

# Impact of tensor and short-range correlations in nuclear physics

J Carlson, LANL

Tensor Force

Deuteron

T20

Nuclear density matrix

Three-nucleon interactions

JLAB/BNL correlated pairs

Inclusive electron/neutrino scattering

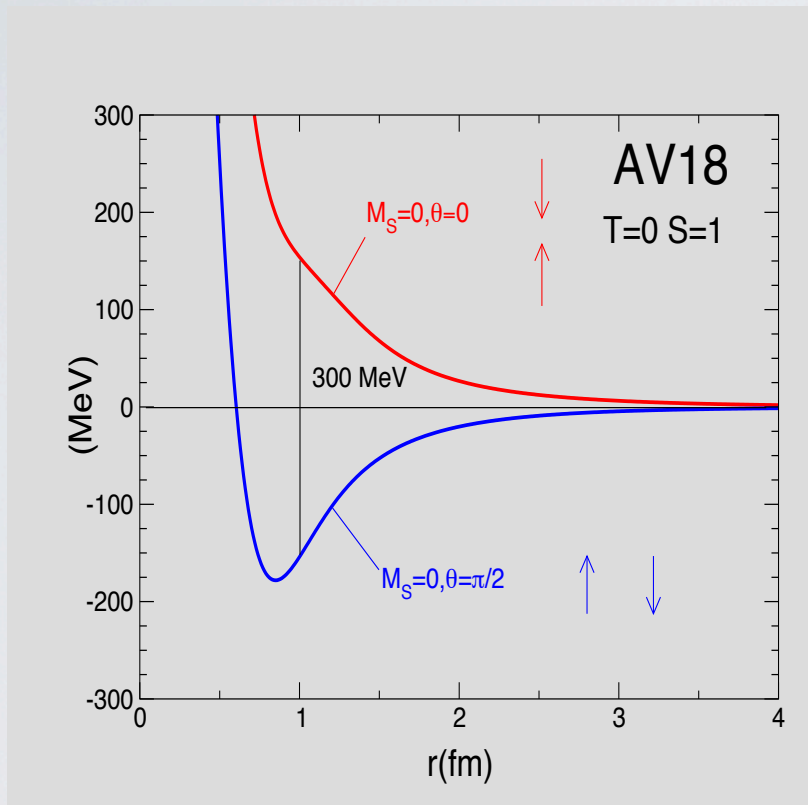
Neutrino emissivity in neutron matter



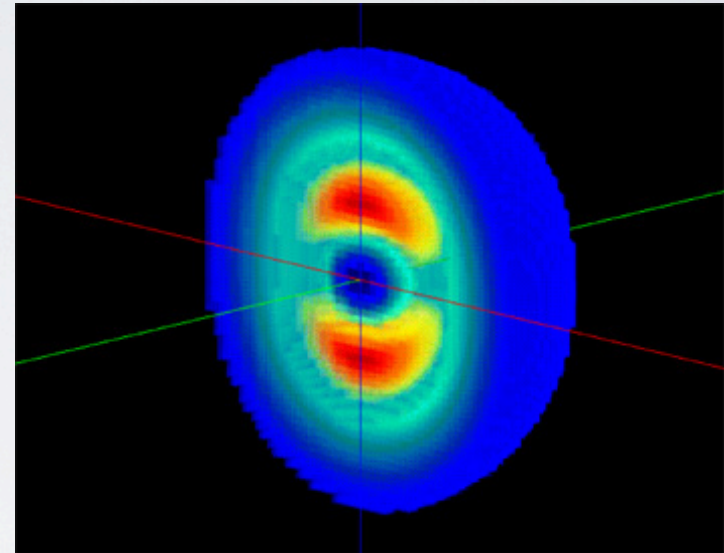
Santa Fe Plaza Dec 9

work with: Wiringa, Schiavilla, Pieper, Shen, Reddy, Gandolfi,...

# Tensor force (and spin-orbit) couple spin to space



diagonal elements of  
force in different spatial  
directions (Forest, PRC, 1996)



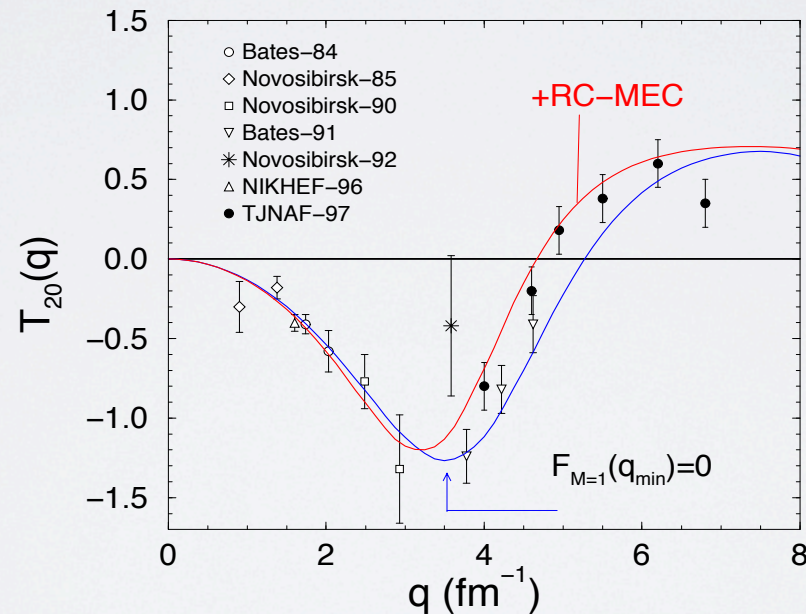
Intrinsic shape of the deuteron (Jlab);  
low density in interior (repulsion) and  
exterior. Intermediate behavior is strongly  
space-dependent

$$Q_d = 0.286 \text{ fm}^2$$

# JLab measurement of T20 in elastic electron scattering

$$A(q) \simeq |F_{M=0}(q)|^2 + 2 |F_{M=1}(q)|^2$$

$$T_{20}(q) \simeq -\sqrt{2} \frac{|F_{M=0}(q)|^2 - |F_{M=1}(q)|^2}{|F_{M=0}(q)|^2 + 2|F_{M=1}(q)|^2}$$

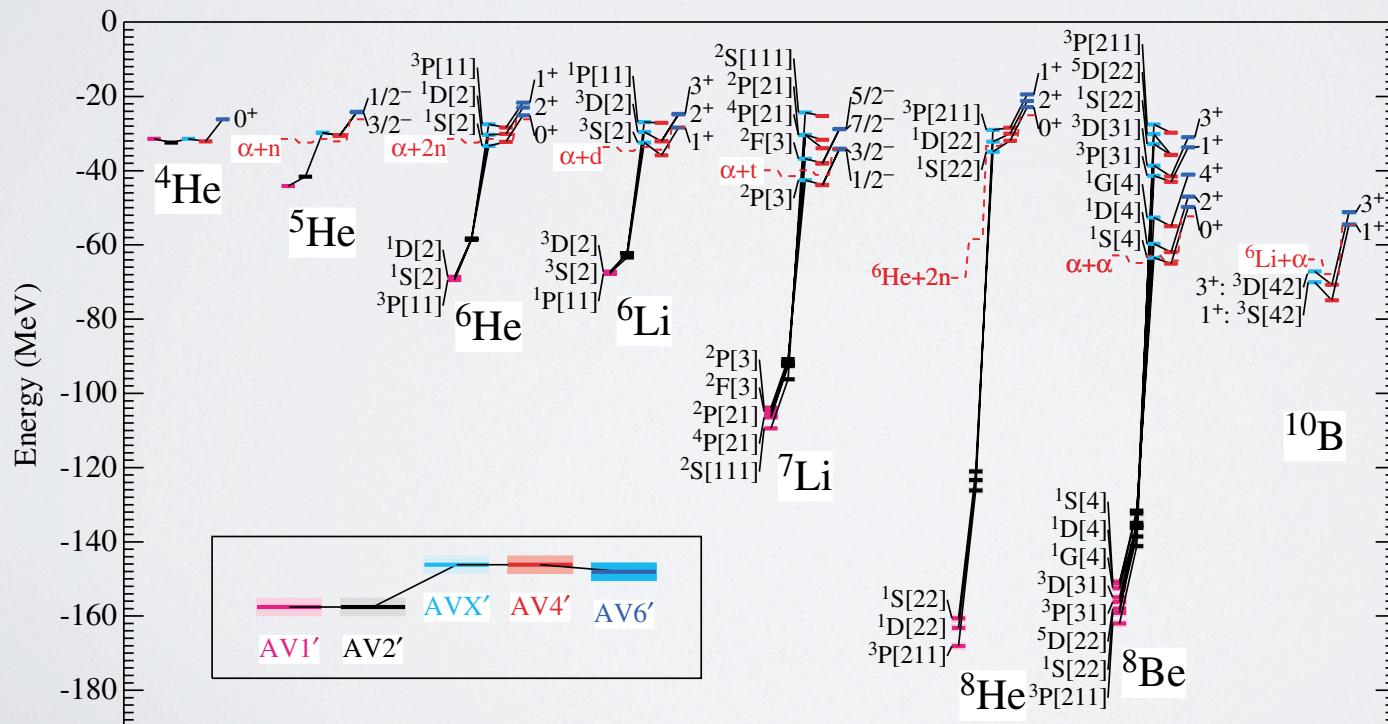


Roy Holt Bonner Prize

# Tensor correlations in $A > 2$

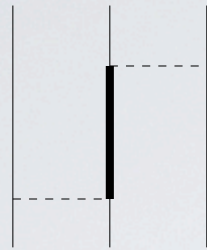
Much more difficult to 'see' in larger nuclei.  
 Typical matrix elements (eg. energy) involve  $S_{ij}^2$

Light nuclei spectra, though, require a 'realistic' force

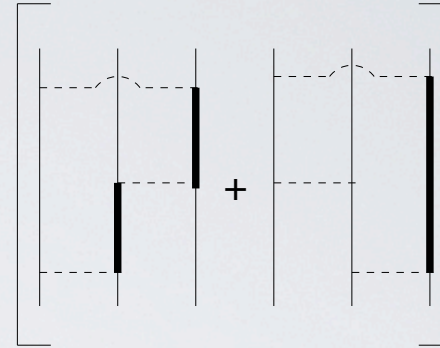


# alpha-n phase shifts

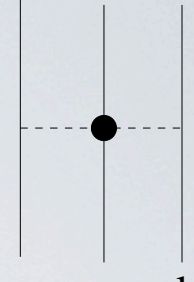
3-nucleon  
interaction: UIX



IL7:  $A^{3\pi}$

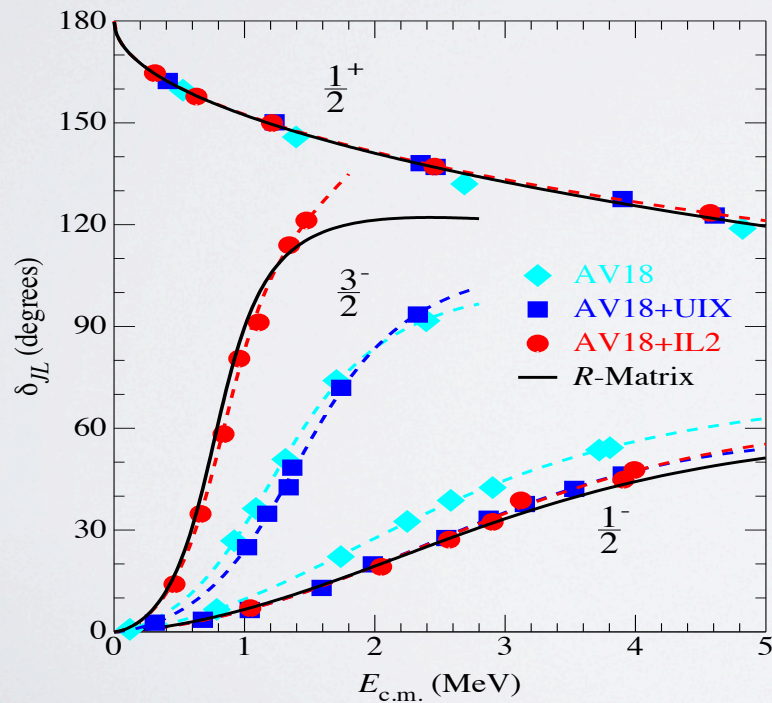


+  $A_{sw}^{2\pi}$

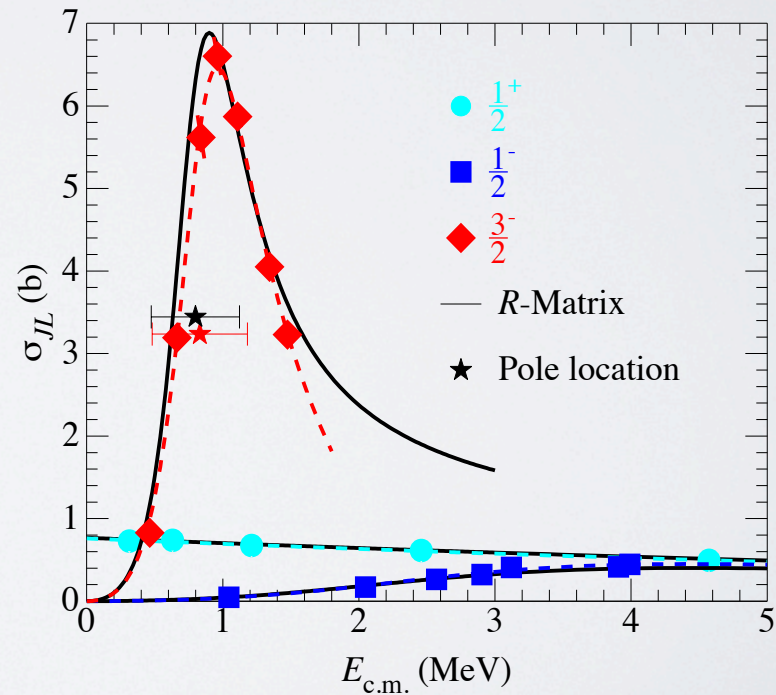


very weak

## Phase Shifts

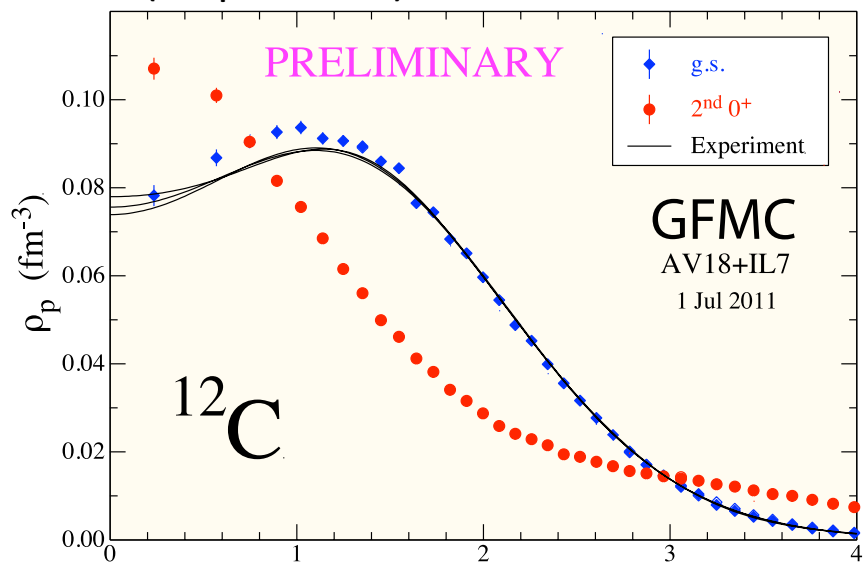


## Cross Sections (AV18/IL2)

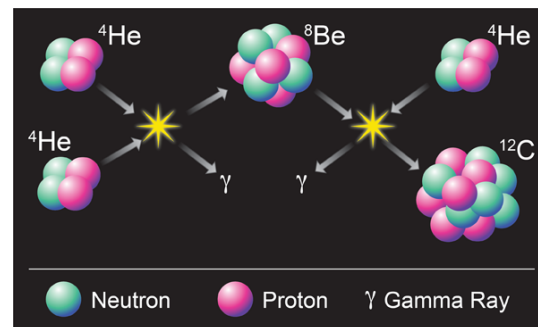
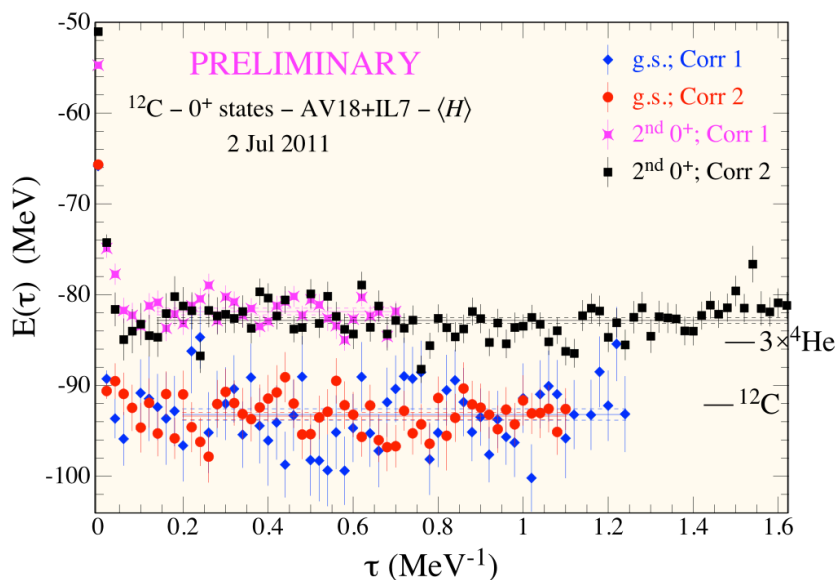


# Long-range correlations (clustering) in Carbon-12

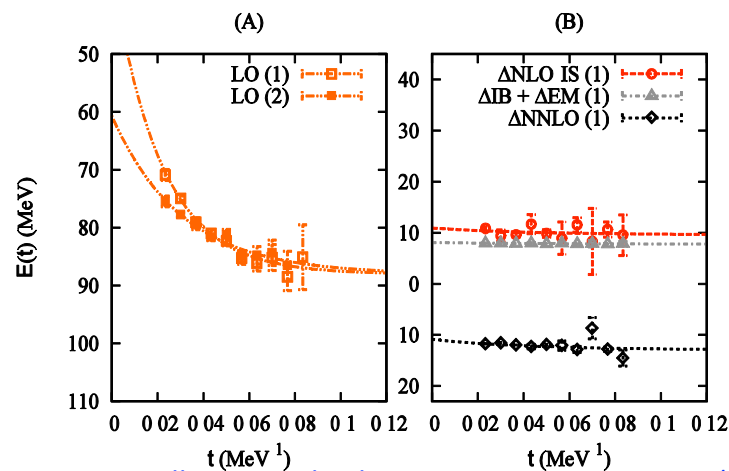
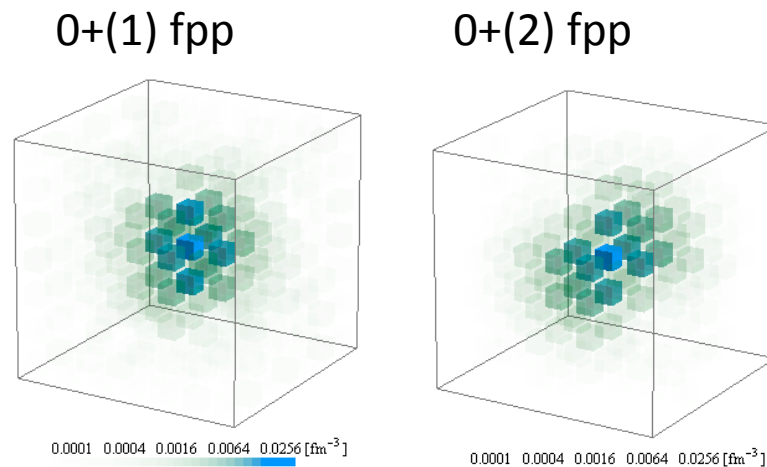
GFMC (Pieper et al.)



The ADLB (Asynchronous Dynamic Load-Balancing) library & GFMC. GFMC energy 93.5(6) MeV; expt. 92.16 MeV. GFMC pp radius 2.35 fm; expt. 2.33 fm.



Lattice EFT (Lee, Epelbaum, Meissner,...)



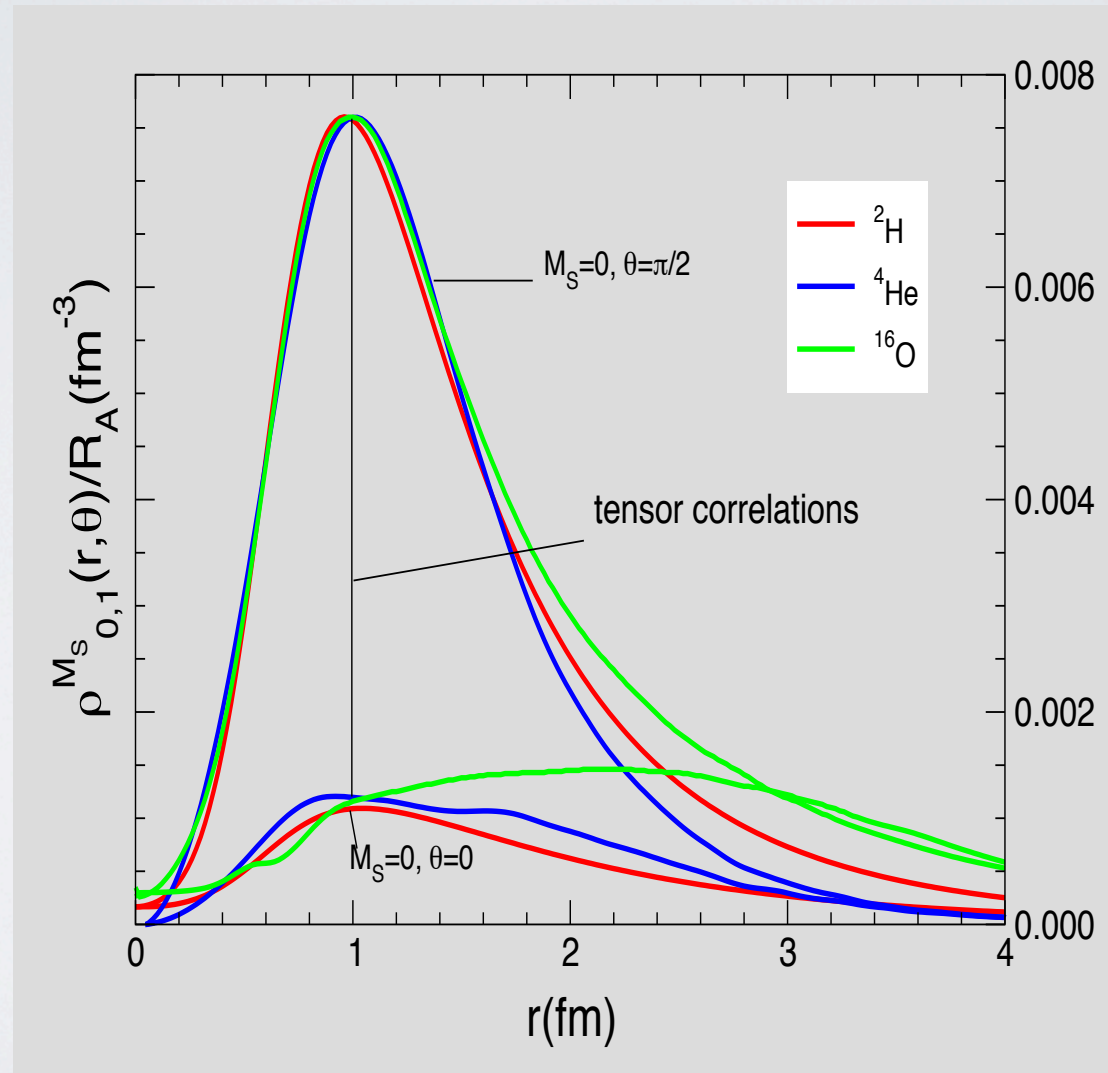
Epelbaum et al., Phys. Rev. Lett. 106, 192501 (2011)

# Shorter-range correlations

Tensor Force: higher momentum transfer:

Two-Nucleon Distribution Functions

$T=0$   $S=1$   
'deuteron'  
pairs



Expectation of short-range operator in  $T=0$ ,  
 $S=1$  pairs: scaling with nucleus

Scaling

	$R_A$	$\langle v^\pi \rangle_A / \langle v^\pi \rangle_d$	$\sigma_A^\pi / \sigma_d^\pi$	$\sigma_A^\gamma / \sigma_d^\gamma$
${}^3\text{He}$	2.0	2.1	2.4(1)	$\simeq 2$
${}^4\text{He}$	4.7	5.1	4.3(6)	$\simeq 4$
${}^6\text{Li}$	6.3	6.3		
${}^7\text{Li}$	7.2	7.8		$\simeq 6.5(5)$



# Two-nucleon momentum distributions:

$$\text{Total } P=0; \quad p_1 \sim -p_2$$

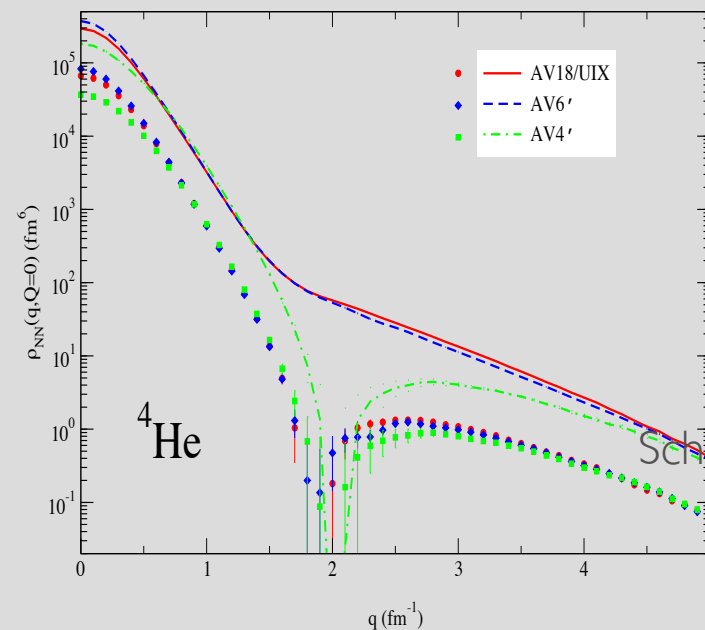
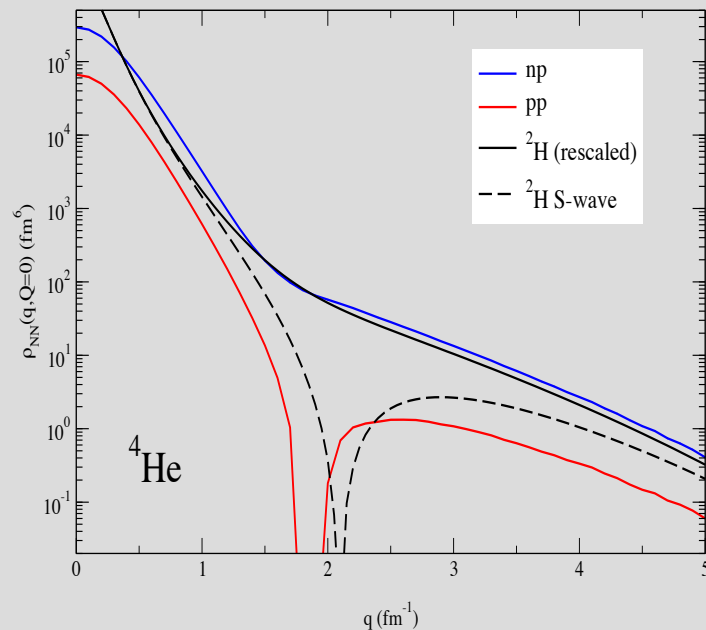
$$\rho^{NN}(\mathbf{q}, \mathbf{Q}) = \frac{1}{2J+1} \sum_{M_J} \langle \psi_{JM_J} | \sum_{i<j} P_{ij}^{NN}(\mathbf{q}, \mathbf{Q}) | \psi_{JM_J} \rangle$$

where  $\mathbf{q}$  and  $\mathbf{Q}$  are respectively the relative and total momenta of the  $NN$  pair, and

$$P_{ij}^{NN}(\mathbf{q}, \mathbf{Q}) \equiv \delta(\mathbf{k}_{ij} - \mathbf{q}) \delta(\mathbf{K}_{ij} - \mathbf{Q}) P_{NN}(ij)$$

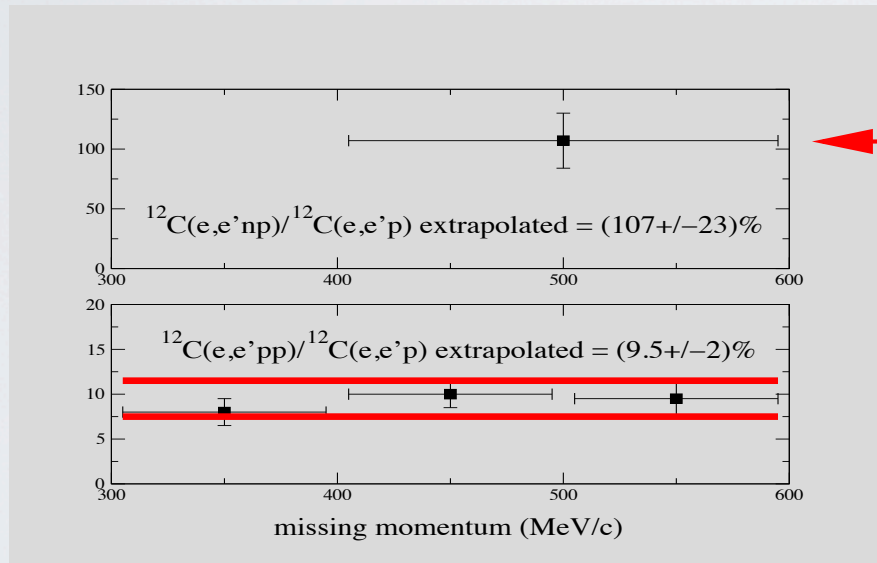
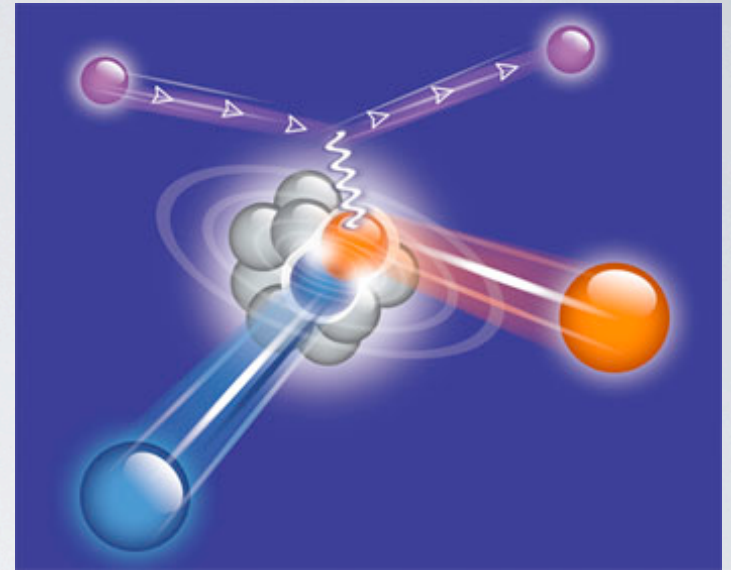
Dominated by  $T=0, S=1$  (np) and  $T=1, S=0$  (pp) pairs

Back-to-back nucleons (total pair momentum vanishes)



Schavilla, PRL 2007

np correlations correlations  
 in back-to-back nucleon  
 knockout:  
 observed at BNL, Jlab



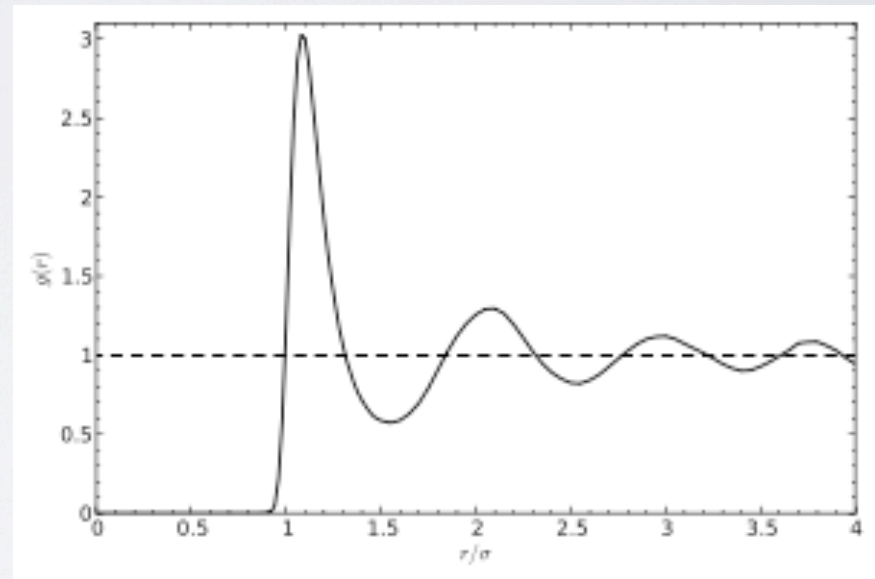
Analysis of BNL data:

$$\frac{P_{pn}}{P_{pX}} = 92^{+8}_{-18}\%$$

<sup>a</sup> Shneor *et al.*, PRL**99**, 072501 (2007); <sup>b</sup> Subedi *et al.*, Science **320**, 1476 (2008); <sup>c</sup> Piassetzky *et al.*, PRL**97**, 162504 (2006); <sup>d</sup> Ashery *et al.*, PRL**47**, 895 (1981)

Typical method to study correlations,  
inclusive scattering

Liquid Helium, ... neutron scattering



pair distribution function  
Lennard-Jones model fluid

# Electron and Neutrino Scattering

$$S(q, \omega) = \sum_f \langle 0 | O^\dagger(q) | f \rangle \langle f | O(q) | 0 \rangle \delta(\omega - (E_f - E_0))$$

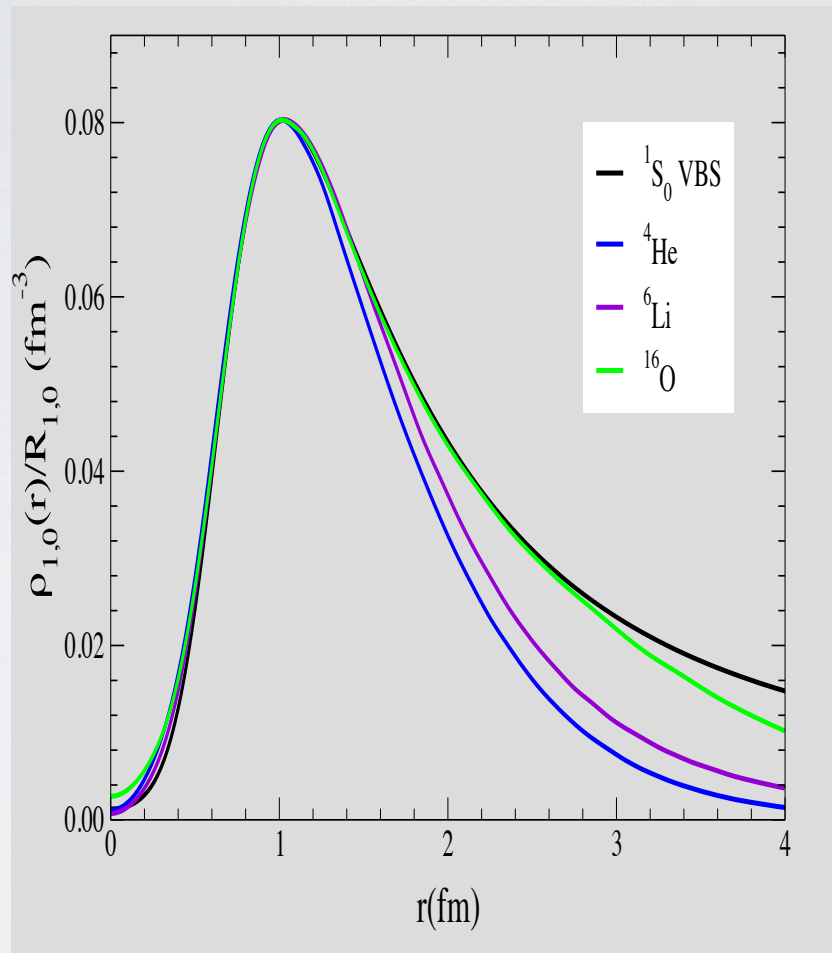
Longitudinal (Charge) scattering

$$O(q) = \sum_i P_p(i) \exp[i\mathbf{q} \cdot \mathbf{r}_i]$$

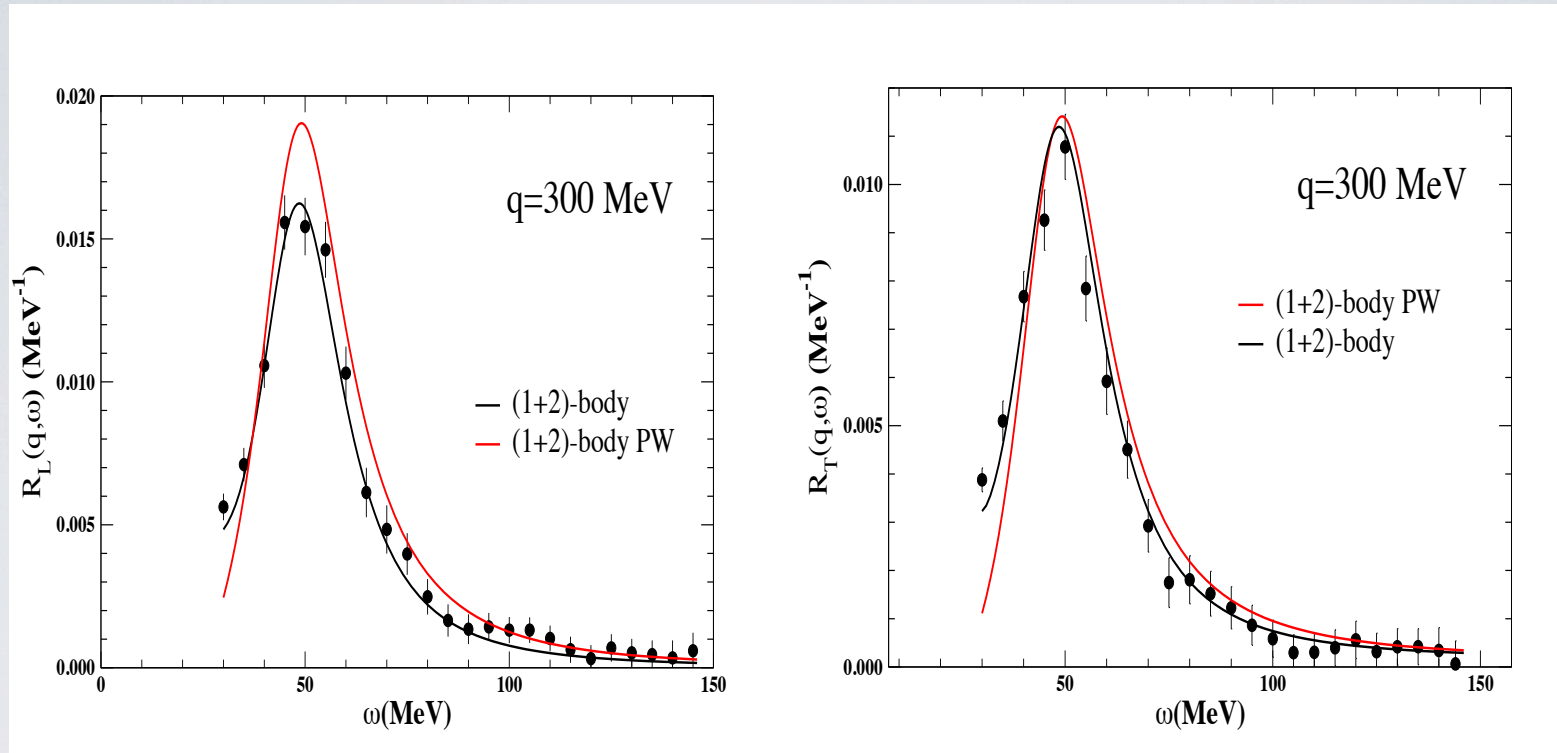
Transverse (Current Scattering)

$$O(q) = \sum_i \mu(i) \exp[i\mathbf{q} \cdot \mathbf{r}_i] + P_p(i) \mathbf{p}_i \exp[i\mathbf{q} \cdot \mathbf{r}_i] + \sum_{i < j} \mathbf{j}_{ij}(\mathbf{q})$$

# Inclusive Electron Scattering and spin 0, $T=1$ pairs (pp, nn)



# Inclusive electron scattering on the deuteron



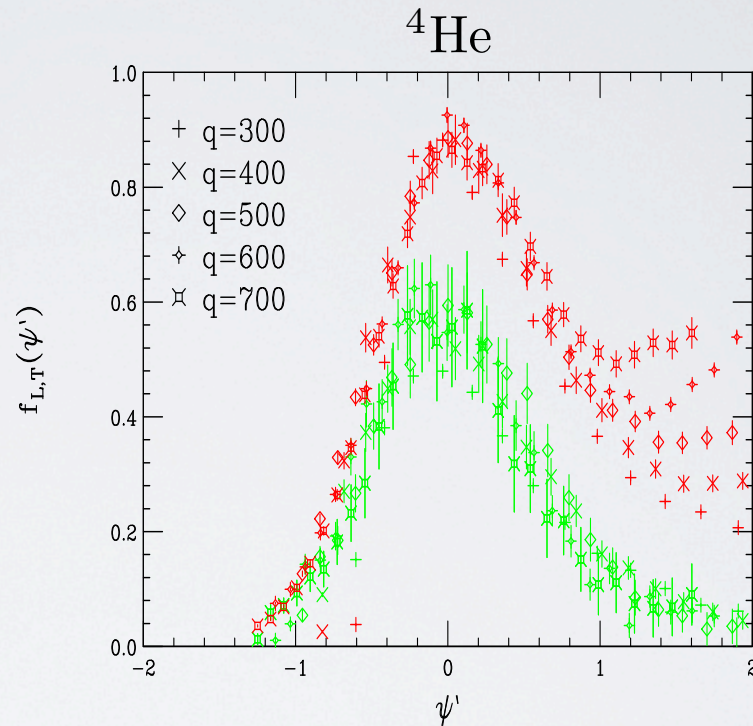
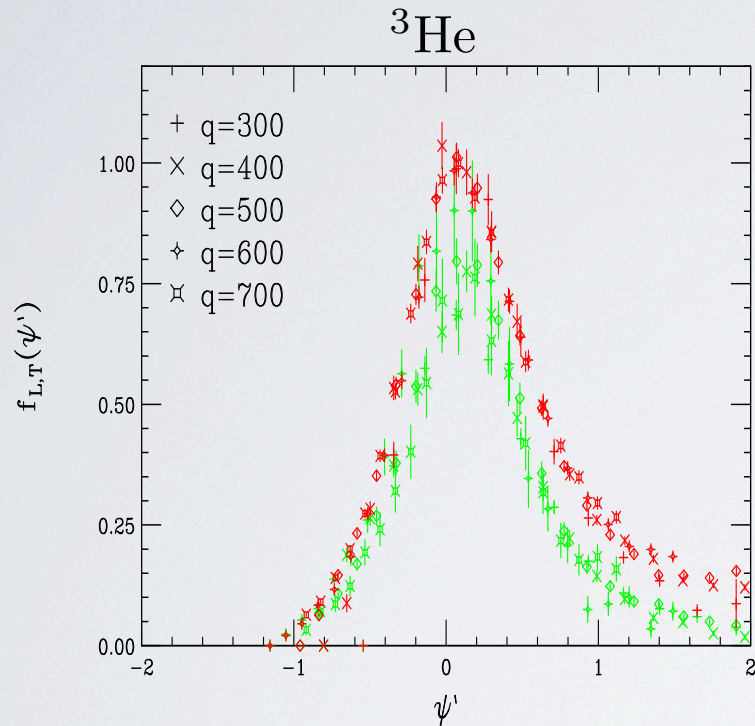
$$R_{\alpha}(q, \omega) = \sum_{f \neq 0} \delta(\omega + E_0 - E_f) | \langle f | O_{\alpha}(\mathbf{q}) | 0 \rangle |^2 \quad \alpha = L, T$$

require knowledge of continuum states: hard to calculate for  $A \geq 3$

- Sum rules: integral properties of response functions
- Integral transform techniques

$$E(q, \tau) = \int_0^{\infty} d\omega K(\tau, \omega) R(q, \omega)$$

# Inclusive Electron Scattering



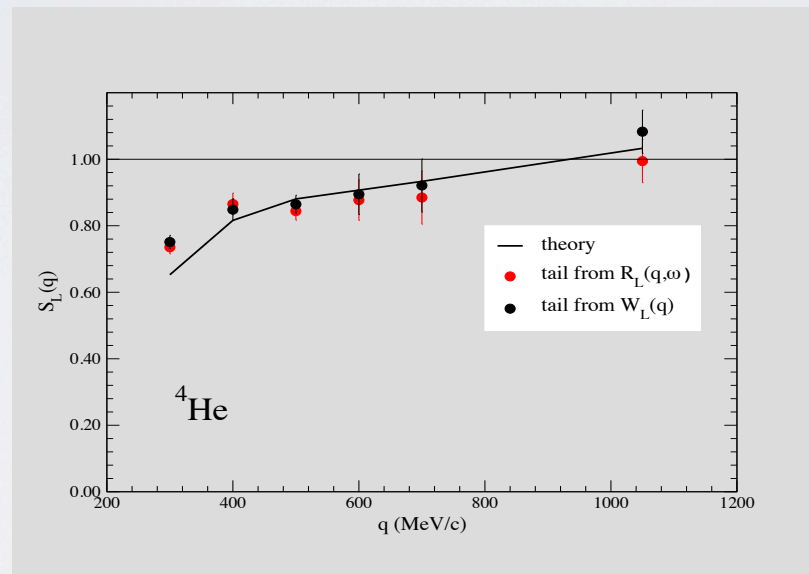
Sick and Donnelly, 1999  
Carlson, Schiavilla, Sick, 2002

Longitudinal sensitive to pp correlations  
Transverse sensitive to np correlations

# Measuring charge-charge ('pp') correlations

## The $^4\text{He}$ Coulomb Sum Rule

- RC/MEC (small) contributions to  $S_L(q)$  tend to cancel out
- Theory and experiment in agreement when using **free**  $G_{Ep}$

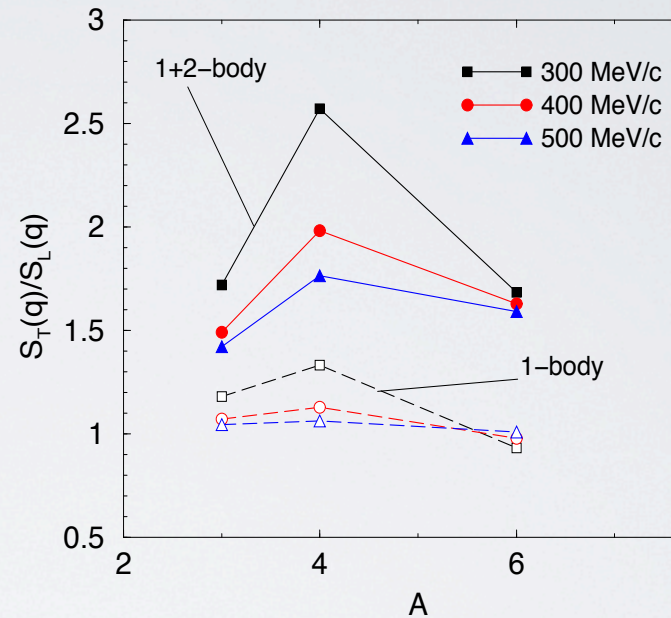
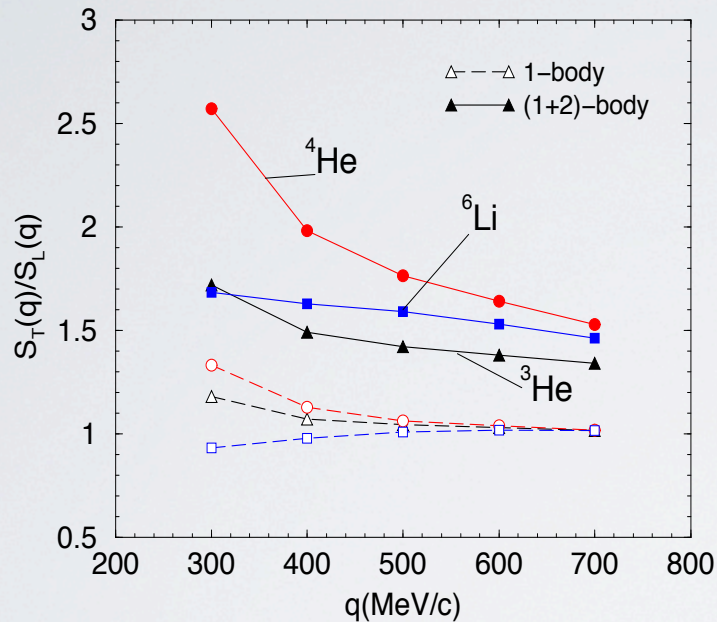


$$W_L(q) = \frac{1}{Z} \int_{\omega_{\text{th}}^+}^{\infty} d\omega \omega \frac{R_L(q, \omega)}{G_{Ep}^2(q, \omega)} = \frac{1}{2Z} \langle 0 | \left[ \rho^\dagger(\mathbf{q}), \left[ H, \rho(\mathbf{q}) \right] \right] | 0 \rangle$$



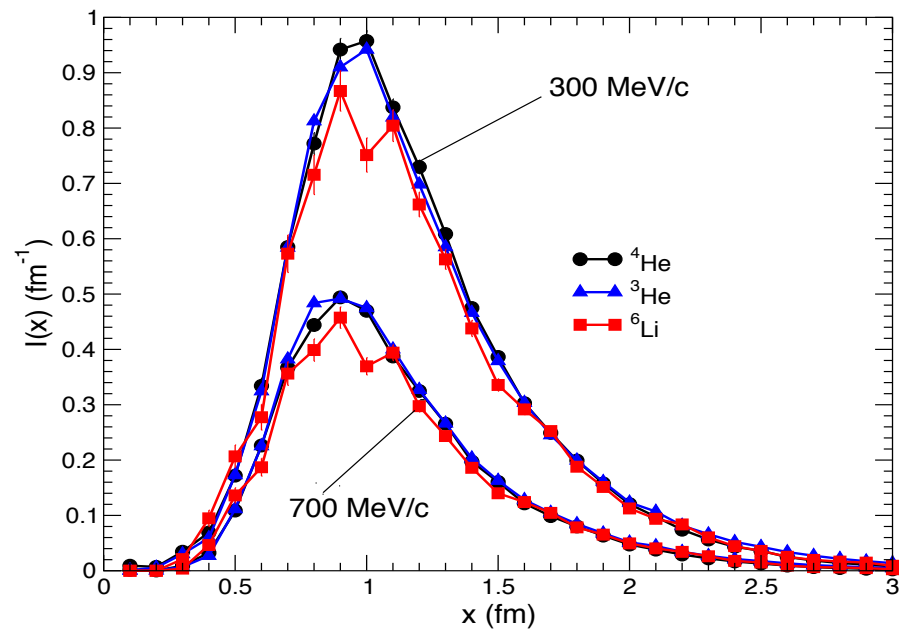
# Transverse Channel: 1 + 2-nucleon currents

## Excess Transverse Strength



- How much of the excess transverse strength  $\Delta S_T = S_T - S_T^{1b}$  in the quasi-elastic peak region?
- Can we understand the  $A$ -dependence of  $\Delta S_T$ ?

## A-Scaling Property



$$\int_0^{\infty} dx I(x) = \Delta S_T \propto R_A / (Z \mu_p^2 + N \mu_n^2)$$

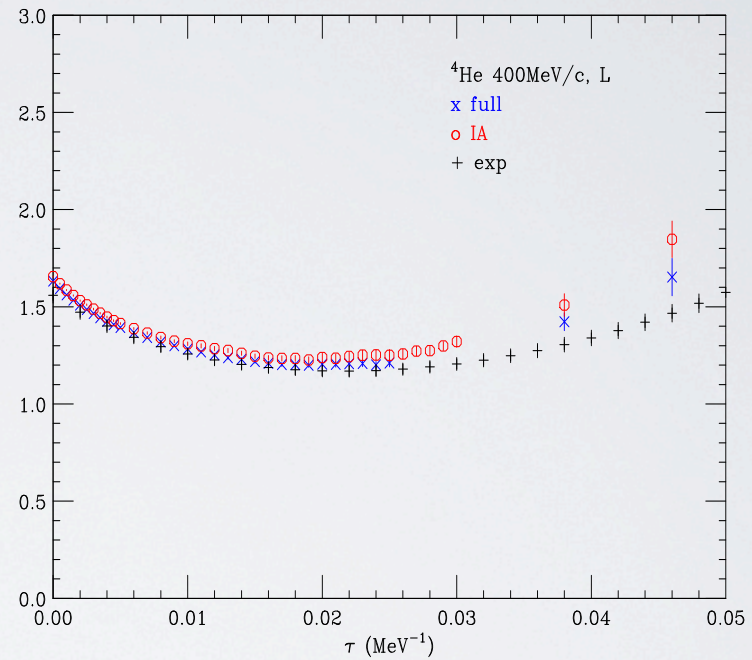
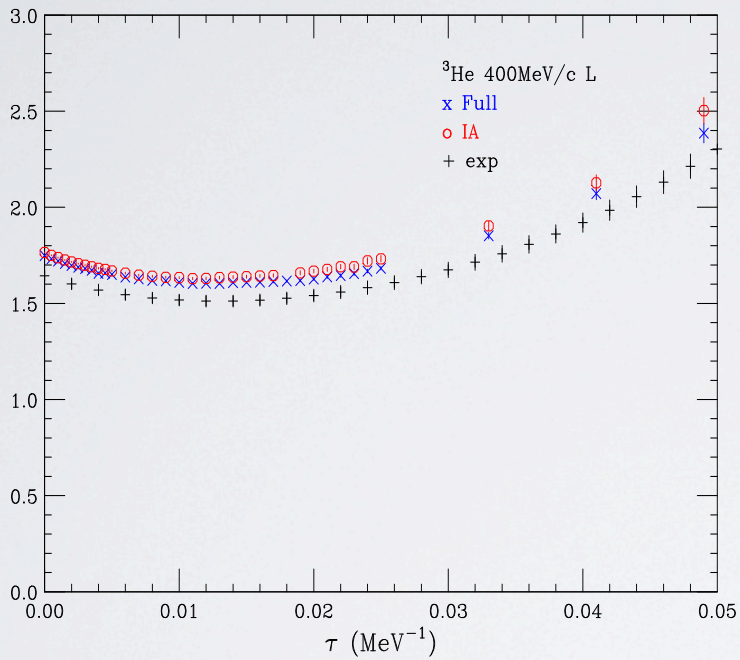
## Euclidean Response Functions

Carlson and Schiavilla (1992,1994)

$$\begin{aligned}\tilde{E}_\alpha(q, \tau) &= \int_{\omega_{\text{th}}^+}^{\infty} d\omega e^{-\tau(\omega - E_0)} \frac{R_\alpha(q, \omega)}{G_{Ep}^2(q, \omega)} \\ &= \langle 0 | O_\alpha^\dagger(\mathbf{q}) e^{-\tau(H - E_0)} O_\alpha(\mathbf{q}) | 0 \rangle - (\text{elastic term})\end{aligned}$$

- $e^{-\tau(H - E_0)}$  evaluated stochastically with QMC
- No approximations made, exact
- At  $\tau = 0$ ,  $\tilde{E}_\alpha(q; 0) \propto S_\alpha(q)$ ; as  $\tau$  increases,  $\tilde{E}_\alpha(q; \tau)$  is more and more sensitive to strength in quasi-elastic region
- Inversion of  $\tilde{E}_\alpha(q; \tau)$  is a numerically ill-posed problem; Laplace-transform data instead

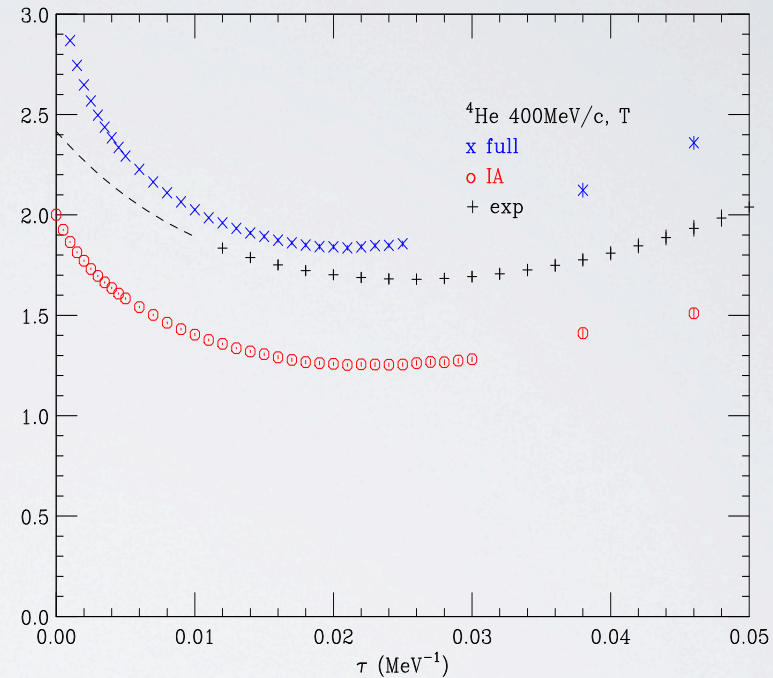
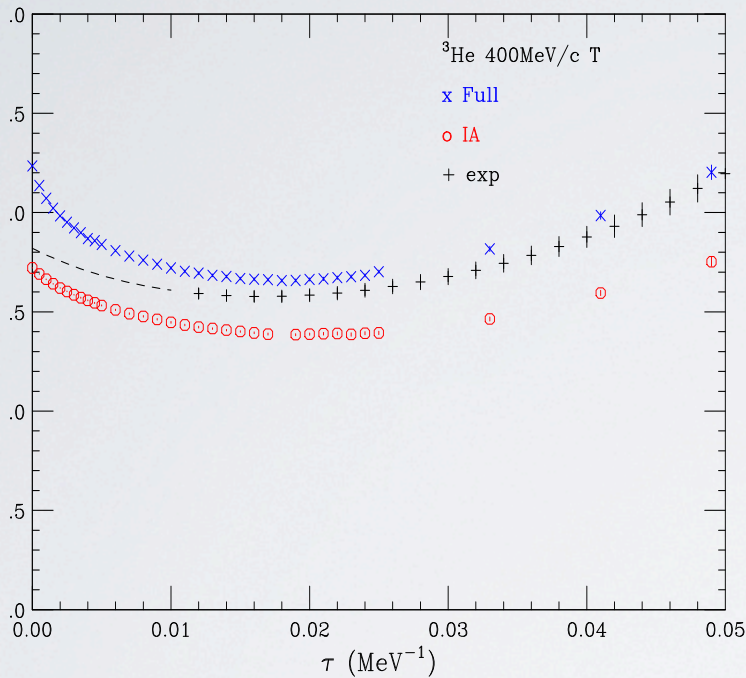
## $^3\text{He}$ and $^4\text{He}$ Longitudinal Euclidean Response Functions



$$E_{\alpha}(q, \tau) = \exp \left[ \tau q^2 / (2 m) \right] \tilde{E}_{\alpha}(q, \tau)$$

and  $E_L(q, \tau) \rightarrow Z$  for a collection of protons initially at rest

## $^3\text{He}$ and $^4\text{He}$ Transverse Euclidean Response Functions

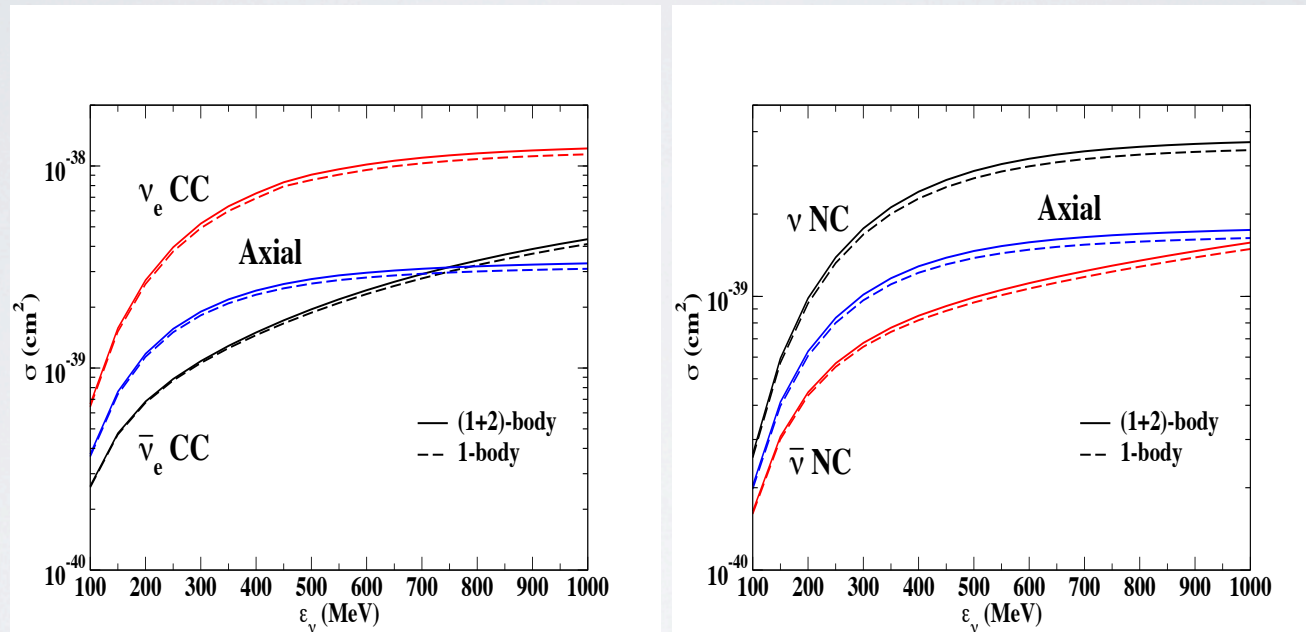


- Excess strength in quasielastic region ( $\tau > 0.01 \text{ MeV}^{-1}$ )
- Larger in  $A = 4$  than in  $A = 3$ , as already inferred from  $S_T$

What happens in  $\sim$ GeV neutrino scattering?

$\nu$ -Deuteron Scattering up to GeV Energy

Shen *et al.* (2012)



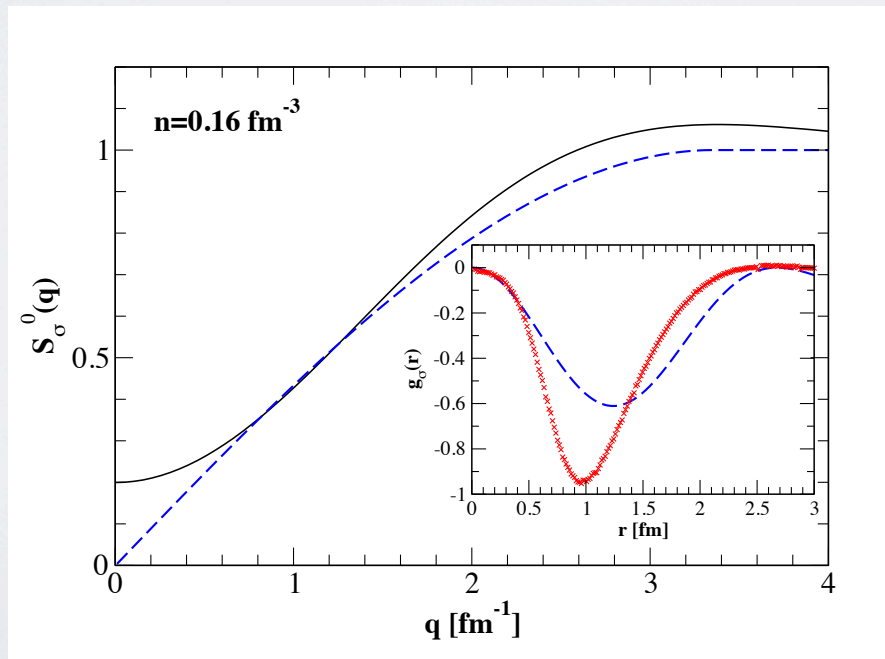
Potentially important in Mini-Boone, LBNE,...

# Spin Response in Neutron Matter

requires tensor and/or spin-orbit interactions at  $q=0$

$$\begin{aligned}
 S_\sigma(\omega, \mathbf{q}) &= \frac{4}{3n} \frac{1}{2\pi} \int_{-\infty}^{\infty} dt e^{i\omega t} \langle \mathbf{s}(t, \mathbf{q}) \cdot \mathbf{s}(0, -\mathbf{q}) \rangle \\
 &= \frac{4}{3n} \sum_f \langle 0 | s(\mathbf{q}) | f \rangle \cdot \langle f | s(-\mathbf{q}) | 0 \rangle \delta(\omega - (E_f - E_0))
 \end{aligned}$$

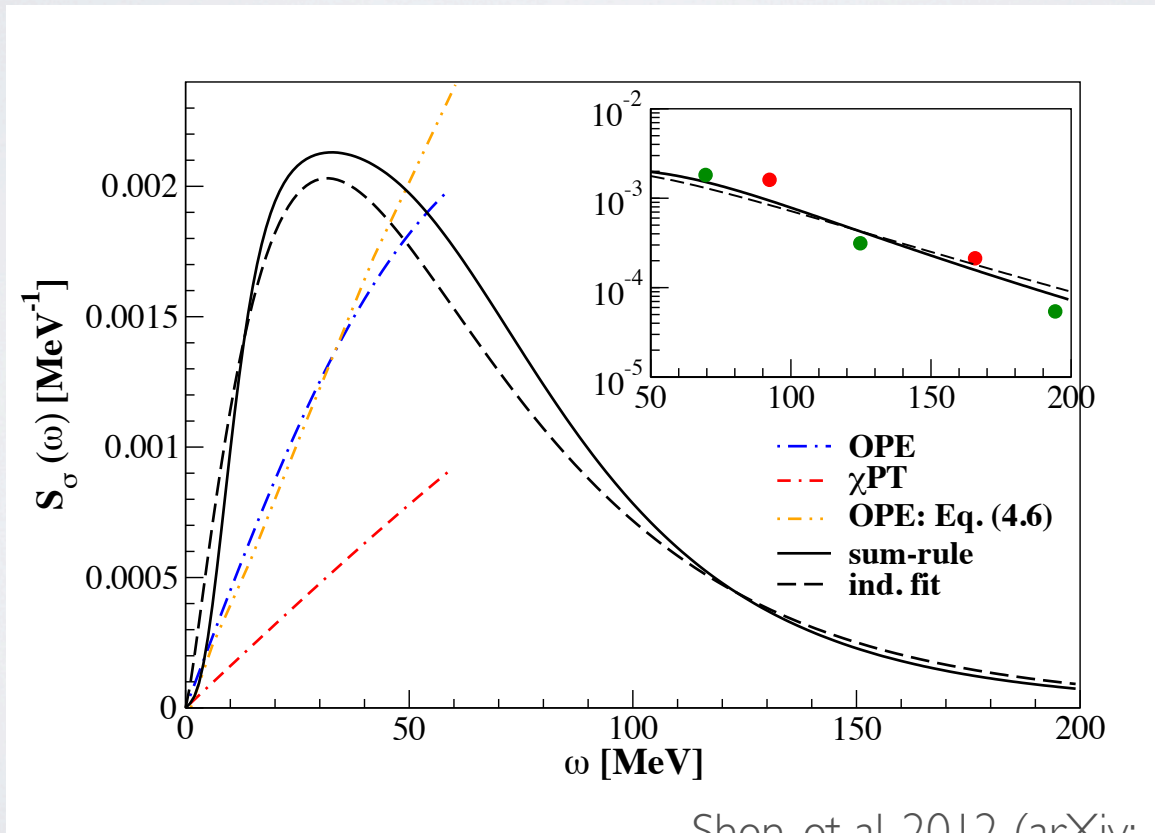
$$Q = \frac{C_A^2 G_F^2 n}{20\pi^3} \int_0^\infty d\omega \omega^6 e^{-\omega/T} S_\sigma(\omega),$$



Sum Rules (inc. spin susceptibility) give  
measure of overall strength, position and width of peak

Table I: AFDMC results for the sum-rules

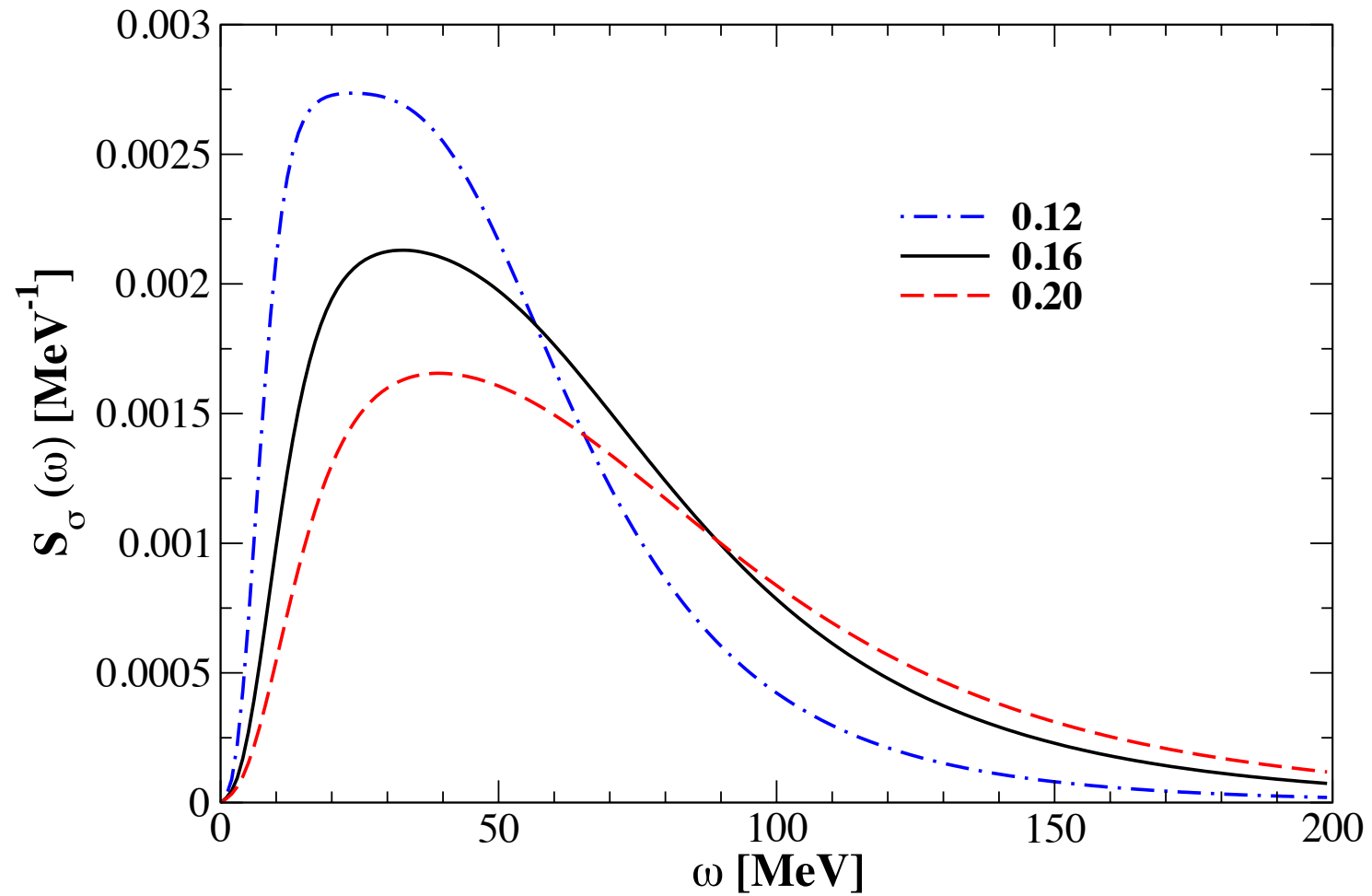
Density ( $\text{fm}^{-3}$ )	$S_{\sigma}^{-1}$ ( $\text{MeV}^{-1}$ )	$S_{\sigma}^0$	$S_{\sigma}^{+1}$ ( $\text{MeV}$ )	$\bar{\omega}_0$ ( $\text{MeV}$ )	$\bar{\omega}_1$ ( $\text{MeV}$ )
$n = 0.12$	0.0057(9)	0.20(1)	8(1)	35(9)	40(8)
$n = 0.16$	0.0044(7)	0.20(1)	11(1)	46(11)	55(8)
$n = 0.20$	0.0038(6)	0.18(1)	14(1)	47(12)	78(10)



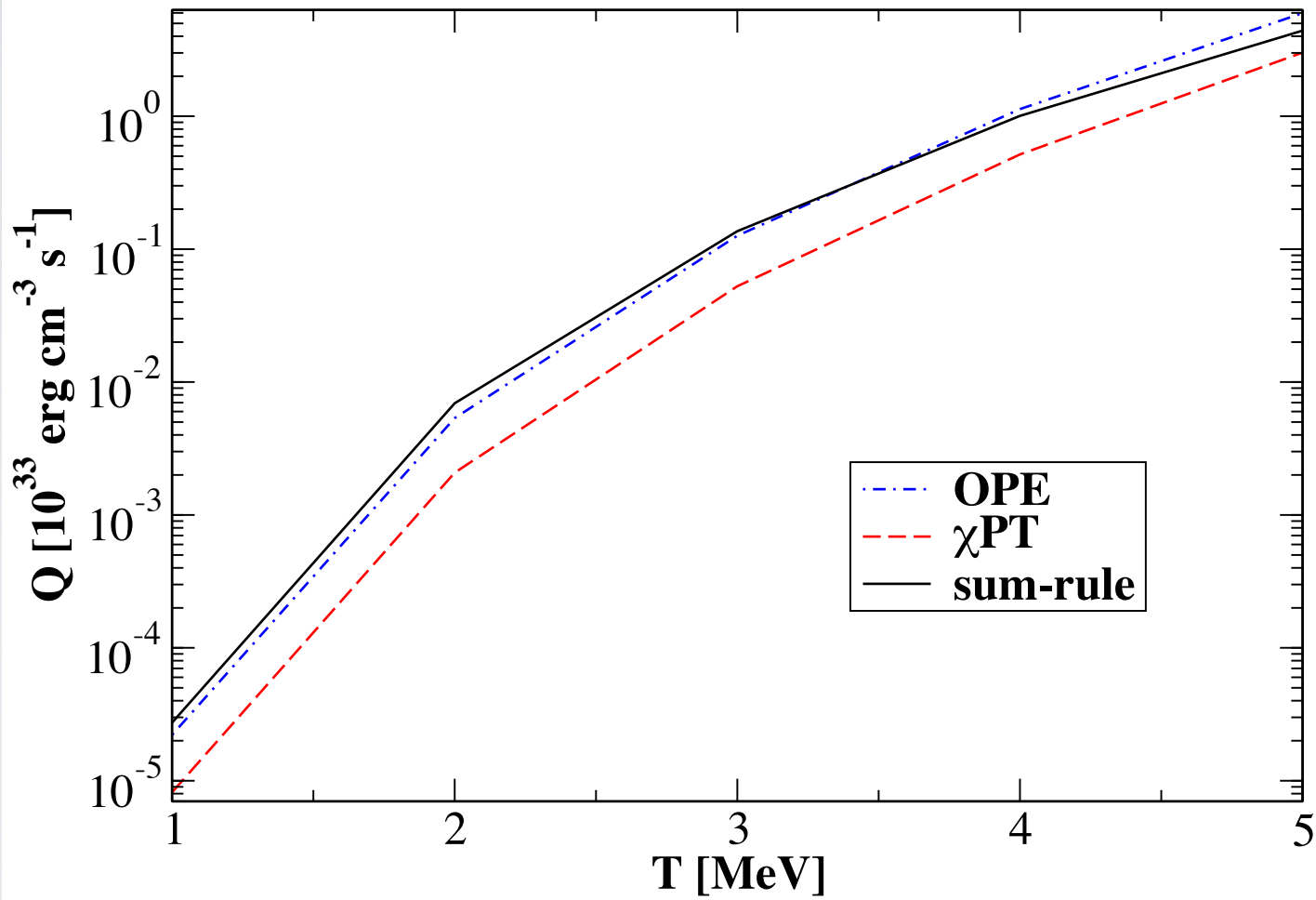
Shen, et al, 2012 (arXiv:



# Density dependence of response



# Energy loss rate $Q$



## Conclusions:

Tensor (and other) correlations critical in nuclear physics

Obvious impacts in the deuteron ( $Q, T_{20}, \dots$ ),  
but for larger nuclei, impact not often seen in  
low-energy observables (spectra, etc.)

Tensor correlations more important in spin observables

More obvious impact at higher momenta:

np vs. pp back-to-back in electron scattering  
inclusive electron and neutrino scattering

Can impact astrophysically relevant behavior:

neutrino propagation,  $3P_2$  pairing, ...