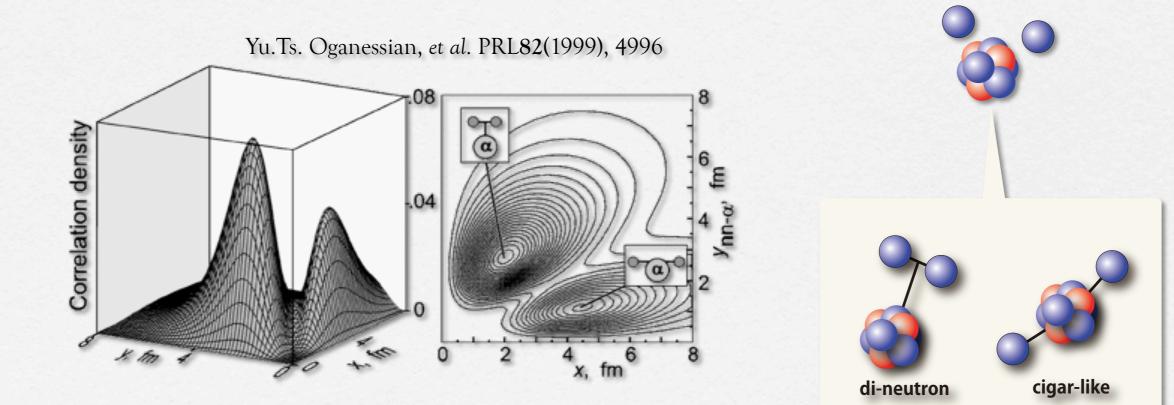
Description of three-body scattering states using complex scaling method and its application to Coulomb breakup reactions of two-neutron halo nuclei

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Two-neutron halo nuclei

- Two-neutron halo nuclei have exotic structure in which the weakly-bound neutrons are spread out beyond the core nucleus.
- From the theoretical calculation using the core+n+n three-body model, the exotic correlation between two halo neutrons, the so-called dineutron, plays a significant role in reproducing the observed two-neutron separation energies and the large matter radii of ⁶He and ¹¹Li.

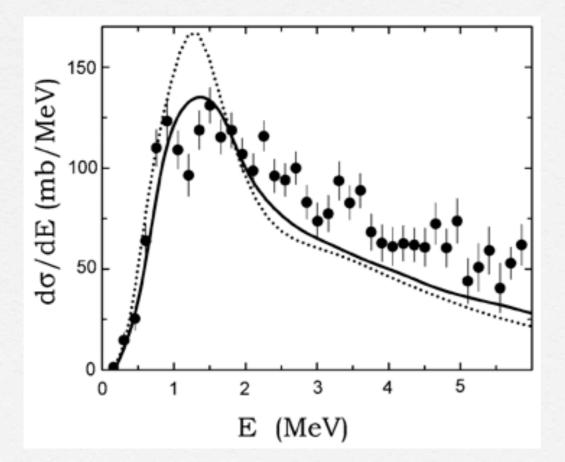


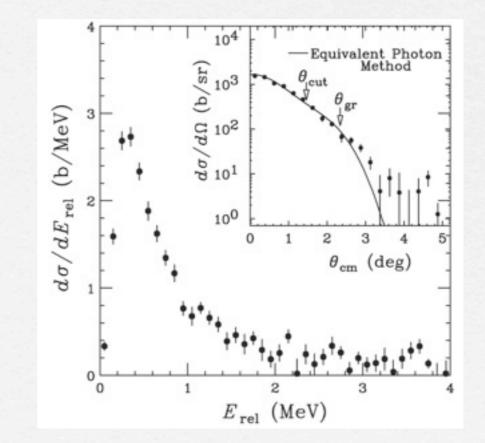
Coulomb breakup reactions of two-neutron halo nuclei

- Coulomb breakup reactions have been performed to investigate the electric response of weakly-bound halo neutrons.
- The observed cross sections show the low-lying enhancements, and those enhancements have been expected to be responsible to the exotic structure of two-neutron halo nuclei.

⁶He breakup: T. Aumann et al., PRC 59, 1252 (1999).

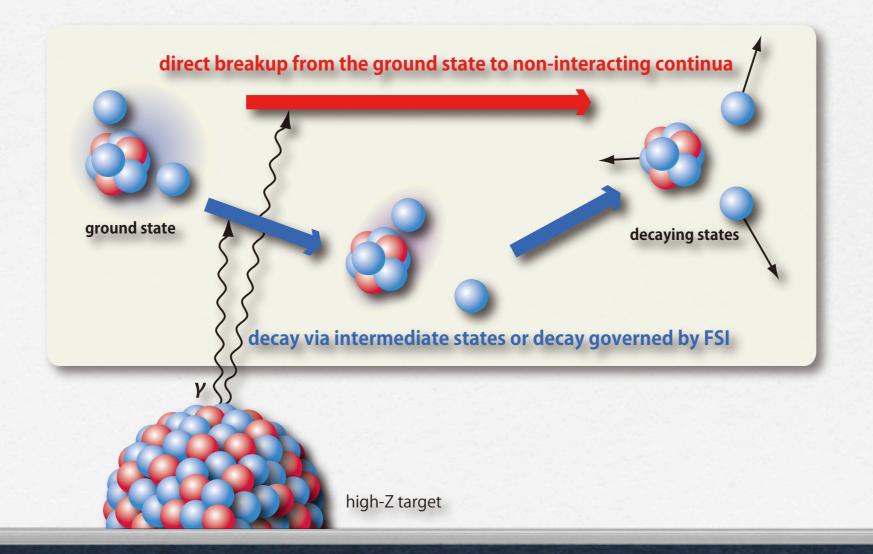
¹¹Li breakup: T. Nakamura et al., PRL 96, 252502 (2006).





To investigate the Coulomb breakups of two-neutron halo nuclei

- Two-neutron halo nuclei are Borromean systems in which any binary subsystems have no bound states, and hence, they are broken up to the core+n+n three-body scattering states.
 - To understand the properties of two-neutron halos from the Coulomb breakups, it is necessary to describe the core+n+n three-body scattering states accurately and to clarify the breakup mechanisms.



In this talk...

- We investigate the Coulomb breakup mechanisms of two-neutron halo nuclei, ⁶He and ¹¹Li, by using the core+n+n three-body models.
 - To describe the final scattering states of the core+n+n system, we use the complexscaled solutions of the Lippmann-Schwinger equation (CSLS).
 - Using CSLS, we discuss the mechanism of the Coulomb breakup reactions of twoneutron halo nuclei.
 - What kinds of processes dominate the Coulomb breakup reactions?
 - □ simultaneous three-body breakup or sequential breakup?
 - What kinds of correlations play roles in the reactions?
 - How the differences in the configurations reflects on the cross section?

 \Box ⁶He : $(p_{3/2})^2$ configuration is dominant

 \Box ¹¹Li: $(p_{1/2})^2$ + $(s_{1/2})^2$ due to the breaking of N=8 magicity

CSM is a powerful tool to investigate the many-body resonances on the same footing as the bound-state cases.

□ In CSM, the relative coordinates and momenta are transformed as follows.

$$U(\theta): \mathbf{r} \to \mathbf{r}e^{i\theta}, \quad \mathbf{k} \to \mathbf{k}e^{-i\theta}$$

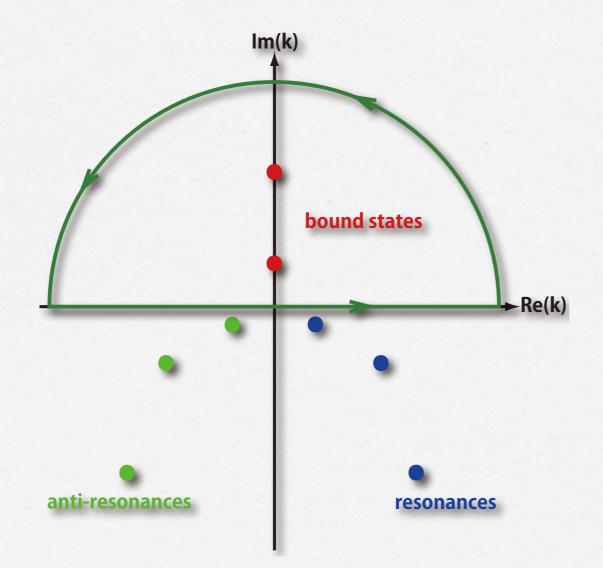
Applying the above transformation to the Schroedinger equation, we obtain the complex-scaled Schroedinger equation as follows.

$$\hat{H}\chi(\mathbf{r}) = E\chi(\mathbf{r}) \to \hat{H}^{\theta}\chi^{\theta}(\mathbf{r}) = E^{\theta}\chi^{\theta}(\mathbf{r})$$

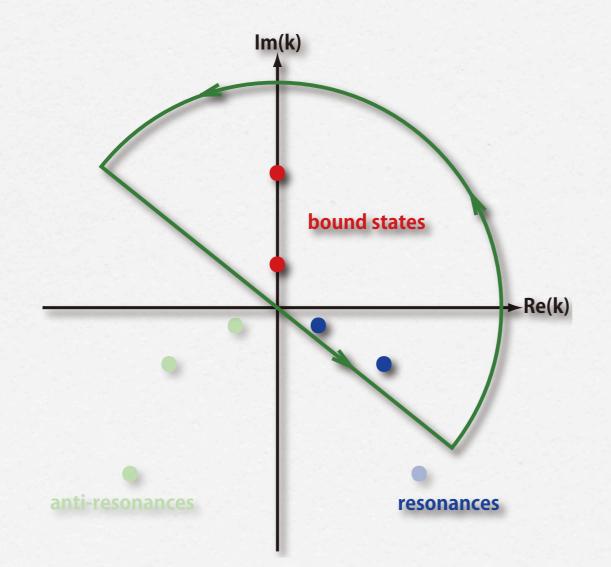
where the complex-scaled wave functions and Hamiltonian are given as

$$\chi^{\theta}(\mathbf{r}) = U(\theta)\chi(\mathbf{r}) = e^{\frac{3}{2}i\theta}\chi(\mathbf{r}e^{i\theta})$$
$$\hat{H}^{\theta} = U(\theta)\hat{H}U^{-1}(\theta)$$

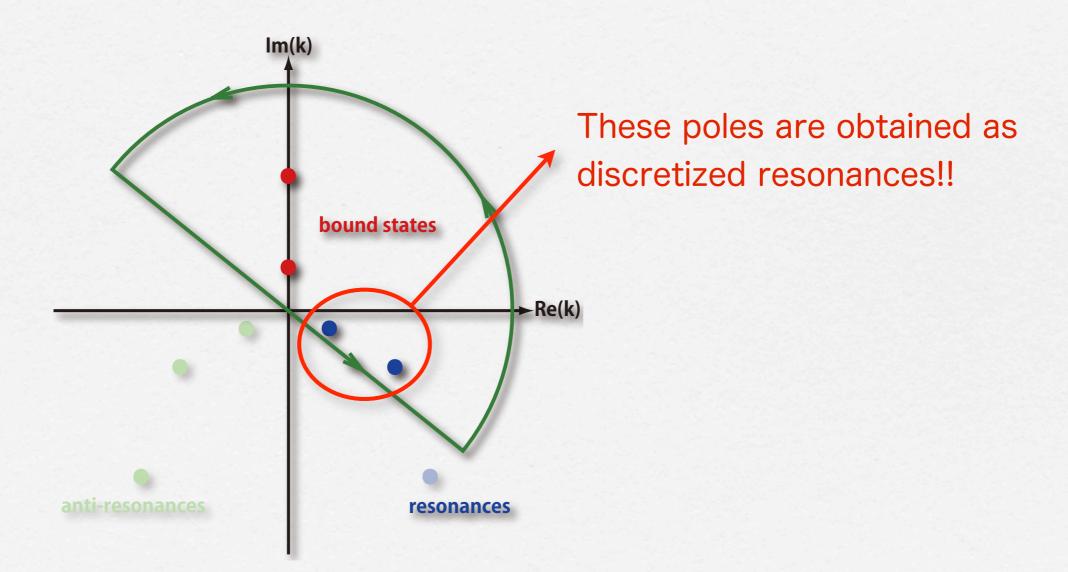
- Under the transformation in CSM, the resonance poles are obtained as the discretized states as well as the bound states.
 - By rotating the contour of the integral pass in the momentum plane with outgoing boundary condition, the resonance poles in the S-matrix are found as the residues.



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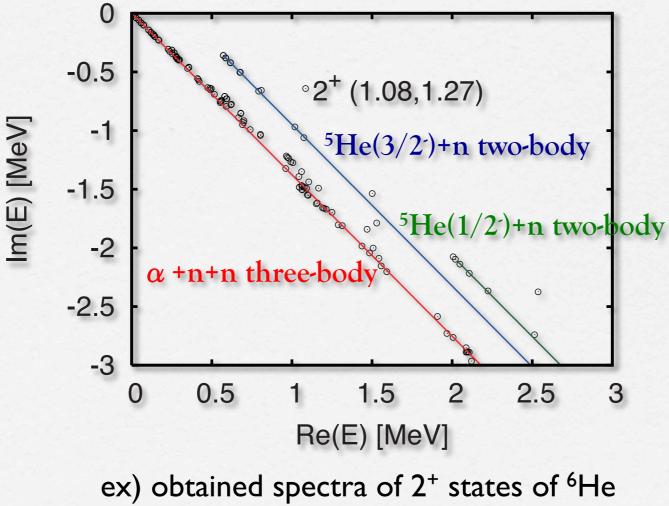


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Energy spectra in CSM

- In CSM, the energy eigenvalues are obtained as complex numbers, and their imaginary parts represent the outgoing boundary conditions.
 - \Box The resonance has the energy of E = E_r-i $\Gamma/2$, where Er and Γ are the resonance energy and decay width, respectively.
 - The continuum states are classified into several families corresponding to decay channels.



- The behaviors of the energy eigenvalues in CSM is useful to describe the scattering of many-body systems since outgoing boundary conditions are imposed.
- Using this characteristic, we describe the many-body scattering states by using the Lippmann-Schwinger formalism.
 - □ We start with the formal solution of the Lippmann-Schwinger equation.

$$\Psi^{(\pm)} = \Phi_0 + \lim_{\varepsilon \to 0} \frac{1}{E - \hat{H} \pm i\varepsilon} \hat{V} \Phi_0$$

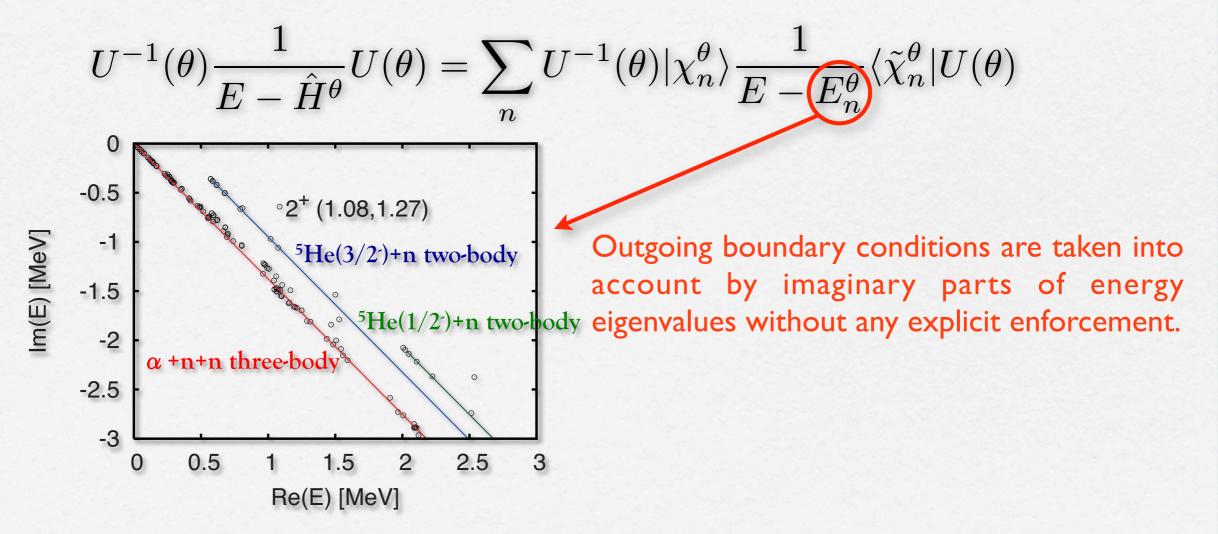
To take into account the outgoing boundary condition in the Green's function, we apply CSM to the Green's function.

$$\lim_{\varepsilon \to 0} \frac{1}{E - \hat{H} + i\varepsilon} = U^{-1}(\theta) \frac{1}{E - \hat{H}^{\theta}} U(\theta)$$

- \Box We expand the complex-scaled Green's function with the complete set constructed with the solved eigenstates of complex-scaled Hamiltonian H^{θ}.
 - \Box Here, we solve the eigenvalue problem of H^{θ} by using the L²-type basis functions as similar to the bound-state cases.

$$U^{-1}(\theta)\frac{1}{E-\hat{H}^{\theta}}U(\theta) = \sum_{n} U^{-1}(\theta)|\chi_{n}^{\theta}\rangle \frac{1}{E-E_{n}^{\theta}}\langle \tilde{\chi}_{n}^{\theta}|U(\theta)$$

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Using the complex-scaled Green's function, we obtain the complex-scaled solutions of the Lippmann-Schwinger equation (CSLS).

$$\begin{split} |\Psi^{(+)}\rangle &= |\Phi_0\rangle + \sum_n U^{-1}(\theta)|\chi_n^{\theta}\rangle \frac{1}{E - E_n^{\theta}} \langle \tilde{\chi}_n^{\theta}|U(\theta)\hat{V}|\Phi_0\rangle \\ \langle \Psi^{(-)}| &= \langle \Phi_0| + \sum_n \langle \Phi_0|\hat{V}U^{-1}(\theta)|\chi_n^{\theta}\rangle \frac{1}{E - E_n^{\theta}} \langle \tilde{\chi}_n^{\theta}|U(\theta) \end{split}$$

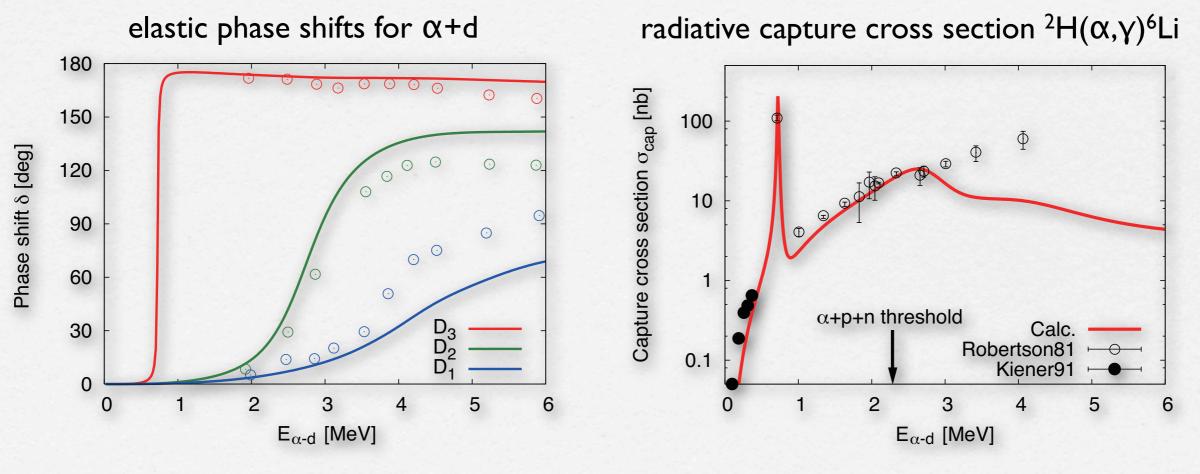
The advantages in CSLS is that we can solve the scattering problems of many-body systems

 \Box in similar way to the bound-state cases and

without explicit enforcement of outgoing boundary conditions.

Reliability of CSLS - three-body scattering problem -

- \Box Application of CSLS to the α +d scattering using α +p+n three-body model.
 - □ The observed elastic phase shift and the capture cross section are nicely reproduced.
 - □ CSLS is reliable method to describe the three-body scattering problem.



R.G.H. Robertson et al., PRL**47**(1981), 1867. J. Kiener et al., PRC**44**(1991), 2195. Core+n+n three-body models

The Hamiltonians are given as follows.

$$\hat{H} = \sum_{i=1}^{3} t_i - T_{\rm cm} + \sum_{i=1}^{2} V_{\rm core-n}(\mathbf{r}_i) + V_{n-n}$$

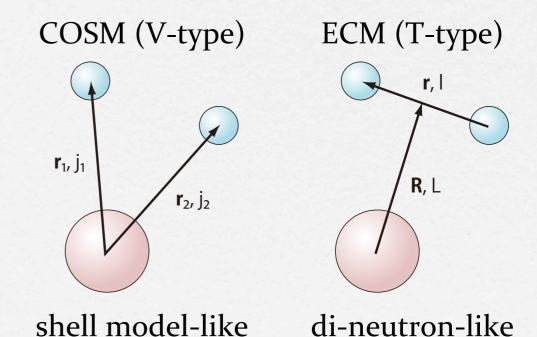
- We use the following interactions in the calculation
 - Core-n interaction
 - \square ⁶He: effective KKNN potential, which reproduce the scattering data of ⁴He+n.
 - ¹¹Li: folding G-matrix int. with the ⁹Li density + phenomenological LS force.
 - n-n interaction
 - ⁶He: Minnesota force
 - □ ¹¹Li:AV8' force
- The wave functions are expanded with the Gaussian basis functions.

T. Myo, et al, PRC**63**, 054313 (2001). T. Myo et al., PRC**76**, 024305 (2007).

H. Kanada, T. Kaneko, S. Nagata, and M. Nomoto, PTP**61**, 1327 (1979).

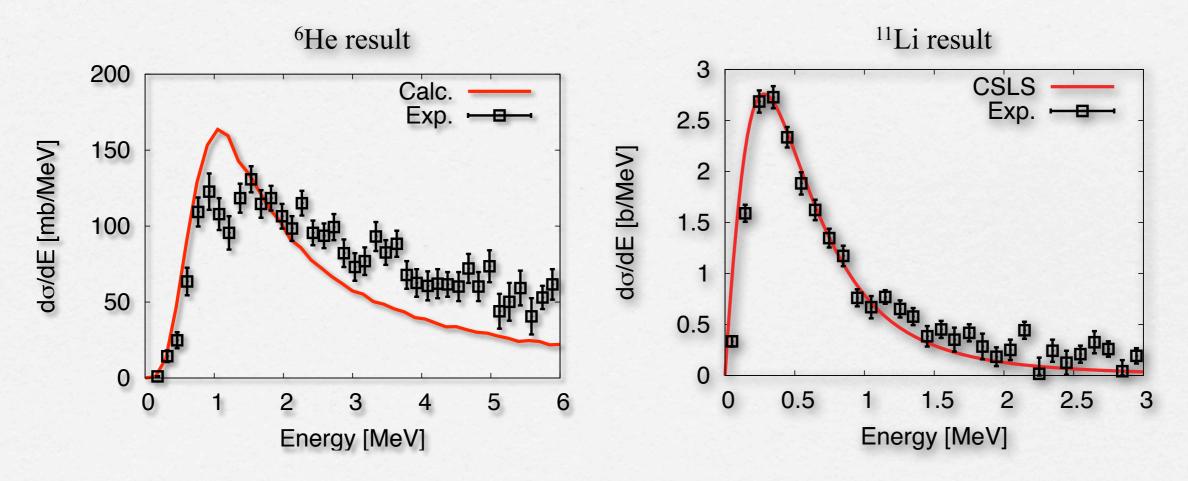
F. Kanada, I. Kaneko, S. Ivagata, and M. Ivomoto, FIFOI, 1527 (1975

E. Hiyama et al., Prog. Part. Nucl. Phys. **51**, 223 (2003).



Coulomb breakup cross sections of ⁶He and ¹¹Li

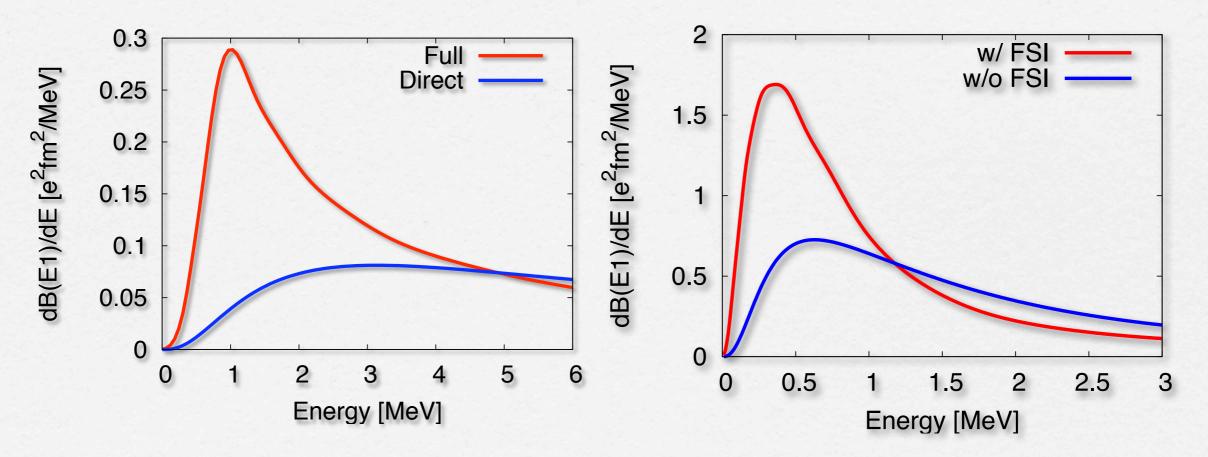
We calculate the Coulomb breakup cross sections of ⁶He and ¹¹Li, and compare the calculated results with experimental data.



- Our results well reproduce the observed data.
- \Box We next investigate the breakup mechanism of ⁶He and ¹¹Li.

Effects of final-state interactions on the cross sections

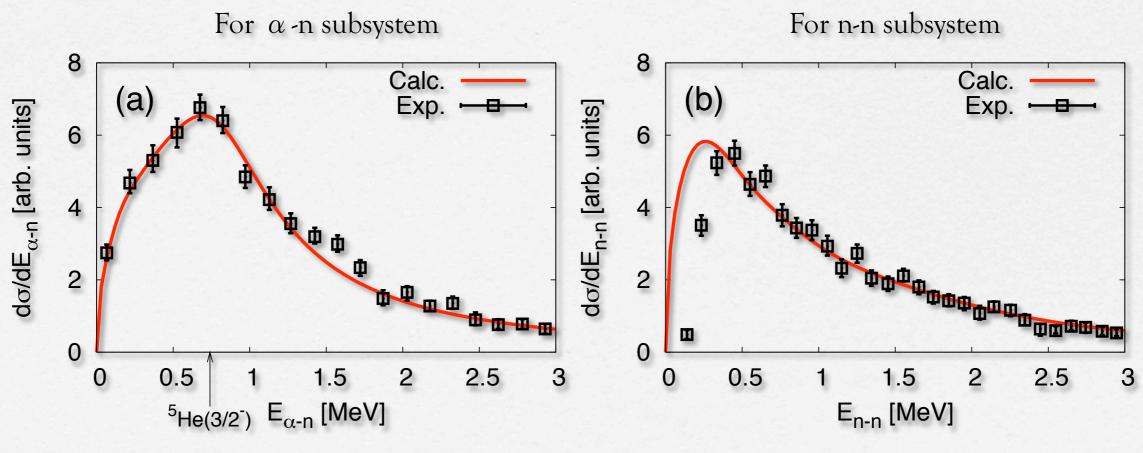
- □ To clarify the breakup mechanism, we investigate the effects of FSI on the cross sections.
 - When the FSI dominate the cross sections, the sequential processes might play roles in the Coulomb breakup reactions.



In both cases of ⁶He and ¹¹Li, FSI have significant effects on the cross sections.
What kinds of FSI does contribute to the breakup cross sections?

Invariant mass spectra for binary subsystems

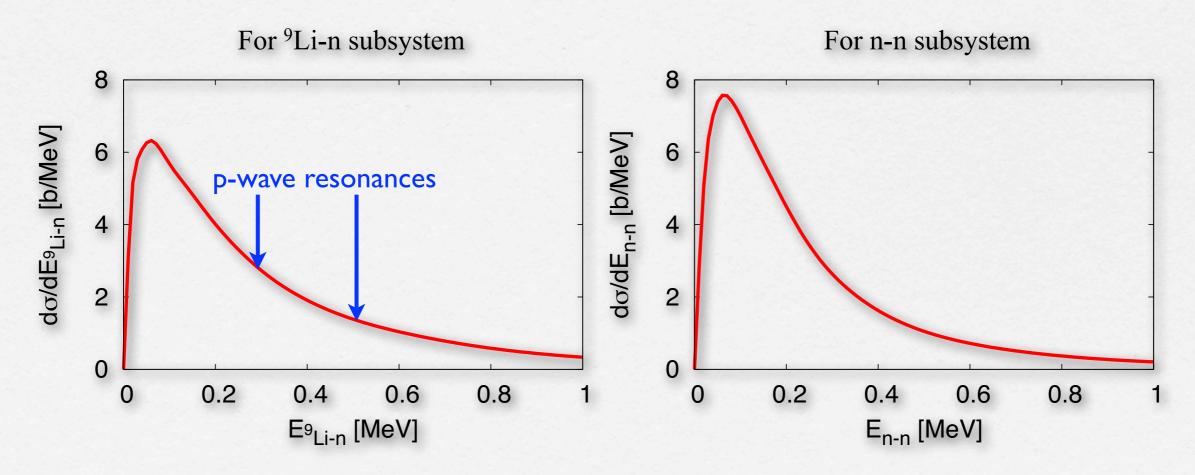
- □ Invariant mass spectra ⁶He case -
 - [□] For α -n, the spectrum shows the peak at around 0.7 MeV, which corresponds to the ${}^{5}\text{He}(3/2^{-})$ resonance energy.
 - \Box The sequential decay via the ⁵He(3/2⁻)+n channel dominates the breakup of ⁶He.
 - □ For n-n, the virtual-state correlation generates the low-lying peak in the spectra.



T.Aumann et al., PRC59, 1252 (1999).

Invariant mass spectra for binary subsystems

- Invariant mass spectra ¹¹Li case -
 - \Box In both cases of ⁹Li-n and n-n, the sharp peaks appear below 0.1 MeV.
 - For n-n, the virtual-state correlation plays a significant role in breakup as similar to ⁶He case.
 - For ⁹Li-n, the peak comes from the virtual state of ¹⁰Li, although the p-wave resonances of ¹⁰Li cannot be found in the spectra.



Summary

- We investigated the mechanisms of Coulomb breakup reactions of two-neutron halo nuclei ⁶He and ¹¹Li based on the core+n+n three-body models.
 - To describe the three-body scattering states of the core+n+n systems, we use the complex-scaled solutions of the Lippmann-Schwinger equation (CSLS).
 - CSLS is capable method of describing the three-body Coulomb breakups.
 - We calculate the Coulomb breakup cross section, and investigate
 - what is the effects of FSI in the Coulomb breakup reaction
 - \rightarrow FSI dominates the cross sections in both cases of ⁶He and ¹¹Li.
 - what kinds of correlations play roles in the Coulomb breakup reactions of twoneutron halo nuclei.
 - We confirmed that binary subsystem correlations have significant contributions to the breakups.
 - \Box For ⁶He, ⁵He(3/2⁻) resonance and n-n virtual state are important.
 - For ¹¹Li, virtual-state correlations in ¹⁰Li and n-n subsystems are clearly seen in the invariant mass spectra, while the p-wave resonances cannot be identified.