## Competing effects of nuclear deformation and density dependence of the $\Lambda N$ interaction in $B_{\Lambda}$ values of hypernuclei

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One of the basic quantities of  $\Lambda$  hypernuclei is the  $\Lambda$  binding energy defined as the energy gain by adding a  $\Lambda$  particle from the core nucleus. Recently, the dependence of  $B_{\Lambda}$  on structures of core nuclei, in particular, nuclear deformations, has been discussed in several hypernuclei theoretically [1, 2, 3]. Values of  $B_{\Lambda}$  are related to nuclear structure in two ways: one is that an increase of deformation reduces the overlap of the densities between the  $\Lambda$  particle and nucleons, which makes  $B_{\Lambda}$  smaller. The other is due to the density dependence of the  $\Lambda N$  effective interaction. In light hypernuclei and/or cluster states, the density overlap between the  $\Lambda$  and nucleons is significantly decreased, which can affect  $B_{\Lambda}$  through the density dependence. When the  $\Lambda N$  effective interaction derived from the *G*-matrix calculation is designed to depend on the nuclear Fermi momentum  $k_F$ , the smaller overlap makes the relevant value of  $k_F$  small, i.e., less Pauli-blocking, resulting in the increase of  $B_{\Lambda}$ . Considering this effect, it is expected that appropriate values of  $k_F$  in finite systems are reduced as overlaps become small with mass numbers, which would affect the mass dependence of  $B_{\Lambda}$ .

The aim of the present work is to reveal how the density dependence of the  $\Lambda N$  effective interaction affects the mass dependence of  $B_{\Lambda}$ . Since *p-sd-pf* shell hypernuclei have various structures, they would affect the values of  $B_{\Lambda}$  through the density dependence of the  $\Lambda N$  interaction. To investigate it, we apply the antisymmetrized molecular dynamics for hypernuclei (HyperAMD) to the  $\Lambda$  hypernuclei with  $9 \leq A \leq 59$ , combined with the generator coordinate method using the HyperAMD wave functions with various nuclear quadrupole deformation  $(\beta, \gamma)$  as the basis. As  $\Lambda N$  effective interactions, we use the  $\Lambda N$  *G*-matrix interactions derived from the Nijmegen potentials, which has the  $k_F$  dependence. In this study, we calculate  $k_F$  from the densities of the  $\Lambda$  particle and nucleons for each hypernucleus based on the averaged density approximation.

The  $B_{\Lambda}$  values are shown in Fig. 1 (solid line) together with observed data. One can see that the present calculation nicely reproduces the observations in wide mass region, which is achieved by taking into account core structure, especially deformation of core nuclei. To see the importance of describing the core structure, we calculate the  $B_{\Lambda}$  values using the spherical basis wave functions only (see dashed line). It shows that the  $B_{\Lambda}$  values are deviated significantly from the observed data, especially in light  $\Lambda$  hypernuclei in Fig. 1(b), if the core nuclei are assumed to be spherical. We also confirmed, for example in  $\frac{1}{\Lambda}B$ , that the present calculation reproduces the intra-band B(E2) of the core nuclei well, implying that core deformation is described properly. The contents of this report have been published [4].



Figure 1: (a)  $B_{\Lambda}$  values calculated with the HyperAMD wave functions with  $(\beta, \gamma)$  deformation (solid) and spherical shape  $(\beta = 0)$  only (dashed). Open circles show observed data with mass numbers from A = 9 up to A = 51. (b) Same as (a), but magnified in the  $5 \le A \le 20$  region.

## References

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