

# Photoproduction of $K^*$ for the study of $\Lambda(1405)$



**Tetsuo Hyodo<sup>a</sup>,**

**A. Hosaka<sup>a</sup>, E. Oset<sup>b</sup>, and M. J. Vicente Vacas<sup>b</sup>**

*RCNP, Osaka<sup>a</sup>      Valencia Univ.<sup>b</sup>*

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## Motivation : Two poles?

There are two poles of the scattering amplitude around nominal  $\Lambda(1405)$  energy region.

- Cloudy bag model

J. Fink, *et al.*, PRC41, 2720

- Chiral unitary model

J. A. Oller, *et al.*, PLB500, 263

E. Oset, *et al.*, PLB527, 99

D. Jido, *et al.*, PRC66, 025203

T. Hyodo, *et al.*, PRC68, 018201

C. Garcia-Recio, *et al.*, PRD67, 076009

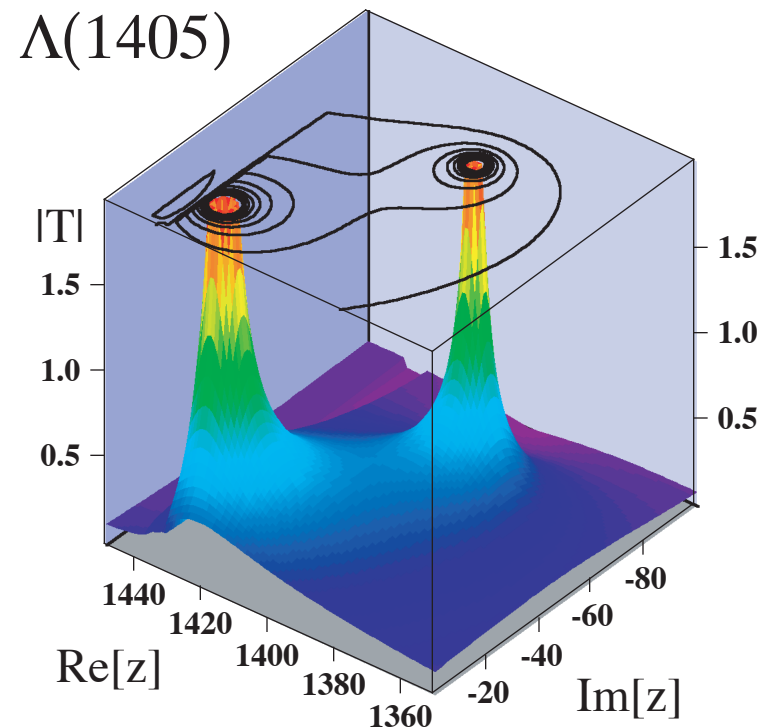
D. Jido, *et al.*, NPA725, 181

T. Hyodo, *et al.*, PRC68, 065203

- Quark model :  $qq-q\bar{q}$

A. Zhang, *et al.*, hep-ph/0403210

$$\Lambda(1405) : J^P = 1/2^-, I = 0$$



ChU model, T. Hyodo

# Chiral unitary model

Flavor SU(3) meson-baryon scatterings (s-wave)

**Chiral symmetry**

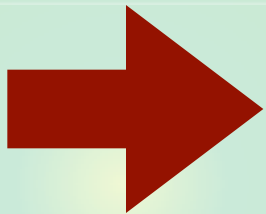
**Low energy  
behavior**



**Unitarity of S-matrix**

**Non-perturbative  
resummation**

**Dynamical  
generation**



$J^P = 1/2^-$  resonances

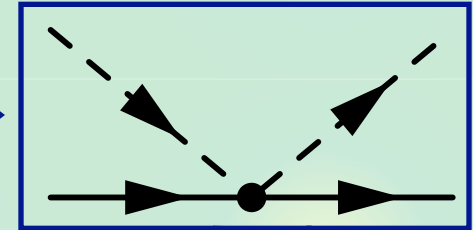
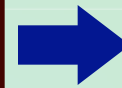
$\Lambda(1405), \Lambda(1670),$   
 $\Sigma(1620), \Xi(1620),$   
 $N(1535)$



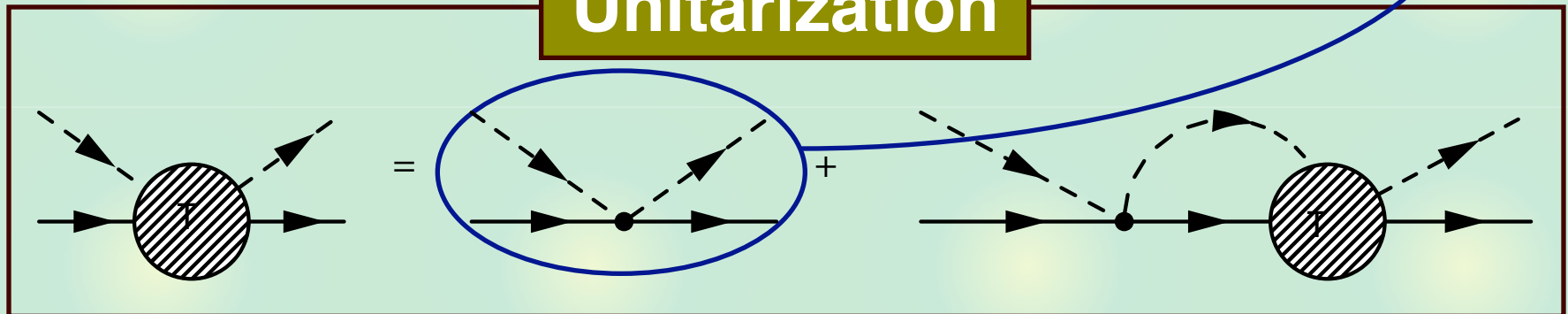
# Framework of the chiral unitary model

## Chiral perturbation theory

$$\mathcal{L}_{WT} = \frac{1}{4f^2} \text{Tr}(\bar{B}i\gamma^\mu[(\Phi\partial_\mu\Phi - \partial_\mu\Phi\Phi), B])$$

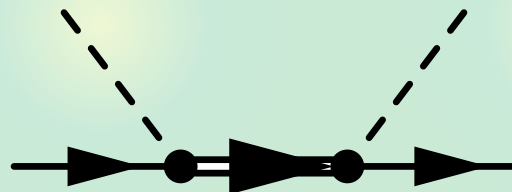


## Unitarization



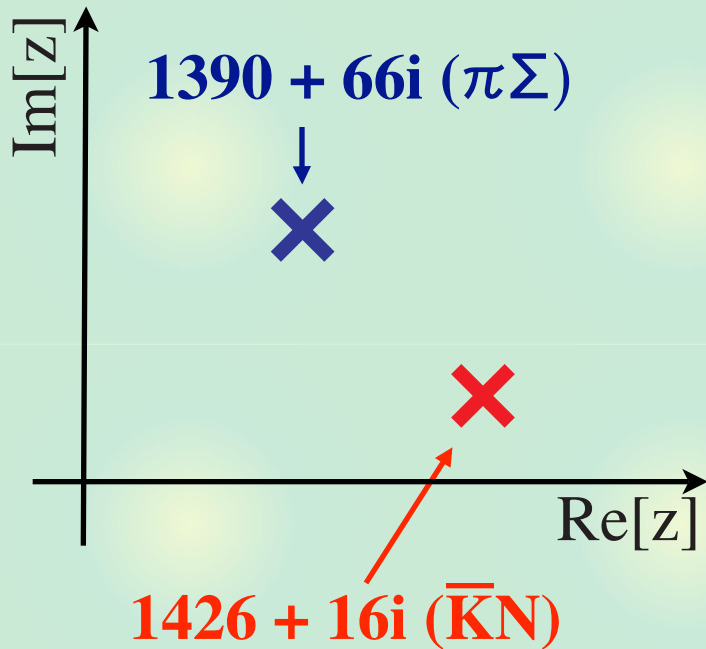
$$T_{ij}(\sqrt{s}) \sim \frac{g_i g_j}{\sqrt{s} - M_R + i\Gamma_R/2} + T_{ij}^{BG}$$

$\sim$

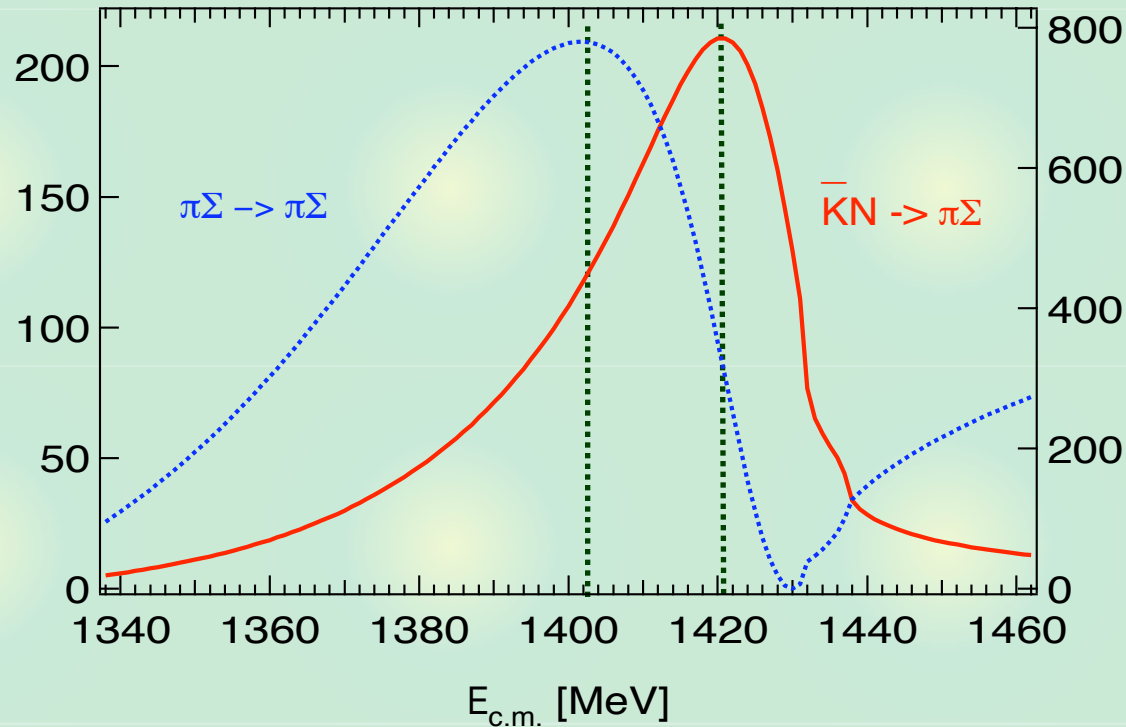


# $\Lambda(1405)$ in the chiral unitary model

## position of poles



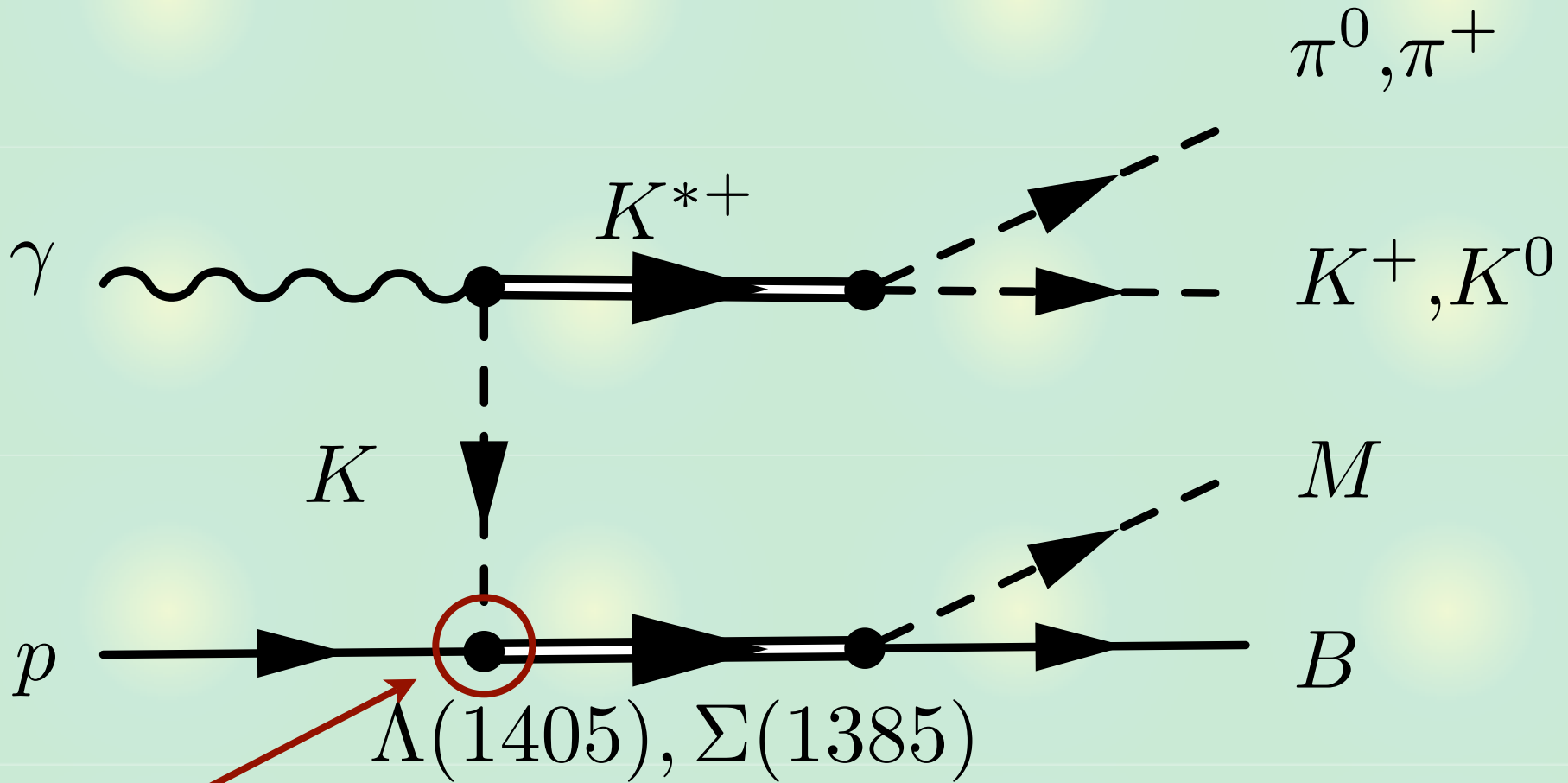
## $\pi\Sigma$ mass distribution



$$\frac{d\sigma}{dM_I} = C |t_{\pi\Sigma \rightarrow \pi\Sigma}|^2 p_{CM} \quad \longrightarrow \quad \frac{d\sigma}{dM_I} = \left| \sum_i C_i t_{i \rightarrow \pi\Sigma} \right|^2 p_{CM}$$

**D. Jido, et al., Nucl. Phys. A 723, 205 (2003)**

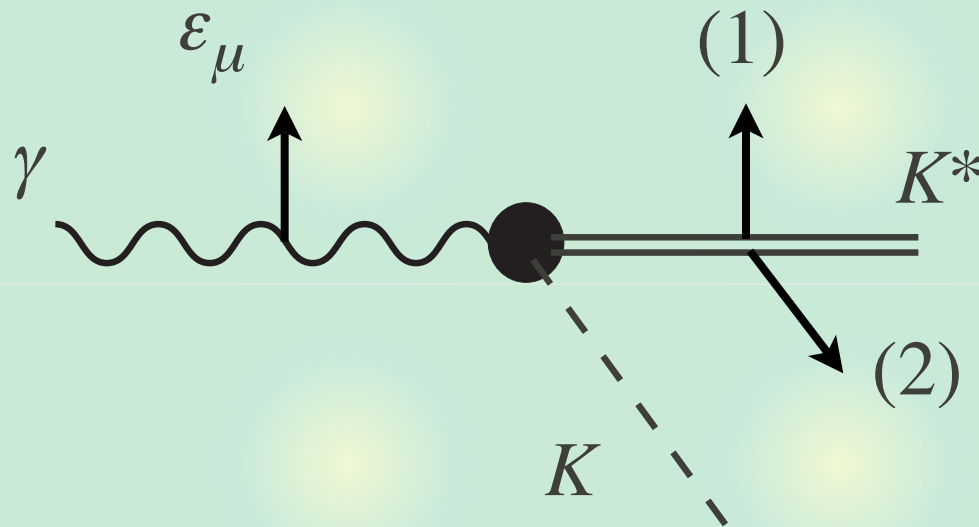
# Photoproduction of $K^*$ and $\Lambda(1405)$



**Only  $K^-p$  channel appears at the initial stage**  
**Higher energy pole ??**

## Advantage of this reaction

There is a correlation between polarization of initial photon and final  $\pi K$  angular distribution



(1)  $\epsilon_\mu(K^*) // \epsilon_\mu(\gamma) : J^P = \text{natural}$

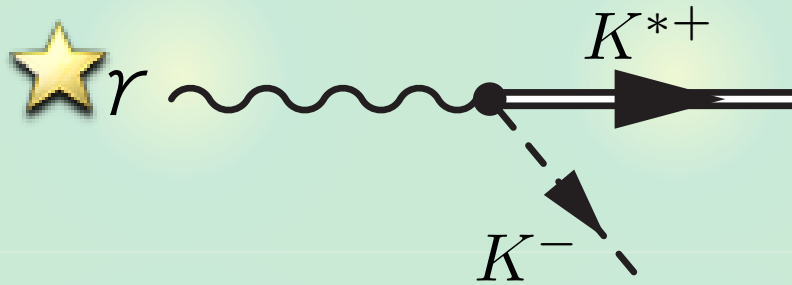
(2)  $\epsilon_\mu(K^*) \perp \epsilon_\mu(\gamma) : J^P = \text{unnatural}$

**Clear mechanism**

# Effective interaction for meson part

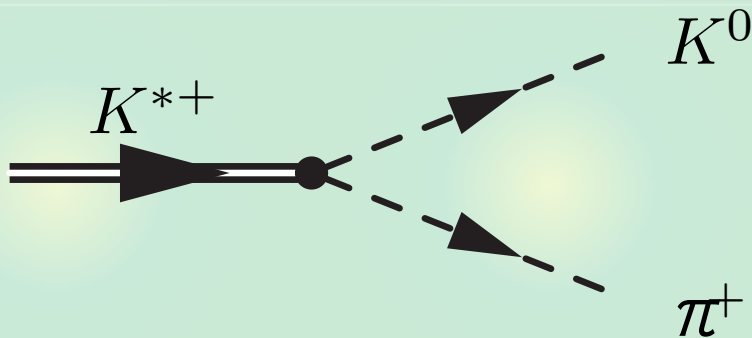
## ★ $\gamma K K^*$ coupling

$$\mathcal{L}_{K^* K \gamma} = g_{K^* K \gamma} \epsilon^{\mu\nu\alpha\beta} \partial_\mu A_\nu (\partial_\alpha K_\beta^{*-} K^+ + \partial_\alpha \bar{K}_\beta^{*0} K^0) + \text{h.c.}$$



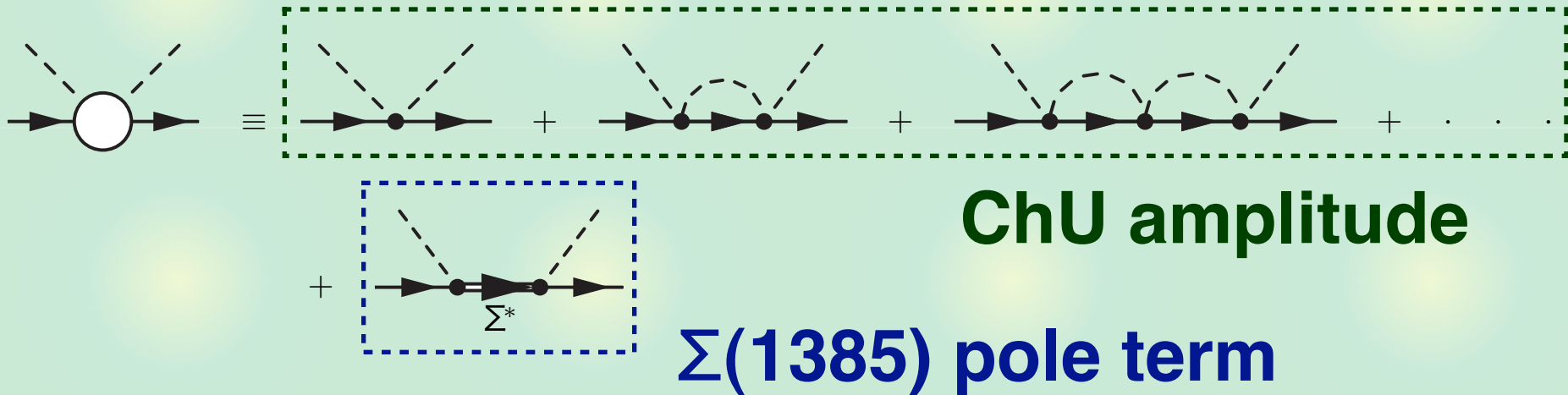
## ★ VPP coupling

$$\mathcal{L}_{VPP} = -\frac{ig_{VPP}}{\sqrt{2}} \text{Tr}(V^\mu [\partial_\mu P, P])$$





# Effective interaction for baryon part



★  **$\Sigma(1385)$ MB coupling**

★ 
$$-it_{\Sigma^* i} = c_i \frac{12}{5} \frac{D + F}{2f} \mathbf{S} \cdot \mathbf{k}_i$$

★ **form factor**

$$F_f(k_1) = \frac{\Lambda^2 - m_K^2}{\Lambda^2 - (k_1)^2}$$

## Isospin decomposition of $\pi\Sigma$ states

Since initial state is  $\bar{K}N$ , we neglect the  $l=2$ .

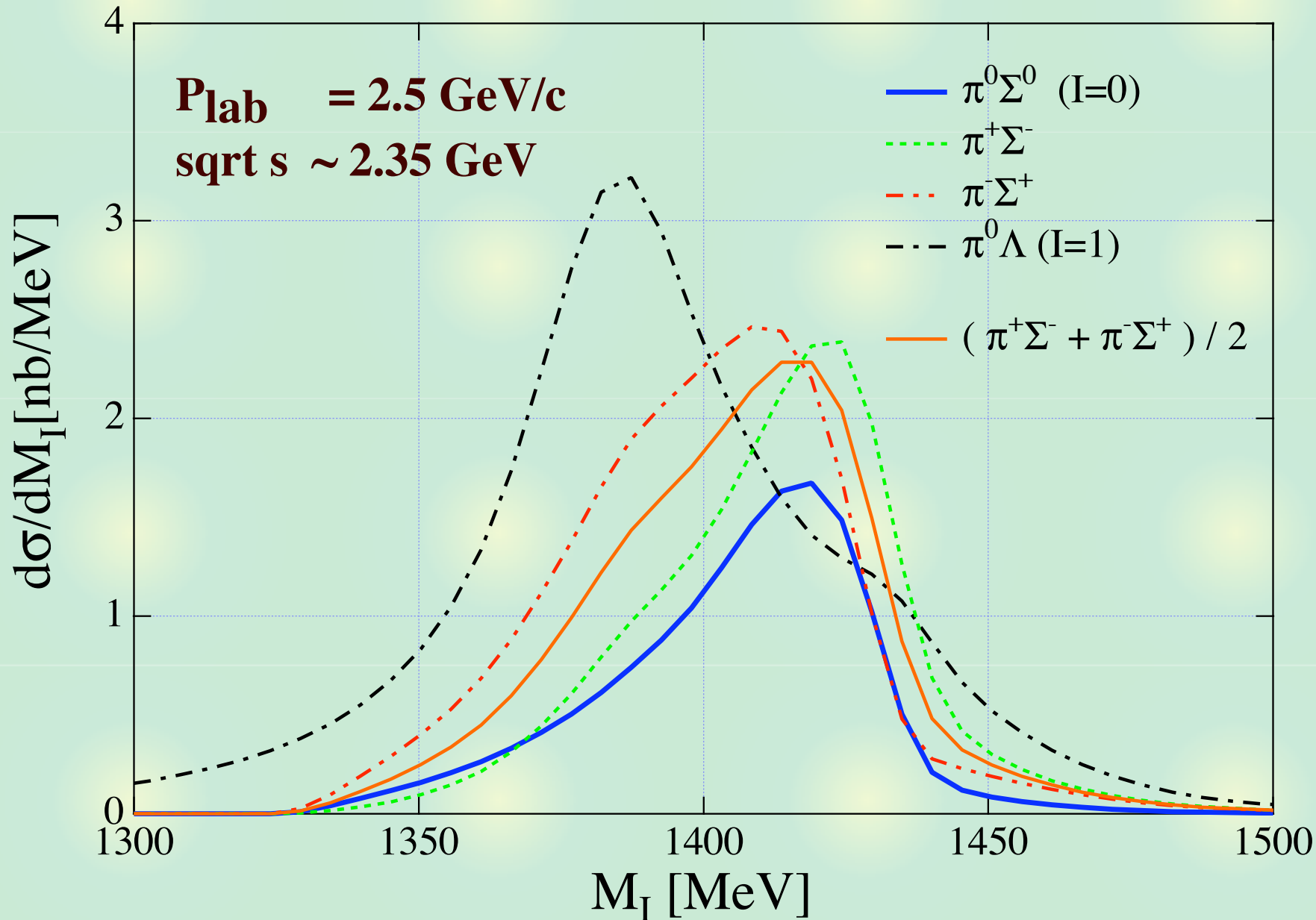
$$\frac{d\sigma(\pi^0\Sigma^0)}{dM_I} \propto \frac{1}{3} |T^{(0)}|^2$$

- Pure  $l=0$  amplitude

$$\frac{d\sigma(\pi^\pm\Sigma^\mp)}{dM_I} \propto \frac{1}{3} |T^{(0)}|^2 + \frac{1}{2} |T^{(1)}|^2 \pm \frac{2}{\sqrt{6}} \text{Re}(T^{(0)}T^{(1)*})$$

- Difference among charged states  
-> when summed up, this term vanishes

# Invariant mass distributions



## Summary

We study the structure of  $\Lambda(1405)$  using the chiral unitary model.

- There are **two poles** of the scattering amplitude around nominal  $\Lambda(1405)$ .
- In the  $\gamma p \rightarrow K^* \Lambda(1405)$  reaction, **Charged  $\pi\Sigma$  states** isolate the **higher energy pole**.
- Neutral  $\pi\Sigma$  state** gives a clear  **$l=0$  distribution**.

[T. H., A. Hosaka, E. Oset, M. J. Vicente Vacas, nucl-th/0401051](#)