Lattice QCD Study for the Penta-Quark Potential

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We perform the first study of the static penta-quark potential [1] in SU(3) lattice QCD with β =6.0 and $16^3 \times 32$ at the quenched level. We investigate the QQ-Q-Q-Q configuration as shown in Fig.1. With the smearing method [2-4] to enhance the ground-state component, we accurately measure the penta-quark potential V_{5Q} from the penta-quark Wilson loop $\langle W_{5Q} \rangle$ as shown in Fig.2 in a gauge-invariant manner.

As the lattice QCD result, we find that the static penta-quark potential V_{5Q} is well described by the sum of one-gluon-exchange (OGE) Coulomb term and multi-Y type linear term, i.e., the OGE plus multi-Y Ansatz [1],

$$V_{5Q} = -A_{5Q} \left\{ \left(\frac{1}{r_{12}} + \frac{1}{r_{34}} \right) + \frac{1}{2} \left(\frac{1}{r_{15}} + \frac{1}{r_{25}} + \frac{1}{r_{35}} + \frac{1}{r_{45}} \right) + \frac{1}{4} \left(\frac{1}{r_{13}} + \frac{1}{r_{14}} + \frac{1}{r_{23}} + \frac{1}{r_{24}} \right) \right\} + \sigma_{5Q} L_{\min} + C_{5Q}$$

$$(1)$$

with $r_{ij} \equiv |\mathbf{r}_i - \mathbf{r}_j|$ and L_{\min} being the minimal length of the flux-tube linking five quarks. Note that there appear three kinds of Coulomb coefficients $(A_{5\mathrm{Q}}, \frac{1}{2}A_{5\mathrm{Q}}, \frac{1}{4}A_{5\mathrm{Q}})$ in the pentaquark system, while only one Coulomb coefficient, $A_{\mathrm{Q}\bar{\mathrm{Q}}}$ or $A_{3\mathrm{Q}}$, appears in the $\mathrm{Q}\bar{\mathrm{Q}}$ or the 3Q system. (In Eq.(1), $A_{5\mathrm{Q}}$ corresponds to $A_{3\mathrm{Q}}$ or $\frac{1}{2}A_{\mathrm{Q}\bar{\mathrm{Q}}}$ in terms of the OGE result.)

Figure 3 shows typical examples of the lattice QCD result for the penta-quark potential $V_{5\mathrm{Q}}$ [1]. The symbols denote the lattice data, and the curves denote the theoretical form of the OGE plus multi-Y Ansatz with $(A_{5\mathrm{Q}}, \sigma_{5\mathrm{Q}})$ fixed to be $(A_{3\mathrm{Q}}, \sigma_{3\mathrm{Q}})$ in the 3Q potential $V_{3\mathrm{Q}}$ in Ref.[3]. (Note that there is no adjustable parameter for the theoretical curves besides an irrelevant constant $C_{5\mathrm{Q}}$, since $A_{5\mathrm{Q}}$ and $\sigma_{5\mathrm{Q}}$ are fixed.) One finds a good agreement between the lattice data and the OGE plus multi-Y Ansatz.

From the comparison with the $Q\bar{Q}$ and the 3Q potentials [2-4], the universality of the string tension and the OGE result are found among $Q\bar{Q}$, 3Q and 5Q systems as

$$\sigma_{Q\bar{Q}} \simeq \sigma_{3Q} \simeq \sigma_{5Q}, \qquad \frac{1}{2} A_{Q\bar{Q}} \simeq A_{3Q} \simeq A_{5Q}.$$
 (2)

This result supports the flux-tube picture [2-5] on the confinement mechanism even for the multi-quark system [1].

Acknowledgements: The lattice-QCD Monte Carlo calculations have been performed on NEC-SX5 supercomputer at Osaka University.

References

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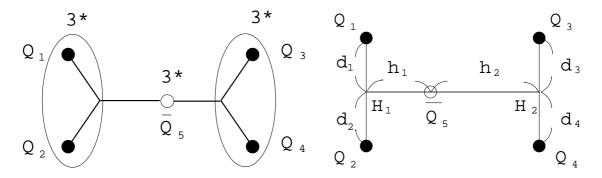


Figure 1: (a) The QQ- \bar{Q} -QQ type configuration for the penta-quark system. The two QQ clusters belong to the $\mathbf{3}^*$ representation of the color SU(3). (b) A planar configuration of the penta-quark system. Here, we take $d_1 = d_2 = d_3 = d_4 \equiv d$ and $h_1 = h_2 \equiv h$.

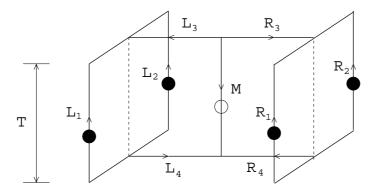


Figure 2: The penta-quark Wilson loop W_{5Q} for the penta-quark potential V_{5Q} .

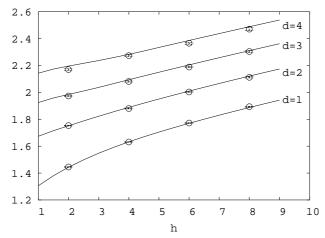


Figure 3: Lattice QCD results for the penta-quark potential V_{5Q} . The symbols denote the lattice data, and the curves denote the theoretical form of the OGE plus multi-Y Ansatz [1]. A good agreement is found between the lattice data and the theoretical curves.