

Photoproduction of K^
for the study of the
structure of $\Sigma(1405)$*



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Motivations : Two poles?

There are two poles of the scattering amplitude around nominal $\Lambda(1405)$ energy region.

- Cloudy bag model
(1990)

J. Fink *et al.* PRC41, 2720

- Chiral unitary model
(2001~)

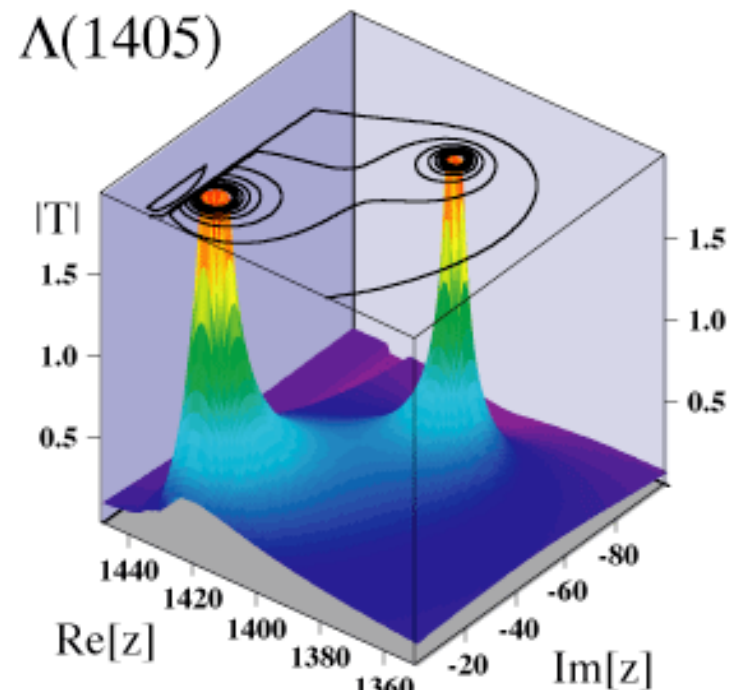
J. A. Oller *et al.* PLB500, 263

E. Oset *et al.* PLB527, 99

D. Jido *et al.* PRC66, 025203

T. Hyodo *et al.* PRC68, 018201

$\Lambda(1405) : J^P=1/2^-, I=0$



ChU model, T. Hyodo

Chiral unitary model

Flavor SU(3) meson-baryon scatterings (s-wave)

Chiral symmetry

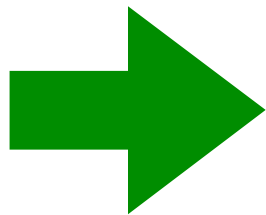
**Low energy
behavior**



Unitarity of S-matrix

**Non-perturbative
resummation**

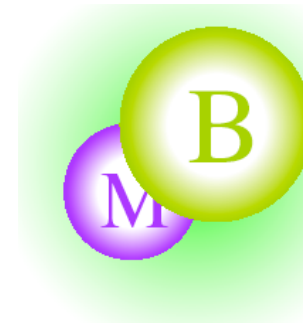
**Dynamical
generation**



$$J^P = 1/2^-$$

Resonances

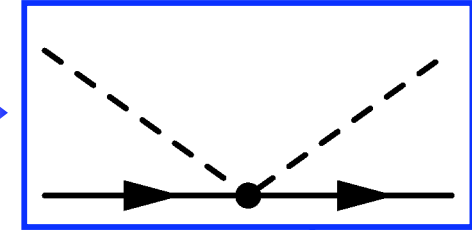
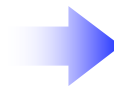
$\Sigma(1405)$, $\Sigma(1670)$, $N(1535)$,
 $\Sigma(1620)$, $\Sigma(1620)$



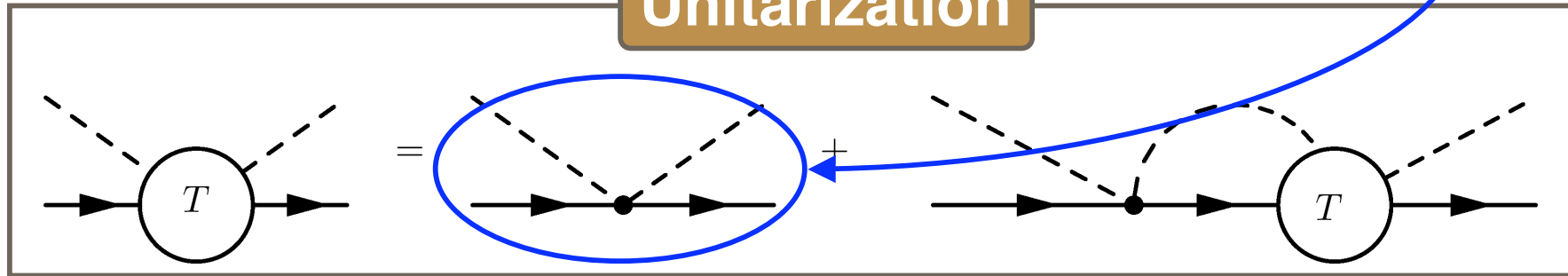
Framework of the chiral unitary model

Chiral perturbation theory

$$\mathcal{L}_{WT} = \frac{1}{4f^2} \text{Tr}(\bar{B}i\gamma^\mu[(\Phi\partial_\mu\Phi - \partial_\mu\Phi\Phi), B])$$

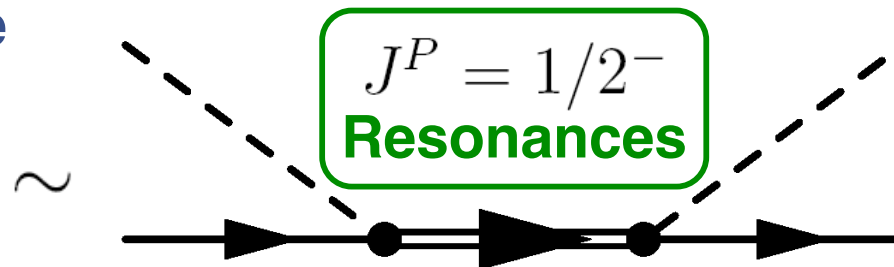


Unitarization

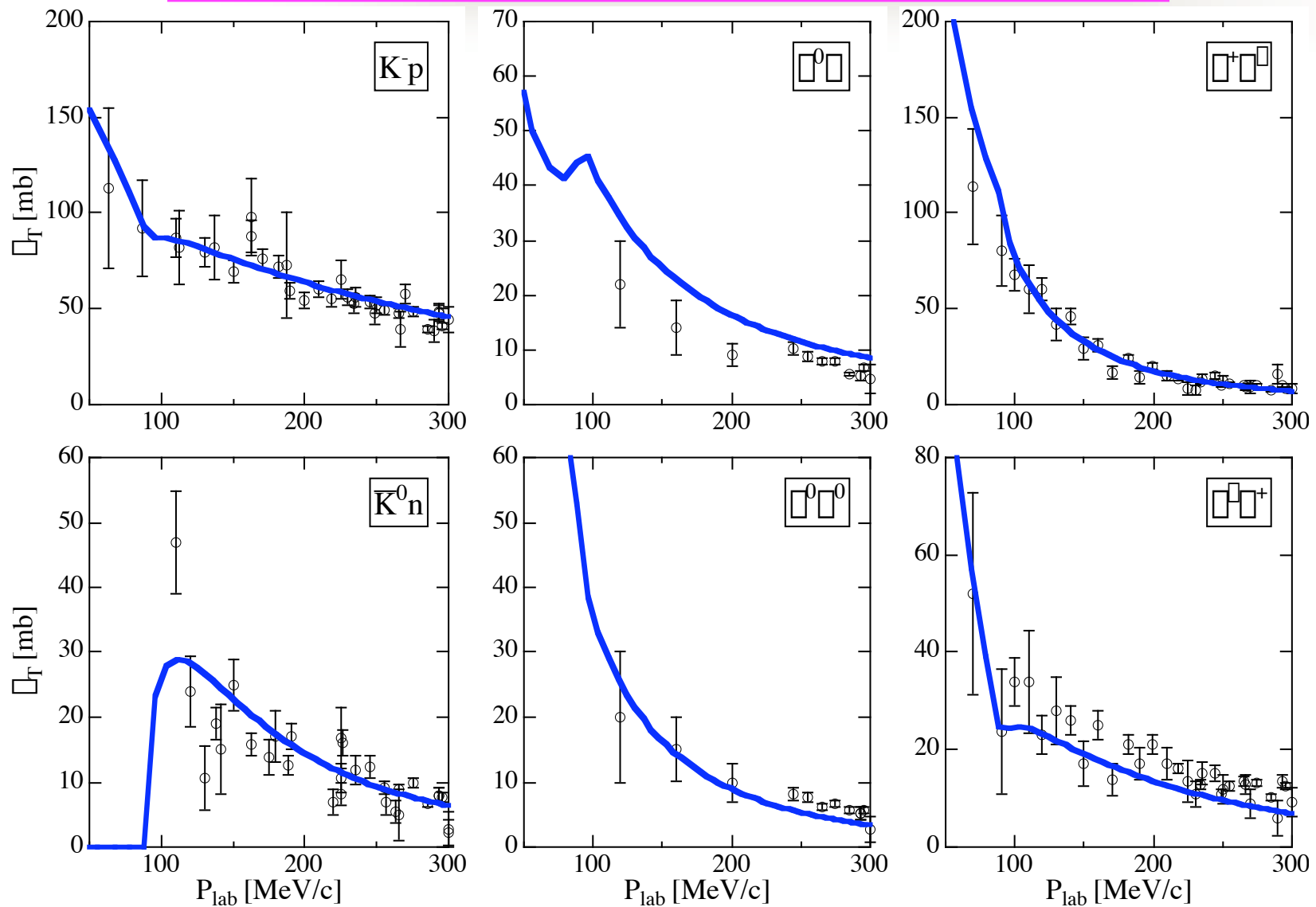


$$T_{ij}(\sqrt{s}) \sim \frac{g_i g_j}{\sqrt{s} - M_R + i\Gamma_R/2} + T_{ij}^{BG}$$

Generated resonances are expressed as poles of the scattering amplitude.



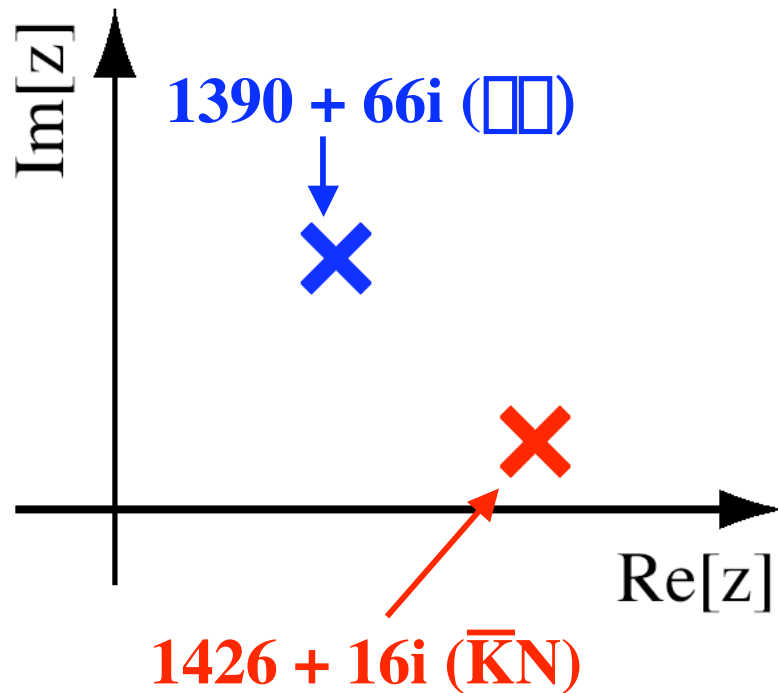
Total cross sections of $K^{\pm}p$ scatterings



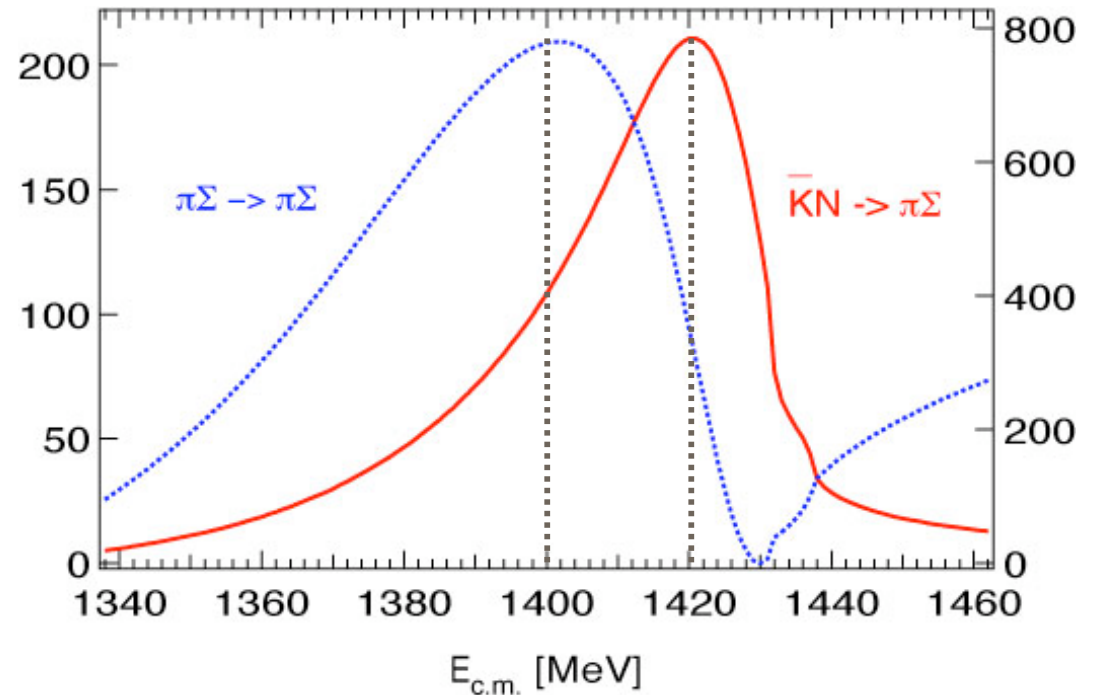
T. Hyodo, et al., Phys. Rev. C 68, 018201 (2003)

$\Xi(1405)$ in the chiral unitary model

position of poles



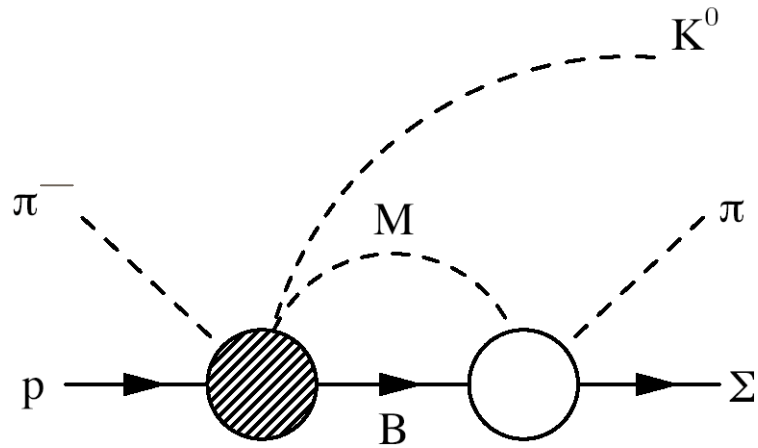
$\Xi\Xi$ mass distribution



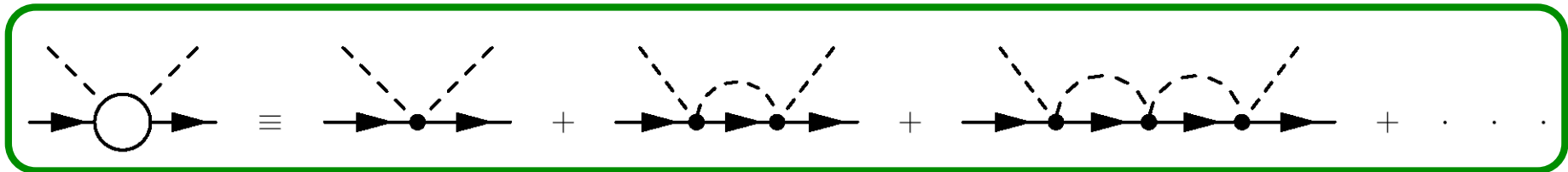
$$\frac{d\sigma}{dM_I} = C |t_{\pi\Sigma \rightarrow \pi\Sigma}|^2 p_{CM} \quad \longrightarrow \quad \frac{d\sigma}{dM_I} = \left| \sum_i C_i t_{i \rightarrow \pi\Sigma} \right|^2 p_{CM}$$

D. Jido, et al., Nucl. Phys. A 723, 205 (2003)

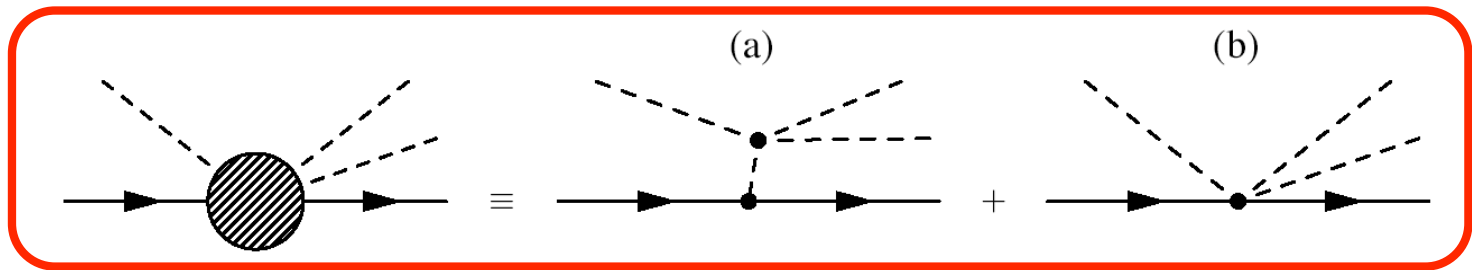
Example : the $\pi^- p \rightarrow K^0 \pi$ reaction



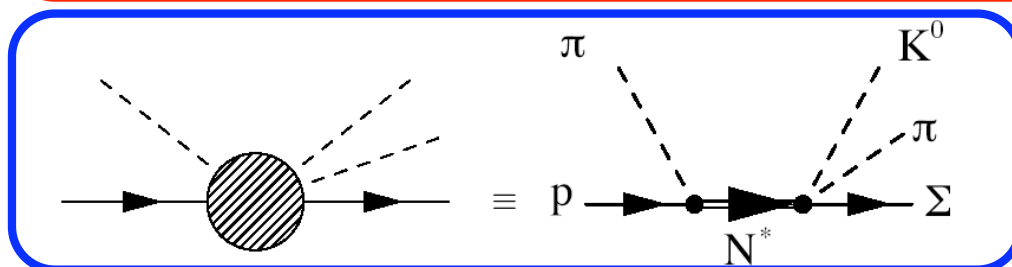
Chiral unitary model



Chiral term

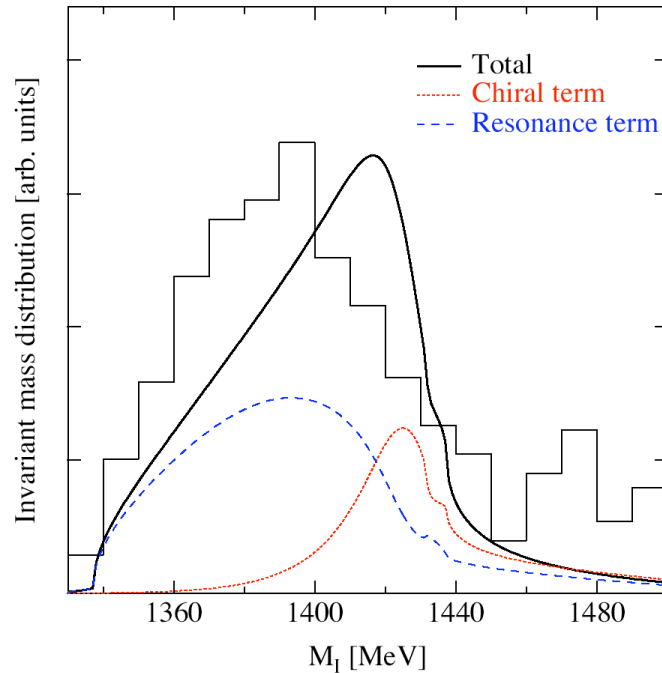


N(1710)



Results for $\pi^+p \rightarrow K^0\pi\pi$

Mass distribution



Total cross sections [mb]

final state	$K^0 K^- p$	$K^0 \bar{K}^0 n$	$K^0 \pi^0 \Lambda$	$K^0 \pi^+ \Sigma^-$	$K^0 \pi^- \Sigma^+$
Exp.	2.9	8.3	104.0	25.1	20.2
total	3.75	5.98	6.02	21.32	20.01
chiral	2.36	2.84	3.14	3.04	6.78
resonance	0.70	0.67	10.85	16.18	5.43

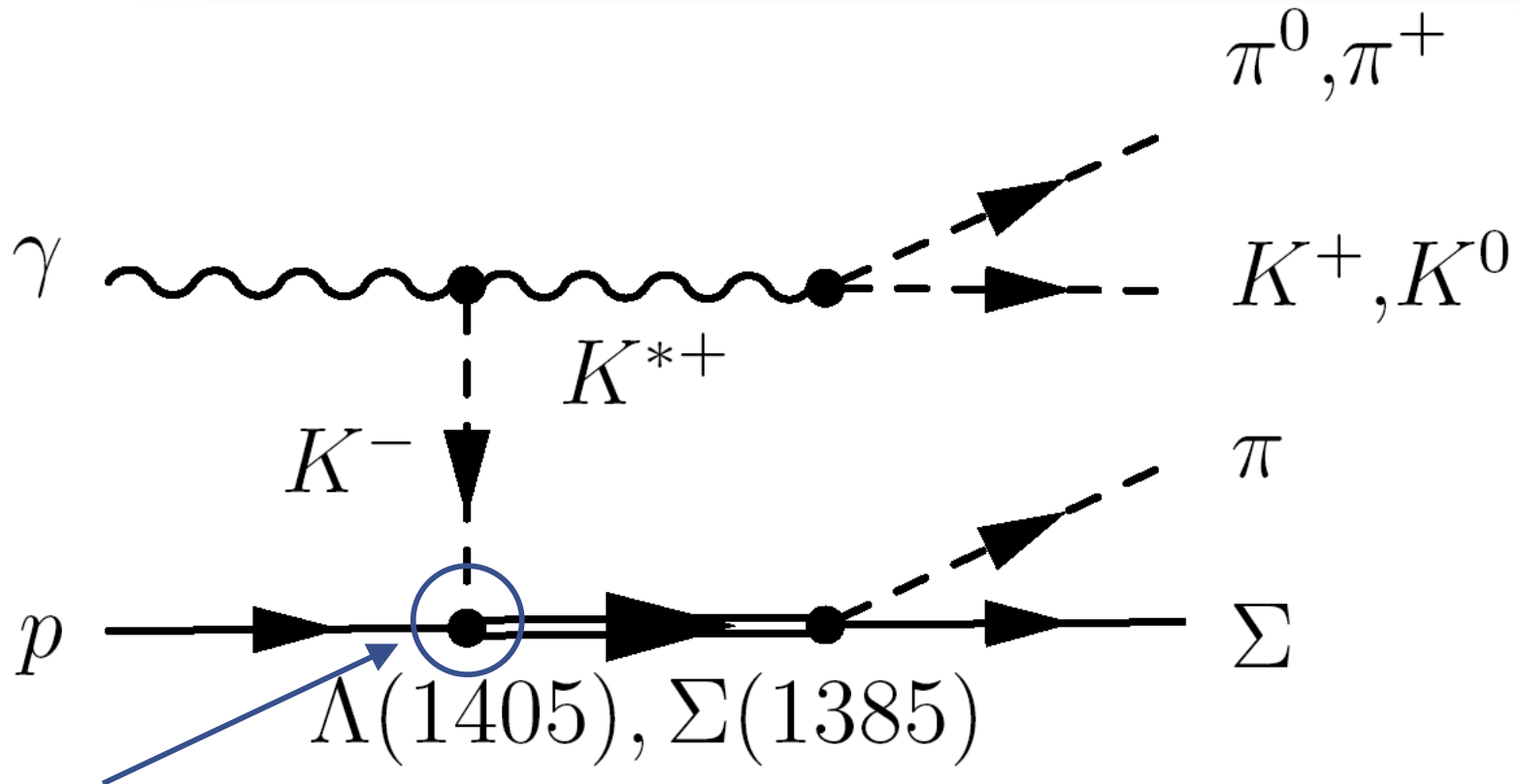
$\Lambda(1385)$ effect

Good agreement

There are two mechanisms in the initial stage interaction, which filter each one of the resonances.

T. Hyodo, et al., nucl-th/0307005, Phys. Rev. C, in press

Photoproduction of $K^* \square (1405)$



Only K^-p channel appears at the initial stage

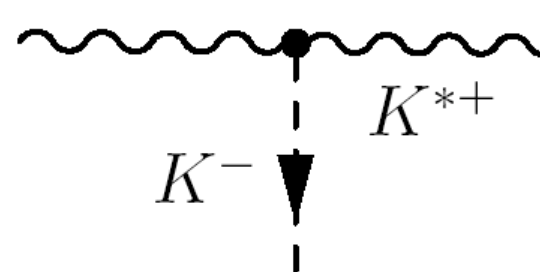
Higher pole ??

Effective interactions for meson part

1. \square VP coupling

$$-it = ig_{\gamma K^* K} \epsilon^{\mu\nu\alpha\beta} P_\mu \epsilon_\nu(K^{*+}) k_\alpha \epsilon_\beta(\gamma), \quad \gamma$$

$$|g_{\gamma K^{*\pm} K^\pm}| = 0.252 \text{ [GeV}^{-1}\text{]},$$

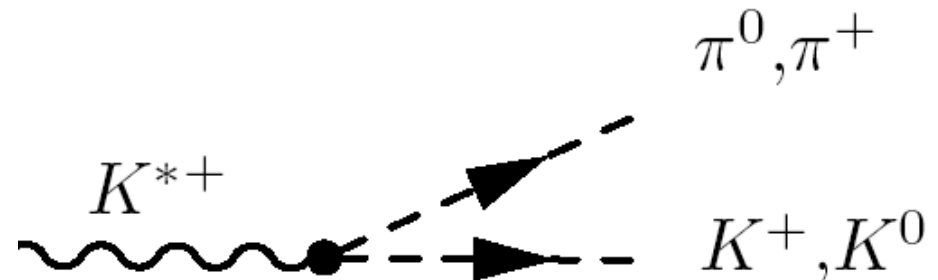
$$|g_{\gamma K^{*0} K^0}| = 0.385 \text{ [GeV}^{-1}\text{]}.$$


2. VPP coupling

$$-it(K^{*+} \rightarrow K^+ \pi^0) = i \frac{g_{VPP}}{\sqrt{2}} \frac{1}{\sqrt{2}} [p_\mu(K^+) - p_\mu(\pi^0)] \epsilon^\mu(K^{*+}),$$

$$-it(K^{*+} \rightarrow K^0 \pi^+) = i \frac{g_{VPP}}{\sqrt{2}} [p_\mu(K^0) - p_\mu(\pi^+)] \epsilon^\mu(K^{*+}),$$

$$g_{VPP} = -6.05$$



Effective interaction for $\Sigma(1385)$

3. $\Sigma(1385)$ MB coupling

$$-it_{\Sigma^*i} = c_i \frac{12D + F}{5} \frac{1}{2f} \mathbf{S} \cdot \mathbf{k}_i$$

SU(6) symmetry



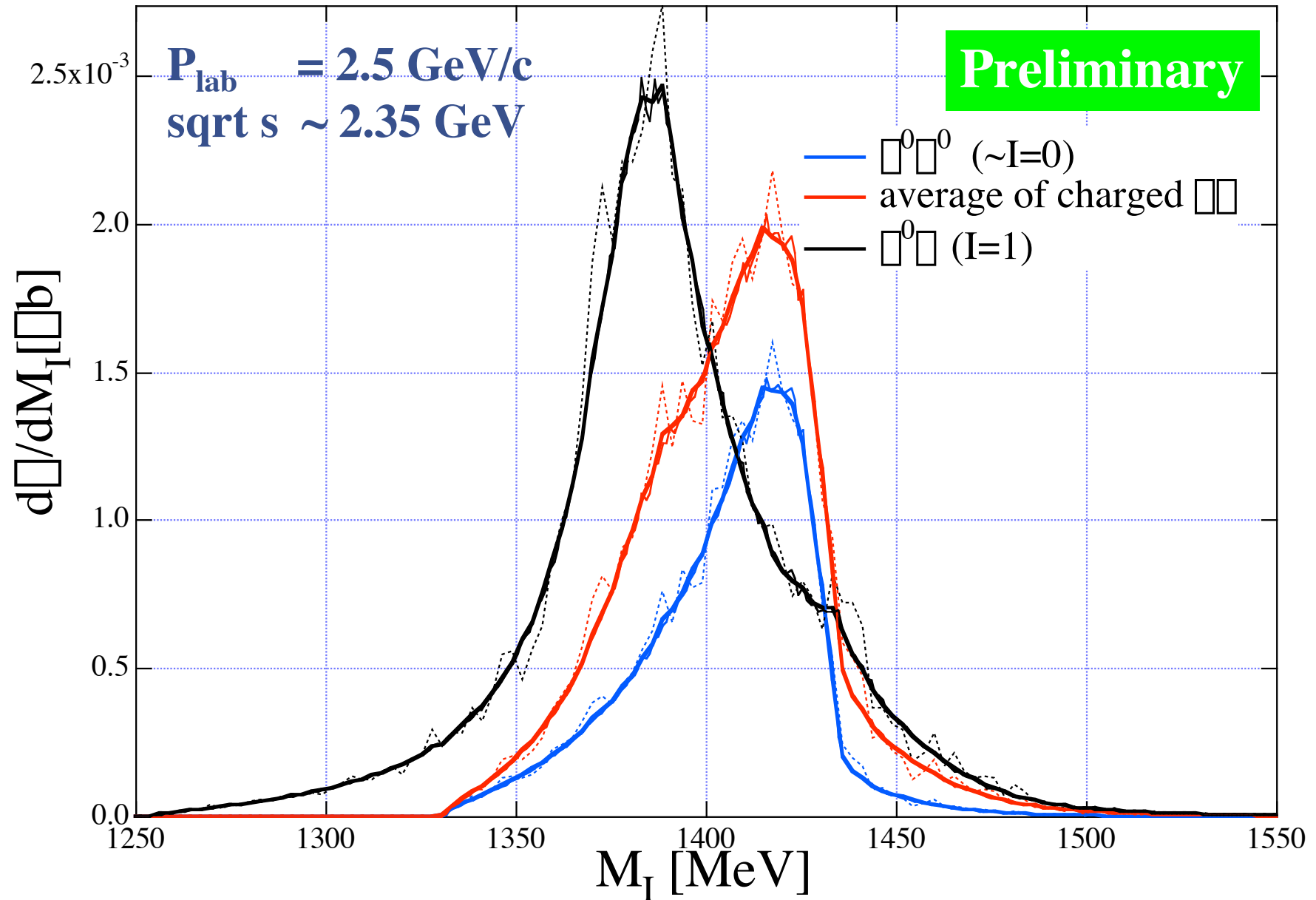
channel i	K^-p	\bar{K}^0n	$\pi^0\Lambda$	$\pi^0\Sigma^0$	$\eta\Lambda$	$\eta\Sigma^0$	$\pi^+\Sigma^-$	$\pi^-\Sigma^+$	$K^+\Xi^-$	$K^0\Xi^0$
c_i	$-\sqrt{\frac{1}{12}}$	$\sqrt{\frac{1}{12}}$	$\sqrt{\frac{1}{4}}$	0	0	$-\sqrt{\frac{1}{4}}$	$-\sqrt{\frac{1}{12}}$	$\sqrt{\frac{1}{12}}$	$\sqrt{\frac{1}{12}}$	$-\sqrt{\frac{1}{12}}$

4. $K^*P \rightarrow \Sigma(1385) \rightarrow$ MB amplitude

$$\begin{aligned}
 -it_{1i} &= c_1 c_i \left(\frac{12D + F}{5} \frac{1}{2f} \right)^2 \mathbf{S} \cdot \mathbf{k}_1 \mathbf{S}^\dagger \cdot \mathbf{k}_i \frac{i}{M_I^{(b)} - M_{\Sigma^*} + i\Gamma_{\Sigma^*}/2} F_f(k_1) \\
 &= c_1 c_i \left(\frac{12D + F}{5} \frac{1}{2f} \right)^2 (k_1)_l (k_i)_m \left(\frac{2}{3} \delta_{lm} - \frac{i}{3} \epsilon_{lmn} \sigma_n \right) \frac{i}{M_I^{(b)} - M_{\Sigma^*} + i\Gamma_{\Sigma^*}/2} F_f(k_1)
 \end{aligned}$$

$$F_f(k_1) = \frac{\Lambda^2 - m_K^2}{\Lambda^2 - (k_1)^2}$$

Experiments : $\omega\omega$ mass distribution



Summary and conclusions

We study the **structure of $\Lambda(1405)$** using the chiral unitary model.

🍏 There are **two poles** of the scattering amplitude around nominal $\Lambda(1405)$.

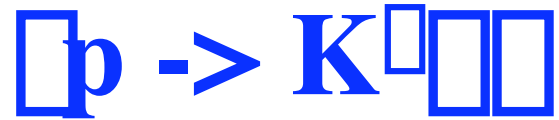
Pole 1 (1426+16i) : strongly couples to KN state

Pole 2 (1390+66i) : strongly couples to $\Lambda\bar{\Lambda}$ state

🍏 By observing the $\Lambda\bar{\Lambda}$ **mass distribution** in the $\bar{p} \rightarrow K^* \Lambda(1405)$ reaction, it could be possible to isolate **higher energy pole**.

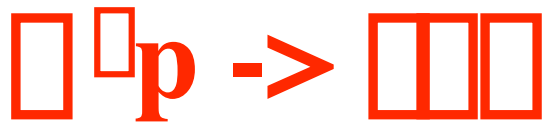
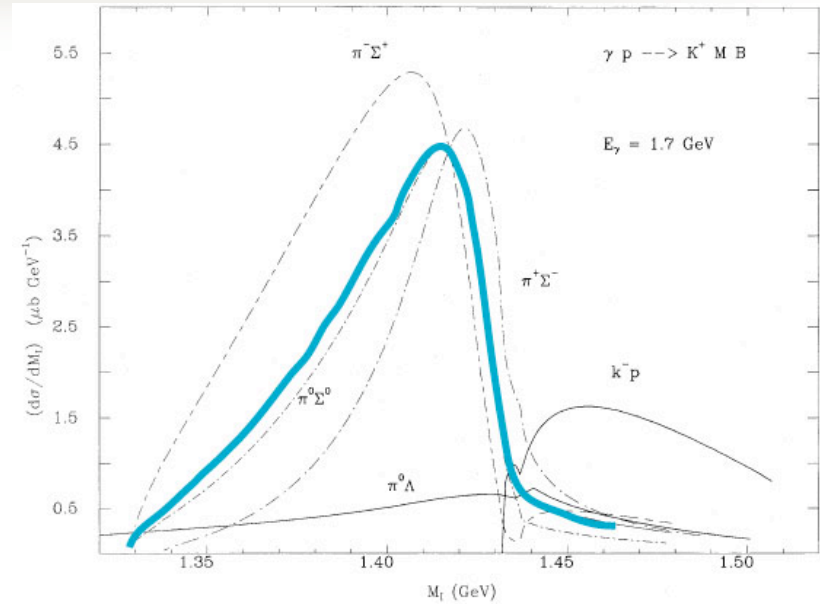
http://www.rcnp.osaka-u.ac.jp/~hyodo/index_e.html

Appendix : other processes



J.C. Nacher, et al., PLB445, 55(1999)

Spring-8



J.C. Nacher, et al., PLB461, 299(1999)

J-PARC?

