## Measurement of $\phi$ decay asymmetry in $\vec{\gamma}d \rightarrow \phi d$ reaction

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The diffractive photoproduction of  $\phi$ -mesons has been used to study the Pomeron exchange process and to look for a exotic daughter Pomeron trajectory associated with a glueball  $(J^{PC} = 0^{++})$  [1]. This is because the baryon and meson exchange amplitudes in the s- and t-channels are suppressed by the Okubo-Zweig-Iizuka (OZI) rule. However, isovector  $\pi$  exchange still contribute to the  $\phi$  photoproduction which makes extraction of the true exotic channel difficult. One of the possible solution to eliminate the  $\pi$ -exchange contribution is to use an isoscalar target such as deuteron. Here, we present a new measurement of decay asymmetry of  $\phi$  meson with respect to the direction of photon polarization in the  $\vec{\gamma}d \rightarrow \phi d$  reaction. The experiment was performed at the SPring8 facility by using linearly polarized photons in conjunction with the LEPS spectrometer system [2, 3].

Photons with a maximum energy of 2.4 GeV were produced from backward-Compton scattering of 351-nm Ar laser off 8-GeV electrons in the SPring-8 storage ring. The  $\phi$  mesons produced in the liquid-deuterium target were analyzed by detecting  $K^+$  and  $K^-$  from the  $\phi$  decay. Chagred patricle momentum was determined by using a silicon-strip vertex detector and three drift chambers. Kaons were identified by the measured time-of-flight between the TOF counters. The missing-mass  $MM_{\gamma\phi}$  calculated by assuming the  $d(\gamma, \phi)X$  kinematics is shown in Fig. 1(a). Backgrounds from the quasi-free  $N(\gamma, \phi)N$  reaction were studied by a Monte Carlo simulation taken into account Fermi motion and off-shell effect in deuteron system. Fig. 1(b) shows the background-subtracted  $MM_{\gamma\phi}$  distribution where peak position and peak width are consistent with deuteron rest mass and the expected value, respectively. The  $\phi$  events were extracted by requiring  $MM_{\gamma\phi} < 1.9 \text{ GeV}/c^2$ .



Figure 1: (a) missing mass spectra for the data (solid) and the simulation(dashed) (b) background subtracted missing mass. Arrows in (a) and (b) correspond to deuteron rest mass. (c) is  $\phi$  decay asymmetry as a function of the photon energy for  $d(\gamma, \phi)d$  (open) and  $p(\gamma, \phi)p$  (closed).

The  $\phi$  decay asymmetry,  $A_{\phi}$ , is defined as an amplitude of the angular distribution of  $\phi$  decay as [3, 1],  $W(\Phi) = 1 + P_{\gamma}A_{\phi}\cos 2\Phi$ , where  $P_{\gamma}$  and  $\Phi$  are magnitude of the photon polarization and  $K^+$  emission angle with respect to the photon polarization direction, respectively. The extracted  $\phi$  events were separated into 8 regions by  $\Phi$  with every  $\pi/4$  step as,  $[N(A)]_{pol} = \Omega(A) \cdot Y_{pol} \int_{\pi/4 \cdot (A-1)}^{\pi/4 \cdot A} W(\Phi) d\Phi$  for A = 1-8, where  $\Omega$  and Y denote the detector acceptance and the number of photons, respectively. pol corresponds to the photon polarization direction (HZ or VT). Double ratio quantities between the HZ and VT events were introduced as,  $R_1 = [\frac{N(3)N(7)}{N(1)N(5)}]_{\rm VT}/[\frac{N(3)N(7)}{N(1)N(5)}]_{\rm HZ}$  and  $R_2 = [\frac{N(2)N(6)}{N(4)N(8)}]_{\rm VT}/[\frac{N(2)N(6)}{N(4)N(8)}]_{\rm HZ}$ . Then, error-weighted average of  $R_1$ and  $R_2$ , R, was calculated. Here  $\Omega$  and Y are automatically canceled out, and  $A_{\phi}$  was thus derived to be  $A_{\phi} = \frac{1}{P_{\gamma}} \cdot \frac{\pi}{2} \cdot \frac{1-R^{1/4}}{1+R^{1/4}}$ . The  $\phi$  decay asymmetry of the  $d(\gamma, \phi)d$  reaction was obtained as a function of the photon energy, as shown in Fig.1(c). It was found that the decay asymmetry is positive, while the  $A_{\phi}$  value of the deuteron reaction is about twice larger than that of the proton reaction.

## References

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