

Photoproduction of K^* for the study of $\Lambda(1405)$



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Introduction : $\Lambda(1405)$

$\Lambda(1405) : J^P = 1/2^-, I = 0$

Mass : 1406.5 ± 4.0 MeV

Width : 50 ± 2 MeV

Decay mode : $\Lambda(1405) \rightarrow (\pi\Sigma)_{I=0} \quad 100\%$

Coupled channel multi-scattering

R.H. Dalitz, T.C. Wong and G. Rajasekaran PR 153, 1617 (1967)

(Naive) Quark model : 1500 MeV?

N. Isgur, and G. Karl, PRD 18, 4187 (1978)

Meson-Baryon or 3-quark?

Introduction : $\Lambda(1405)$

Recent works

Quark model : LS partner of $\Lambda(1520)$

C.L. Schat, J.L. Goity and N.N. Scoccola PRL 88, 100202 (2002)

Lattice QCD with 3-quark operator

W. Melnitchouk, *et. al.*, PRD 67, 114506 (2003)

--> not a 3-quark

Y. Nemoto, *et. al.*, PRD 68, 094505 (2003)

F.X. Lee, *et. al.*, NP 119, 296 (2003)

--> 3-quark?

Chiral unitary approaches

Valencia, GSI, ...

--> Meson-baryon picture

Motivation : Two poles?

There are two poles of the scattering amplitude around nominal $\Lambda(1405)$ energy region.

- **Cloudy bag model**

J. Fink, *et al.*, PRC41, 2720

- **Chiral unitary model**

J. A. Oller, *et al.*, PLB500, 263

E. Oset, *et al.*, PLB527, 99

D. Jido, *et al.*, PRC66, 025203

T. Hyodo, *et al.*, PRC68, 018201

T. Hyodo, *et al.*, PTP112, 73

C. Garcia-Recio, *et al.*, PRD67, 076009

D. Jido, *et al.*, NPA725, 181

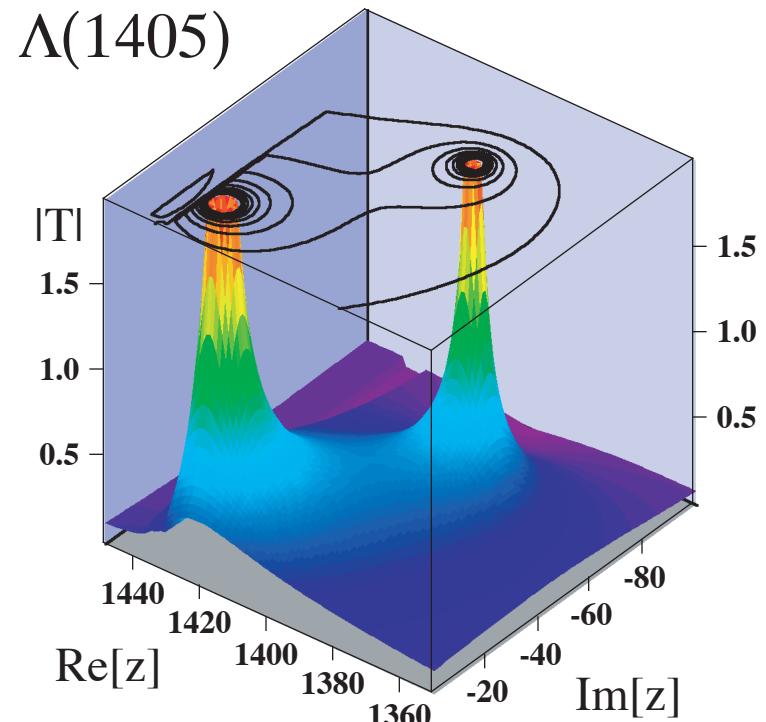
T. Hyodo, *et al.*, PRC68, 065203

C. Garcia-Recio, *et al.*, PLB582, 49

- **Correlated quark model**

A. Zhang, *et al.*, hep-ph/0403210

$$\Lambda(1405) : J^P = 1/2^-, I = 0$$



ChU model, T. Hyodo

Chiral unitary model

Flavor SU(3) meson-baryon scatterings (s-wave)

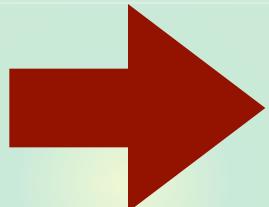
Chiral symmetry

Low energy behavior

Unitarity of S-matrix

Non-perturbative resummation

Dynamical generation



$J^P = 1/2^-$ resonances

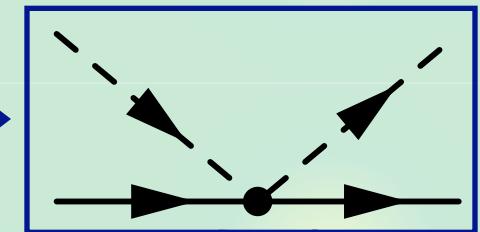
$\Lambda(1405), \Lambda(1670),$
 $\Sigma(1620), \Xi(1620),$
 $N(1535)$



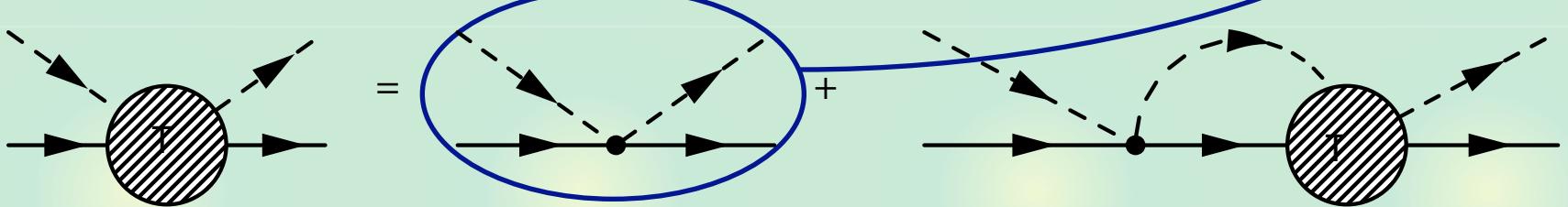
Framework of the chiral unitary model

Chiral perturbation theory

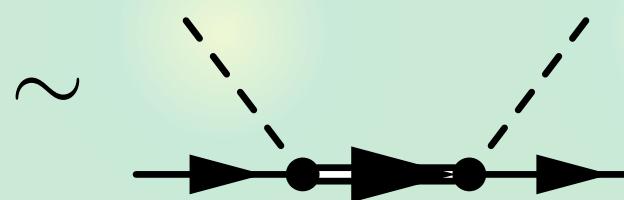
$$\mathcal{L}_{WT} = \frac{1}{4f^2} \text{Tr}(\bar{B}i\gamma^\mu [(\Phi\partial_\mu\Phi - \partial_\mu\Phi\Phi), B])$$



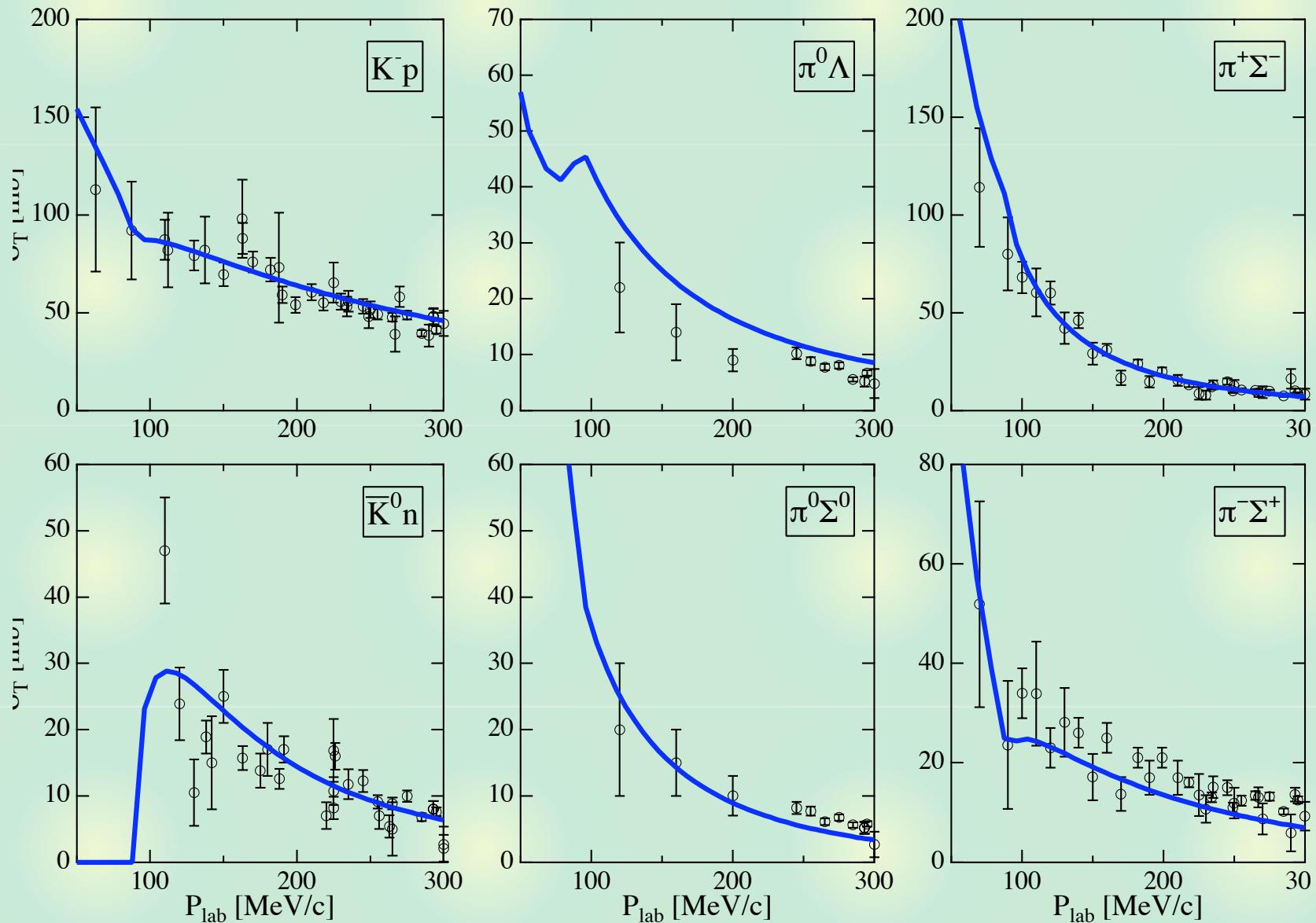
Unitarization



$$T_{ij}(\sqrt{s}) \sim \frac{g_i g_j}{\sqrt{s} - M_R + i\Gamma_R/2} + T_{ij}^{BG}$$



Total cross sections of K-p scattering

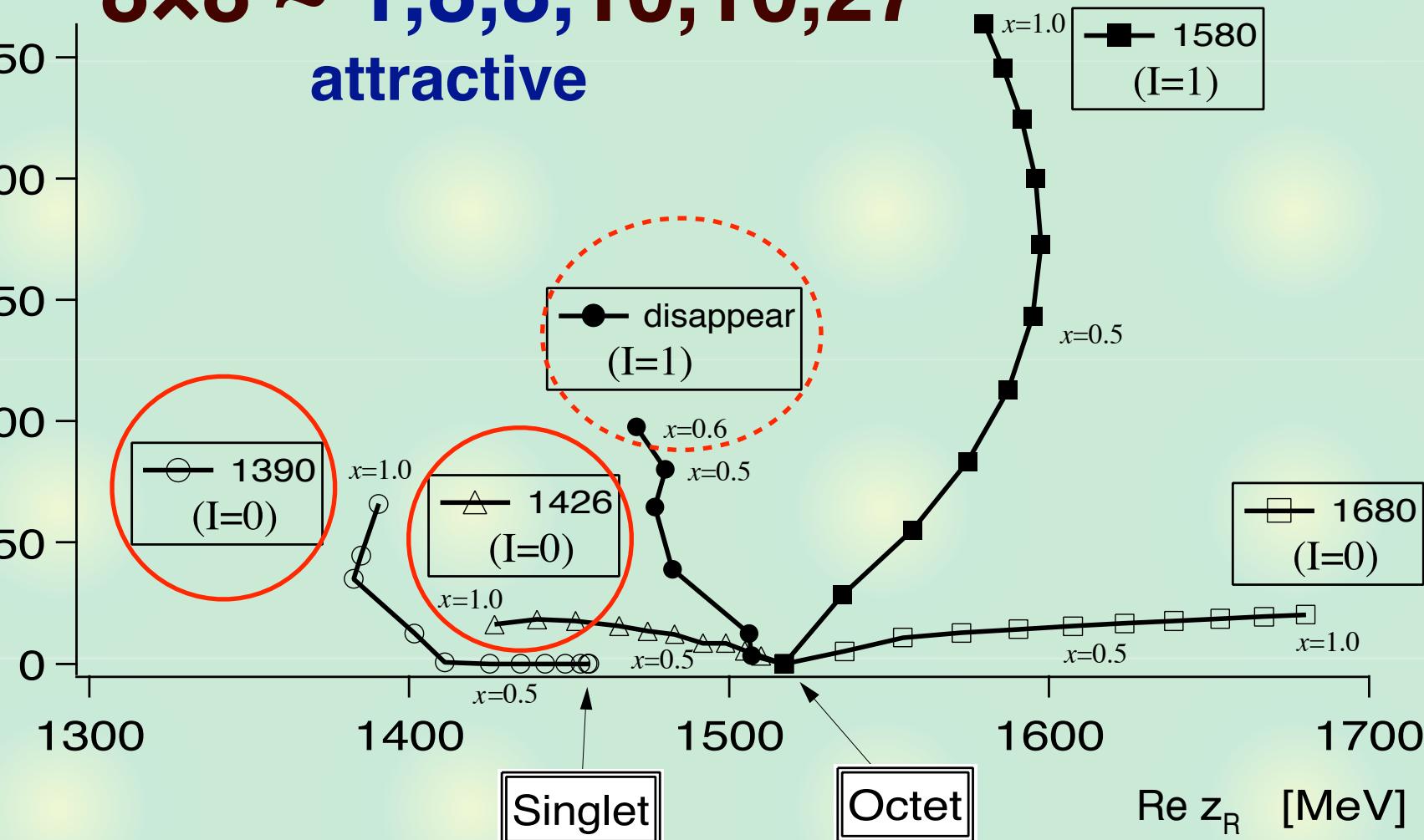


T. Hyodo, *et al.*, Phys. Rev. C 68, 018201 (2003)

Trajectories of the poles with SU(3) breaking ($S = -1$)

$8 \times 8 \sim 1, 8, 8, 10, \overline{10}, 27$
attractive

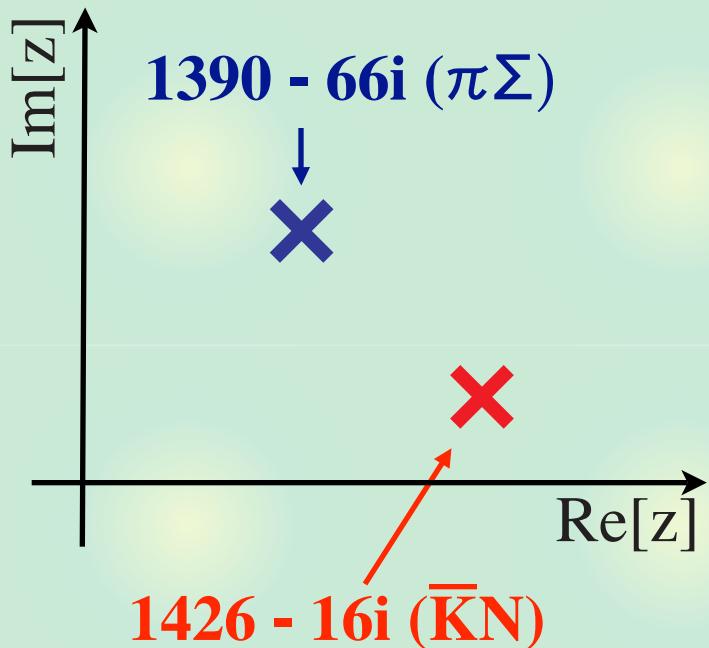
$\text{Im } z_R$ [MeV]



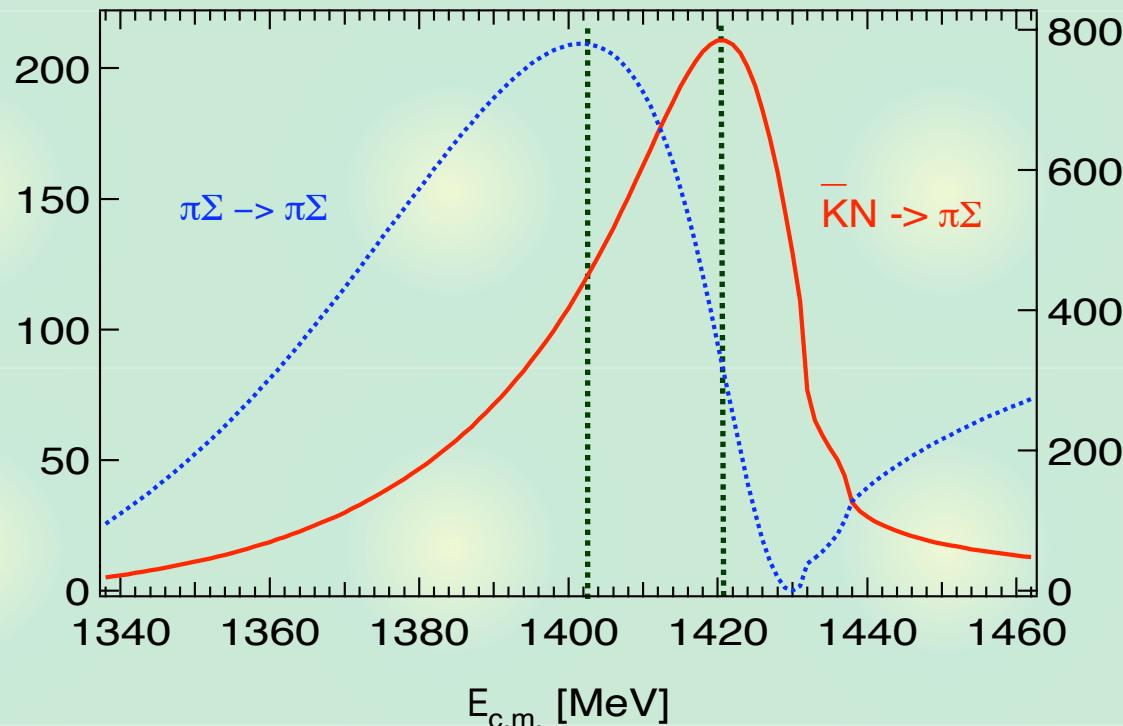
D. Jido, et al., Nucl. Phys. A 723, 205 (2003)

$\Lambda(1405)$ in the chiral unitary model

position of poles



$\pi\Sigma$ mass distribution



$$\frac{d\sigma}{dM_I} = C |t_{\pi\Sigma \rightarrow \pi\Sigma}|^2 p_{CM} \rightarrow \frac{d\sigma}{dM_I} = \left| \sum_i C_i t_{i \rightarrow \pi\Sigma} \right|^2 p_{CM}$$

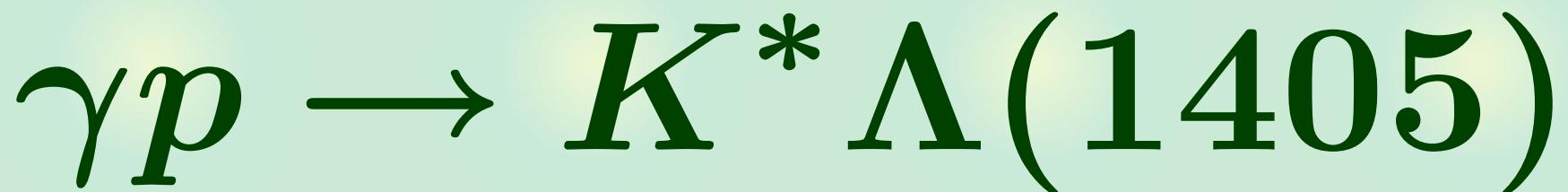
D. Jido, et al., Nucl. Phys. A 723, 205 (2003)

Photoproduction of K^* and $\Lambda(1405)$

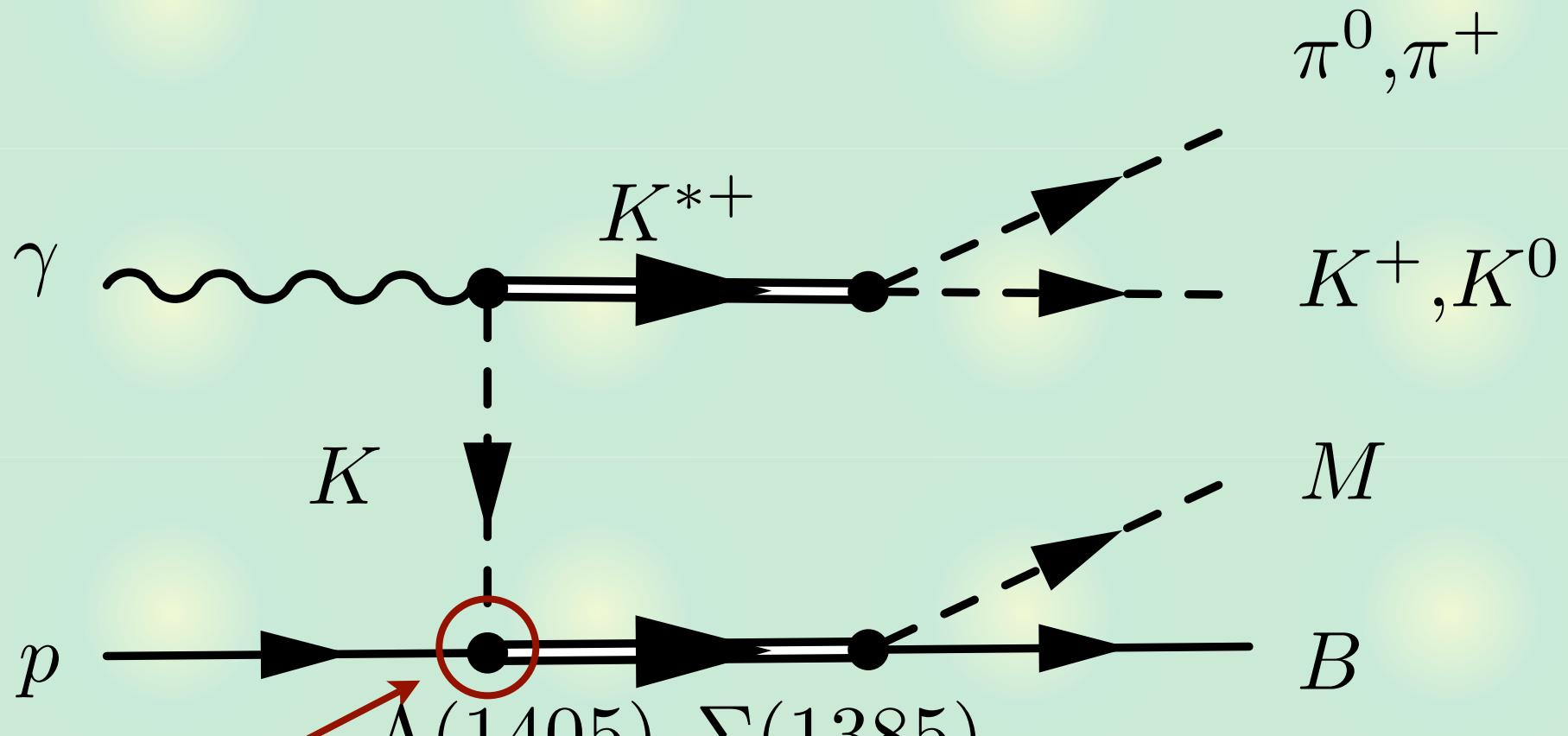
In order to study

- ★ $S = -1, l = 0, s\text{-wave} : \Lambda(1405)$
 - ★ two poles?
 - ★ $1426 - 16i : \bar{K}N$
 - ★ $1390 - 66i : \pi\Sigma$
- ★ $S = -1, l = 1, s\text{-wave}$
 - ★ pole? --> see paper

we calculate

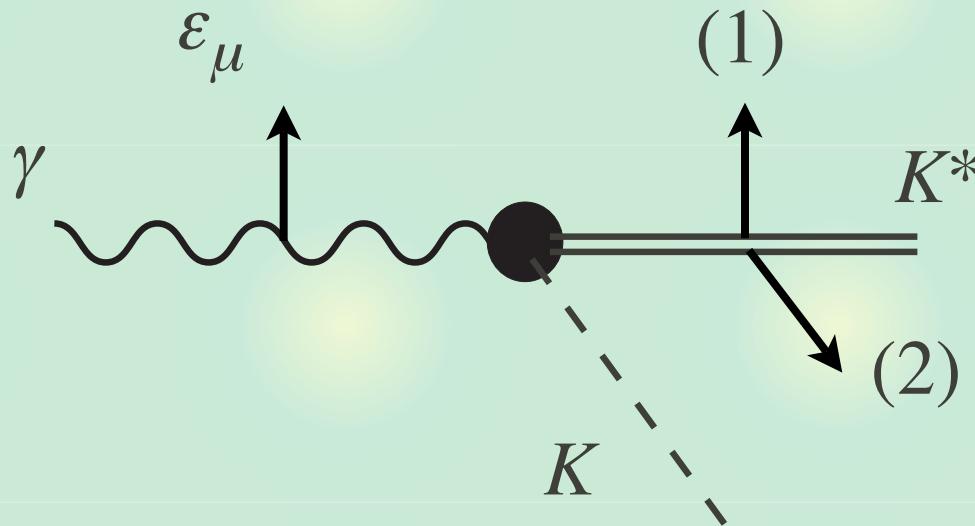


Photoproduction of K^* and $\Lambda(1405)$



**Only K^-p channel appears at the initial stage
Higher energy pole ??**

Advantage of this reaction



(1) $\varepsilon_\mu(K^*) \parallel \varepsilon_\mu(\gamma) : J^P = \text{natural}$

(2) $\varepsilon_\mu(K^*) \perp \varepsilon_\mu(\gamma) : J^P = \text{unnatural}$

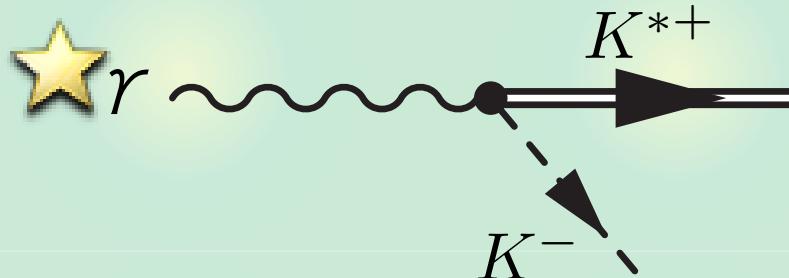
With polarized photon beam, the exchanged particle can be identified.

Clear mechanism

Effective interaction for meson part

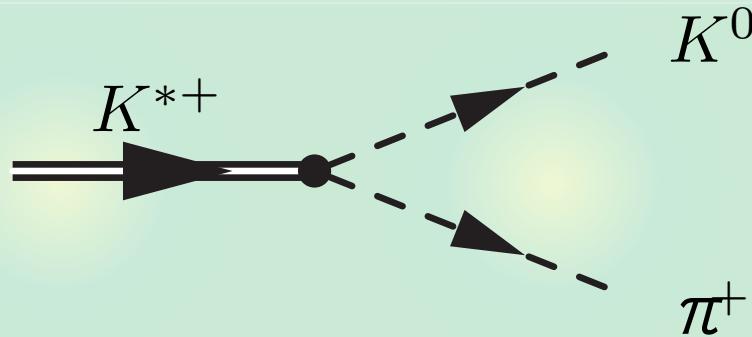
★ $\gamma K K^*$ coupling

$$\mathcal{L}_{K^* K \gamma} = g_{K^* K \gamma} \epsilon^{\mu\nu\alpha\beta} \partial_\mu A_\nu (\partial_\alpha K_\beta^{*-} K^+ + \partial_\alpha \bar{K}_\beta^{*0} K^0) + \text{h.c.}$$

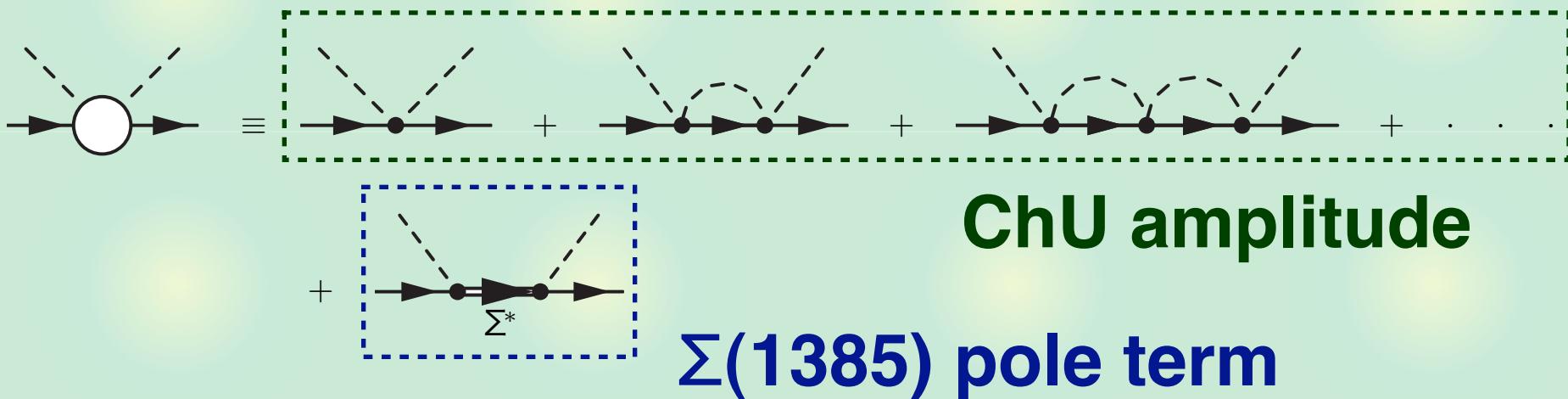


★ VPP coupling

$$\mathcal{L}_{VPP} = -\frac{ig_{VPP}}{\sqrt{2}} \text{Tr}(V^\mu [\partial_\mu P, P])$$



Effective interaction for baryon part



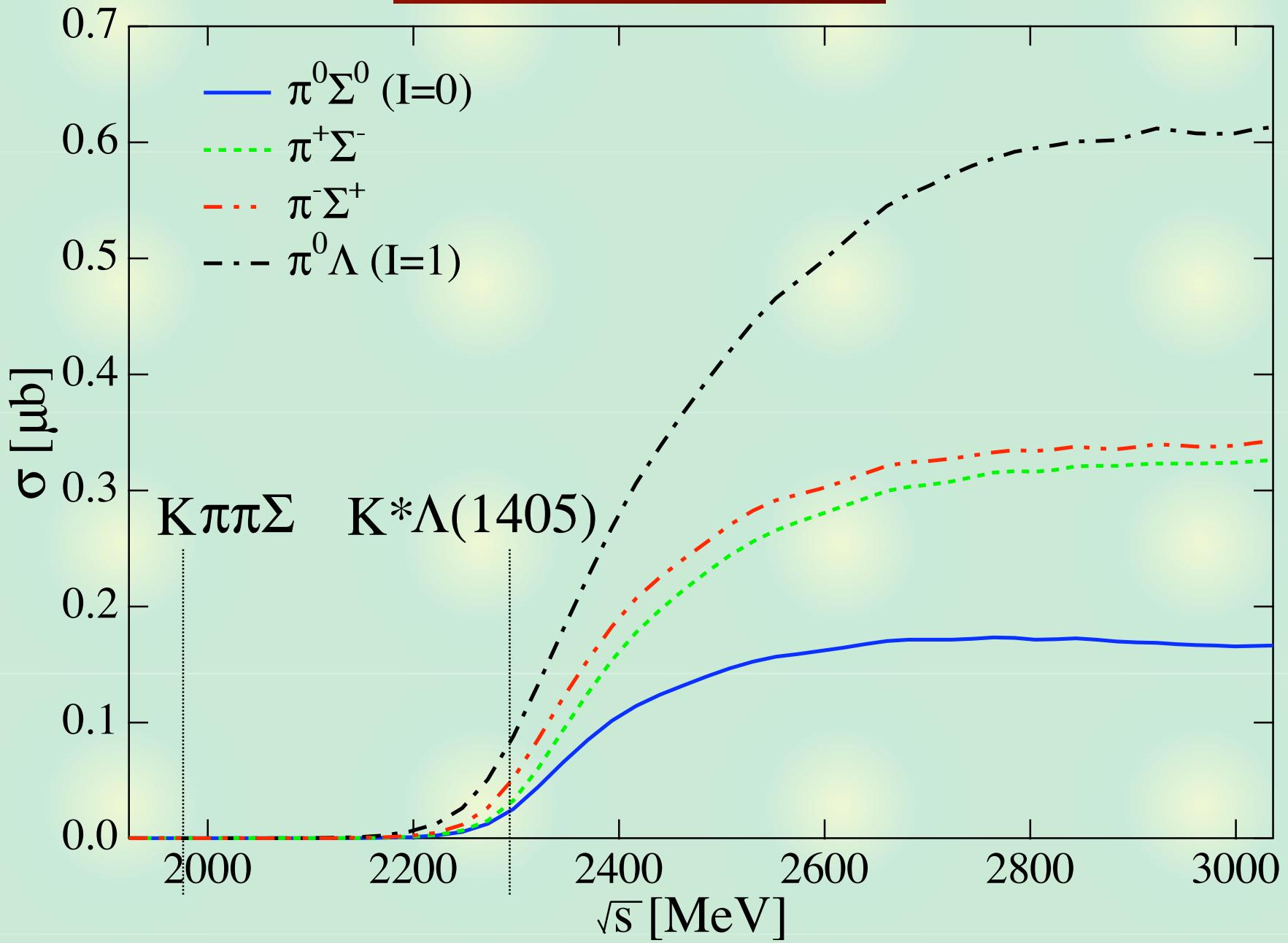
★ **$\Sigma(1385)$ MB coupling**

★ $-it_{\Sigma^* i} = c_i \frac{12}{5} \frac{D + F}{2f} S \cdot k_i$

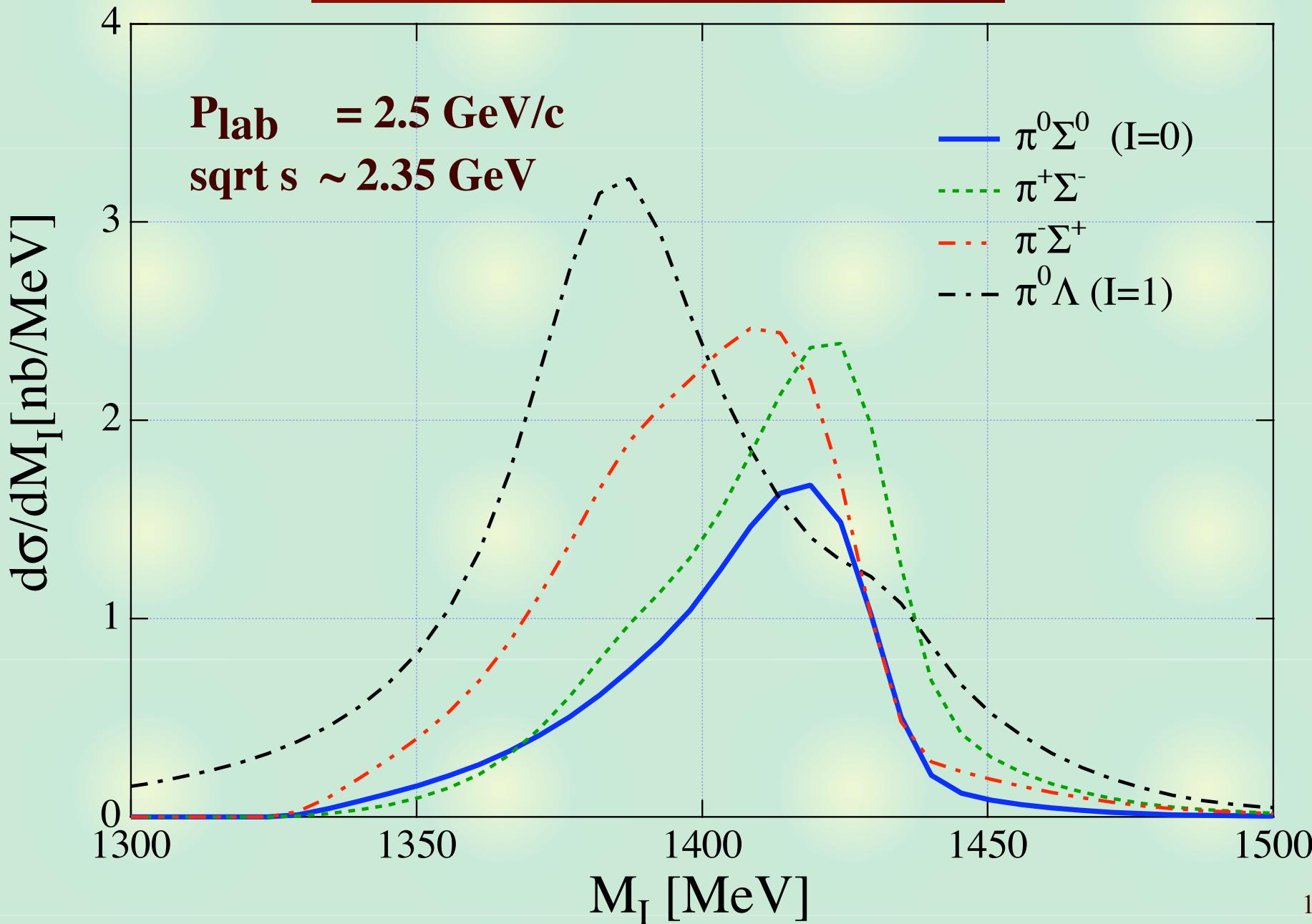
★ **form factor**

$$F_f(k_1) = \frac{\Lambda^2 - m_K^2}{\Lambda^2 - (k_1)^2}$$

Total cross sections



Invariant mass distributions



Isospin decomposition of $\pi\Sigma$ states

Since initial state is KN, we neglect the $I=2$.

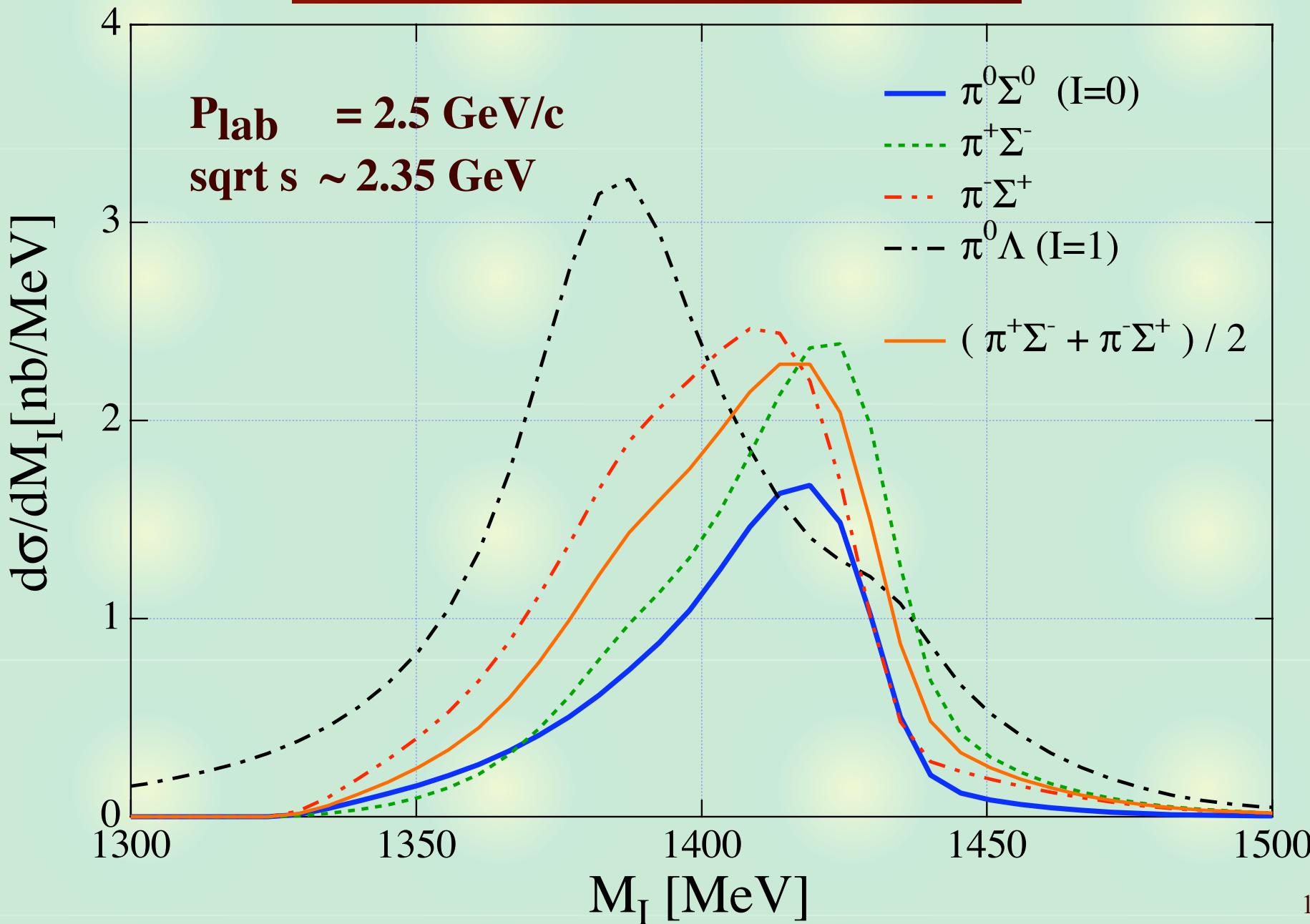
$$\frac{d\sigma(\pi^0\Sigma^0)}{dM_I} \propto \frac{1}{3}|T^{(0)}|^2$$

- Pure $I=0$ amplitude

$$\frac{d\sigma(\pi^\pm\Sigma^\mp)}{dM_I} \propto \frac{1}{3}|T^{(0)}|^2 + \frac{1}{2}|T^{(1)}|^2 \pm \frac{2}{\sqrt{6}}\text{Re}(T^{(0)}T^{(1)*})$$

- Difference among charged states
-> when summed up, this term vanishes
- No p-wave contribution
-> $I=1$ s-wave amplitude

Invariant mass distributions 2



Summary and conclusions 1

We study the structure of $\Lambda(1405)$ using the chiral unitary model.

There are two poles of the scattering amplitude around nominal $\Lambda(1405)$.

Pole 1 (1426–16i) : strongly couples to $\bar{K}N$ state

Pole 2 (1390–66i) : strongly couples to $\pi\Sigma$ state

By observing the charged $\pi\Sigma$ states in the $\gamma p \rightarrow K^* \Lambda(1405)$ reaction, it is possible to isolate the higher energy pole.

Summary and conclusions 2

- If we observe neutral $\pi\Sigma$ state, clear $|l=0$ distribution is obtained.
- Combining three $\pi\Sigma$ states, we can also study the s-wave $|l=1$ amplitude, where the existence of another pole is argued.

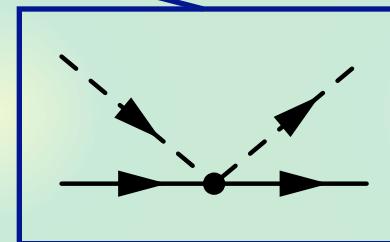
T. H., A. Hosaka, E. Oset, M. J. Vicente Vacas, PLB593, 75 (2004)

<http://www.rcnp.osaka-u.ac.jp/~hyodo/>

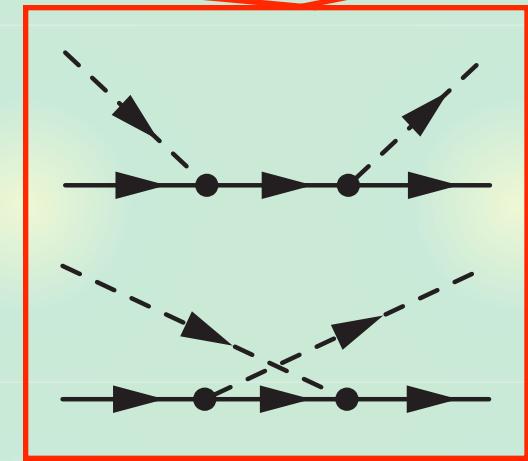
Appendix : ChPT Lagrangian

$$\mathcal{L}^{(1)} = \text{Tr} \left(\bar{B} (i \cancel{D} - M_0) B - D (\bar{B} \gamma^\mu \gamma_5 \{A_\mu, B\}) - F (\bar{B} \gamma^\mu \gamma_5 [A_\mu, B]) \right)$$

$$\mathcal{D}_\mu B = \partial_\mu B + i [\underline{V}_\mu, B]$$



$$\xi(\Phi) = \exp\{i\Phi/\sqrt{2}f\}$$



$$D + F = g_A$$

$$B = \begin{pmatrix} \frac{1}{\sqrt{2}}\Sigma^0 + \frac{1}{\sqrt{6}}\Lambda & \Sigma^+ & p \\ \Sigma^- & -\frac{1}{\sqrt{2}}\Sigma^0 + \frac{1}{\sqrt{6}}\Lambda & n \\ \Xi^- & \Xi^0 & -\frac{2}{\sqrt{6}}\Lambda \end{pmatrix} \quad \Phi = \begin{pmatrix} \frac{1}{\sqrt{2}}\pi^0 + \frac{1}{\sqrt{6}}\eta & \pi^+ & K^+ \\ \pi^- & -\frac{1}{\sqrt{2}}\pi^0 + \frac{1}{\sqrt{6}}\eta & K^0 \\ K^- & \bar{K}^0 & -\frac{2}{\sqrt{6}}\eta \end{pmatrix}$$

$$\underline{V}_\mu = -\frac{i}{2} (\xi^\dagger \partial_\mu \xi + \xi \partial_\mu \xi^\dagger) = \frac{i}{4f^2} \underline{(\Phi \partial_\mu \Phi - \partial_\mu \Phi \Phi)} + \dots$$

$$\underline{A}_\mu = -\frac{i}{2} (\xi^\dagger \partial_\mu \xi - \xi \partial_\mu \xi^\dagger) = -\frac{1}{f} \underline{\partial_\mu \Phi} + \dots$$

Appendix : Several treatments 1

N. Kaiser, P. B. Siegel and W. Weise, NPA594, 325, PLB362, 23 (1995)

(HB)ChPT p^2 , Form factor (channel dep.)

S = -1, 0, $\Lambda(1405)$, $N(1535)$

E. Oset and A. Ramos, NPA635, 99 (1998)

WT term, 3-momentum cutoff (channel indep.)

S = -1, $\Lambda(1405)$

J. A. Oller and U. G. Meissner, PLB500, 263 (2001)

ChPT p , dimensional reg. (channel indep.)

S = -1, $\Lambda(1405)$

Analytic solution for BS eq.

-> pole structure in complex plane

Appendix : Several treatments 2

E. Oset, A. Ramos and C. Bennhold, PLB527, 99 (2002)

T. Inoue, E. Oset, and M. J. Vicente Vacas, PRC 65, 035204 (2002)

A. Ramos, E. Oset and C. Bennhold, PRL 89, 252001 (2002)

WT term, dimensional reg. (channel dep.)

$\Lambda(1405)$, $\Lambda(1670)$, $\Sigma(1620)$, $N(1535)$, $\Xi(1620)$

Scattering observables, $S = -1, 0$, analytic

M. F. M. Lutz and E. Kolomeitsev, NPA700, 193 (2002)

ChPT p^3 , Optimal reno. (channel indep.)

Scattering observables, Numerical solution

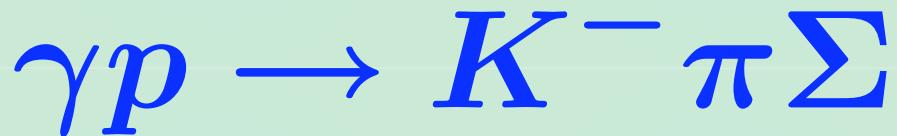
C. Garcia-Recio, M. F. M. Lutz and J. Nieves, PLB582, 49 (2004)

WT term, Optimal reno. (channel indep.)

$\Lambda(1405)$, $\Lambda(1670)$, $\Sigma(1620)$, $N(1535)$, $\Xi(1620)$, $\Xi(1690)$

Scattering observables?

Appendix : other processes



J.C. Nacher, et al., PLB445, 55



J.C. Nacher, et al., PLB461, 299

