## Lattice Study of Low-lying Nonet Scalar Mesons in Quenched Approximation

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Recent experimental and theoretical evidence for the existence of the  $\sigma$  and the  $\kappa$  [1], indicates that light scalar mesons constitute the nonet state. They cannot be ordinary  $q\bar{q}$  mesons as described in the nonrelativistic constituent quark model since in such a quark model, the  $J^{PC}=0^{++}$  meson is realized in the  ${}^{3}P_{0}$  state, which implies that the mass of the  $\sigma$  and the  $\kappa$  must be as high as  $1.2 \sim 1.6$  GeV. Thus, the low-lying scalar mesons below 1 GeV have been a source of various ideas of exotic structures, as mentioned above: they may be fourquark states such as  $qq\bar{q}\bar{q}$ , or  $\pi\pi$  or K $\pi$  molecules.

We present a simulation with weaker couplings on a larger lattice than any other previous simulations although in the quenched level. We perform quenched level simulations on the  $\kappa$  meson so as to clarify the structure of the mysterious scalar meson rather than to reproduce the experimental value of the mass; a precise quenched-level simulation should give a rather clear perspective on whether the system can fit with the simple constituent-quark model picture or not.

We perform a quenched QCD calculation using the Wilson fermions, with the plaquette gauge action, on a relatively large lattice  $(20^3 \times 24)$ . The values of the hopping parameter for the u/d quark are  $h_{u/d} =$ 0.1589, 0.1583 and 0.1574, while  $h_s = 0.1566$  and 0.1557 for the *s* quark. Using these hopping parameters except for  $h_s = 0.1557$ , CP-PACS collaboration performed a quenched QCD calculation of the light meson spectrum with a larger lattice  $(32^3 \times 56)$  [2], which we refer to for comparison. The gauge configurations are generated by the heat bath algorithm at  $\beta = 5.9$ . After 20000 thermalization iterations, we start to calculate the meson propagators. On every 2000 configurations, 80 configurations are used for the ensemble average.

We emply the point-like source and sink for the  $\kappa^+$  meson

$$\hat{\kappa}(x) \equiv \sum_{c=1}^{3} \sum_{\alpha=1}^{4} \bar{s}^{c}_{\alpha}(x) u^{c}_{\alpha}(x) \quad , \tag{1}$$

where u(x) and s(x) are the Dirac operators for the u/d and s quarks, and the indices c and  $\alpha$  denote the color and Dirac-spinor indices, respectively. The point source and sink in Eq.(1) lead a positive spectral function  $\rho(m^2)$  in the correlation function  $\langle \hat{\kappa}(t)\hat{\kappa}(0) \rangle = \int dm\rho(m^2)\exp(-mt)$ . The result obtained here is thus an upper bound of  $\kappa$  mass, because our result should include excited states.

Our estimated value of the mass of the  $\kappa$  is ~ 1.7 GeV, which is larger than twice the experimental mass ~ 800 MeV. This result was expected on the basis of our experience in calculating the  $\sigma$  meson. The relatively heavy mass of the  $\kappa$  may be at least partly attributed to the absence of the disconnected diagram in the  $\kappa$  propagator; the  $\kappa$  propagator is composed of only a connected diagram. While the disconnected diagram was essential for realizing the low-mass  $\sigma$ , it does not exist for the  $\kappa$ ; therefore, the mass of the  $\kappa$  is not made lighter by the disconnected diagram. Indeed, the mass of the valence  $\sigma_v$  described solely with the connected propagator is far larger than the experimental value ~ 500-600 MeV. Our lattice study and the quark model analysis suggest that the simple two-body constituent-quark picture of the  $\kappa$  meson does not agree well with the experimentally observed  $\kappa$ . Note that the quench simulation is a clean theoretical experiment in which a virtual intermediate like  $qq\bar{q}\bar{q}$  is highly suppressed. Therefore, if its existence with the reported low mass is experimentally established, the dynamical quarks may play an essential role for making the  $\kappa$  mass so lighter or the  $\kappa$  may contain an unconventional state such as a  $qq\bar{q}\bar{q}$  or  $K\pi$  molecular state, which are missing in the calculation here.

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## References

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- [2] CP-PACS Collaboration, S. Aoki, et al., Phys. Rev. D67 (2003) 034503.