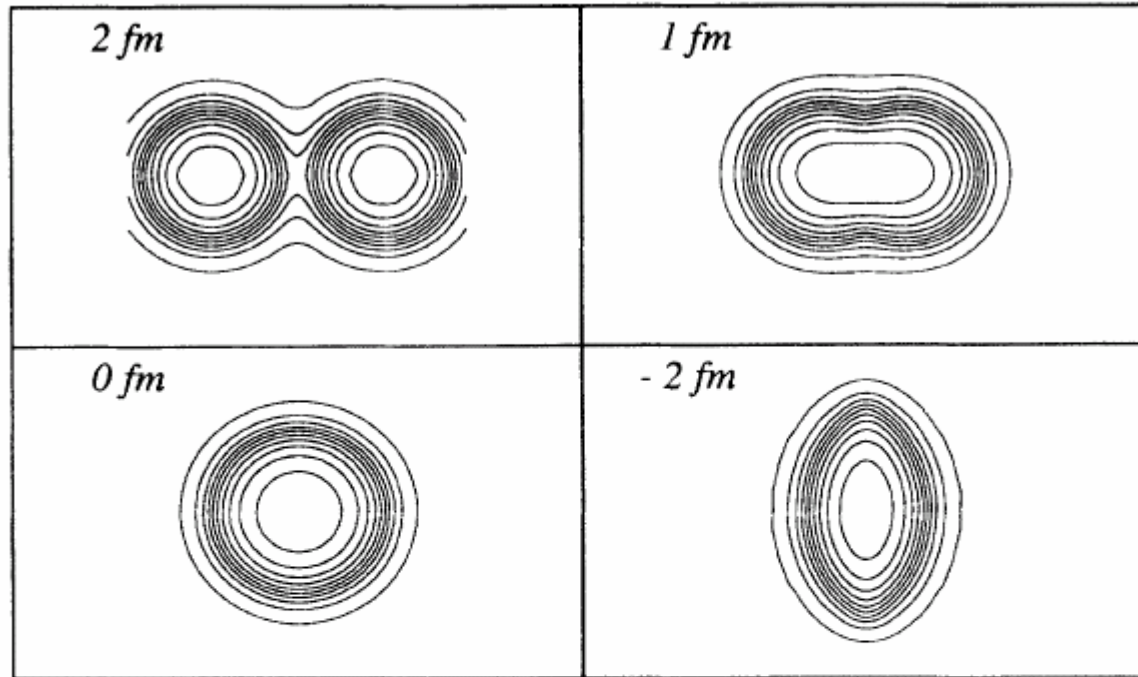


FIG. 1. The scalar field, $\sigma_\alpha(\mathbf{r})$ of Eq. (16), from which the single-quark wave functions are generated, for four different values of the deformation parameter α between 2 fm and -2 fm. The fields correspond to the parameter set with $f = 3$ and $c = 10\,000$, and are shown with equal increments between adjacent contours.

Variable rms proton charge radius

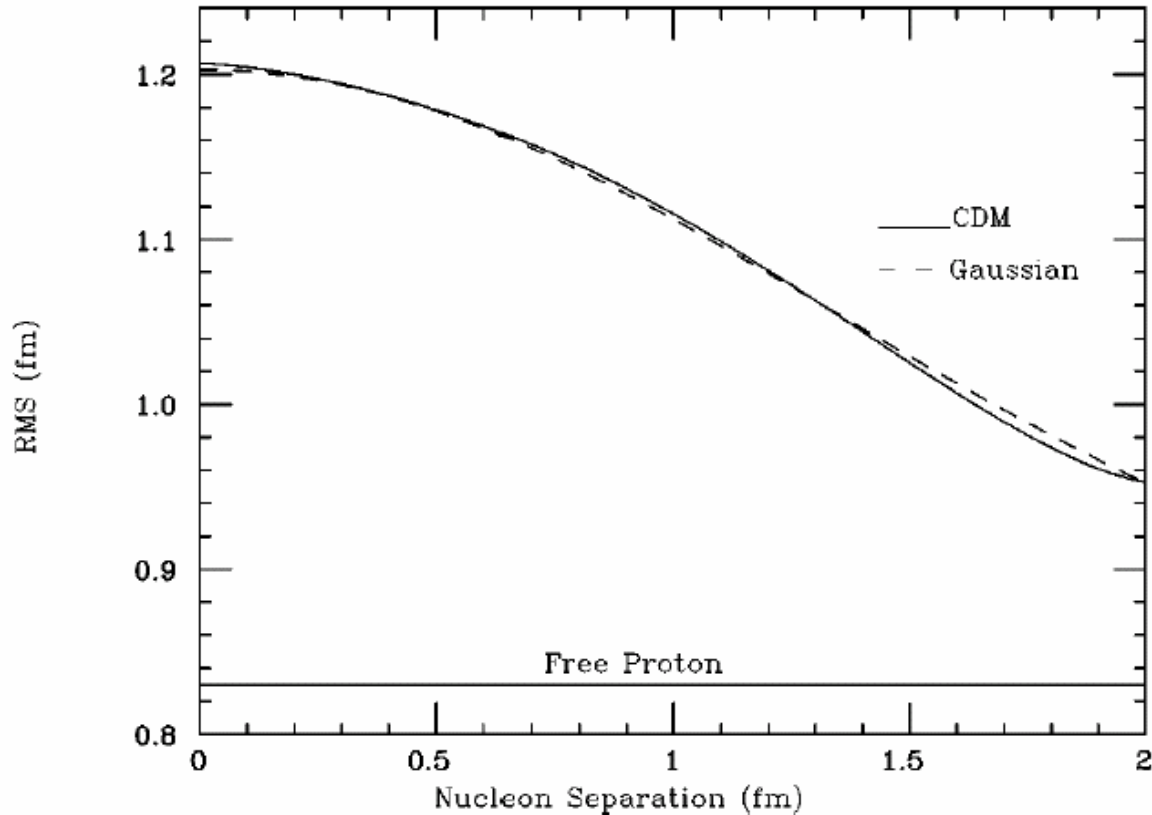


FIG. 3. Proton rms charge radius r_p of Eq. (1) as a function of internucleon separation. The line labeled CDM is the calculated chromodielectric model result. The dashed line is a Gaussian approximation, normalized to the free value, with a size parameter given by Eq. (3)

Folding variable proton size to the point proton density distribution

- assume Gaussian distribution $b(r') = \sqrt{2/3}r_p(r')$.
- 2-nucleon pair charge density

$$\rho_{\text{pair}}(\mathbf{r}_i, \mathbf{r}_j; \mathbf{r}) = \{ \delta_{ip} \exp[-|\mathbf{r} - \mathbf{r}_i|^2 / b^2(r_{ij})] + \delta_{jp} \exp[-|\mathbf{r} - \mathbf{r}_j|^2 / b^2(r_{ij})] \} / \pi^{3/2} b^3(r_{ij}),$$

- independent pair approximation

$$\rho_{\text{ch}}(r) = \frac{1}{3} \sum_{i < j} \int d^3 r_i \int d^3 r_j f_2(\mathbf{r}_i, \mathbf{r}_j) \rho_{\text{pair}}(\mathbf{r}_i, \mathbf{r}_j; \mathbf{r}).$$
$$f_2(\mathbf{r}_1, \mathbf{r}_2) = \int |\psi(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3, \mathbf{r}_4)|^2 d^3 r_3 d^3 r_4,$$

^4He density distribution

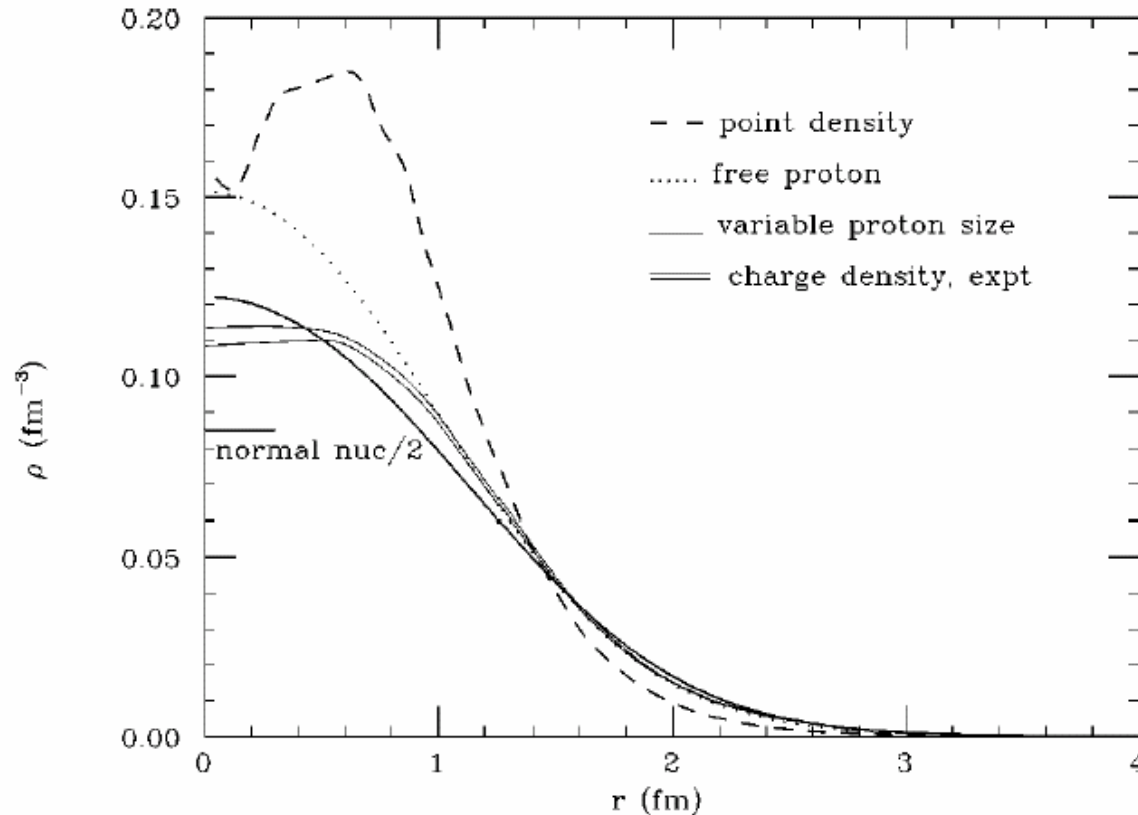
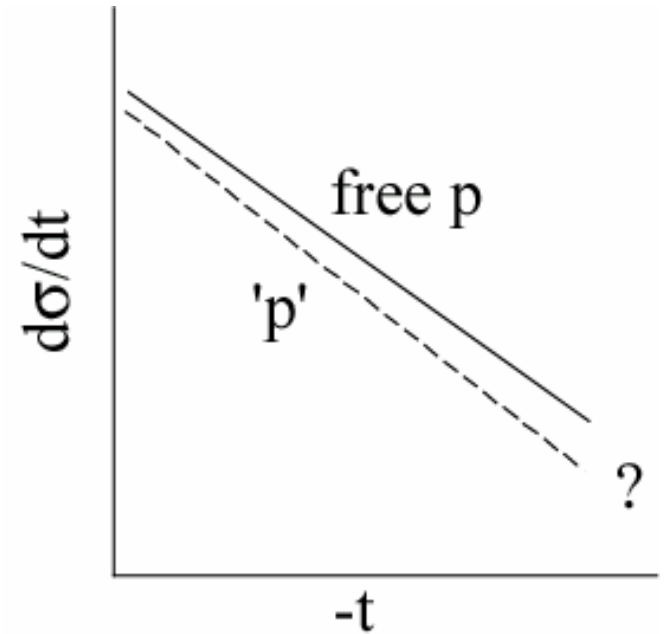
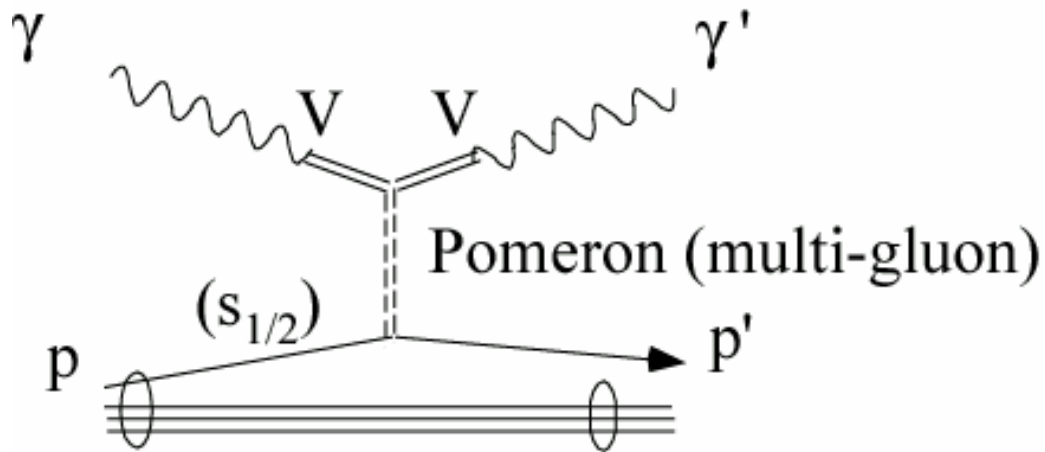


FIG. 4. ^4He density distributions: The dashed line is the point density from a parametrized Green's function Monte Carlo calculation. The curve labeled "free proton" is the charge distribution obtained from a Gaussian proton charge distribution with a fixed size parameter (as is usually done). The curve labeled "variable proton size" uses the Gaussian fit of Fig. 3. We also indicate half the normal nuclear density $0.17/2 \text{ fm}^{-3}$.

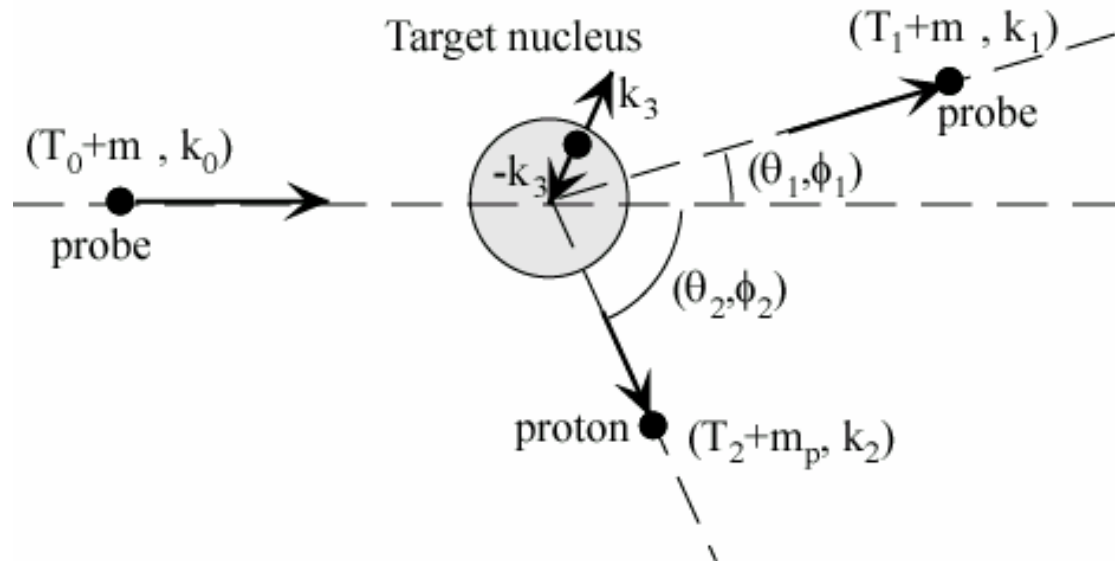
${}^4\text{He}(\gamma, \gamma p){}^3\text{H}$ in VD energy region



ISI, FSI: small
 no MEC ?
 $E_\gamma \gg E_{\text{sep}}$

+
 Polarization observables

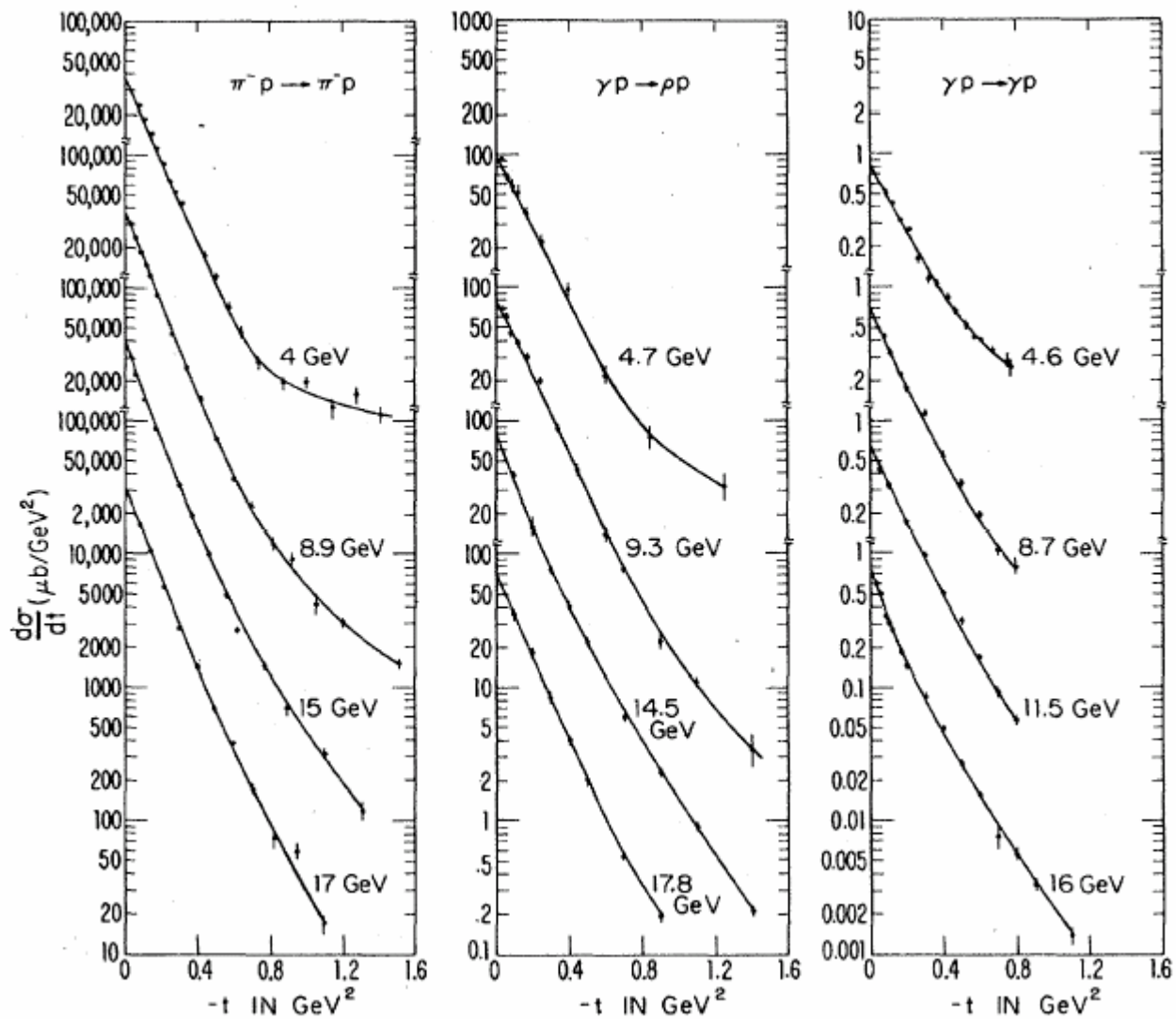
Quasifree kinematics



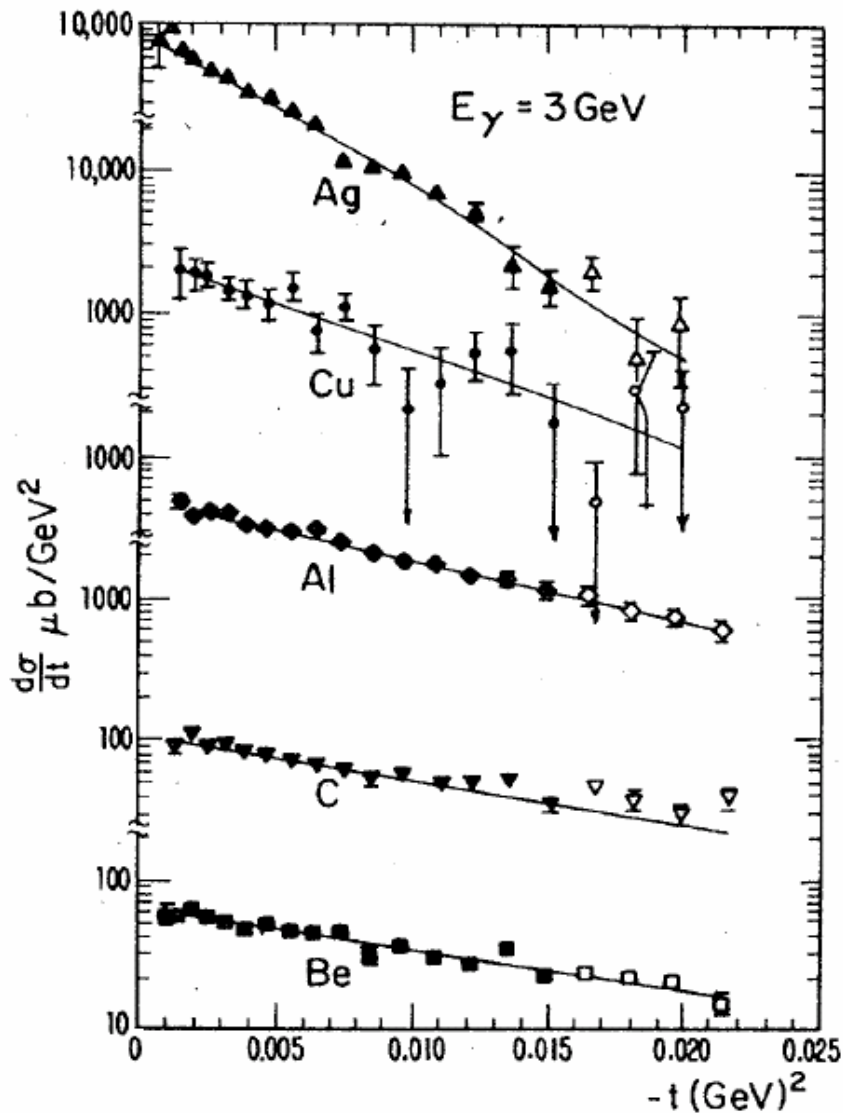
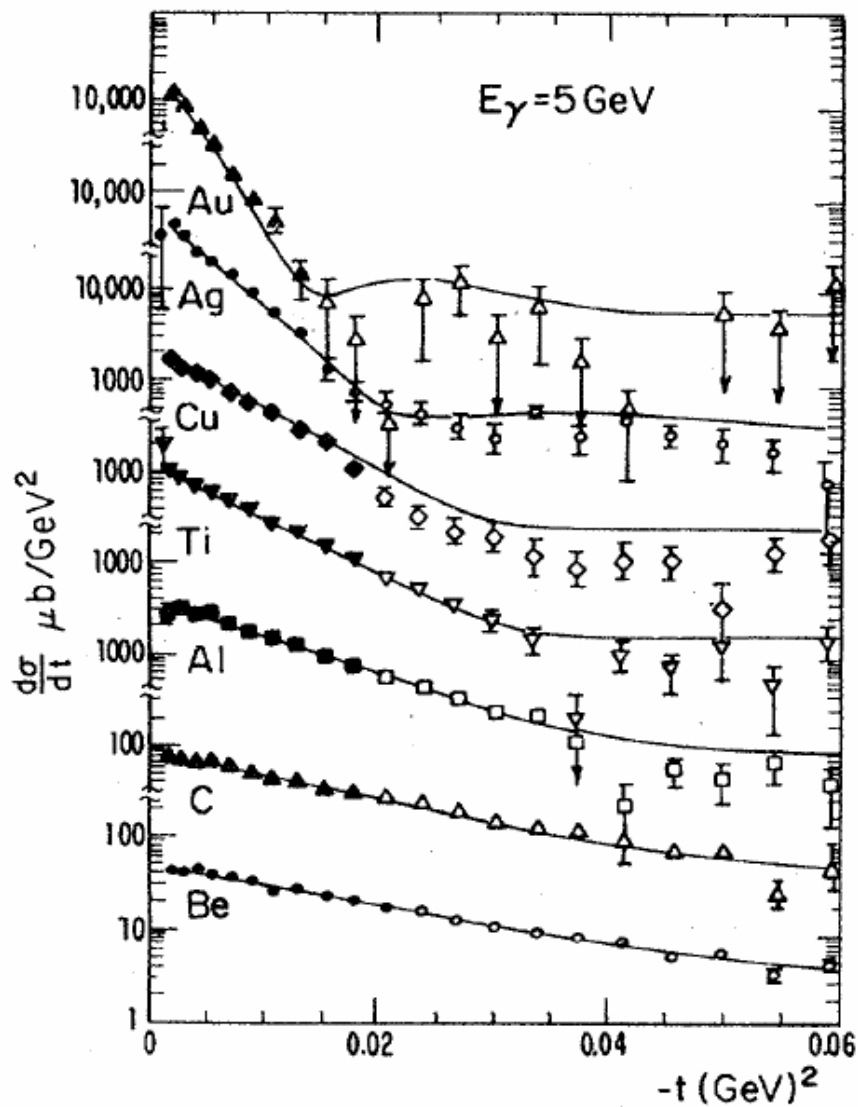
$$\vec{k}_3 = \vec{k}_0 - \vec{k}_1 - \vec{k}_2$$

$$E_{sep} = T_0 - (T_1 + T_2 + T_3)$$

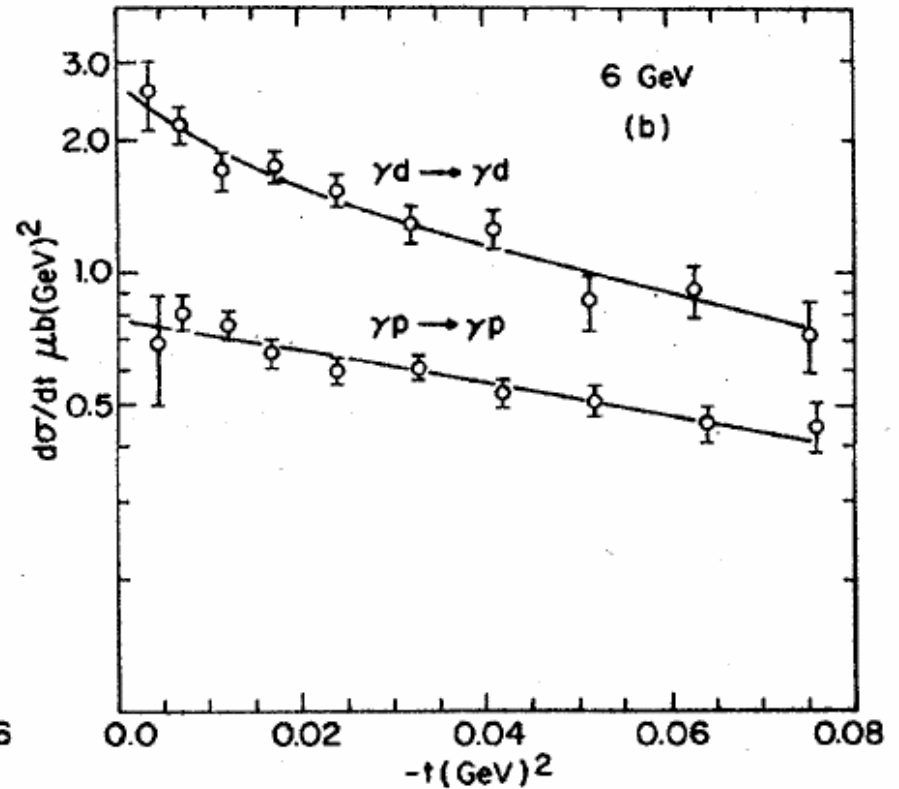
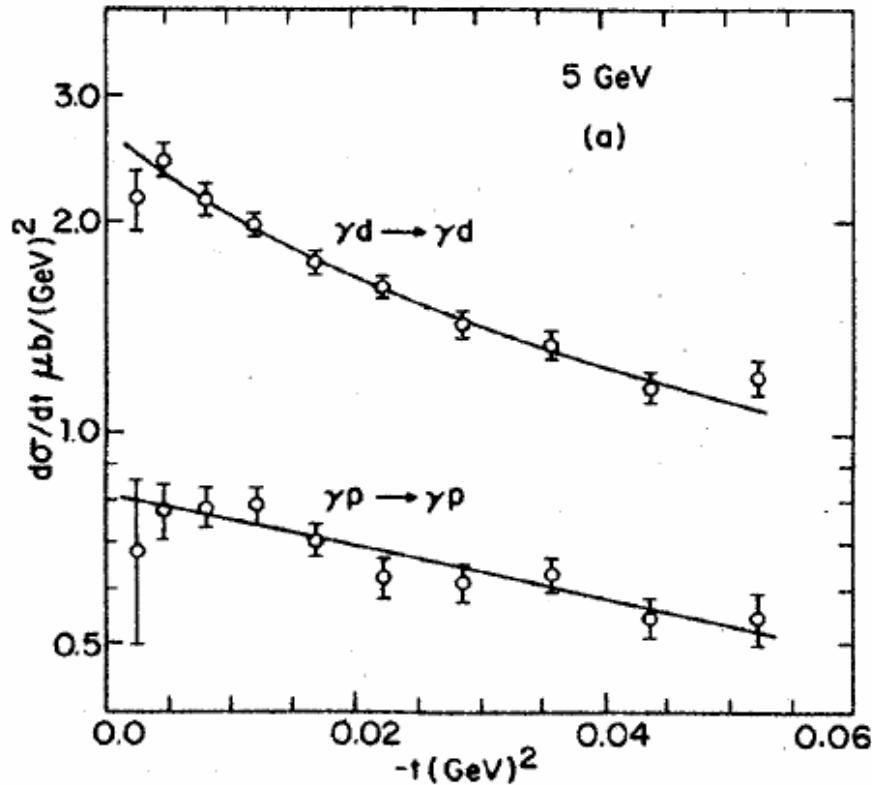
γp γp data above GeV energy (1960's-1970's)



Compton scattering for nuclear target (coherent + incoherent)



γd γd data



after 1990, many precise γp γp below 1 GeV
→ Nucleon polarizability

レーザー電子光によるクォーク核物理専用施設実行計画書 (平成8年)

GeV領域偏極ガンマ線による陽子コンプトン散乱実験

1・実験提案者

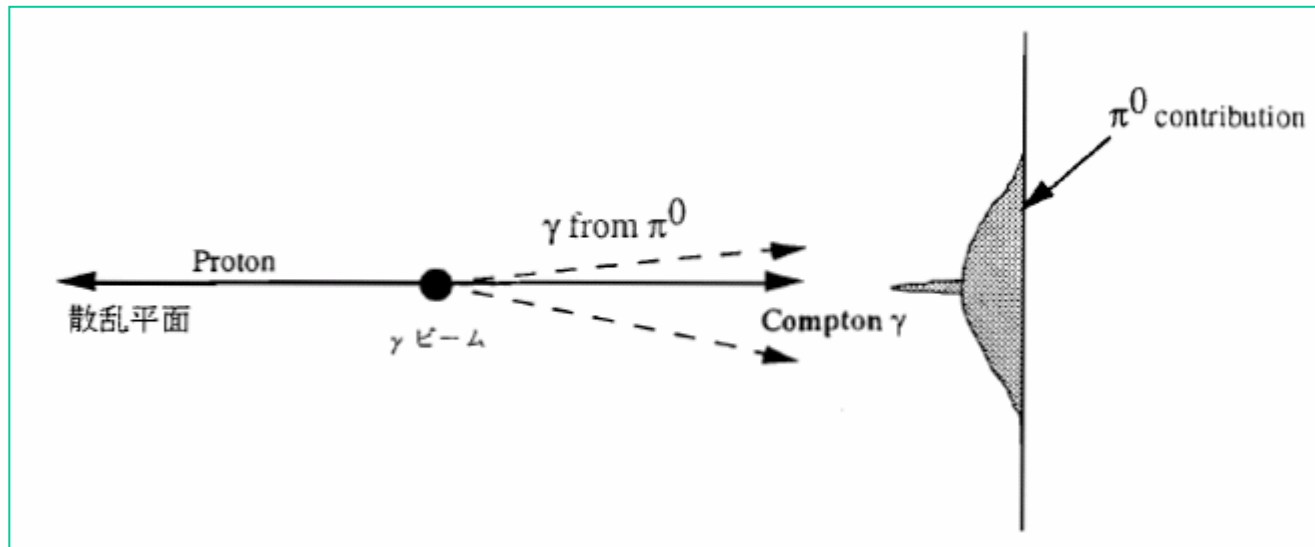
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2・目的:

GeV領域の偏極 γ 線による陽子コンプトン散乱反応の断面積及び偏極分解能を測定し、核子構造を調べる。

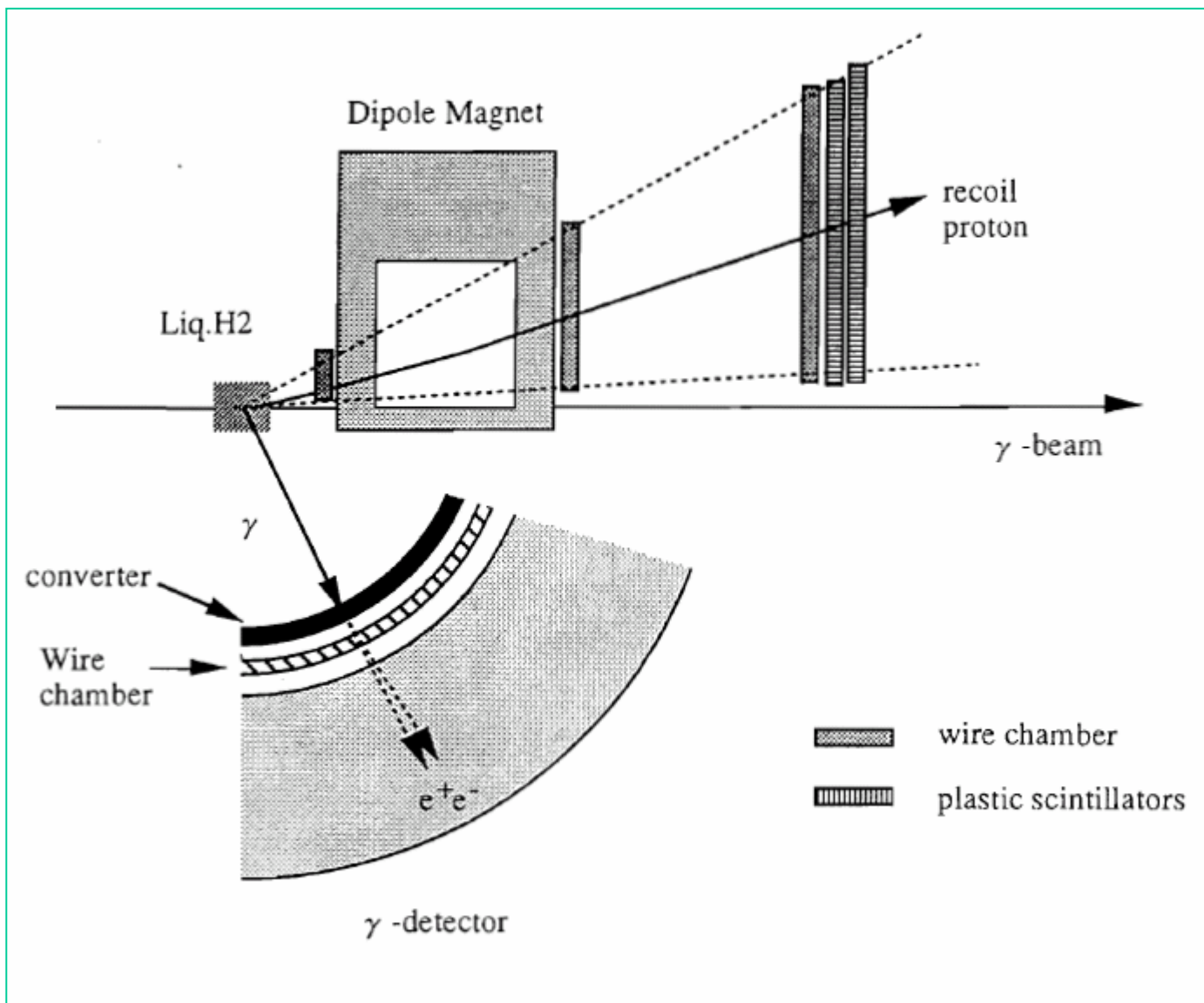
Identification of Compton events

coincidence measurements of scatted- γ and recoil-p



Co-planarity is important for Compton and π_0 separation

Detection system (by T.Suda)



Yield estimation

- $d\sigma/d\Omega \sim 1 \text{ nb/sr}$
- $N_\gamma \sim 10^7 / \text{sec}$
- $Nt \sim 1 \text{ mol}$

$$\rightarrow Y = 0.006 \Delta\Omega \quad (\Delta\Omega \sim 1 \text{ sr})$$

need one order high intensity to divide E_γ into 10 energy bins