

チャンネル結合 ${}^9\text{Li}+n+n$ 模型を用いた  
 ${}^{11}\text{Li}$ のクーロン分解反応の分析

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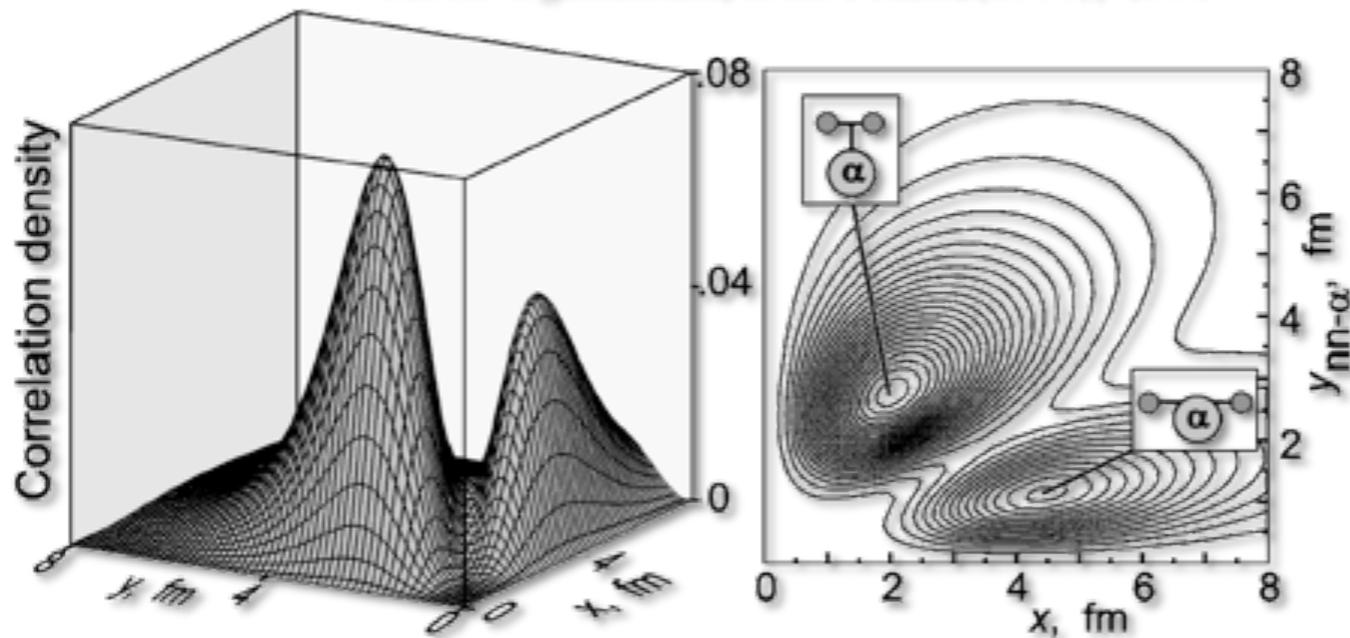
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published in PRC87, 034606 (2013).

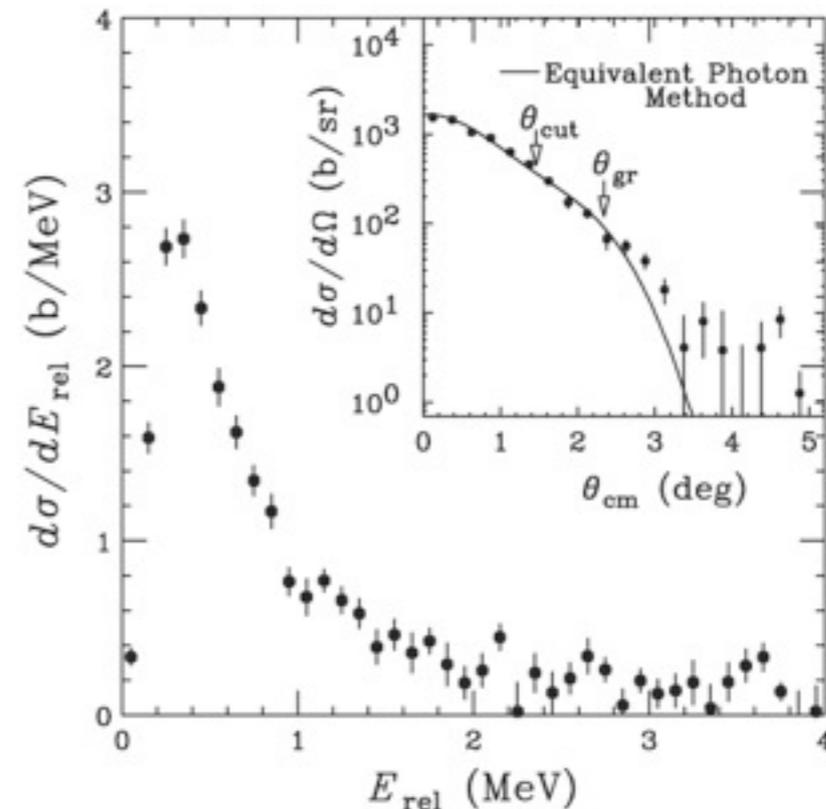
## 2n halo nuclei and their Coulomb breakups

- Theoretical studies based on core+2n picture
  - “dineutron” correlation
- Experiments using Coulomb breakup reactions
  - Low-lying enhancement above the breakup threshold
  - Related to dineutron?

Yu.Ts. Oganessian, *et al.* PRL82(1999), 4996



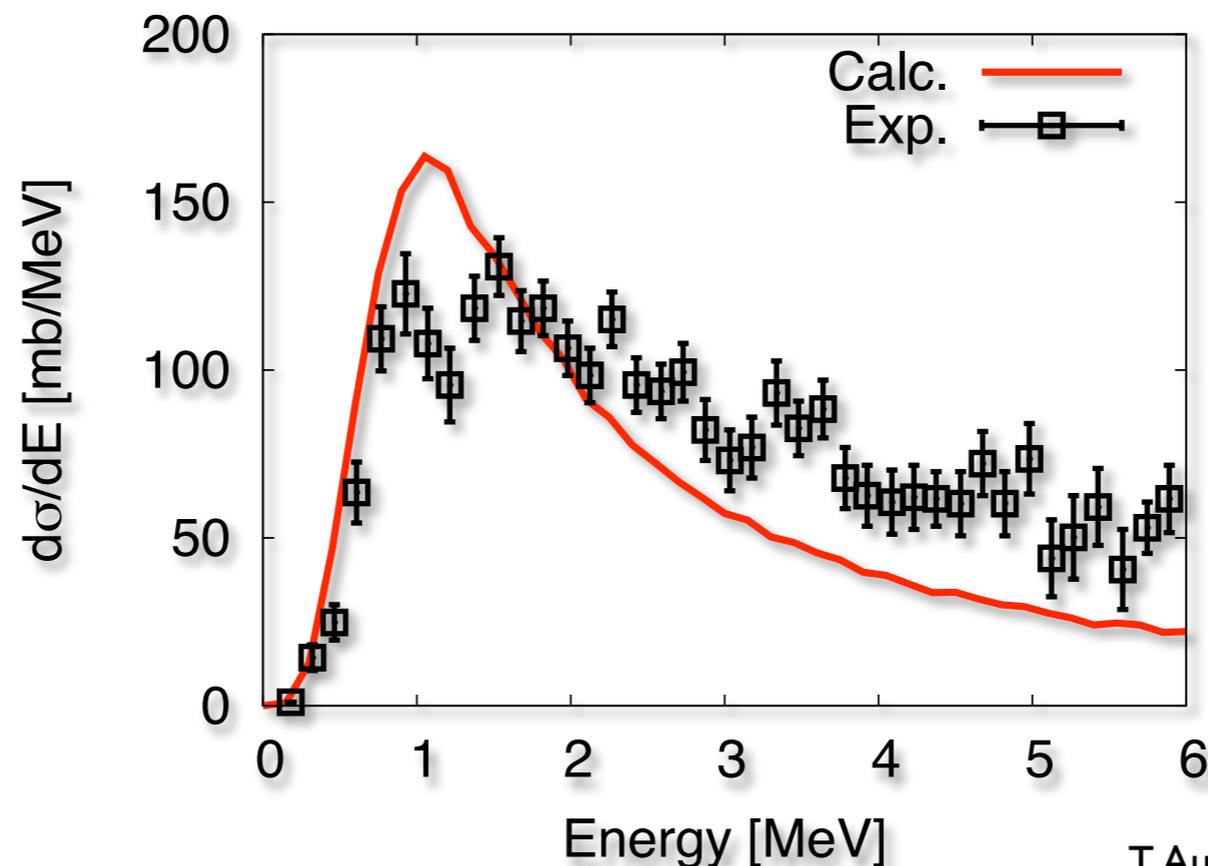
${}^6\text{Li}$  breakup: T. Nakamura *et al.*, PRL 96, 252502 (2006).



# Coulomb breakup reaction of ${}^6\text{He}$

YK et al., PTP 122, 499 (2009).  
YK et al., PRC 81, 044308 (2010).

- Our previous analysis based on  $\alpha+n+n$  three-body model
  - Good agreement with the observed cross section
  - Dominance of FSI
    - ${}^5\text{He}(3/2^-)$  resonance & n-n virtual state
    - The information on g.s. structure is masked



T.Aumann et al., PRC 59, 1252 (1999).

# Exotic properties of $^{11}\text{Li}$

- Breaking of the  $N=8$  magic number
  - Large  $(s_{1/2})^2$  component of halo neutron comparable to the  $(p_{1/2})^2$  one
  - Such a large s-wave mixing causes the well-developed halo structure.
- Suggested point from the simple model assuming the inert  $^9\text{Li}$  core
  - Difficult to explain the charge radius and the dipole strength in  $^{11}\text{Li}$  simultaneously\*.

K. Hagino et al., PRC80, 031301(R) (2009).

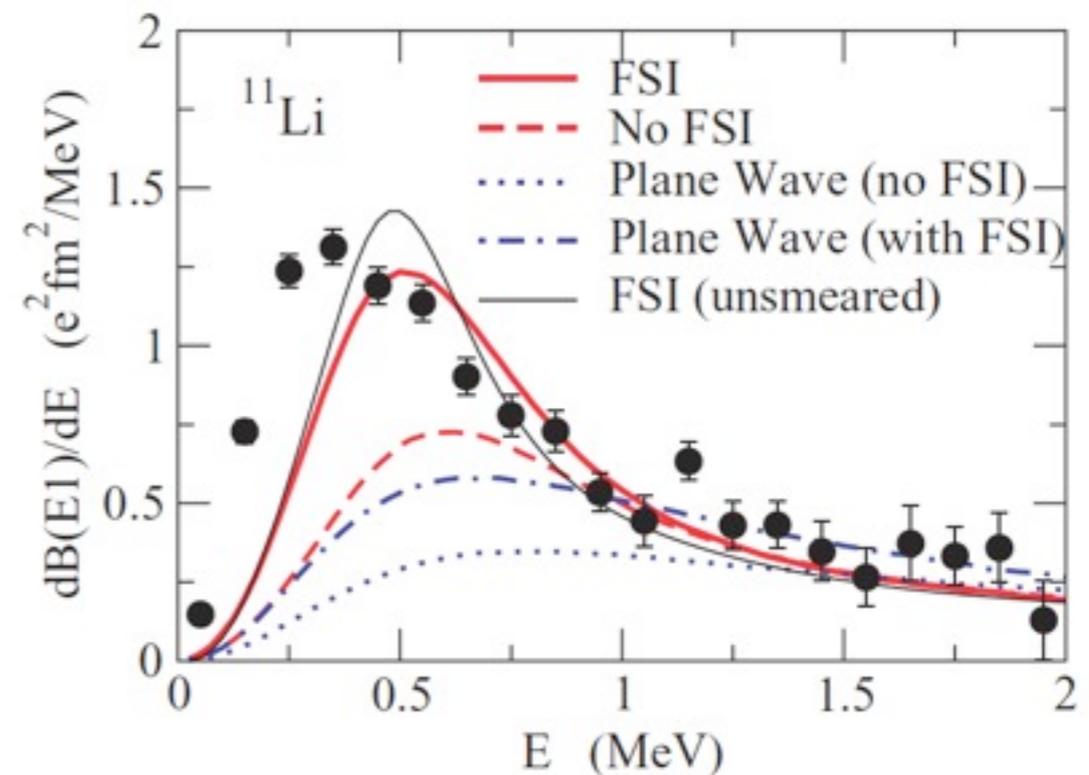
\*H. Esbensen et al., PRC76, 024302 (2007).

## Theory

$(1s_{1/2})^2$  component = 20.6 %  
scattering length for  $^9\text{Li-n}$  = -5.6 fm

## Expt.

$(1s_{1/2})^2$  component =  $45 \pm 10\%$   
scattering length for  $^9\text{Li-n}$   $\sim$  -30 fm



## In this talk...

- We investigate the Coulomb breakup mechanism of  $^{11}\text{Li}$ .
- We calculate the cross section by using the coupled-channel  $^9\text{Li}+n+n$  three-body model, in which the coupling between the relative motion of halo neutrons and the excitations of  $^9\text{Li}$  core is taken into account.
- We discuss the following points:
  - Effects of FSI
  - Binary subsystem correlations
  - Effects of  $^9\text{Li}$  core excitation
- Our approach T. Myo et al., PTP119, 561 (2008).
  - 2p-2h excitations due to tensor and pairing correlations in  $^9\text{Li}$  (TOSM)
    - ➡ well reproduce the observed quantities in g.s. and EI strength

# Coupled-channel ${}^9\text{Li}+n+n$ three-body model

- We take into account the coupling between the relative motion of halo neutrons and the excitation of the  ${}^9\text{Li}$  core.

- Hamiltonian

$$\hat{H} = \sum_{i=1}^3 t_i - T_{\text{cm}} + \sum_{i=1}^2 V_{{}^9\text{Li}-n}(\mathbf{r}_i) + V_{n-n} + \lambda |\Phi_{\text{PF}}\rangle \langle \Phi_{\text{PF}}| + h({}^9\text{Li})$$

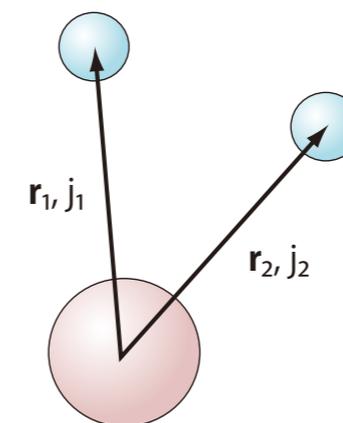
$V_{{}^9\text{Li}-n}$  : folded MHN potential     $V_{n-n}$  : Argonne v8'

- Wave function

$$\Phi^{J\pi}({}^{11}\text{Li}) = \sum_i C_i [\Phi_i({}^9\text{Li}) \otimes \chi_{nn}]_{J\pi}$$

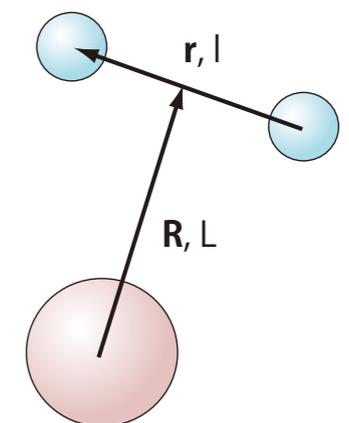
0p-0h & 2p-2h configurations

COSM (V-type)



shell model-like

ECM (T-type)

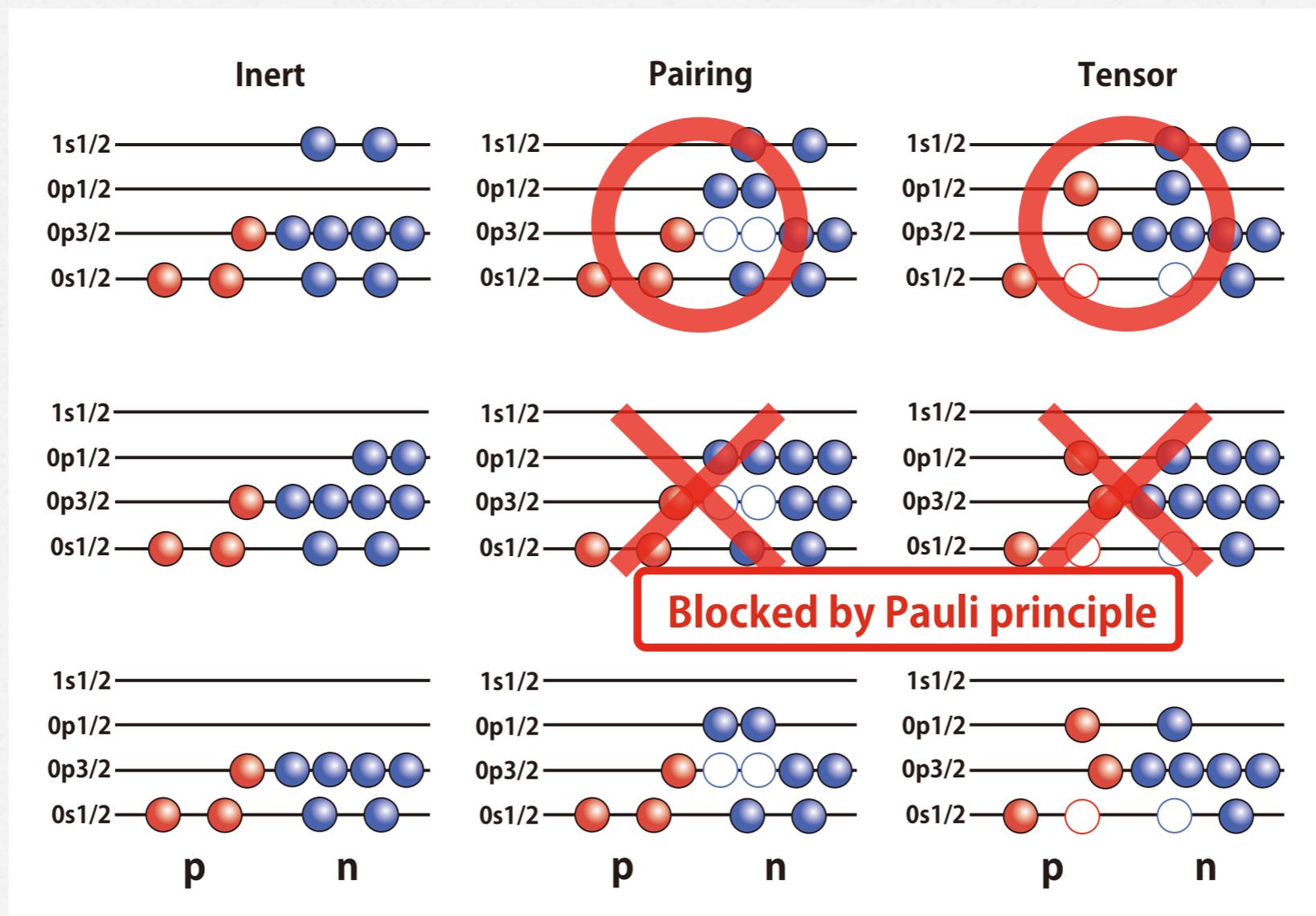


di-neutron-like

## 2p-2h excitations in ${}^9\text{Li}$ core

- The 2p-2h configurations make the single-particle energy of the  $p_{1/2}$  orbit pushed up by the Pauli blocking effect.
  - Large  $(s_{1/2})^2$  mixing & well-developed halo structure

T. Myo et al., PTP119, 561 (2008).



# Properties of $^{11}\text{Li}$ and $^{10}\text{Li}$ with coupled-channels

- For  $^{11}\text{Li}$ 
  - Large s-wave mixing of 44.0% ( $45 \pm 10\%$  in Exp.)
- For  $^{10}\text{Li}$ 
  - large s-wave scattering length  $\rightarrow$   $^{10}\text{Li}$  virtual state
  - Two p-wave resonances @  $E_r = 275$  keV and 506 keV

For  $^{11}\text{Li}$

For  $^{10}\text{Li}$

|                      | Theor. | Exp.                |              | 1 <sup>+</sup> | 2 <sup>+</sup> | 1 <sup>-</sup> | 2 <sup>-</sup> | Exp.              |
|----------------------|--------|---------------------|--------------|----------------|----------------|----------------|----------------|-------------------|
| $S_{2n}$ (keV)       | 377    | $378 \pm 5^a$       | Energy (keV) | 275            | 506            |                |                |                   |
| $R_{\text{m}}$ (fm)  | 3.39   | $3.12 \pm 0.16^b$   | Width (keV)  | 150            | 388            |                |                |                   |
|                      |        | $3.53 \pm 0.06^c$   | $a_s$ (fm)   |                |                | -6.8           | -45.0          | $-30^{+12}_-31^a$ |
|                      |        | $3.71 \pm 0.20^d$   |              |                |                |                |                |                   |
| $R_{\text{ch}}$ (fm) | 2.43   | $2.467 \pm 0.037^e$ |              |                |                |                |                |                   |
|                      |        | $2.423 \pm 0.034^f$ |              |                |                |                |                |                   |
| $P((s_{1/2})^2)(\%)$ | 44.0   | $45 \pm 10^g$       |              |                |                |                |                |                   |
| $P((p_{3/2})^2)(\%)$ | 2.5    | -                   |              |                |                |                |                |                   |
| $P((p_{1/2})^2)(\%)$ | 46.9   | -                   |              |                |                |                |                |                   |
| $P((d_{5/2})^2)(\%)$ | 3.1    | -                   |              |                |                |                |                |                   |
| $P((d_{3/2})^2)(\%)$ | 1.7    | -                   |              |                |                |                |                |                   |

# Coulomb breakup cross section

- Equivalent photon method

$$\frac{d^6\sigma}{d\mathbf{k}d\mathbf{K}} = \frac{16\pi^3}{9\hbar c} \cdot N_{E1}(E_\gamma) \cdot \frac{d^6B(E1)}{d\mathbf{k}d\mathbf{K}}$$
$$\frac{d^6B(E1)}{d\mathbf{k}d\mathbf{K}} = \frac{1}{2J_{\text{gs}} + 1} \left| \langle \Psi^{(+)}(\mathbf{k}, \mathbf{K}) || \hat{O}(E1) || \Phi_{\text{gs}} \rangle \right|^2$$

- For scattering states:  
Complex-scaled solutions of the Lippmann-Schwinger equation (CSLS)
  - Interaction region
    - Few-body technique similar to the bound states
  - Boundary conditions
    - Complex scaling method

# Complex-scaled solutions of LS equation (CSLS)

- Formal solution of LS equation

$$\Psi^{(\pm)} = \Phi_0 + \lim_{\varepsilon \rightarrow 0} \frac{1}{E - \hat{H} \pm i\varepsilon} \hat{V} \Phi_0$$

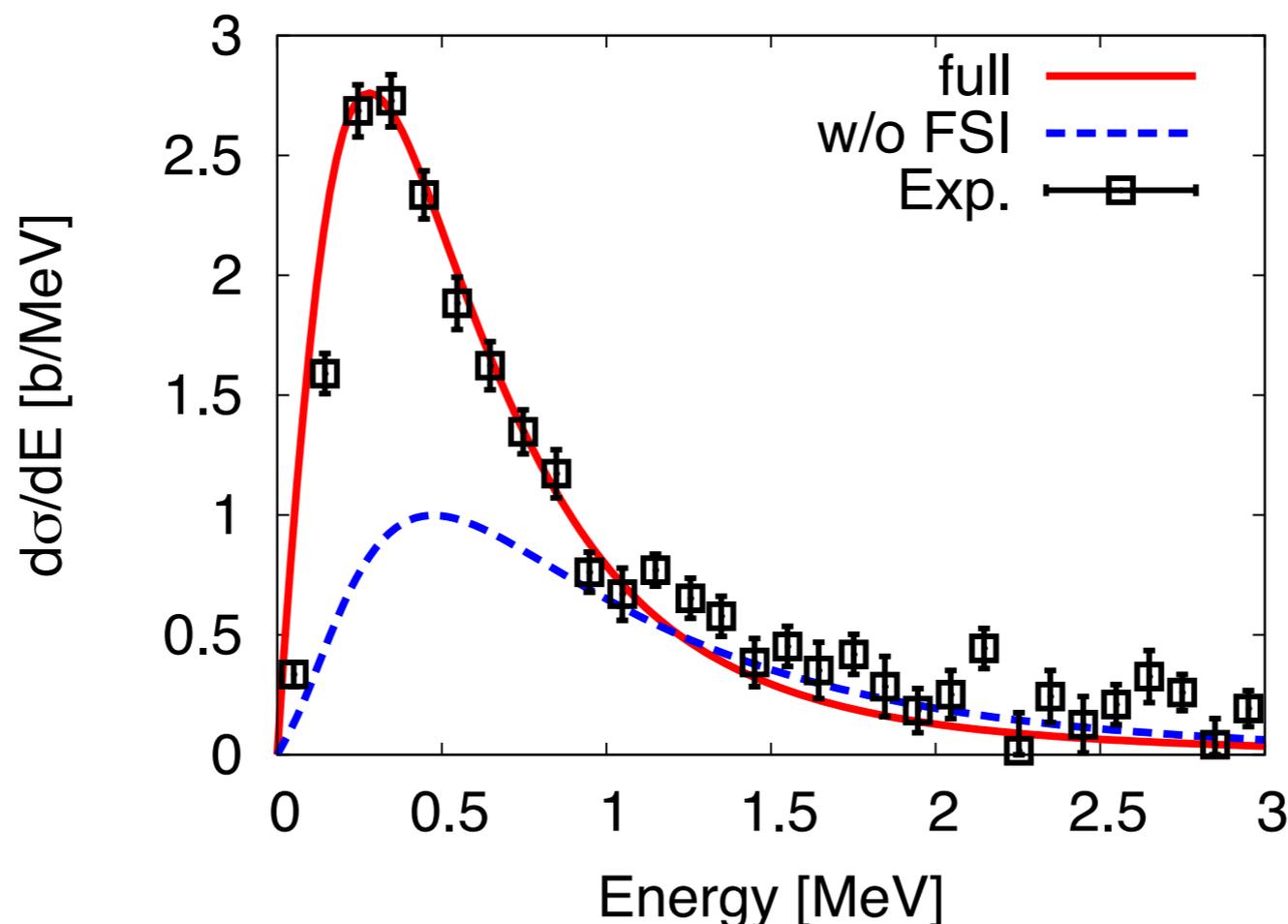
- Complex-scaled Green's function

$$\lim_{\varepsilon \rightarrow 0} \frac{1}{E - \hat{H} + i\varepsilon} = U^{-1}(\theta) \frac{1}{E - \hat{H}^\theta} U(\theta) = \sum_n U^{-1}(\theta) |\chi_n^\theta\rangle \frac{1}{E - E_n^\theta} \langle \tilde{\chi}_n^\theta | U(\theta)$$

- Expand with complete set constructed by the eigenstates of  $H^\theta$
- Coupling b/w relative motion of halo neutrons and excitation of  ${}^9\text{Li}$  core

# Coulomb breakup cross section of $^7\text{Li}$

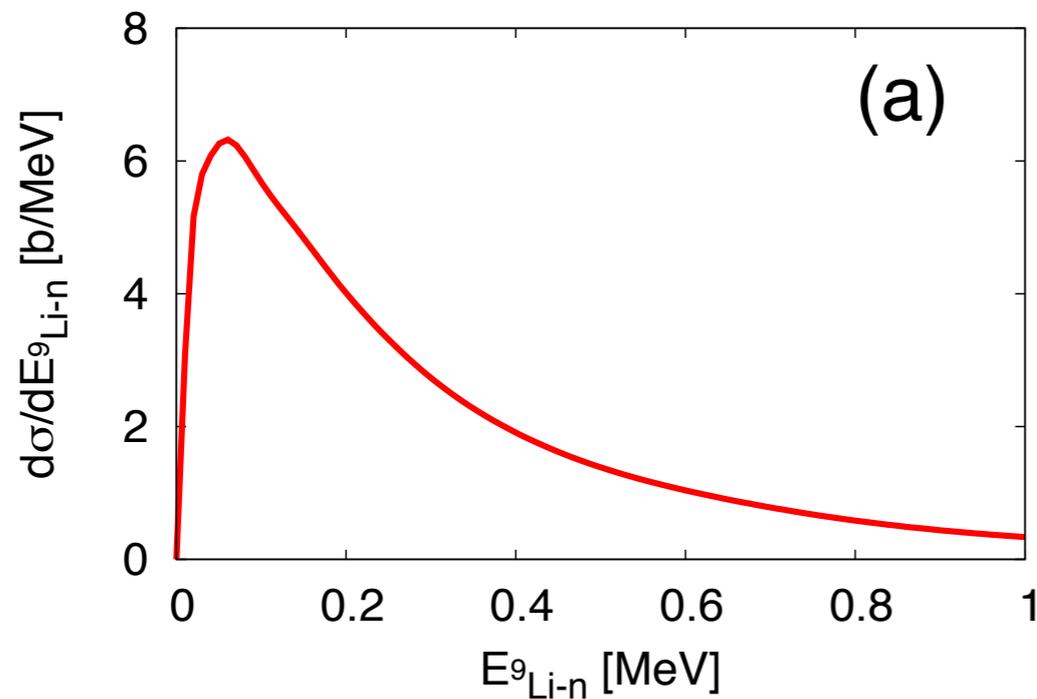
- Calculated differential cross section with respect to total excitation energy.
- Perfect agreement with the experiment in whole energy region.
- Low-lying enhancement comes from FSI  $\rightarrow$  similar to the  $^6\text{He}$  case.



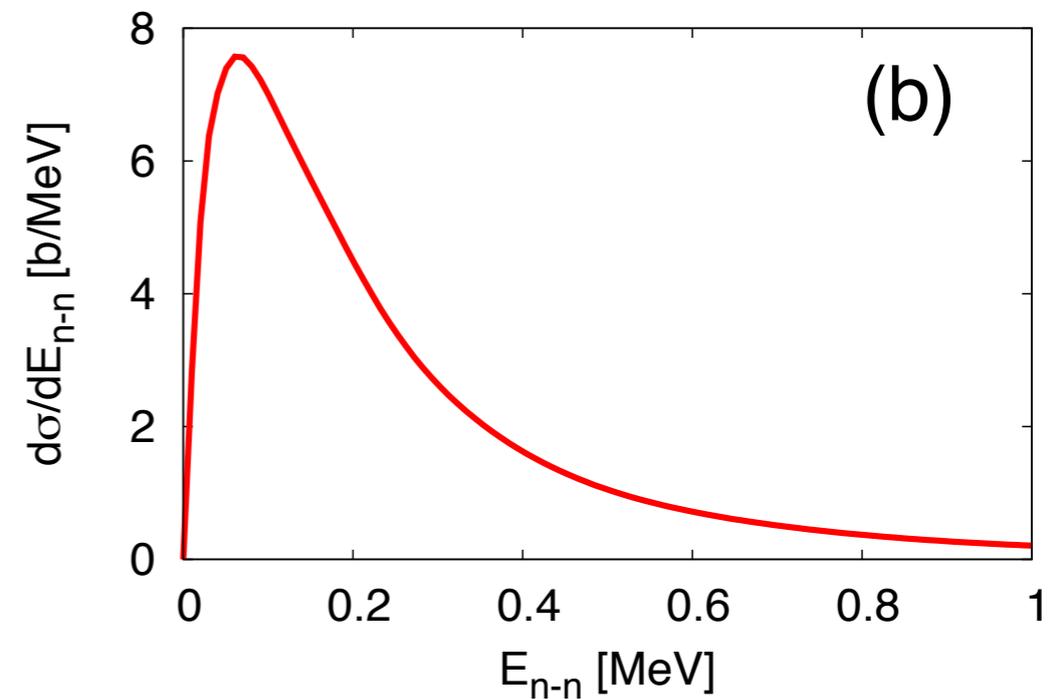
# Breakup mechanism of $^{11}\text{Li}$ Coulomb breakup

- Invariant mass spectra for binary subsystems
  - $^{10}\text{Li}$  and n-n virtual states
  - p-wave resonances of  $^{10}\text{Li}$  → not identified

For  $^9\text{Li}$ -n subsystem



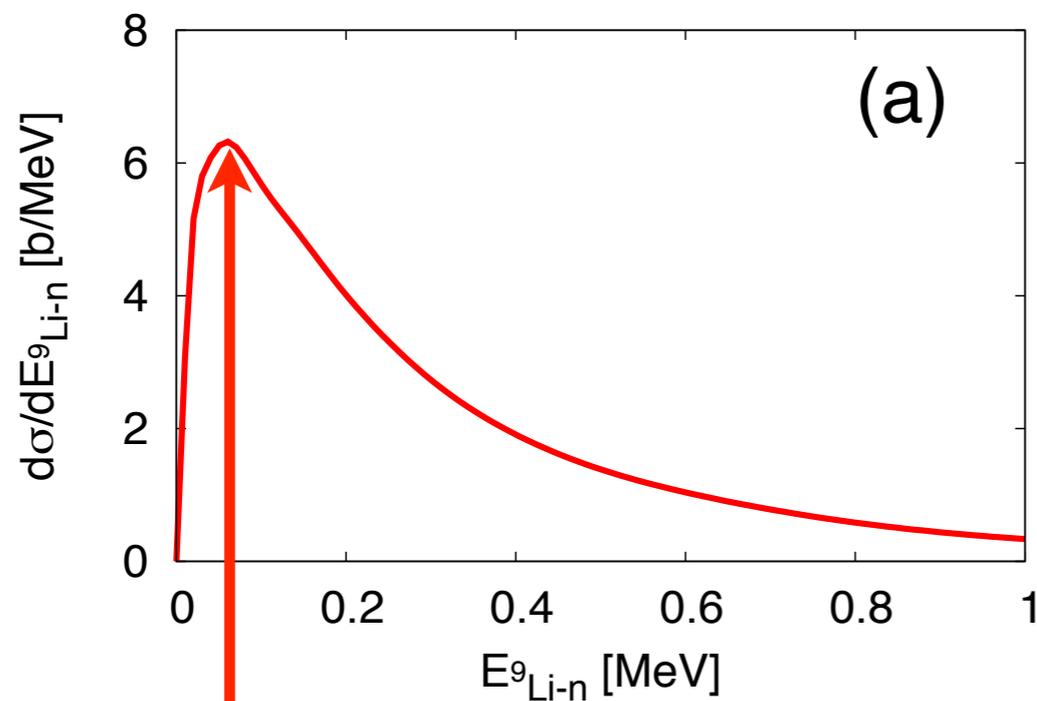
For n-n subsystem



# Breakup mechanism of ${}^7\text{Li}$ Coulomb breakup

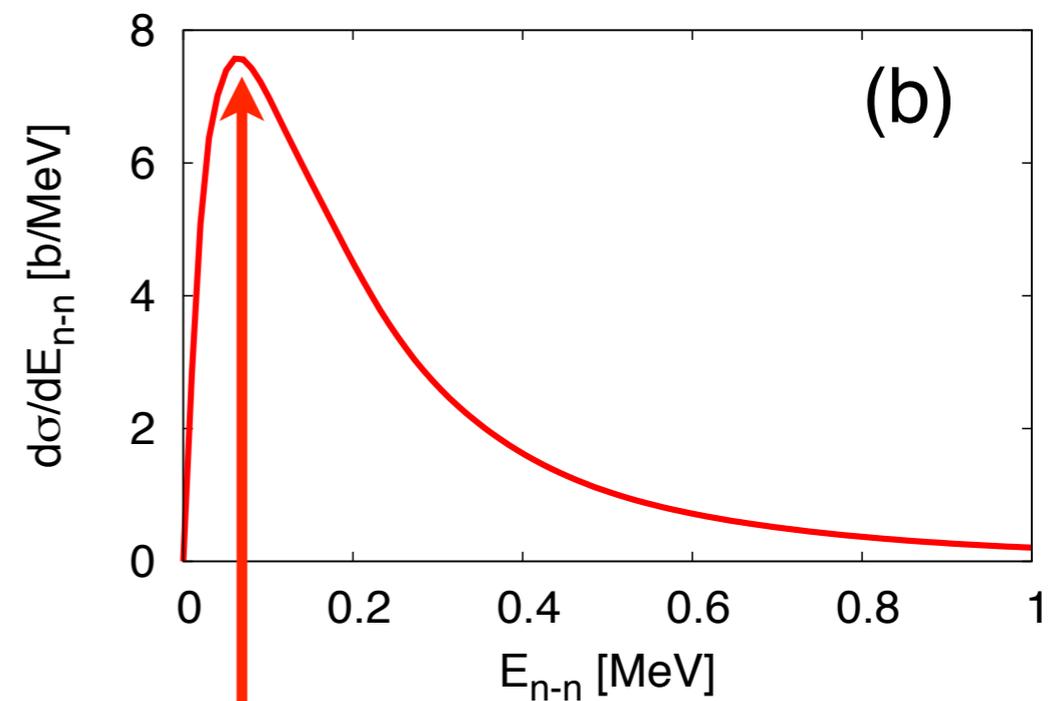
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For  ${}^9\text{Li}$ -n subsystem



**${}^{10}\text{Li}$  virtual state**

For n-n subsystem

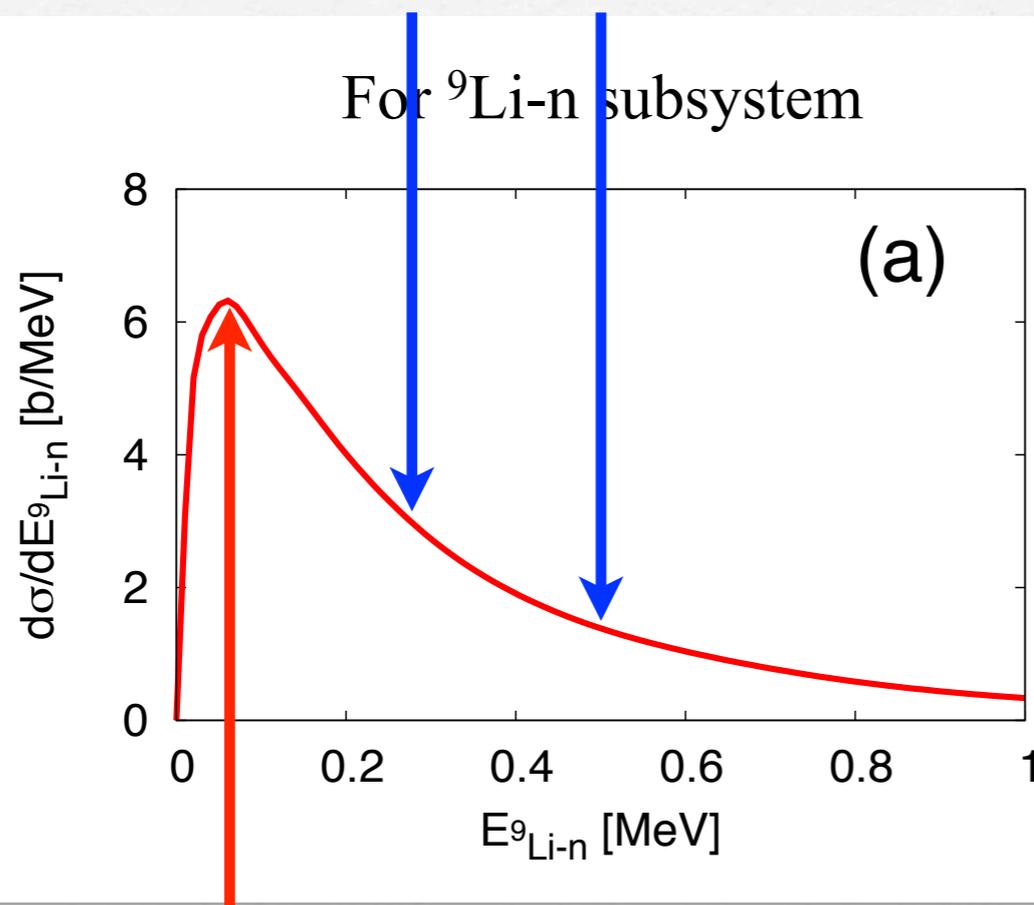


**n-n virtual state**

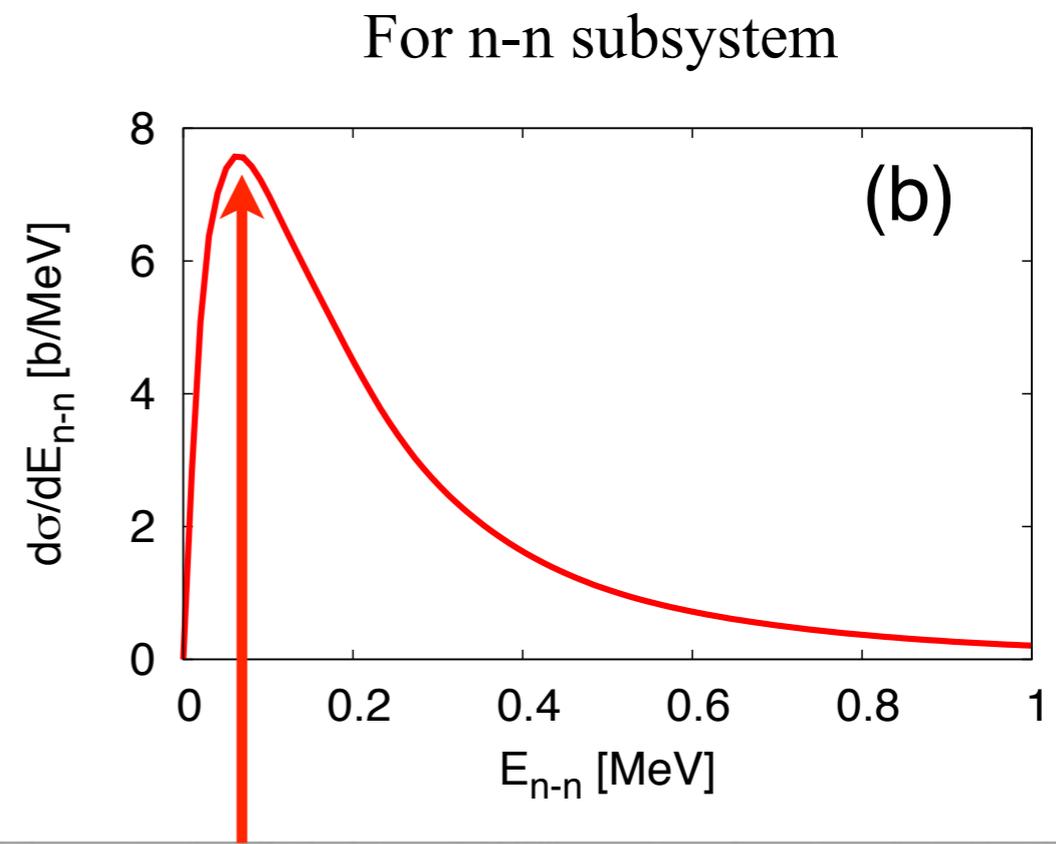
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## p-wave resonances



$^{10}\text{Li}$  virtual state



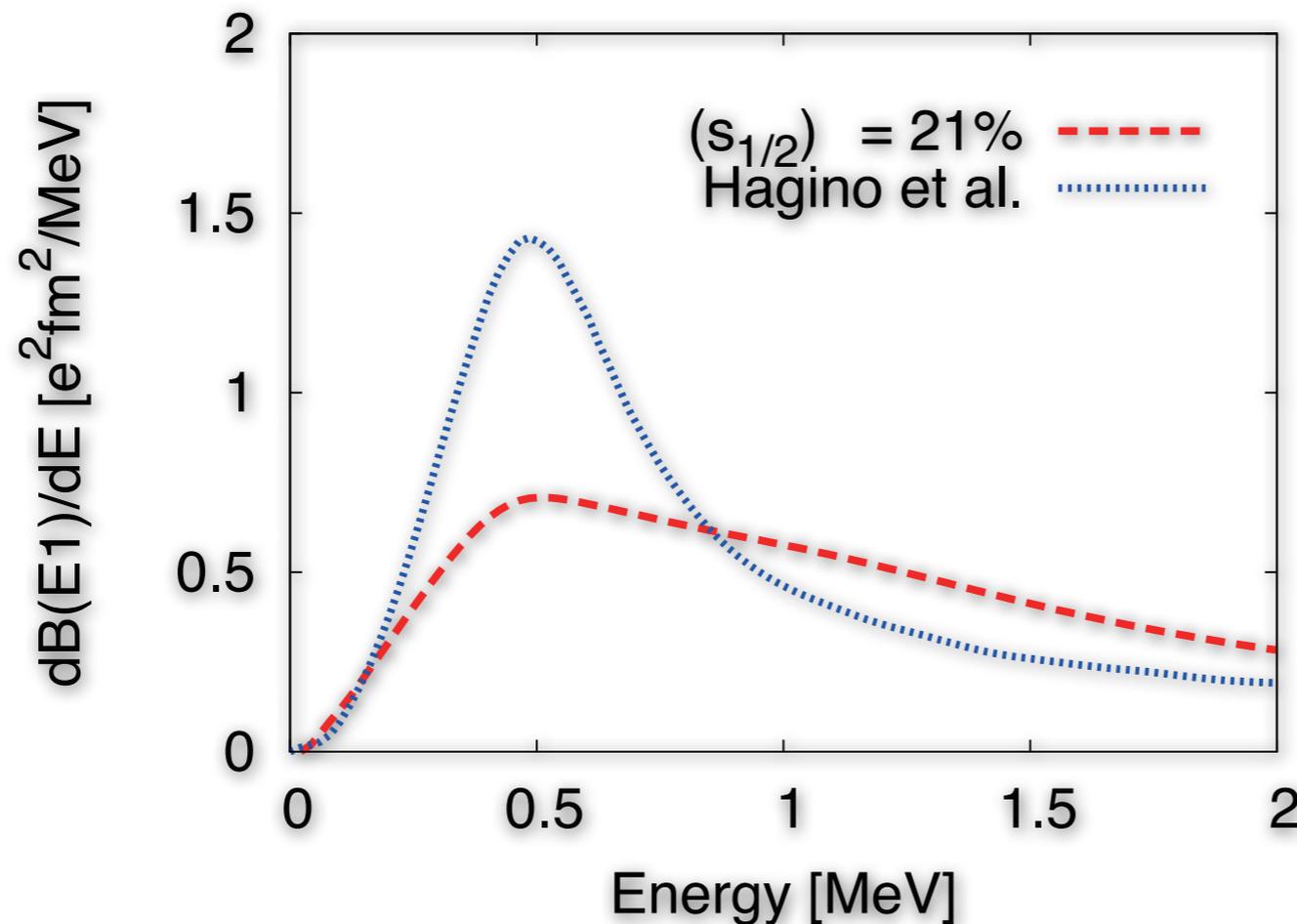
n-n virtual state

# Effects of ${}^9\text{Li}$ core excitation on the cross section

- To see the effects of the excitation of the  ${}^9\text{Li}$  core
  - For comparison...
    1. Our model including only pairing-type 2p-2h excitation
      - $V_{\text{core-n}}$  is tuned to reproduce the g.s. binding energy
      - $(s_{1/2})^2$  component = 21.0 %
    2. Simple model assuming inert  ${}^9\text{Li}$  core Hagino et al., PRC80, 031301(R) (2009).
      - Using different  $V_{\text{core-n}}$  for even- and odd-parity states
      - $(s_{1/2})^2$  component = 20.6 %

# Effects of ${}^9\text{Li}$ core excitation on the cross section

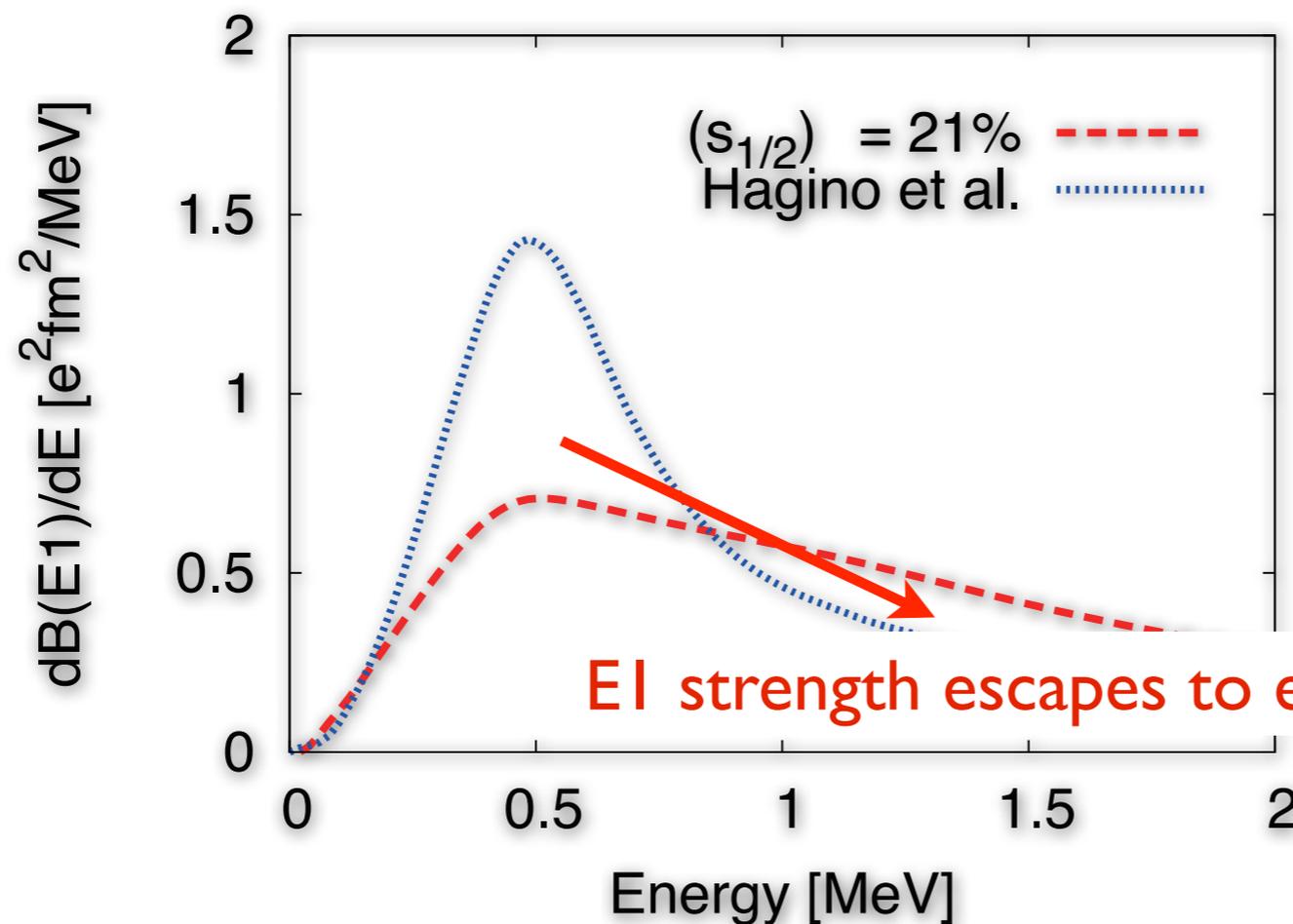
- ${}^9\text{Li}$  core excitation
  - ➔ Effect on shape & sum rule reduction  $\sim 15\%$



K. Hagino et al., PRC80, 031301(R) (2009).

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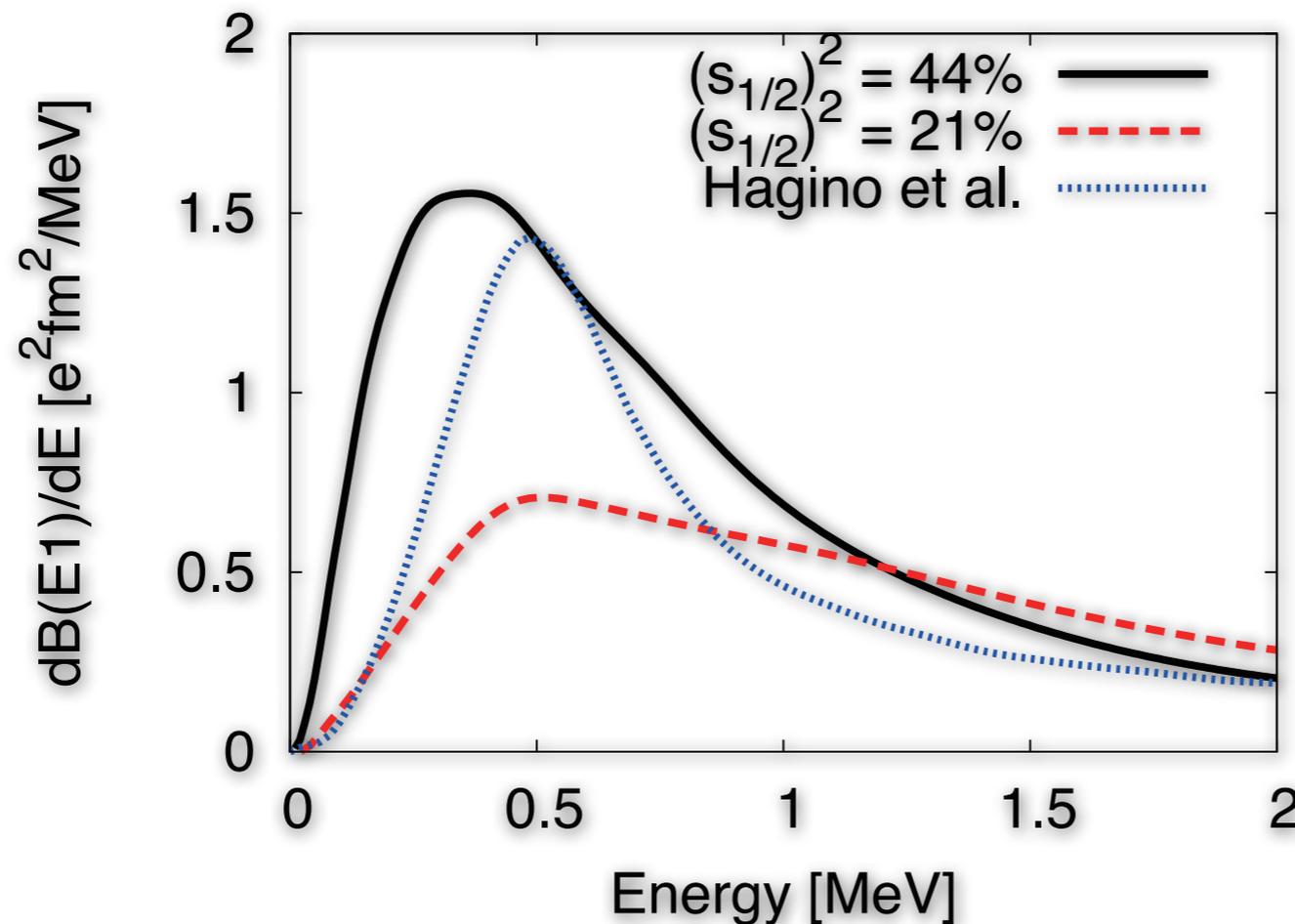


K. Hagino et al., PRC80, 031301(R) (2009).

El strength escapes to excited component of  ${}^9\text{Li}$

# Effects of $^9\text{Li}$ core excitation on the cross section

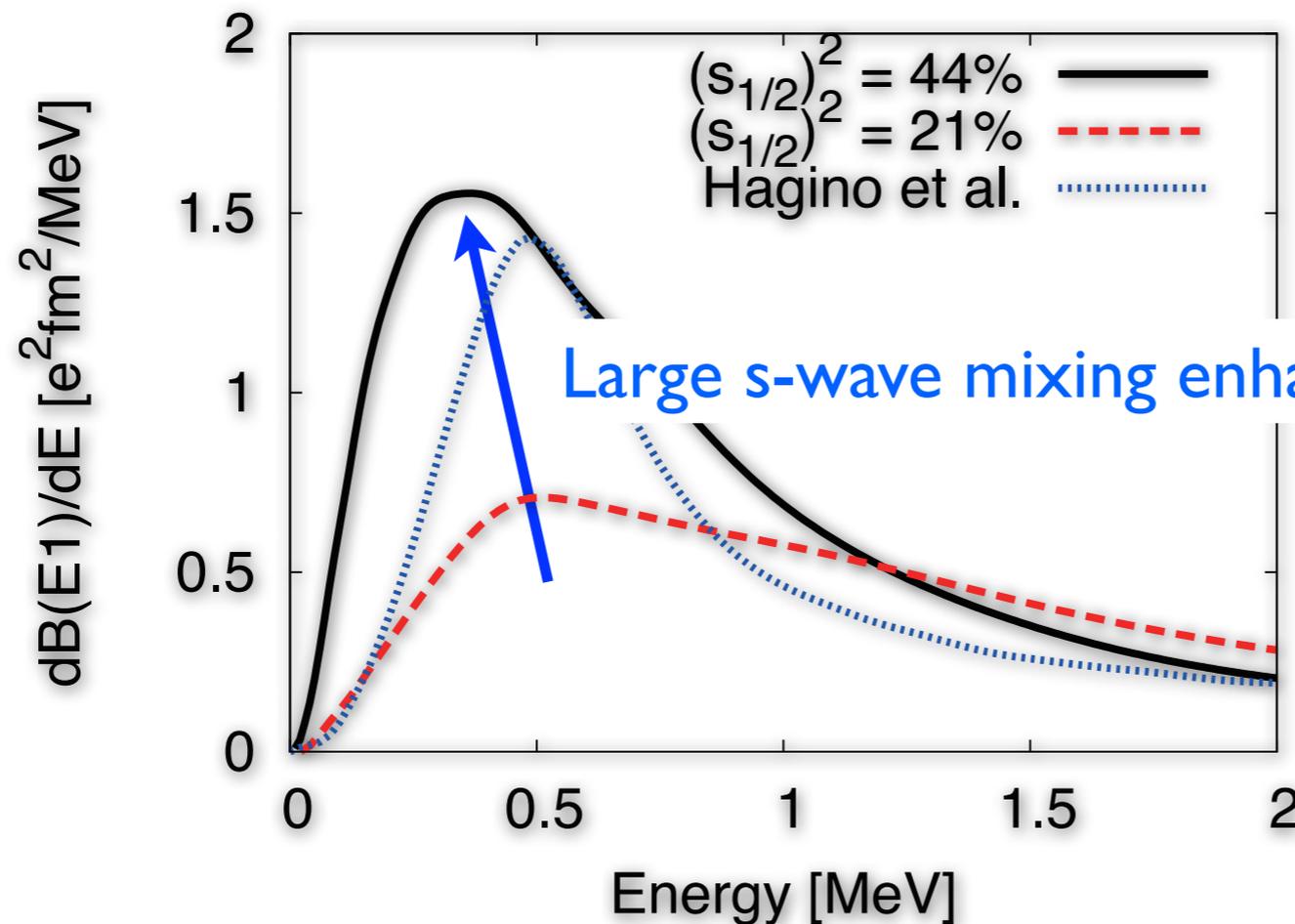
- $^9\text{Li}$  core excitation
  - ➔ Effect on shape & sum rule reduction  $\sim 15\%$
- Large s-wave mixing
  - ➔ Enhances the strength at low energy



K. Hagino et al., PRC80, 031301(R) (2009).

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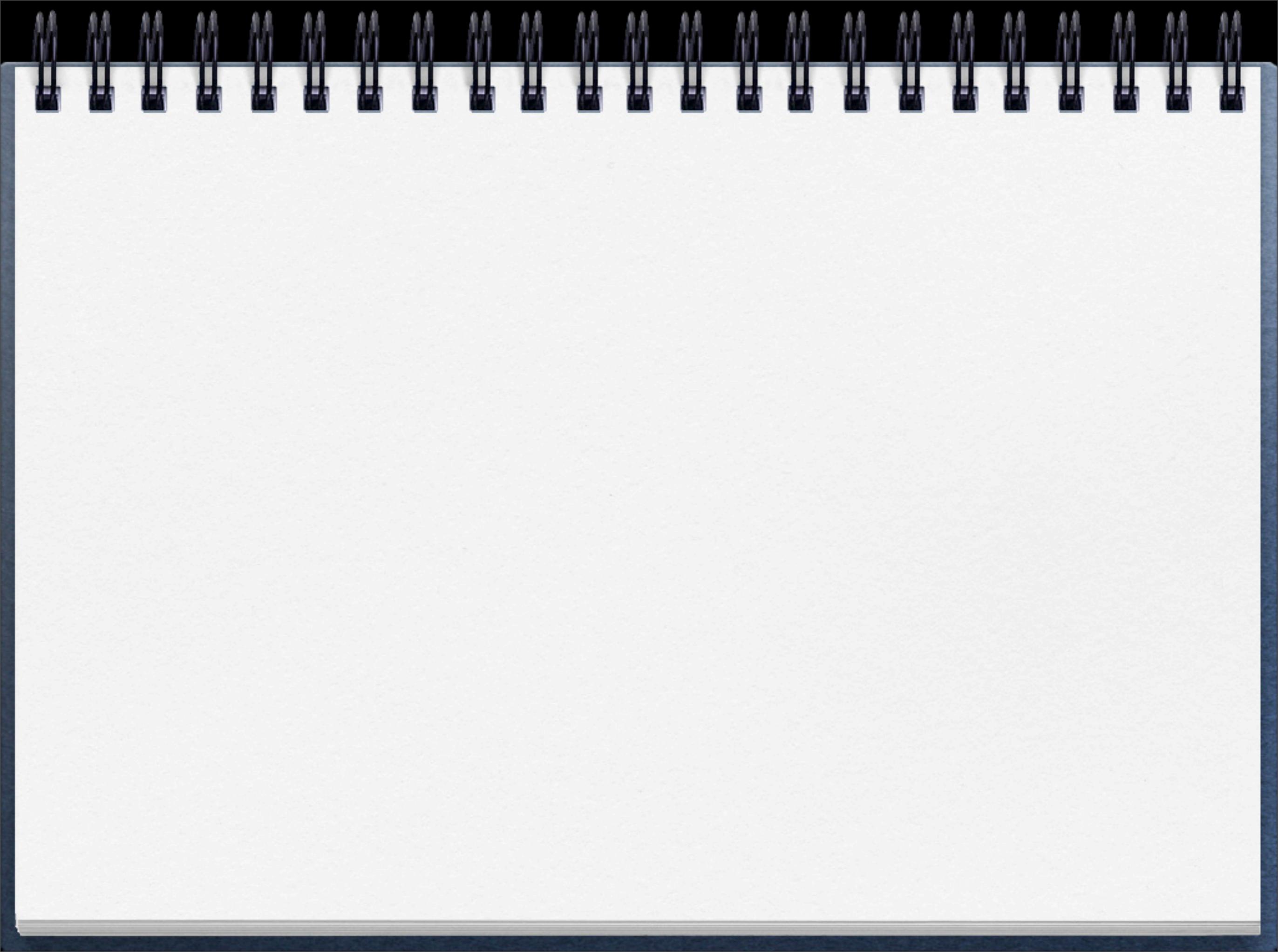


K. Hagino et al., PRC80, 031301(R) (2009).

# Summary

- We investigated the Coulomb breakup mechanism of  $^{11}\text{Li}$ .
- We calculated the cross section by using the coupled-channel  $^9\text{Li}+n+n$  three-body model.
  - ➔ Perfect agreement with the observed cross section
- We discussed the following points:
  - Effects of FSI ➔ Dominated the cross section
  - Binary subsystem correlations ➔  $^{10}\text{Li}$  & n-n virtual states  
 $^{10}\text{Li}$  p-wave resonances are not identified
  - Effects of  $^9\text{Li}$  core excitation ➔ 15 % reduction of EI sum rule

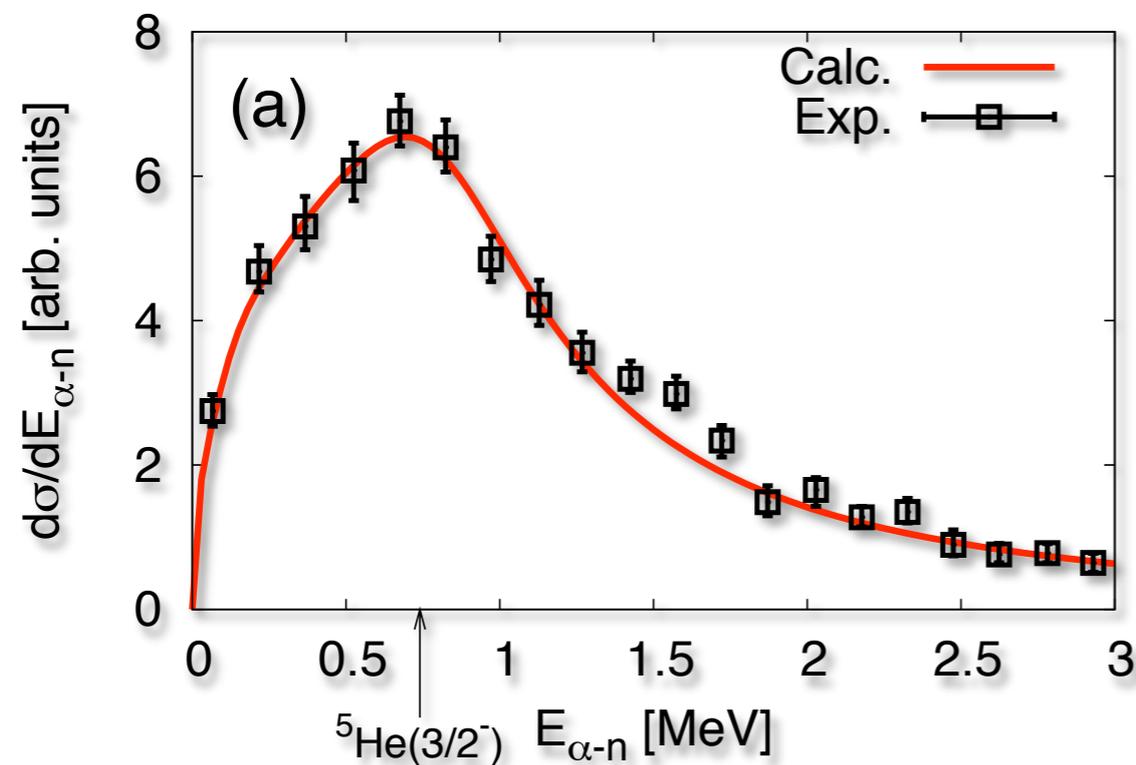
**2p-2h excitation due to the tensor and pairing correlations in  $^9\text{Li}$  is a key to reproduce the observed quantities.**



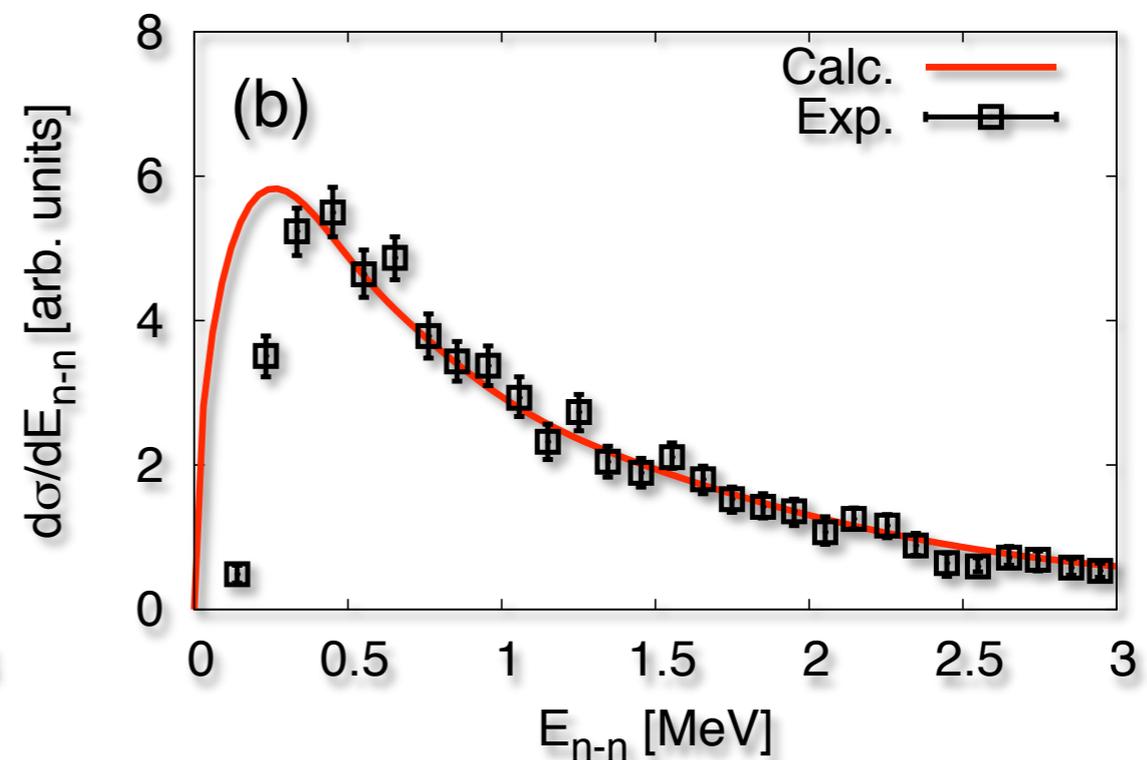
# Results - Application to ${}^6\text{He}$ -

- What kinds of FSI are important?
  - We calculate invariant mass spectra as functions of energies of binary subsystems.
  - For  ${}^4\text{He}$ -n, the spectra has a peak, whose position corresponds to resonance energy of  ${}^5\text{He}(3/2^-)$ .
  - For n-n, the spectra has a sharp peak in low-energy region, which comes from the virtual s-state of n-n system.

For  $\alpha$ -n subsystem



For n-n subsystem



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