

# Baryon resonances as hadronic molecular states

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in collaboration with

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$\Lambda(1405)$  in few-body systems  
three-body quasibound states

$N^*$



# Introduction

## Baryon resonances : decay with strong interactions

in understanding the structure of baryon resonances

**dynamical aspect** (*hadron dynamics*)

decaying resonance  $\rightarrow$  large hadronic components

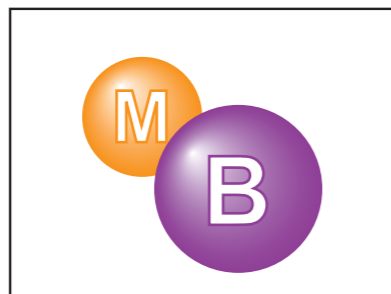
hadron dynamics is important

**symmetry aspect** (*quark dynamics*)

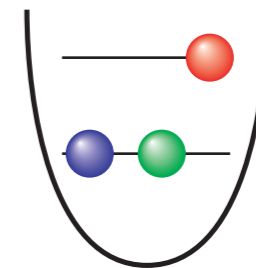
symmetry in terms of quarks

ex. chiral partners: N(1535) chiral partner of nucleon ??

cluster picture



shell model picture



- difference in range of dynamics
- meson cloud effects

# Hadronic molecular states

- *system of multiple hadrons described by hadron dynamics*

typical constituents are ground states hadrons

octet baryons:  $N, \Lambda, \Sigma, \Xi$     octet meson:  $\pi, K, \eta$

**mesons are key ingredients** (potential picture ok ?)

- *constituents keep their identity*

typical binding energy  $\sim 10\text{-}20$  MeV    weakly bound system

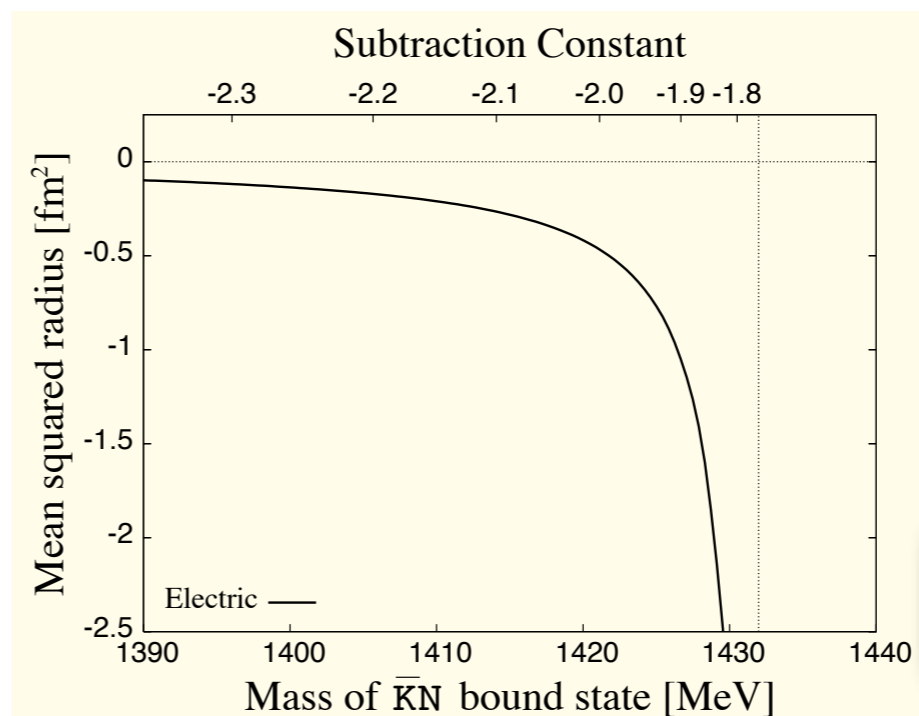
decay width  $\sim 50$  MeV (strong interactions)    **quasi-bound state**

- *spatially extended (large size)*

typically more than 1 fm

- *softer form factors*

strong energy dependence  
in production



Sekihara, Hyodo, DJ  
PLB669, 133 (2008).

**quark degrees of freedom may be less important**

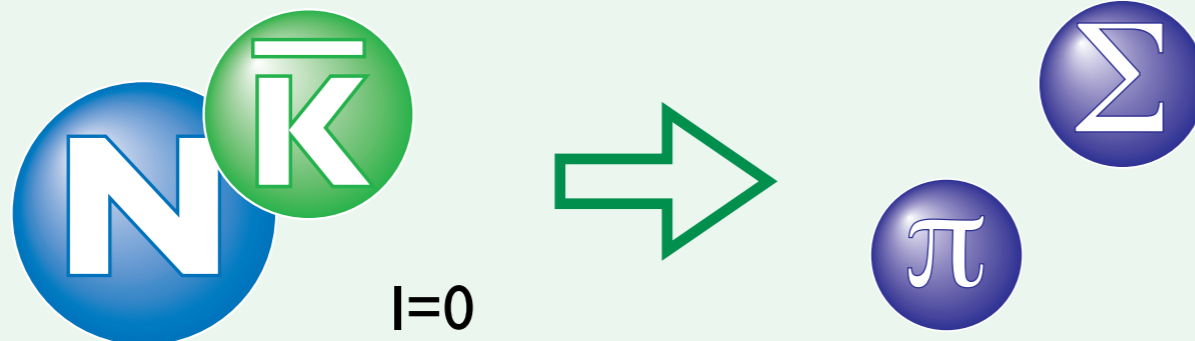
# $\Lambda(1405)$ as Quasi-bound state of $K^{\text{bar}}N$

$\Lambda(1405) \quad J^\pi = 1/2^-, \quad I = 0, \quad S = -1, \quad Q = 0$

- most established resonance, seen in many exp. clearly
- mass :  $1406.5 \pm 4.0$  MeV, (below  $K^-p$  threshold)
- width :  $50 \pm 2$  MeV (PDG)
- decay mode  $\Lambda(1405) \rightarrow (\Sigma\pi)_{I=0} \quad 100\% \quad \mathbf{s\text{-wave}}$

$\bar{K}N$	<u>1435 MeV</u>	$\Lambda(1405)$
$\pi\Sigma$	<u>1331 MeV</u>	

## $\Lambda(1405)$ : quasi-bound states of $K^{\text{bar}}N$



have been considered as a quasi-bound state of  $K^{\text{bar}}N$  since late 50's

Recently this idea has been developed to light nuclear system

$\Lambda(1405)$  is described as a superposition of two states.

**important to understand kaons in nuclei**

Dalitz, Tuan, PRL 2, 425 ('59)

Akaishi, Yamazaki, PRC 65, 044005 ('02)

DJ, Oller, Oset, Ramos, Meissner  
NPA725, 181 ('03)

**$\Lambda(1405)$  is composed mostly of meson-baryon components.**

Hyodo, Jido, Hosaka, PRC78, 025203 ('08)

# Double pole structure of $\Lambda(1405)$

DJ, Oller, Oset, Ramos, Meissner  
NPA725, 181 ('03)

$\Lambda(1405)$  is a superposition of two states **having different properties.**

there are two attractive channels

group theoretically

SU(3) singlet and octet

physically

$K^{\text{bar}}N$  and  $\pi\Sigma$  [Hyodo, Weise, PRC77, 035204 (08)]

$\Lambda(1405)$

below threshold of  $\bar{K}N$

pole 1 : 1390 - 66i

- wider width

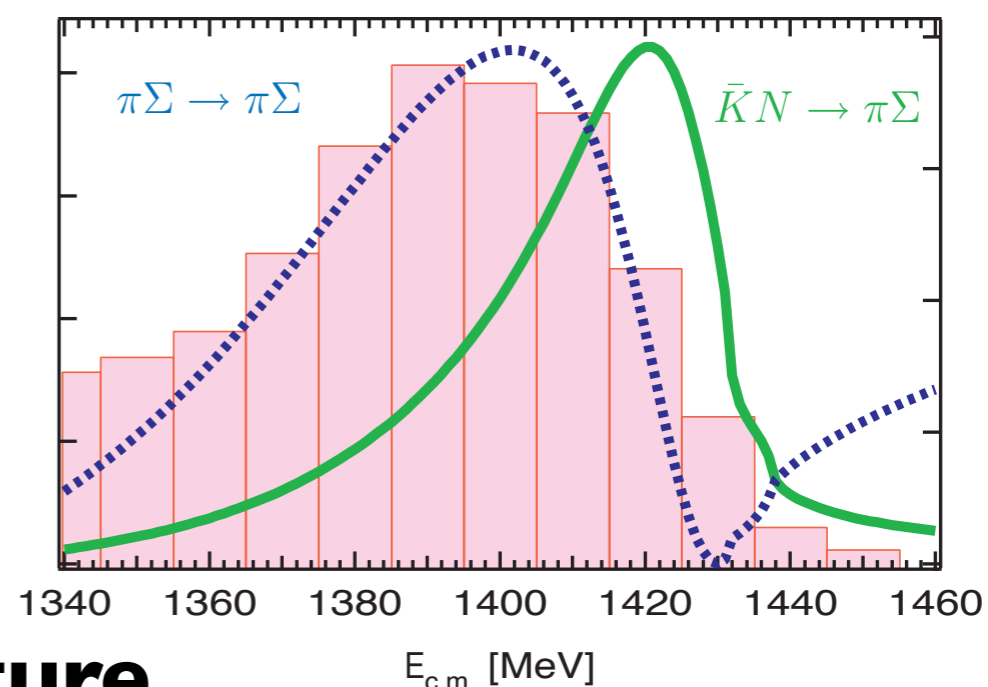
- strongly couples to  $\pi\Sigma$  state

pole 2 : 1426 - 16i

- narrower width

- dominantly couples to  $\bar{K}N$  state

$\pi\Sigma$  Mass distribution



## Implication of double pole structure

$\Lambda(1405)$  spectrum is dependent on channels

Resonance position in  $K^{\text{bar}}N$  channel  $\sim 1420$  MeV with narrower width  
instead of 1405 MeV

This 15 MeV difference is important for  $K^{\text{bar}}N$  interactions

# $\Lambda(1405)$ in few-body systems ~hadronic molecule states with kaons~

**three-body quasibound states**  $I=1/2, J^P=1/2^+$



in collaboration with **Y. Kanada-En'yo** (YITP, Kyoto)

## References

- Y. Kanada-En'yo and D. Jido, **Phys. Rev. C78, 025212 (2008)**
- D. Jido and Y. Kanada-En'yo, **Phys. Rev. C78, 035203 (2008)**

# Peculiarities of K meson

small binding energy  $\sim 10\text{-}30$  MeV    small kinetic energy

- **heavy particle** compared with kinetic energy

half of nucleon mass

cf. pion     $m_\pi \approx 140$  MeV

non-relativistic potential model

isospin averaged mass

$$m_K = 495.7 \text{ MeV}$$

$$m_N = 938.9 \text{ MeV}$$

- **Nambu-Goldstone boson**

smaller mass compared with typical hadron mass scale

strong s-wave attraction in  $K^{\text{bar}}N$

chiral effective theory    momentum expansion

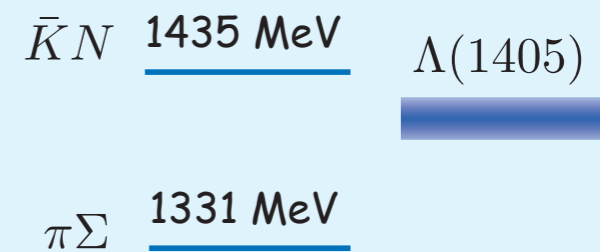
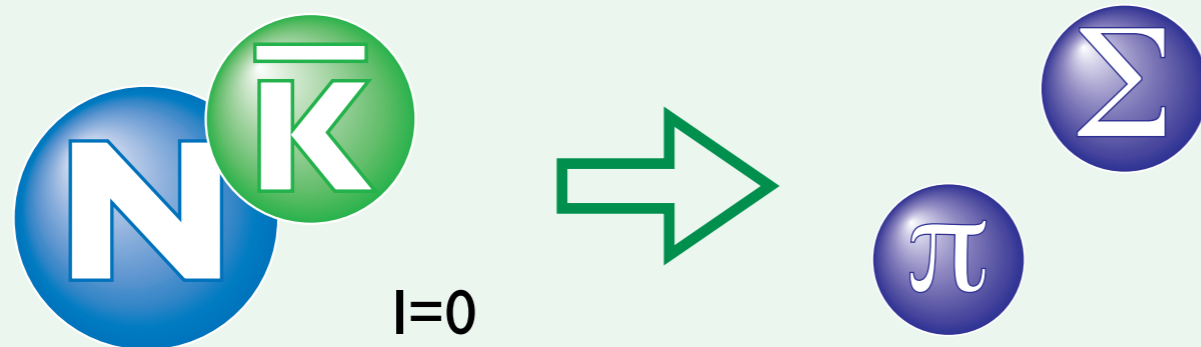
s-wave int. proportional to K energy

**Kaons are different from pions** in the energies of our interest !!

# Hadronic molecular states in two-body sys.

## Assumption:

$\Lambda(1405)$  : one of the historical examples



have been considered as a quasi-bound state of  $K^{\text{bar}}N$  since late 50's

Recently this idea has been developed to light nuclear system

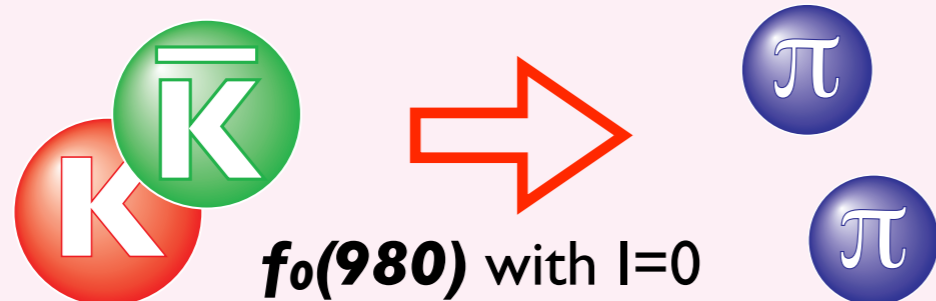
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DJ, Oller, Oset, Ramos, Meissner  
NPA725, 181 ('03)

$f_0(980)$ ,  $a_0(980)$  : the candidates of  $KK^{\text{bar}}$  QBS  
scalar mesons



Weinstein, Isgur, PRL 48, 659 ('90)



$\Lambda(1405)$  $f_0(980), a_0(980)$ 

### Interactions in $KK^{\text{bar}}N$ system

		$I=0$	$I=1$	threshold
$N^*$ 	$\bar{K}N$	attraction $\Lambda(1405)$	weak attraction	1434.6 MeV
	$K\bar{K}$	$f_0(980)$	$a_0(980)$	991.4 MeV
	$KN$	repulsion very weak	strong repulsion	1434.6 MeV

### Interactions in $K^{\text{bar}}K^{\text{bar}}N$ system

		$I=0$	$I=1$	threshold
$\Xi^*$ 	$\bar{K}N$	attraction $\Lambda(1405)$	weak attraction	1434.6 MeV
	$\bar{K}\bar{K}$	repulsion weak	strong repulsion	991.4 MeV

**Formulation** for three-body system:  $K\bar{K}N$  and  $\bar{K}\bar{K}N$

## non-relativistic potential model

### Hamiltonian

$$H = T + V_{\bar{K}N}(r_1) + V_{KN}(r_2) + V_{K\bar{K}}(r_3),$$

two-body effective interactions

**local potentials** obtained by s-wave two-body scattering

**Gaussian potential**  $V(r) = U \exp[-(r/b)^2]$

**complex potentials** to implement coupled-channels effects

no three-body interactions ex:  $\bar{K}N \rightarrow \pi\Sigma$

no transitions to two hadrons  $\bar{K}\bar{K}N \rightarrow MB$

*will be suppressed in hadronic molecular states*

### recipe

1st: solve Schrödinger eq. without imaginary potential

obtain wavefunction  $\Psi$  and real part of energy

2nd: estimate imaginary part of energy  $E^{\text{Im}} = \langle \Psi | \text{Im}V | \Psi \rangle.$

# Effective interactions

$$\bar{K}N \quad V_{\bar{K}N} = U_{\bar{K}N}^{I=0} \exp[-(r/b)^2] + U_{\bar{K}N}^{I=1} \exp[-(r/b)^2]$$

## Hyodo-Weise potential (HW-HNJH)

PRC77,035204 (08)

derived from chiral unitary approach

energy dependent, but small in energy of interest

resonance position ~ **1420 MeV** (double pole structure)

interaction range  $b = 0.47$  fm

## Akaishi-Yamazaki potential (AY)

PRC64,044005 (02)

obtained phenomenologically

$I=0$  : reproduce  $\Lambda(1405)$  as quasi-bound state of  $K^{\text{bar}}N$

mass: **1405 MeV**, width: 40 MeV

$I=1$ : scattering and Konic atom data

interaction range  $b = 0.66$  fm

# Effective interactions

Gaussian potential  $V(r) = U \exp[-(r/b)^2]$

$\bar{K}N$   
attractive

HW-HNJH and AY potentials

binding energy : 11 MeV (HW), 31 MeV (AY)

$K^{\text{bar}}$ -N distance : **1.9 fm** (HW), **1.4 fm** (AY)

$K\bar{K}$   
attractive

reproduce masses and widths of  $f_0$  and  $a_0$

mass: 980 MeV, width: 60 MeV

reproduced

binding energy : 11 MeV

$K$ - $K^{\text{bar}}$  distance : **2.1 fm**

**PDG**

mass:  $980 \pm 10$  MeV

width: 40~100 MeV

mass:  $984 \pm 1.2$  MeV

width: 50~100 MeV

$KN$   
repulsive

reproduce scattering lengths

experimental data

$$a_{KN}^{I=0} = -0.035 \text{ fm}$$

$$a_{KN}^{I=1} = -0.310 \pm 0.003 \text{ fm}$$

# Result $KK^{\text{bar}}N$

$N^*$  around 1900 MeV

$K\bar{K}N$  is bound below thresholds of  $\Lambda(1405)+K$ ,  $a_0(f_0)+N$

## - loosely bound system

binding energy	width
HW: 19 MeV	88 MeV
AY: 39 MeV	98 MeV

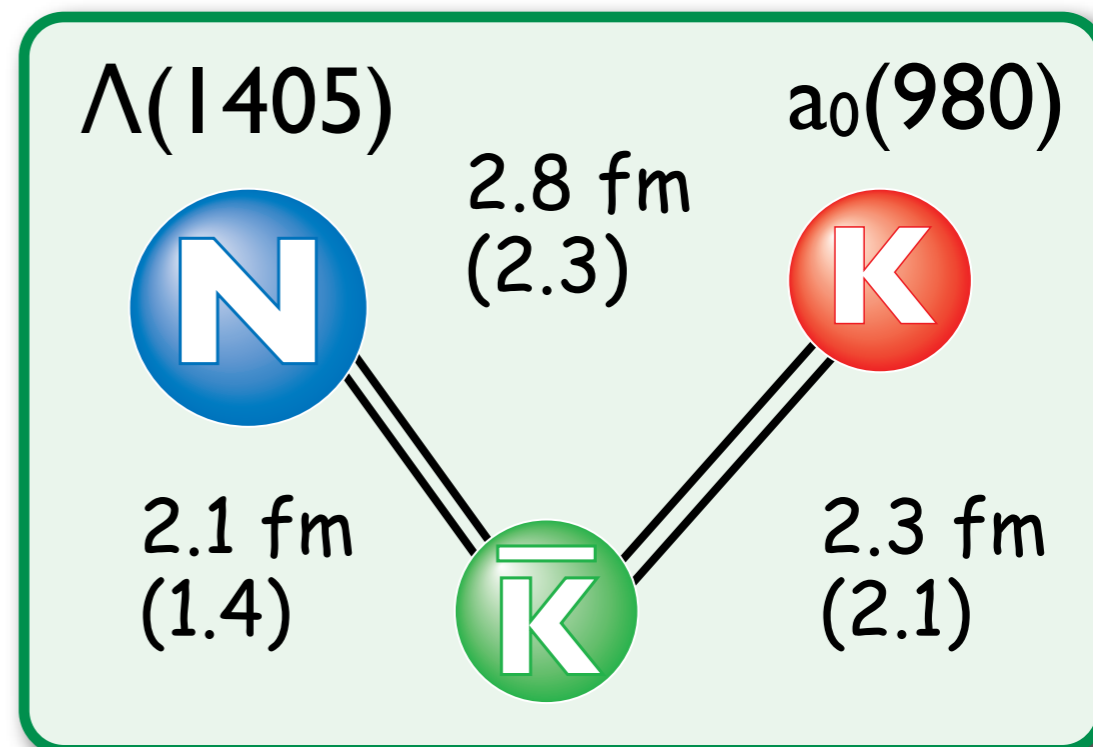
sum of those of isolated two-particle systems

## - coexistence of two quasi-bound states keeping their characters

$\Lambda(1405)+K$      $a_0(980)+N$

## - main decay modes

$\pi\Sigma K$	from $\Lambda(1405)$
$\pi\eta N$	from $a_0(980)$



two-body systems	HW	AY
$\bar{K}-N$ B.E. (MeV)	11	31
width (MeV)	44	20
$K-\bar{K}$ B.E. (MeV)	11	11
width (MeV)	60	60
$\bar{K}-N$ distance (fm)	1.9	1.4
$K-\bar{K}$ distance (fm)	2.1	2.2

DJ, Y. Kanada-En'yo, PRC78, 035203 (2008)

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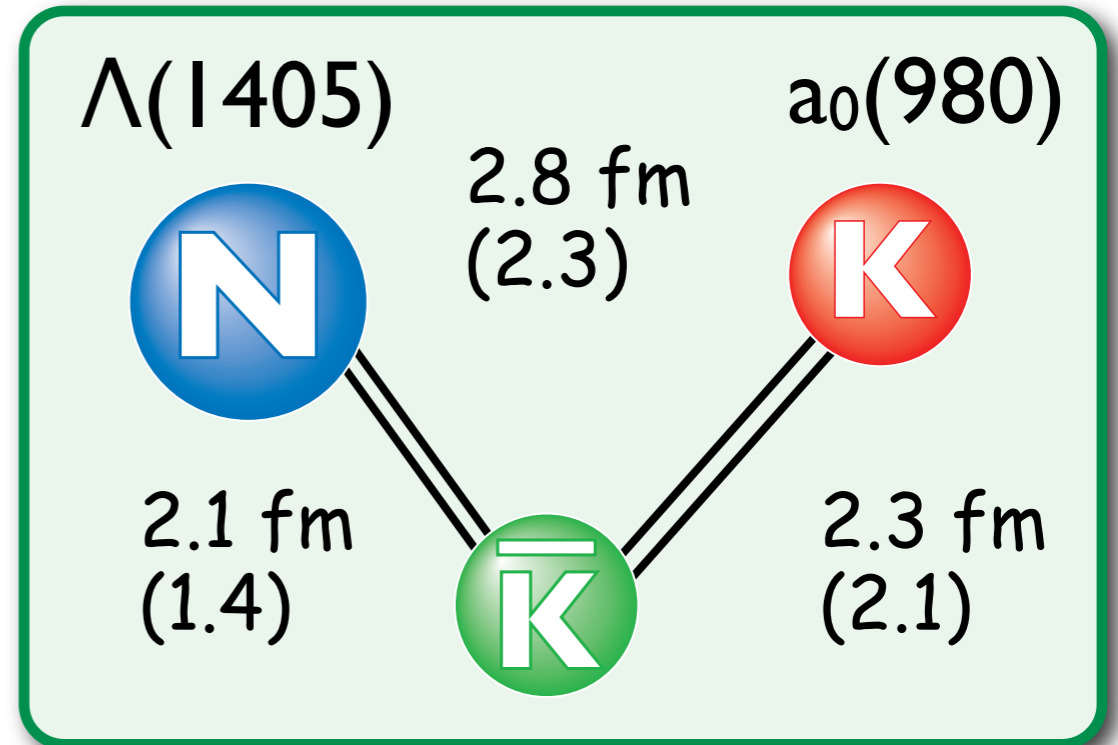
sum of those of isolated two-particle systems

## spacial structure

hadron-hadron distances are comparable with nucleon-nucleon distances in nuclei

r.m.s radius (1.7 fm) is larger than that of  $^4\text{He}$  (1.4 fm)

mean hadron density: 0.07 hadrons/(fm<sup>3</sup>)



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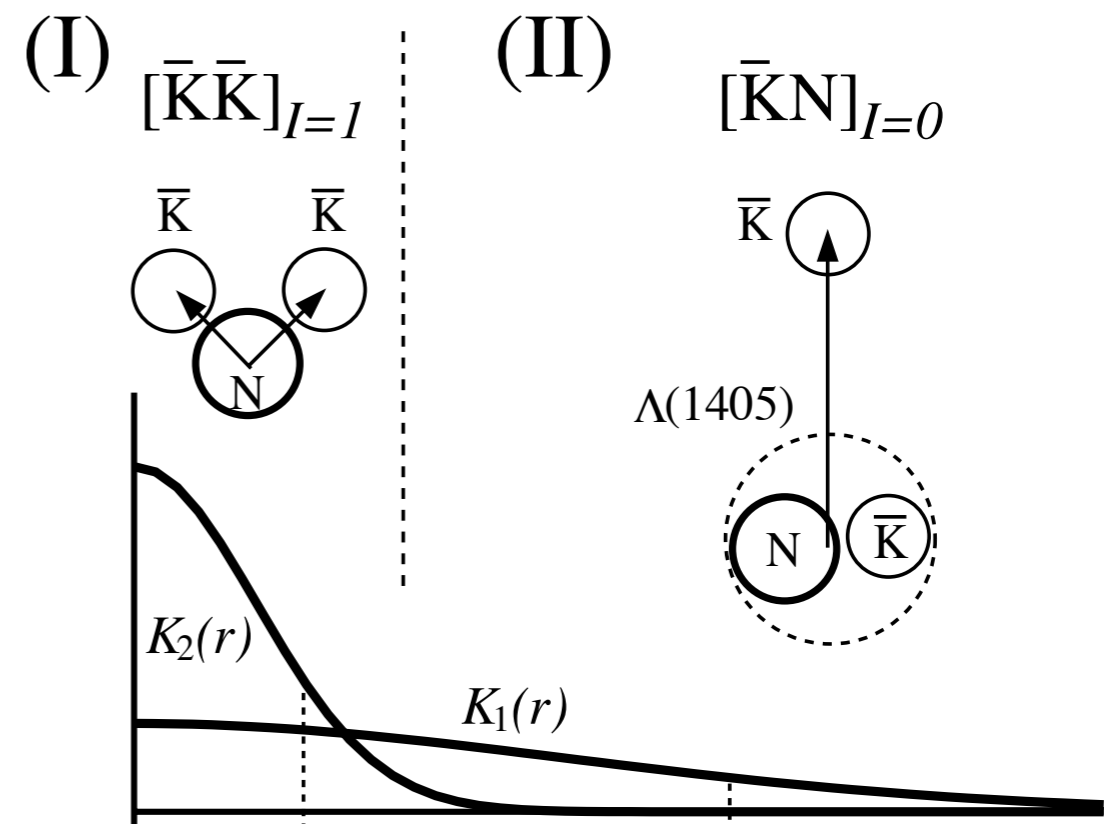
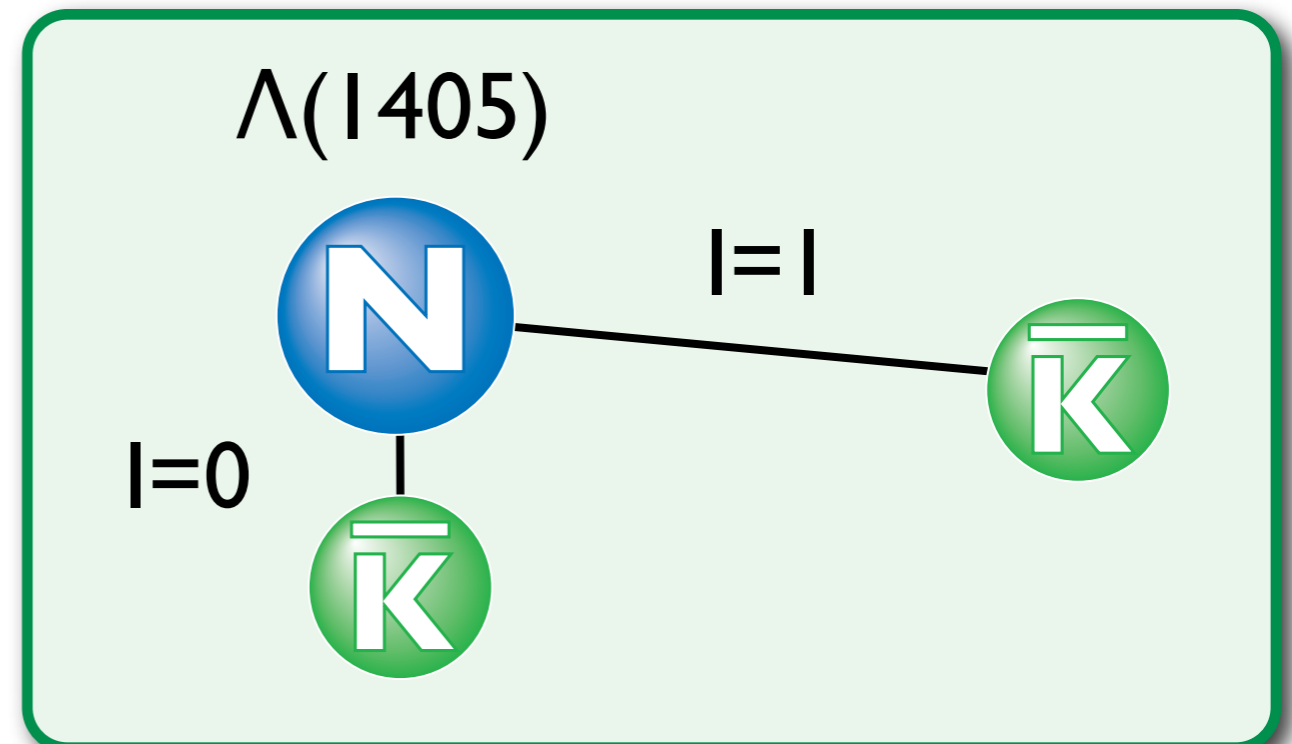
DJ, Y. Kanada-En'yo, PRC78, 035203 (2008)

# $\bar{K}\bar{K}N$ system with $I=1/2, J^P=1/2^+$ ( $\Xi^*$ )

Once  $\Lambda(1405)$  forms in a  $K^{\text{bar}}N$  system with  $I=0$ , another  $K^{\text{bar}}$  and  $N$  has dominantly  $I=1$  component, which is weak attraction. This is not enough to overcome the repulsive  $K^{\text{bar}}K^{\text{bar}}$  interaction.

**very weak binding**

binding energy  $\sim 2$  MeV



Y. Kanada-En'yo, DJ, PRC78, 025212 (2008)

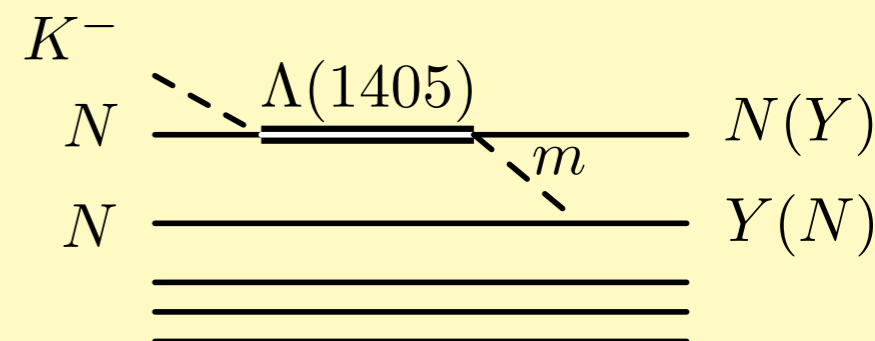
# Decay properties of Kaonic nuclei

Sekihara, DJ, Y. Kanada-En'yo,  
in progress

活発な K 中間子原子核構造・生成研究

## picture

$K^{\text{bar}}$  と核子で  $\Lambda(1405)$  が作られ、  
核内で  $\Lambda(1405)$  が崩壊

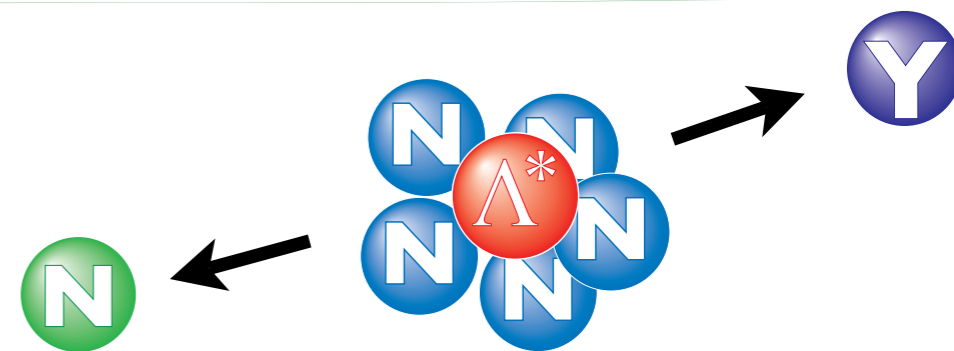


$\Lambda(1405)$  が K 中間子吸収の doorway state

## K 中間子原子核を $\Lambda^*$ ハイパー核として見る

この描像が正しいかどうかは、核中で  $\Lambda^*$  がどれだけ個性を保っているかによる

- ハドロン分子状態の束縛エネルギー：数十 MeV  
より深い束縛なら、別のピクチャーが必要
- 理論的には、ハイパー核のテクニクが使える



崩壊モードから性質を調べることができないか？

特に、nonmesonic (nonpionic) decay (最も基本的な多体効果)



# Summary

keyword : **hadronic molecular state**

**$\Lambda(1405)$  = Quasi-bound state of  $K^{\text{bar}}N$**

## - $\Lambda(1405)$ in few-body systems

$\Lambda(1405)$  keeps its identity in few-body systems

nuclear physics

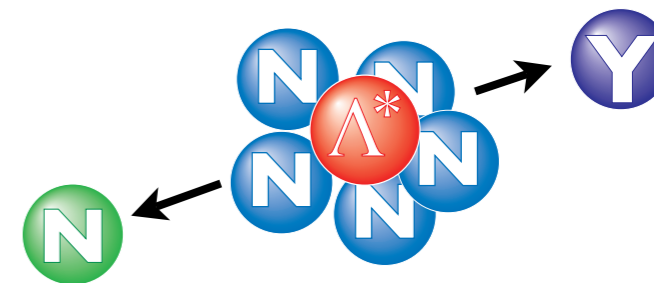
potential picture, coupled channel approach



## - Decay properties of Kaonic nuclei

**Kaonic nuclei =  $\Lambda^*$  hypernuclei**

decay pattern of non-mesonic decay



# BACKUP SLIDES

# $\Lambda(1405)$ as Quasi-bound state of $\bar{K}N$

$\Lambda(1405)$   $J^\pi = 1/2^-, I = 0, S = -1, Q = 0$

- most established resonance, seen in many exp. clearly
- mass :  $1406.5 \pm 4.0$  MeV, (below  $K^-p$  threshold)
- width :  $50 \pm 2$  MeV (PDG)
- decay mode  $\Lambda(1405) \rightarrow (\Sigma\pi)_{I=0}$  100 % **s-wave**

$\bar{K}N$	<u>1435 MeV</u>	$\Lambda(1405)$
$\pi\Sigma$	<u>1331 MeV</u>	

## chiral unitary model

model parameters tuned so as to

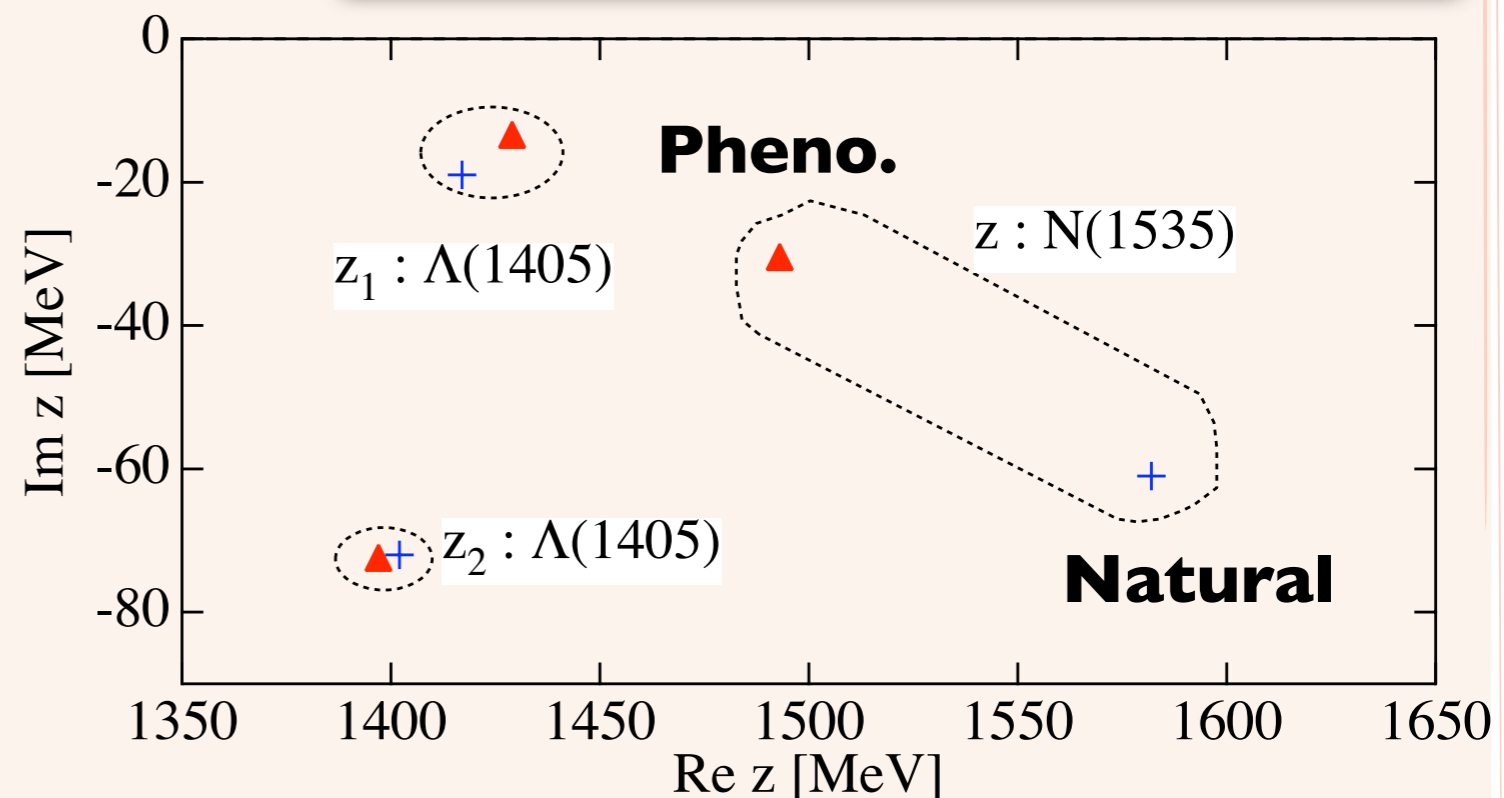
- exclude quark-originated states theoretically

**+ Natural**

- reproduce scattering data

**▲ Pheno.**

Hyodo, Jido, Hosaka, PRC78, 025203 ('08)



**$\Lambda(1405)$  has mostly meson-baryon components.**

# Result $KK^{\text{bar}}N$

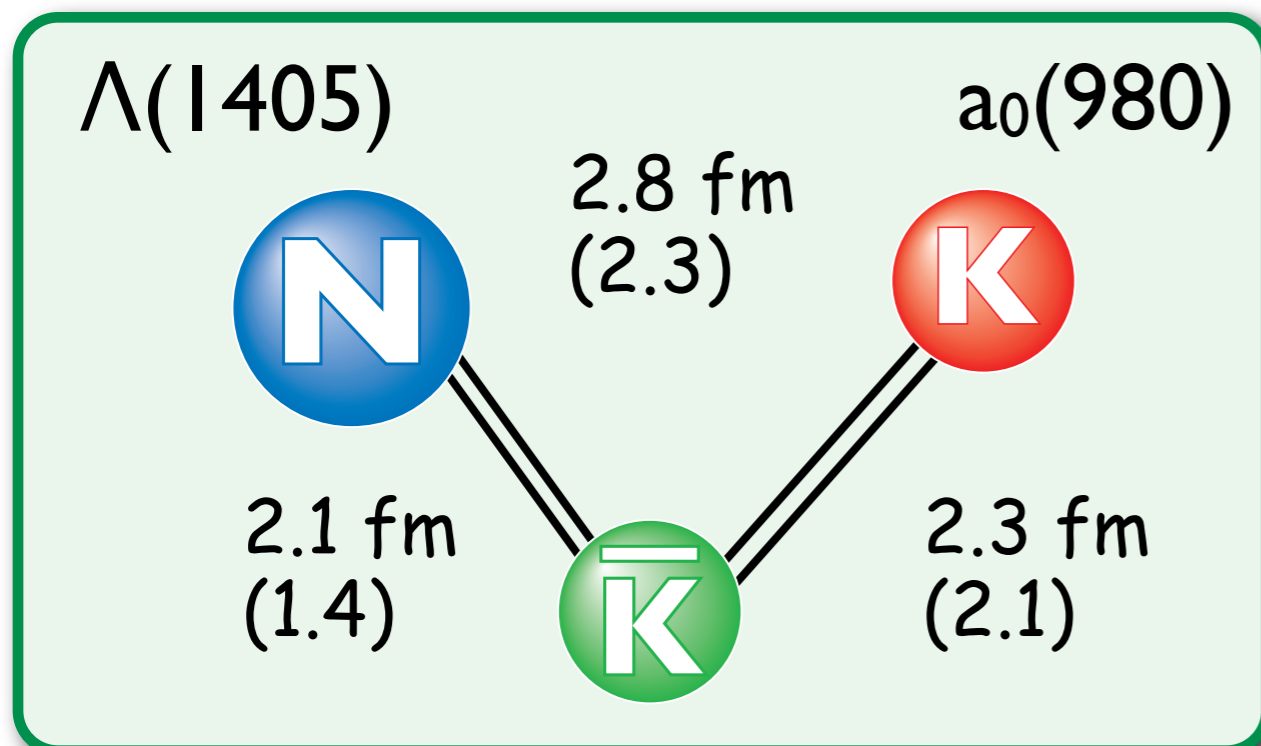
## - main decay modes

$\pi\Sigma K$  from  $\Lambda(1405)$

$\pi\eta N$  from  $a_0(980)$

## - KN repulsion important

Without the KN repulsion, the quasi-bound state would be more bound and have much larger width.



Parameter set	(A)	(A)	(B)	(B)
$V_{\bar{K}N}$	HW-HNJH	HW-HNJH	AY	AY
$V_{KN}$	On	Off	On	Off
ReE	-19	-39	-41	-57
$\langle T \rangle$	169	282	175	227
$\langle \text{Re}V \rangle$	-188	-320	-216	-284
ImE	-44	-72	-49	-63
$\langle \text{Im}V_{KN}^{I=0} \rangle$	-17	-30	-19	-23
$\langle \text{Im}V_{KN}^{I=1} \rangle$	-1	0	0	0
$\langle \text{Im}V_{K\bar{K}}^{I=0} \rangle$	-1	-10	-4	-10
$\langle \text{Im}V_{K\bar{K}}^{I=1} \rangle$	-25	-31	-25	-31
$\langle T_{KN} \rangle$	113	185	131	157
$\langle \text{Re}V_{KN}^{I=0} \rangle$	-87	-152	-139	-162
$\langle \text{Re}V_{KN}^{I=1} \rangle$	-2	0	0	0
Re $\mathcal{E}_{KN}$	25	33	-9	-4
Re $\mathcal{E}_{\bar{K}N}^{I=0}$	-6	-4	-28	-27
$\langle T_{K\bar{K}} \rangle$	104	162	86	115
$\langle \text{Re}V_{K\bar{K}}^{I=0} \rangle$	-4	-42	-11	-31
$\langle \text{Re}V_{K\bar{K}}^{I=1} \rangle$	-101	-127	-75	-92
Re $\mathcal{E}_{K\bar{K}}$	-1	-7	-1	-7
$\langle T_{KN} \rangle$	59	108	55	83
$\langle \text{Re}V_{KN}^{I=0} \rangle$	0	0	0	0
$\langle \text{Re}V_{KN}^{I=1} \rangle$	6	0	10	0
Re $\mathcal{E}_{KN}$	65	108	65	83

# Subthreshold properties of $K^{\text{bar}}N$

$K^- d \rightarrow \Lambda(1405)n$  in chiral unitary model

in collaboration with

**E. Oset** (Valencia) and **T. Sekihara** (Kyoto)

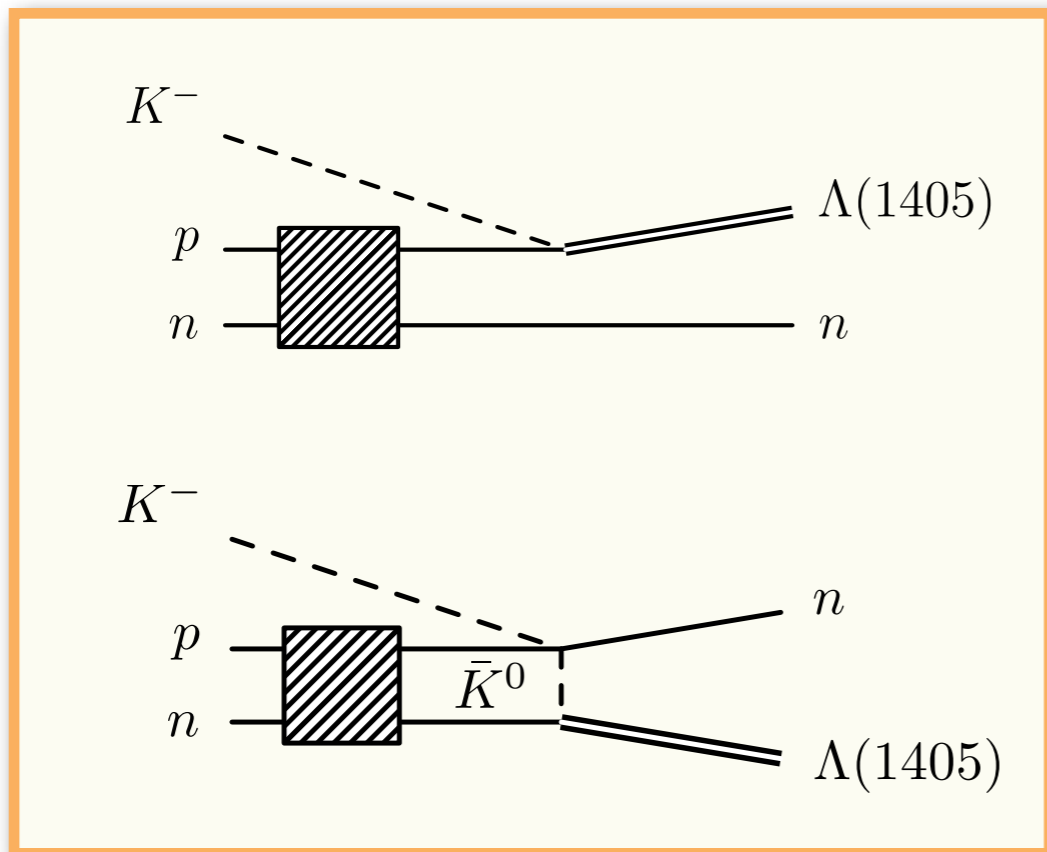
# Subthreshold properties of $K^{\text{bar}}N$

to confirm the  $\Lambda(1405)$  resonance position in  $K^{\text{bar}}N$  channel

$\Lambda(1405)$  spectrum in  $K^{\text{bar}}N$  channel

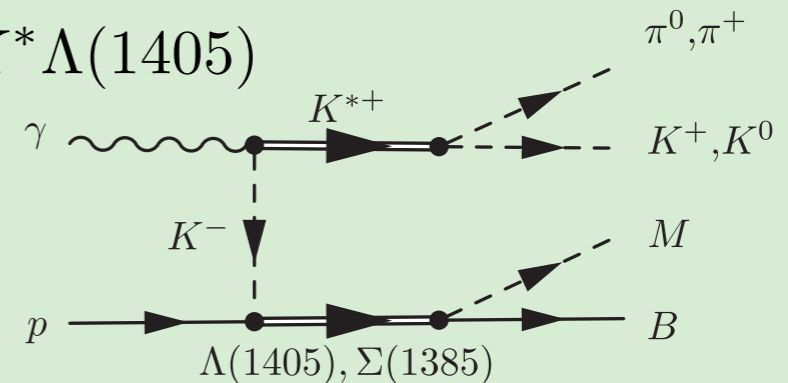
$$\bar{K}N \rightarrow \Lambda(1405)$$

$$K^- d \rightarrow \Lambda(1405)n$$



$K^*$  photoproduction

$$\gamma p \rightarrow K^* \Lambda(1405)$$



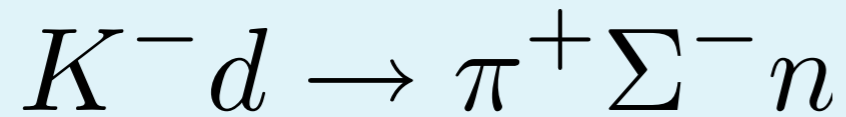
Hyodo, Hosaka, Vacas, Oset, PLB593, 75 (04)

# Subthreshold properties of $K^{\text{bar}}N$

DJ, Oset, Sekihara, in preparation

bubble chamber

## Experiment



$\pi\Sigma$  invariant mass spectrum

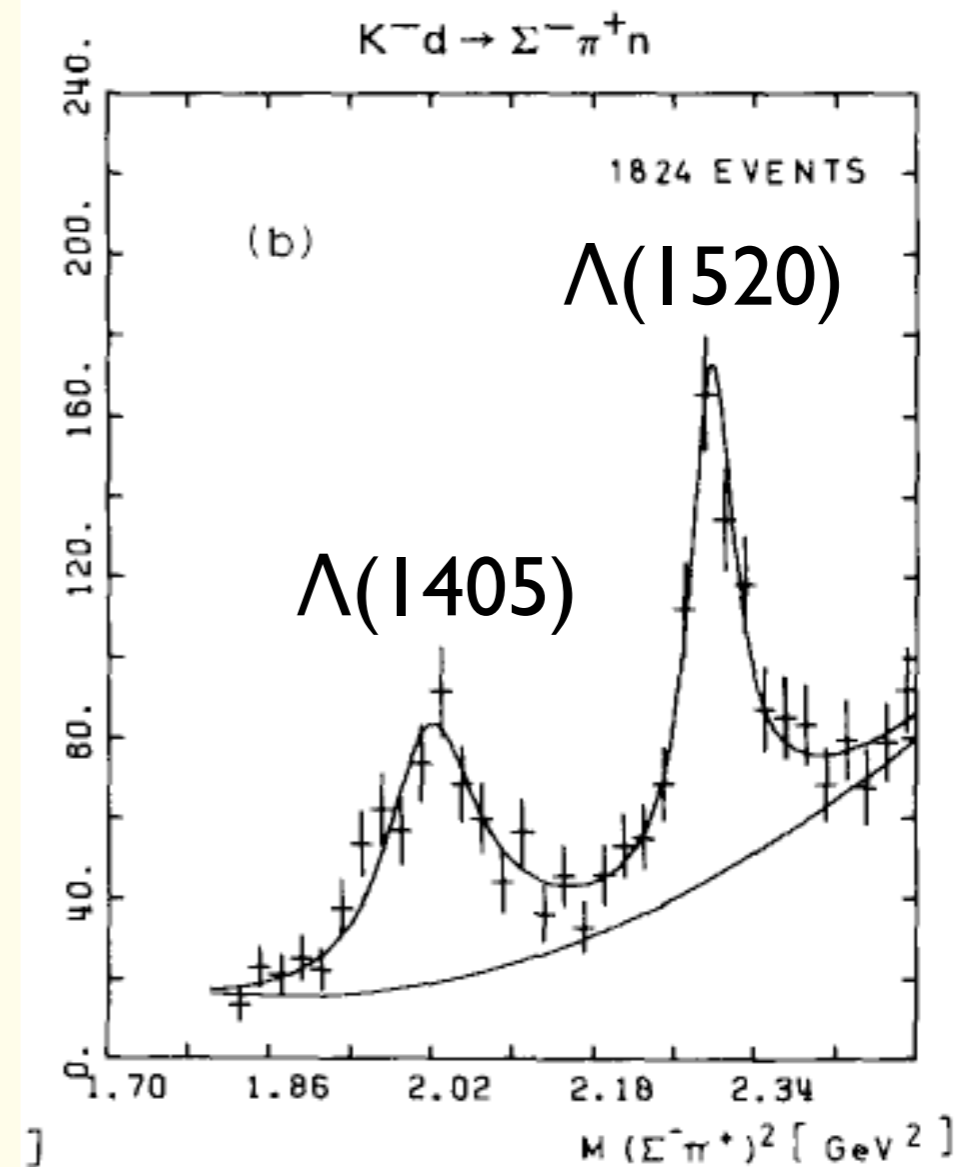
initial K momentum

686 ~ 844 MeV/c



cross section

$410 \pm 100 \mu\text{b}$



**peak position 1420 MeV**

Braun et al. NPB129,1,(77)

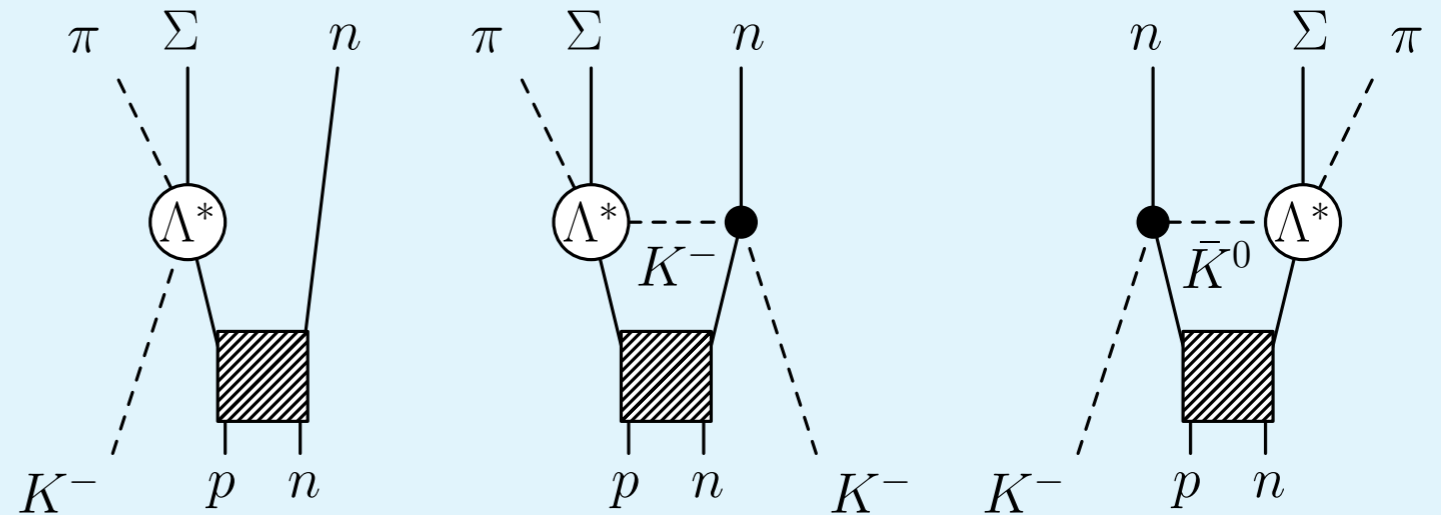
# Subthreshold properties of $K^{\text{bar}}N$

DJ, Oset, Sekihara, in preparation

## chiral unitary model

$\Lambda^*$ : T-matrix by ChUM

K momentum: 800 MeV/c



$$K^- d \rightarrow \pi^+ \Sigma^- n$$

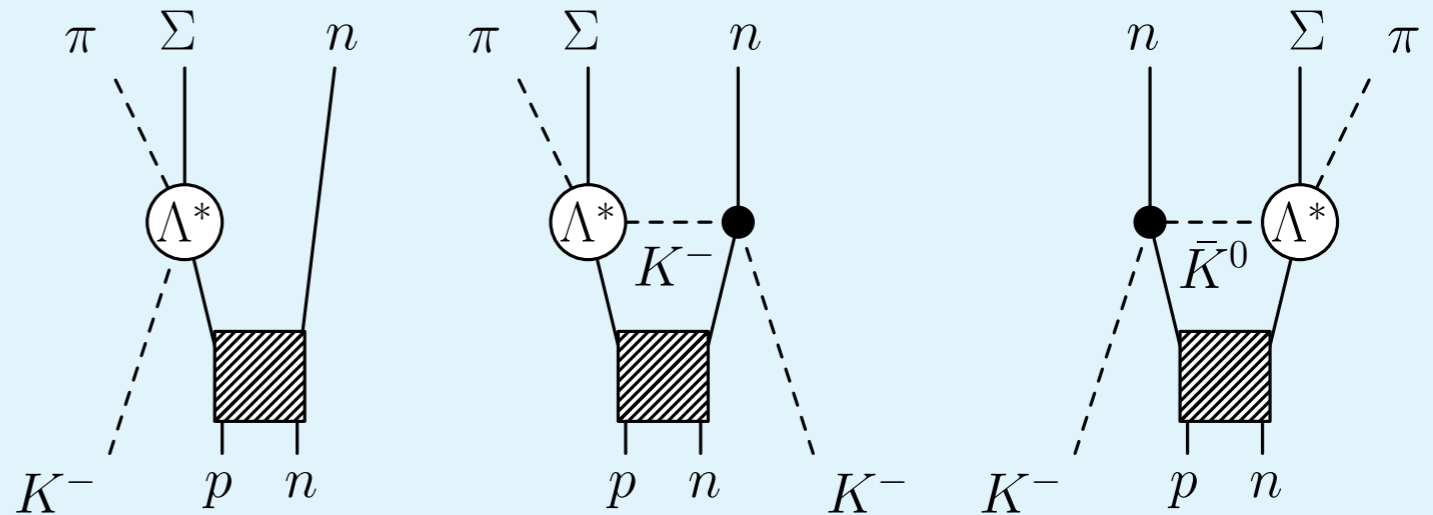


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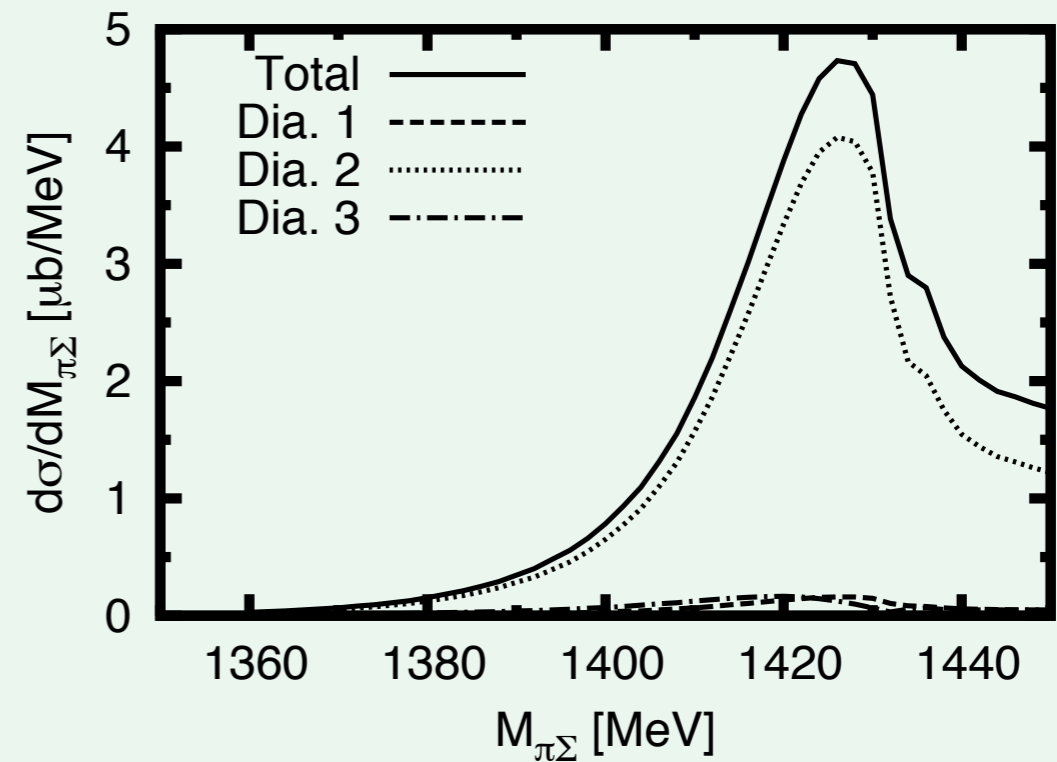
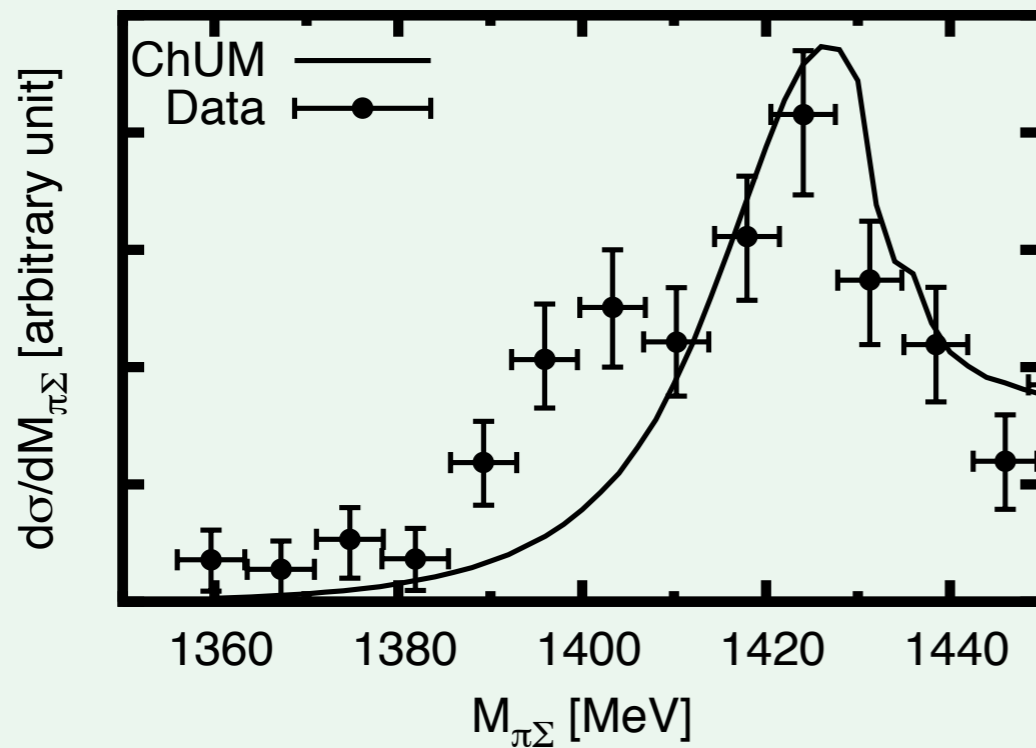
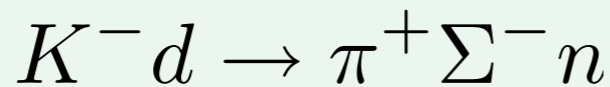
## chiral unitary model

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K momentum: 800 MeV/c



## result



$K^- d \rightarrow \Lambda(1405)n$  cross section

$\sim 380 \mu\text{b}$  (exp.  $410 \pm 100 \mu\text{b}$ )

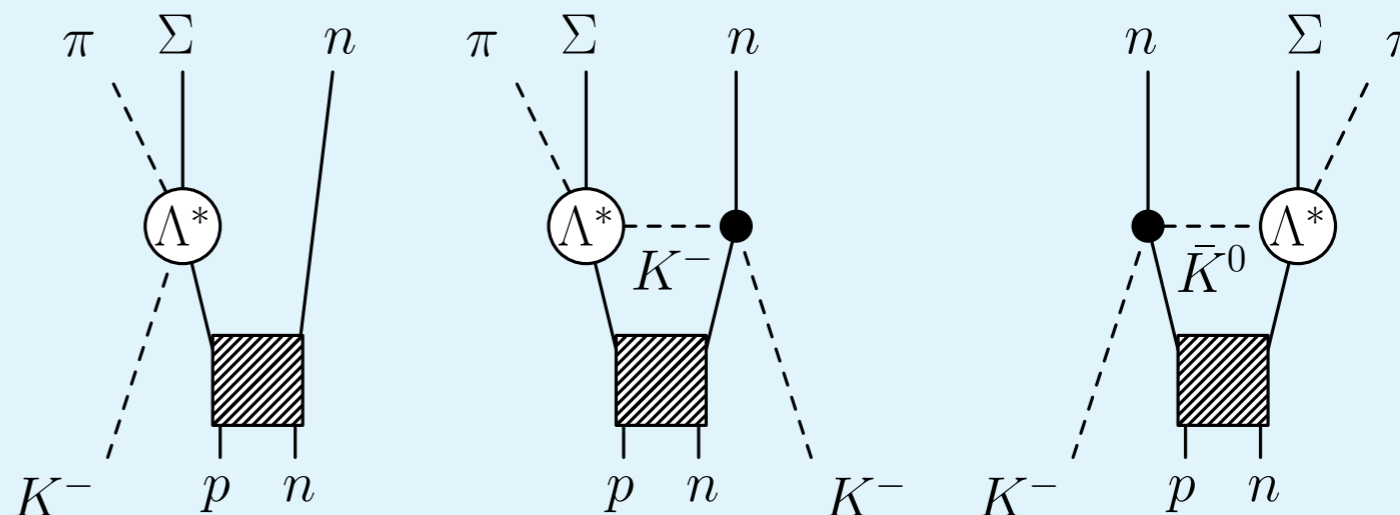
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DJ, Oset, Sekihara, in preparation

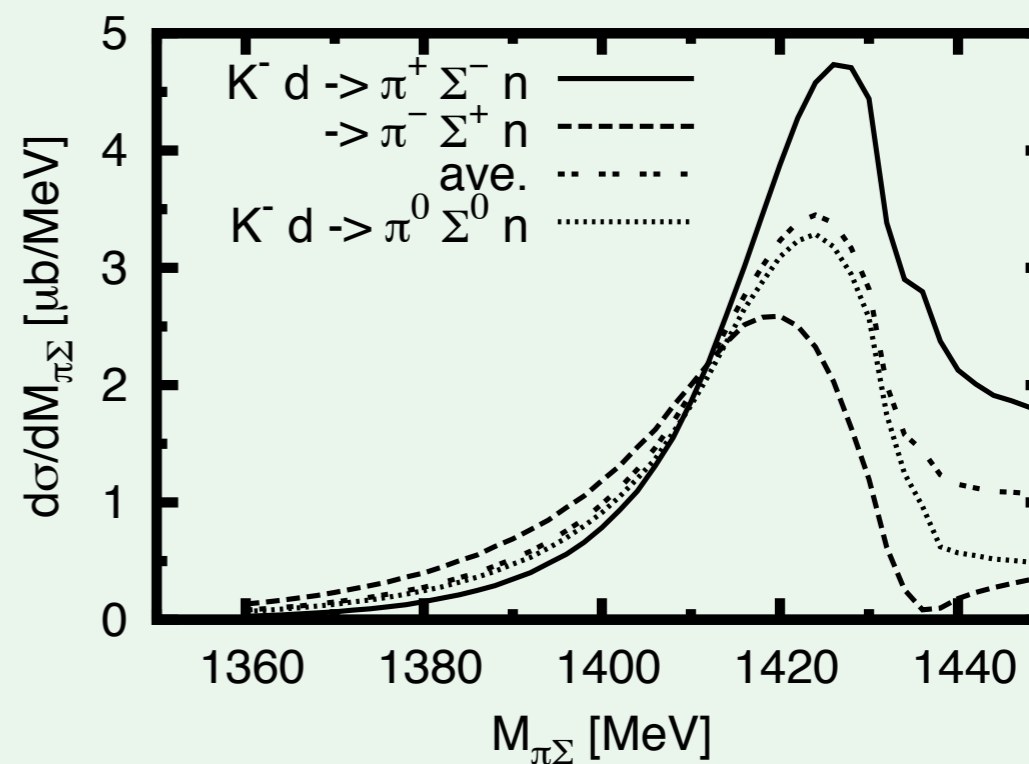
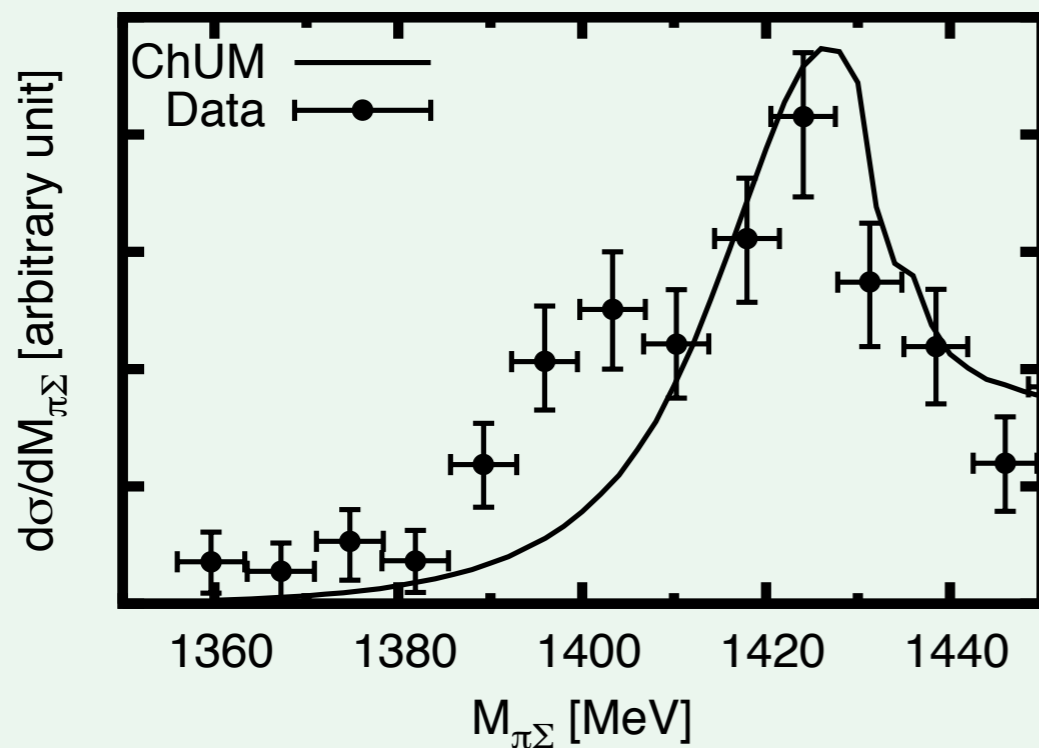
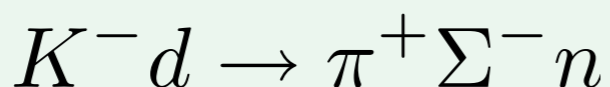
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