

Role of tensor force in light nuclei with tensor-optimized shell model

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Outline

- **Role of V_{tensor}** in light nuclei
 - He, Li, Be isotopes
- Tensor Optimized Shell Model (**TOSM**) to describe tensor correlation.
- Unitary Correlation Operator Method (**UCOM**) to describe short-range correlation.

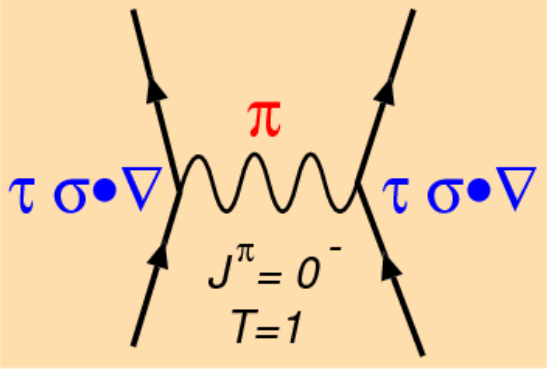
TM, A. Umeya, H. Toki, K. Ikeda PRC84 (2011) 034315

TM, A. Umeya, H. Toki, K. Ikeda PRC86 (2012) 024318

TM, A. Umeya, K. Horii, H. Toki, K. Ikeda PTEP (2014) 033D01

Pion exchange interaction vs. V_{tensor}

$$3(\vec{\sigma}_1 \cdot \hat{q})(\vec{\sigma}_2 \cdot \hat{q}) \frac{q^2}{m^2 + q^2} = (\vec{\sigma}_1 \cdot \vec{\sigma}_2) \frac{q^2}{m^2 + q^2} + S_{12} \frac{q^2}{m^2 + q^2}$$

$$= (\vec{\sigma}_1 \cdot \vec{\sigma}_2) \left[\frac{m^2 + q^2}{m^2 + q^2} - \frac{m^2}{m^2 + q^2} \right] + S_{12} \frac{q^2}{m^2 + q^2}$$


δ interaction Yukawa interaction

$S_{12} \frac{q^2}{m^2 + q^2}$

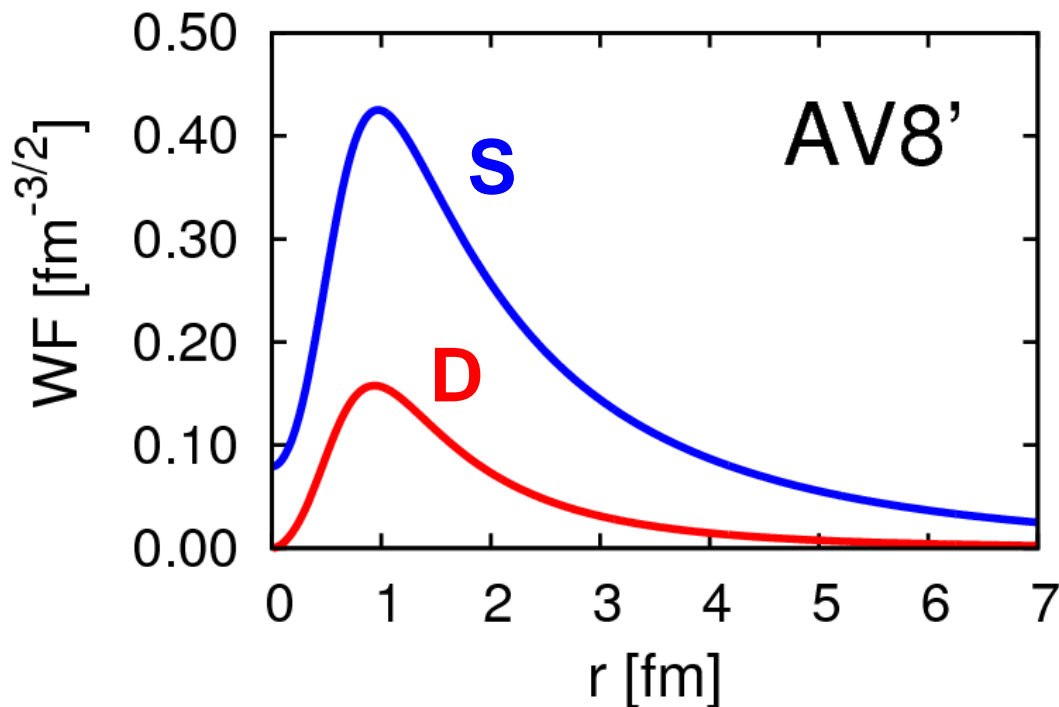
involve large momentum

Tensor operator

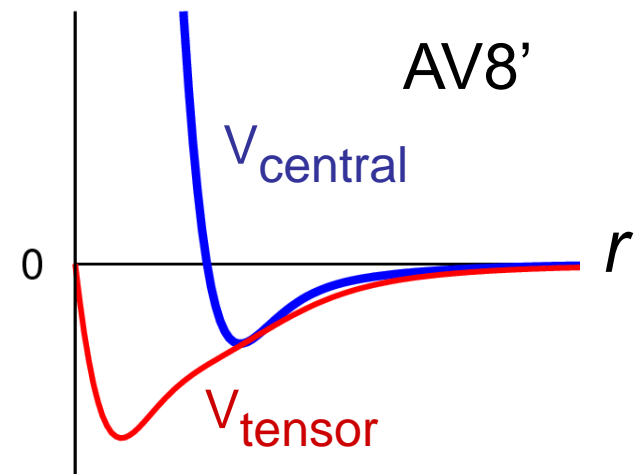
$$S_{12} = 3(\vec{\sigma}_1 \cdot \hat{q})(\vec{\sigma}_2 \cdot \hat{q}) - (\vec{\sigma}_1 \cdot \vec{\sigma}_2)$$

- V_{tensor} produces the high momentum component.
- **Horii** (4P-2B) treats $N \rightarrow \Delta$ transition in light nuclei.

Deuteron properties & tensor force



Energy	-2.24 MeV
Kinetic	19.88
Central	-4.46
Tensor	-16.64
LS	-1.02
P(L=2)	5.77%
Radius	1.96 fm



$$R_m(s) = 2.00 \text{ fm}$$

$$R_m(d) = 1.22 \text{ fm}$$

d-wave is
“spatially compact”
 (high momentum)

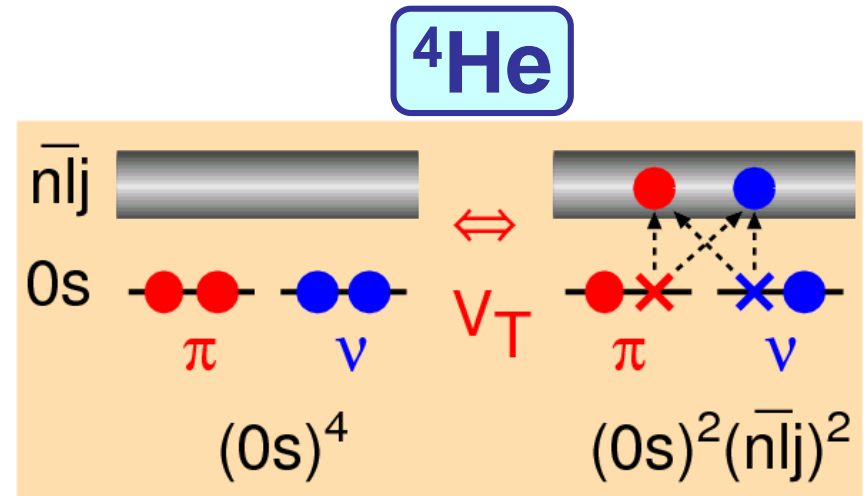
Tensor-optimized shell model (TOSM)

TM, Sugimoto, Kato, Toki, Ikeda PTP117(2007)257

- **2p2h excitations**
with high- L orbits.

- V_{tensor} is **NOT** treated as residual interactions

cf. $\frac{V_{\pi}}{V_{NN}} \sim 80\%$ in GFMC



- Length parameters such as b_{0s} , b_{0p} , ... are optimized **independently**, or **superposed by many Gaussian bases**.
 - **spatial shrinkage** of **D-wave** as seen in deuteron.
HF (Sugimoto, NPA740) , RMF (Ogawa, PRC73), AMD (Dote et al., PTP115)
- Satisfy few-body results with Minnesota central force (${}^4, {}^6\text{He}$)

Hamiltonian and variational equations in TOSM

$$H = \sum_{i=1}^A t_i - T_G + \sum_{i<j}^A v_{ij},$$

(0p0h+1p1h+2p2h)

$$\Phi(A) = \sum_k C_k \cdot \psi_k(A)$$

Shell model type configuration
with mass number A

Particle state : Gaussian expansion for each orbit

$$\varphi_{lj}^{n'}(\mathbf{r}) = \sum_{n=1}^N C_{lj,n}^{n'} \cdot \phi_{lj,n}(\mathbf{r})$$

$$\phi_{lj,n}(\mathbf{r}) \propto r^l \exp\left[-\frac{1}{2}\left(\frac{r}{b_{lj,n}}\right)^2\right] \left[Y_l(\hat{\mathbf{r}}), \chi_{1/2}^\sigma \right]_j$$

$$\langle \varphi_{lj}^{n'} | \varphi_{lj}^{n''} \rangle = \delta_{n',n''}$$

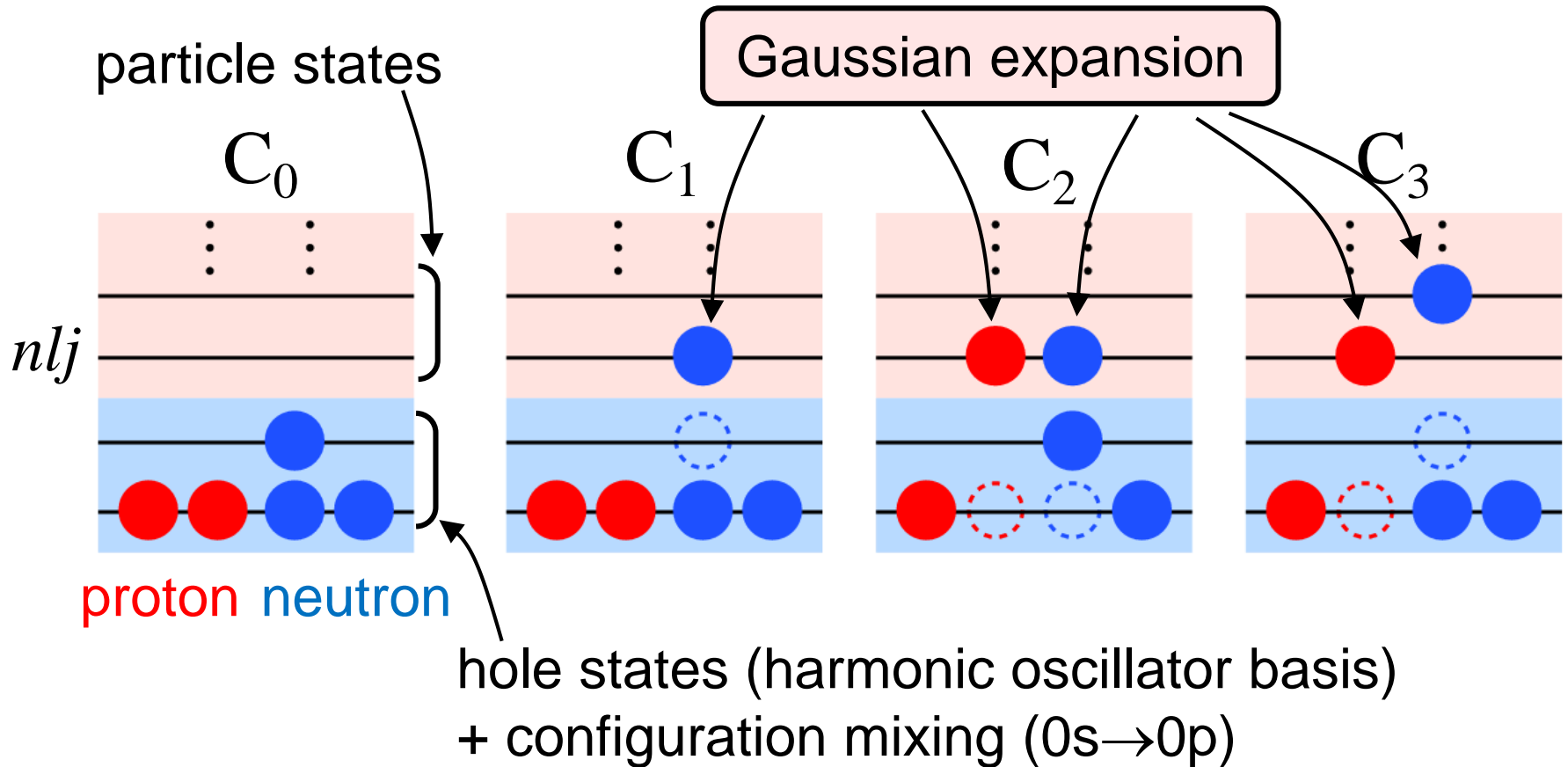
Gaussian basis function

Hiyama, Kino, Kamimura
PPNP51(2003)223

$$\frac{\partial \langle H - E \rangle}{\partial C_k} = 0, \quad \frac{\partial \langle H - E \rangle}{\partial b_{lj,n}} = 0$$

c.m. excitation is excluded
by Lawson's method

Configurations in TOSM

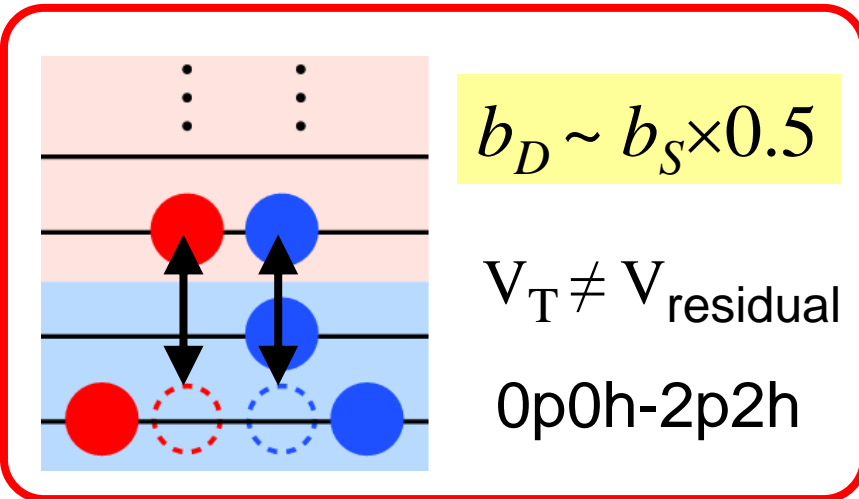
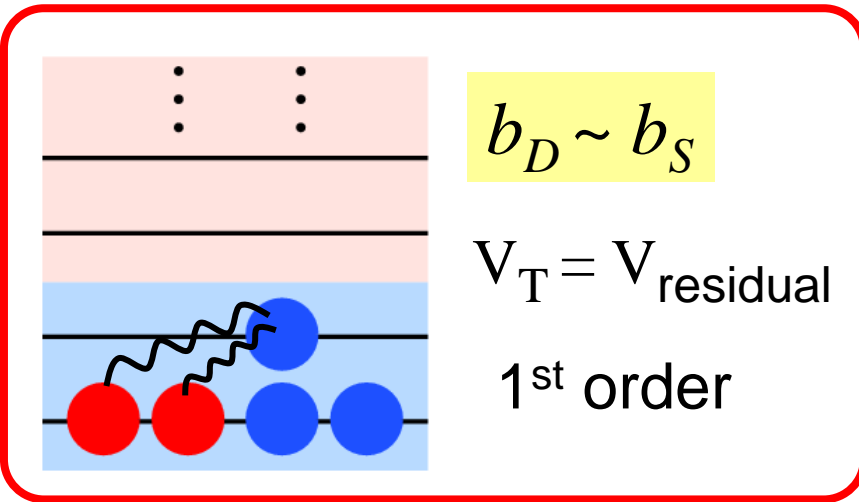


Application to Hypernuclei to investigate ΛN - ΣN coupling by **Umeya** (NIT), **Hiyama** (RIKEN)

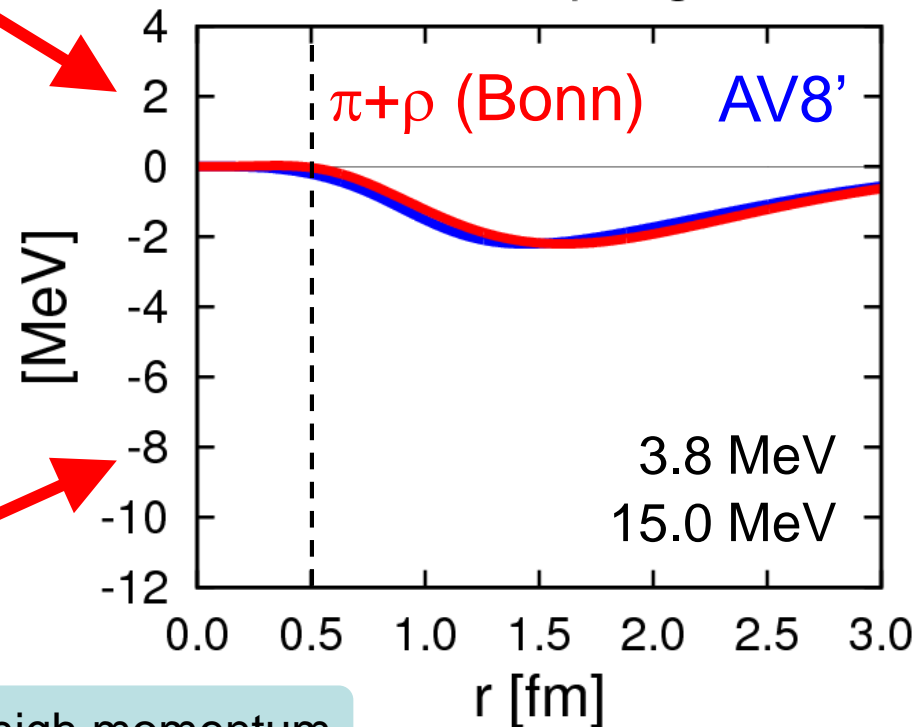
Tensor force matrix elements

$$M_{SD}(r) = r^2 \phi_S(r, b_S) \cdot V_T \cdot \phi_D(r, b_D)$$

Integrand of Tensor ME



SD coupling



- Centrifugal potential (1GeV@0.5fm) pushes away *D*-wave.

Unitary Correlation Operator Method

(short-range part)

$$\Psi_{\text{corr.}} = C \cdot \Phi_{\text{uncorr.}}$$

TOSM

short-range correlator

$$C^\dagger = C^{-1} \quad (\text{Unitary trans.})$$

$$H\Psi = E\Psi \rightarrow C^\dagger H C \Phi \equiv \hat{H}\Phi = E\Phi$$

Bare Hamiltonian

Shift operator depending on the relative distance

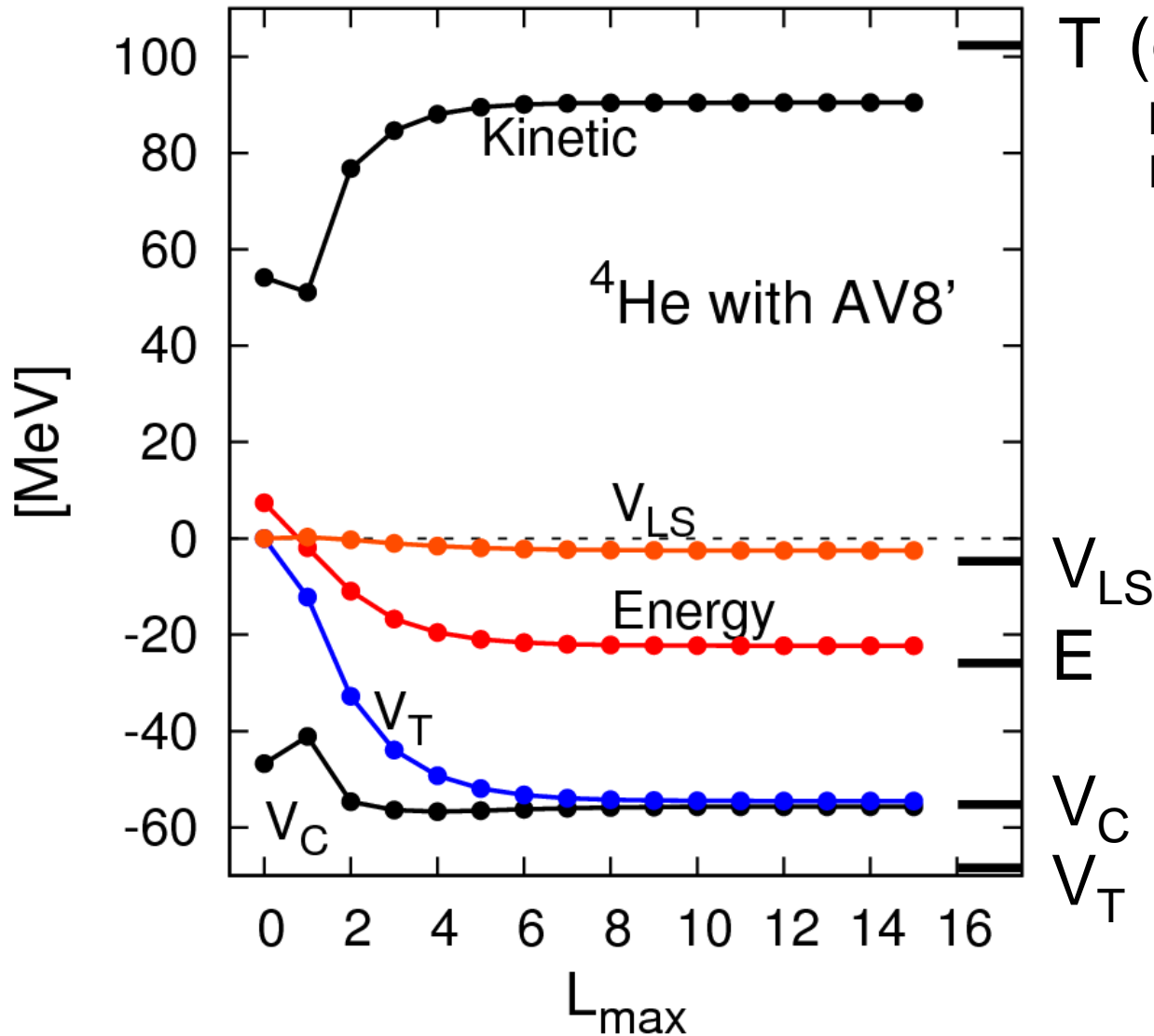
$$C = \exp(-i \sum_{i < j} g_{ij}), \quad g_{ij} = \frac{1}{2} \{ p_r \underline{s(r_{ij})} + \underline{s(r_{ij})} p_r \} \quad \vec{p} = \vec{p}_r + \vec{p}_\Omega$$

Amount of shift, variationally determined.

$$C^\dagger r C = r + s(r) + \frac{1}{2} s(r) s'(r) \dots$$

2-body cluster expansion

^4He in TOSM + short-range UCOM



T (exact)

Kamada et al.
PRC64 (Jacobi)

TM, H. Toki, K. Ikeda
PTP121(2009)511

- variational calculation

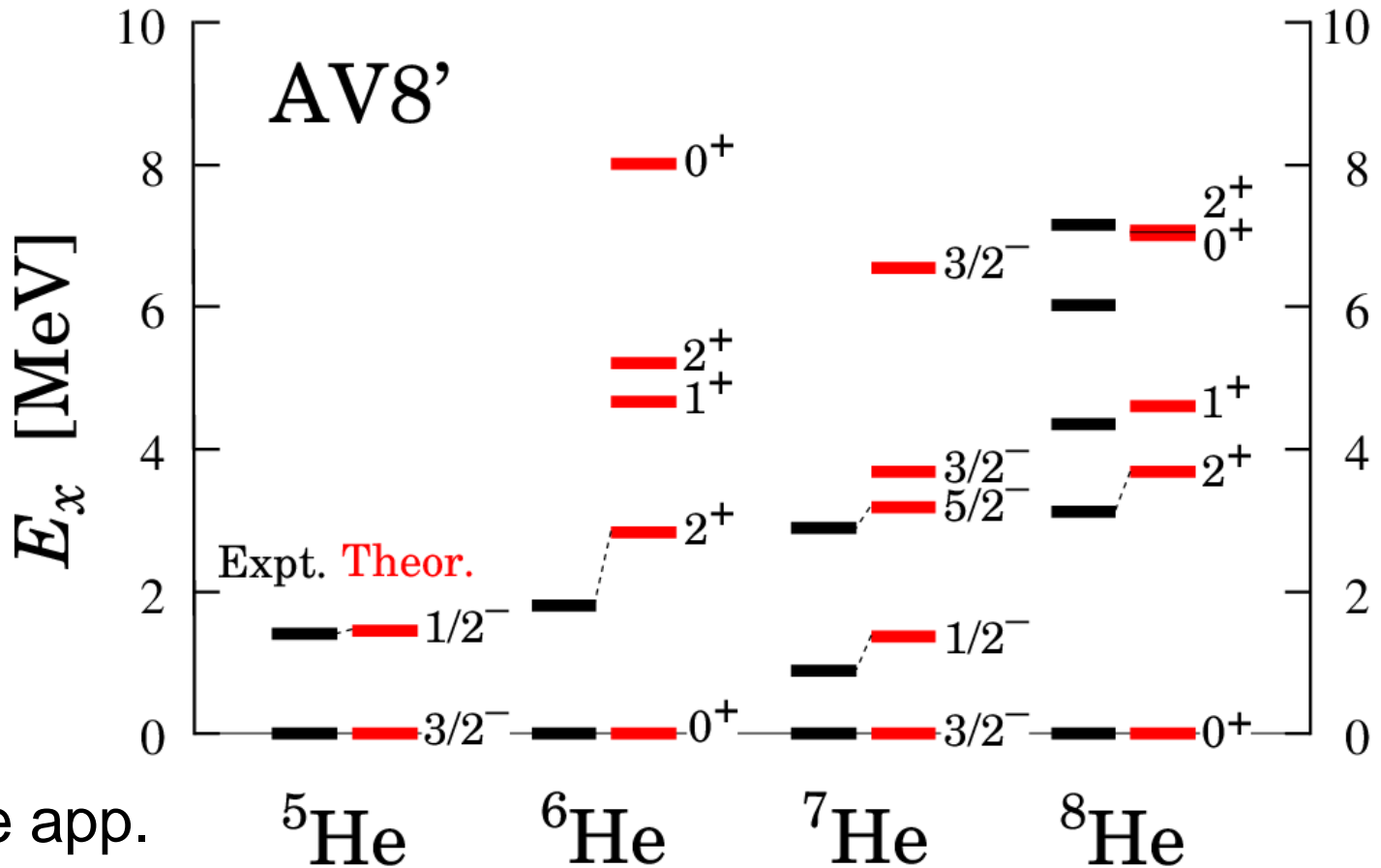
- Gaussian expansion with 9 bases

good convergence

$5\text{-}8\text{He}$ with TOSM+UCOM

TM, A. Umeya, H. Toki, K. Ikeda
 PRC84 (2011) 034315

- Excitation energies in MeV



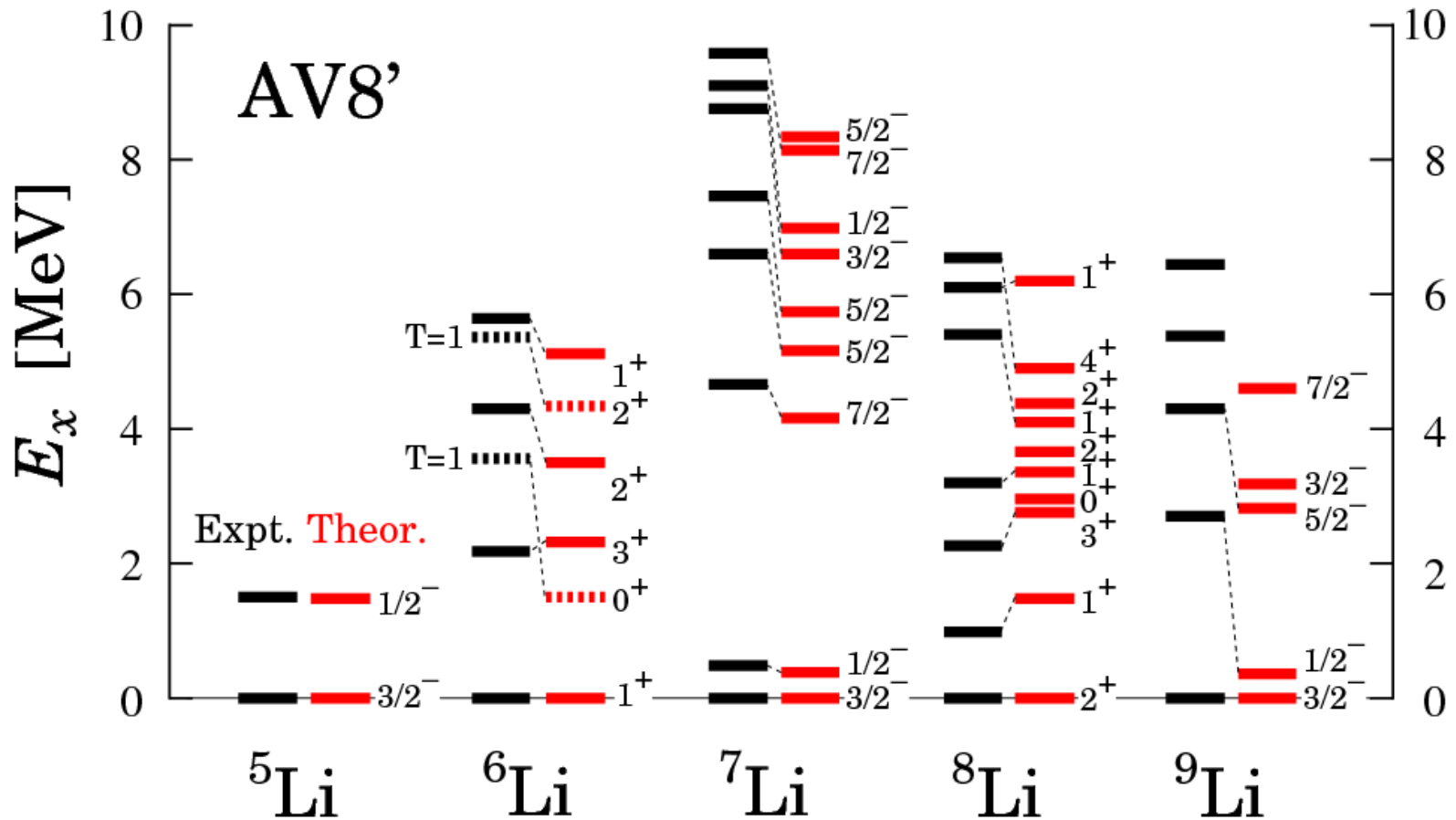
- Bound state app.
- No continuum
- No V_{NNN}

Excitation energy spectra are reproduced well

^{5-9}Li with TOSM+UCOM

TM, A. Umeya, H. Toki, K. Ikeda
PRC86(2012) 024318

- Excitation energies in MeV

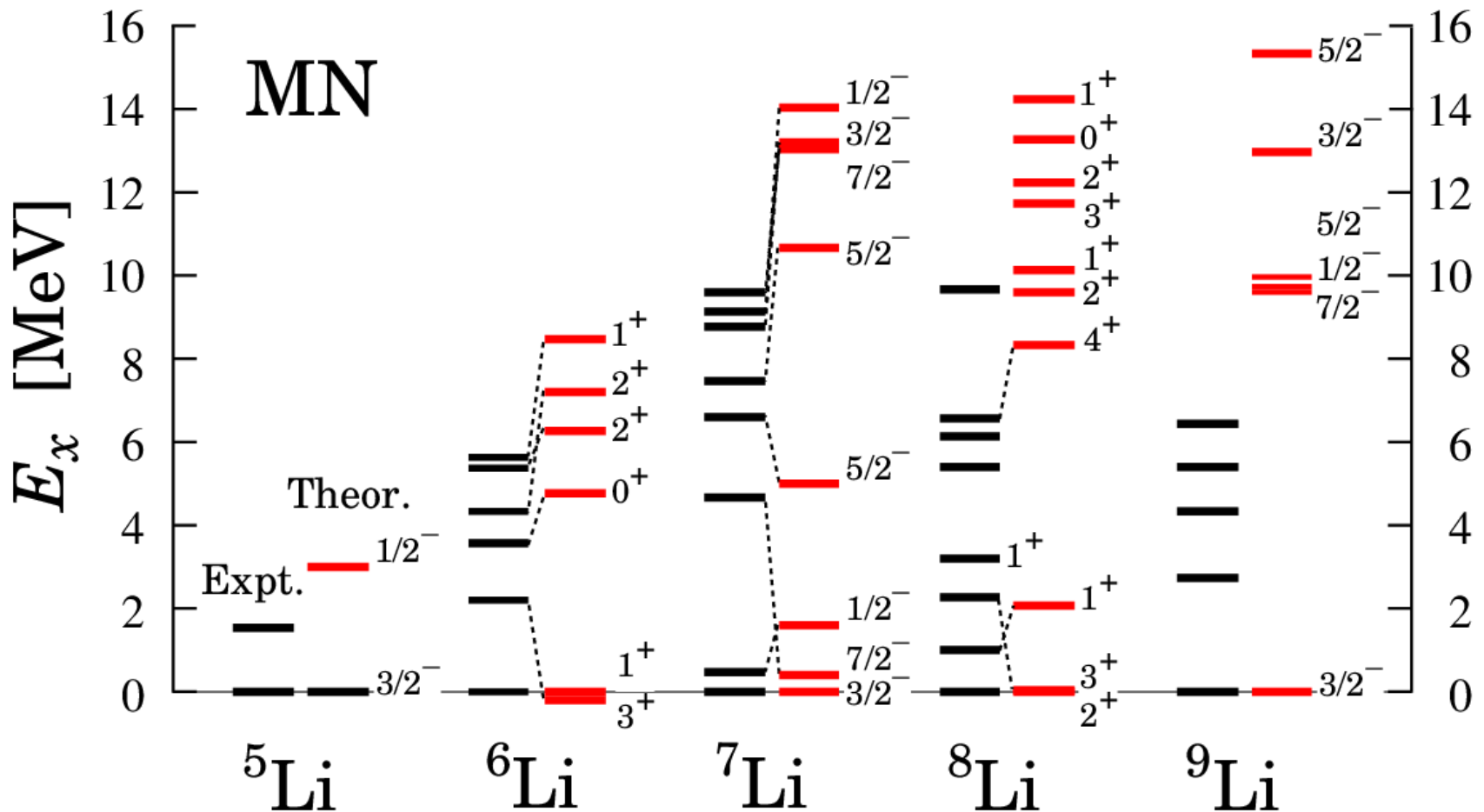


- Excitation energy spectra are reproduced well

^{5-9}Li with TOSM

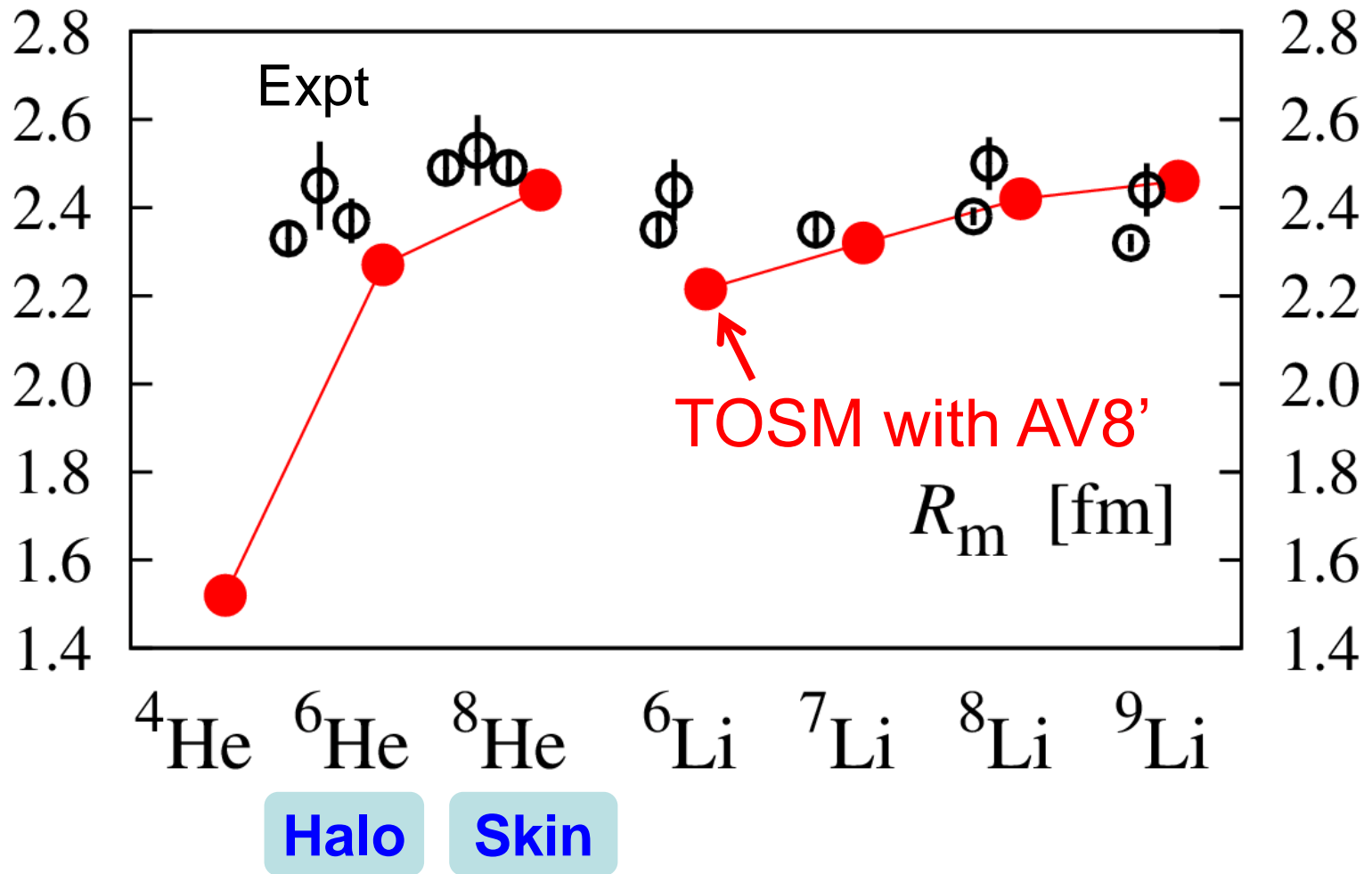
Minnesota force
NO tensor

- Excitation energies in MeV



- Too large excitation energy

Radius of He & Li isotopes



I. Tanihata et al., PLB289('92)261

O. A. Kiselev et al., EPJA 25, Suppl. 1('05)215.

A. Dobrovolsky, NPA 766(2006)1

G. D. Alkhazov et al., PRL78('97)2313

P. Mueller et al., PRL99(2007)252501

Configurations of ${}^4\text{He}$ with AV8'

$(0s_{1/2})^4$	83.0 %
$(0s_{1/2})^{-2}_{JT}(p_{1/2})^2_{JT}$ $JT=10$	2.6
$JT=01$	0.1
$(0s_{1/2})^{-2}_{10}(1s_{1/2})(d_{3/2})_{10}$	2.3
$(0s_{1/2})^{-2}_{10}(p_{3/2})(f_{5/2})_{10}$	1.9
Radius [fm]	1.54

TM, H. Toki, K. Ikeda
PTP121(2009)511

• deuteron correlation
with $(J, T)=(1, 0)$

Cf. R.Schiavilla et al. (VMC)
PRL98(2007)132501
R. Subedi et al. (JLab)
Science320(2008)1476

${}^{12}\text{C}(e, e' pN)$

S.C.Simpson, J.A.Tostevin
PRC83(2011)014605

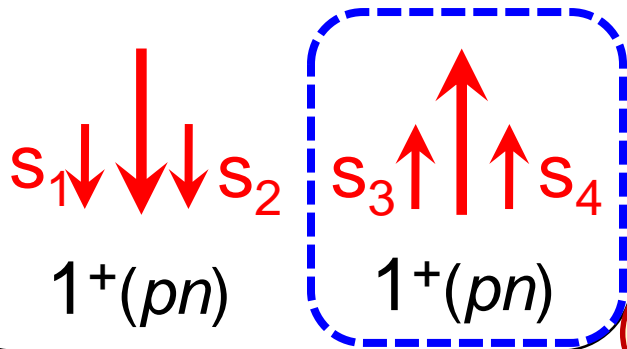
${}^{12}\text{C} \rightarrow {}^{10}\text{B} + pn$

- ${}^4\text{He}$ contains $p_{1/2}$ of “ pn -pair”
 - Same feature in ${}^5\text{He}$ - ${}^8\text{He}$ ground state

Selectivity of the tensor coupling in ${}^4\text{He}$

$$0p0h : (0s)_{00}^4 \\ \supset (0s)_{10}^2 (0s)_{10}^2$$

$$l_1 = l_2 = l_3 = l_4 = 0$$



V_T

V_T

Selectivity of S_{12}

$$\Delta L=2, \Delta S=-2$$

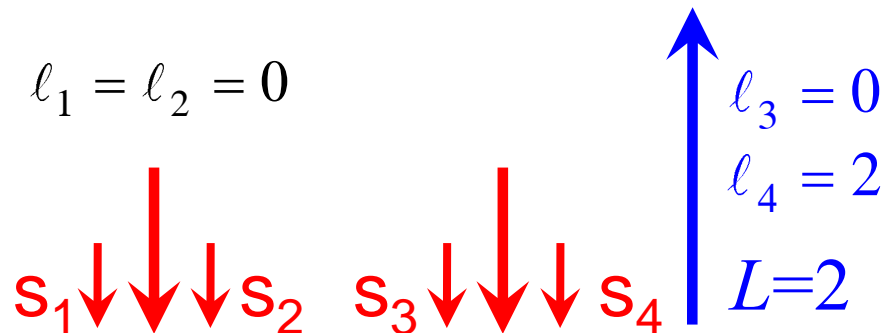
$$2p2h : (0s)_{10}^2 (0p_{1/2})_{10}^2$$

$$l_1 = l_2 = 0$$

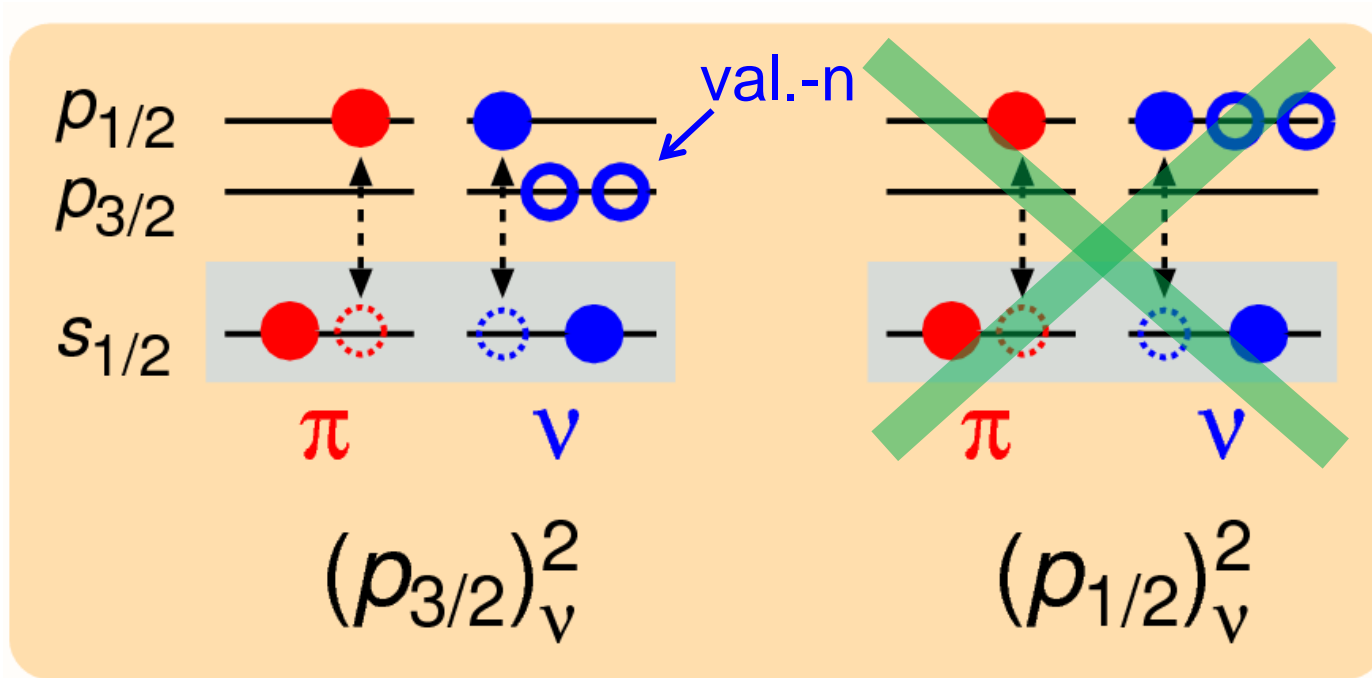


$$2p2h : (0s)_{10}^2 [(1s)(0d_{3/2})]_{10}$$

$$l_1 = l_2 = 0$$



Tensor correlation in ${}^6\text{He}$



Ground state

halo state (0^+)

Excited state

↑
Tensor correlation is **suppressed**
due to Pauli-Blocking

${}^6\text{He}$: Hamiltonian components

- Difference from ${}^4\text{He}$ in MeV

${}^6\text{He}$	0^+_1	0^+_2
n^2 config	$(p_{3/2})^2$	$(p_{1/2})^2$
$\Delta\text{Kin.}$		
$\Delta\text{Central}$		
ΔTensor		
ΔLS		

$$b_{\text{hole}} = 1.5 \text{ fm}$$

$$\hbar\omega = 18.4 \text{ MeV}$$

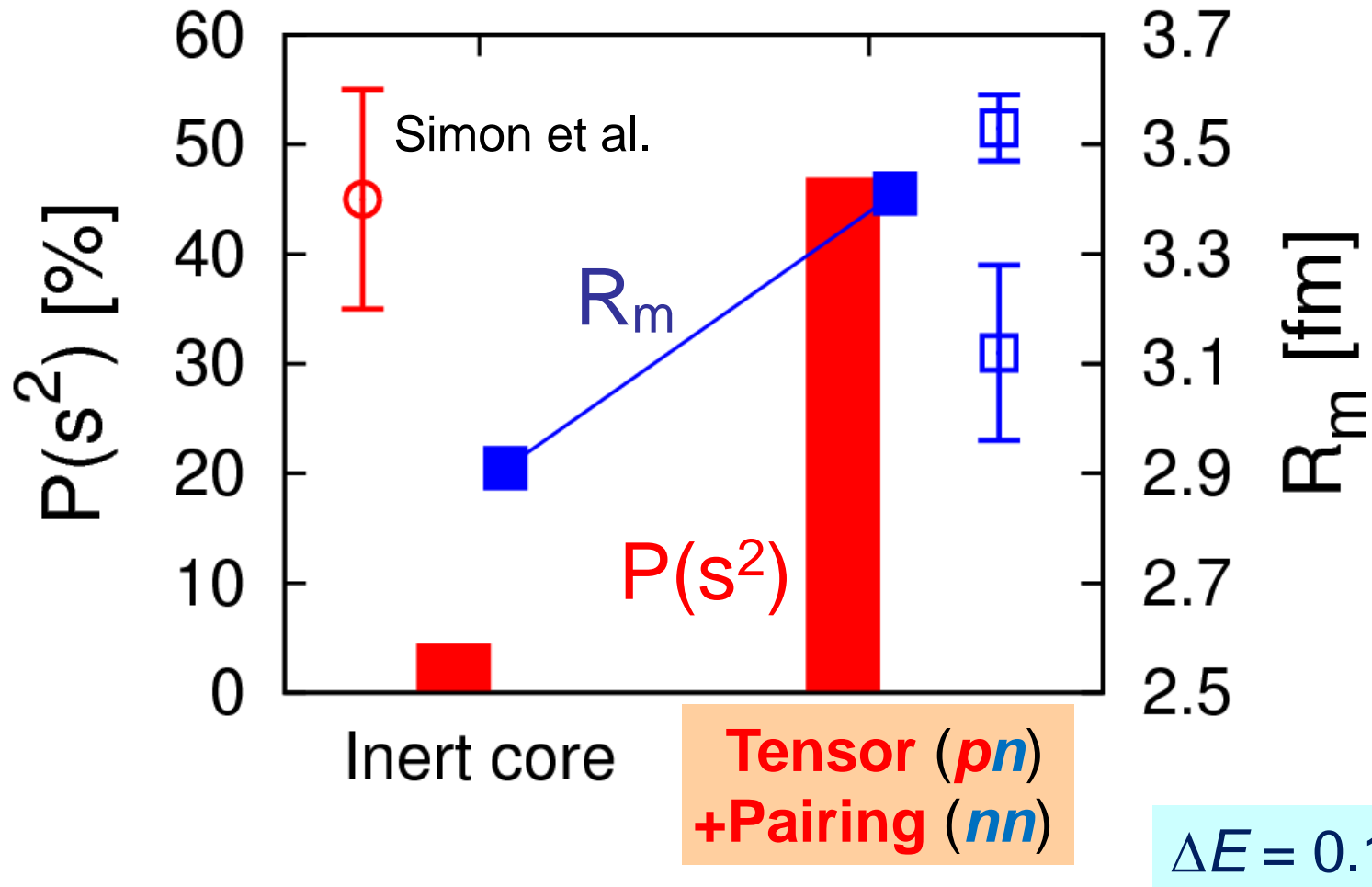
LS splitting
energy in ${}^6\text{He}$

same trend
in ${}^5\text{-}{}^8\text{He}$, ${}^{10,11}\text{Li}$

Halo formation with
large s-wave in ${}^{11}\text{Li}$

^{11}Li with $^9\text{Li}(\text{TOSM})+n+n$

$$S_{2n}=0.31 \text{ MeV}$$



Pairing correlation between halo neutrons couples $(0p)^2$ and $(1s)^2$

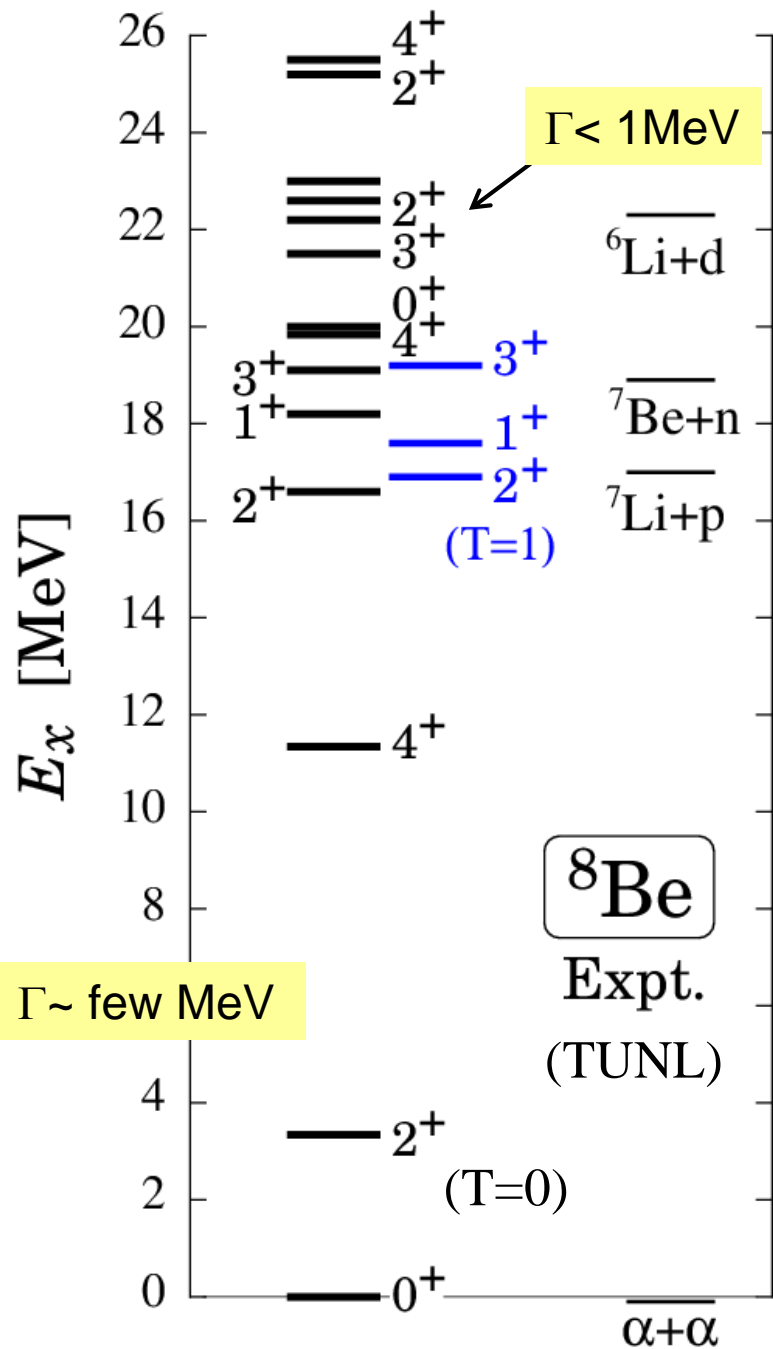
TM, K.Kato, H.Toki, K.Ikeda, PRC76(2007)024305

TM, Y.Kikuchi, K.Kato, H.Toki, K.Ikeda, PTP119(2008)561

Li isotopes: ground state configurations

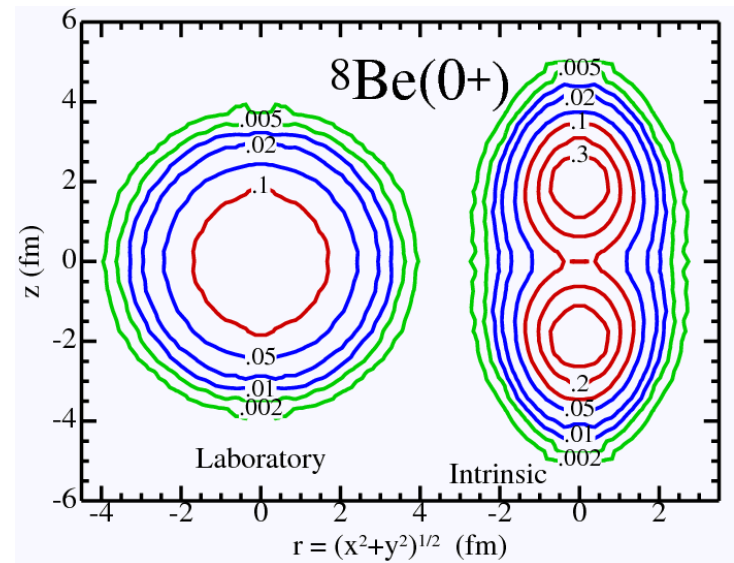
		<i>p</i> -shell config.	weight	
<i>LS</i>	${}^6\text{Li} (1^+, T=0)$	$(0p_{1/2})(0p_{3/2})$ $\Delta V_T = -12 \text{ MeV}$	43%	$S=1$
<i>jj</i>	${}^6\text{Li} (0^+, T=1)$	$(0p_{3/2})^2$ $\Delta V_T = -4 \text{ MeV}$	72%	IAS of ${}^6\text{He}$
<i>jj</i>	${}^7\text{Li} (3/2^-)$	$(0p_{3/2})^3$	48%	
<i>jj</i>	${}^8\text{Li} (2^+)$	$(0p_{3/2})^4$	41%	
<i>jj</i>	${}^9\text{Li} (3/2^-)$	$(0p_{3/2})^5$	46%	

- ${}^6\text{Li}_{\text{gs}}$... *LS* coupling → Indication of $\alpha+d$ clustering deformation



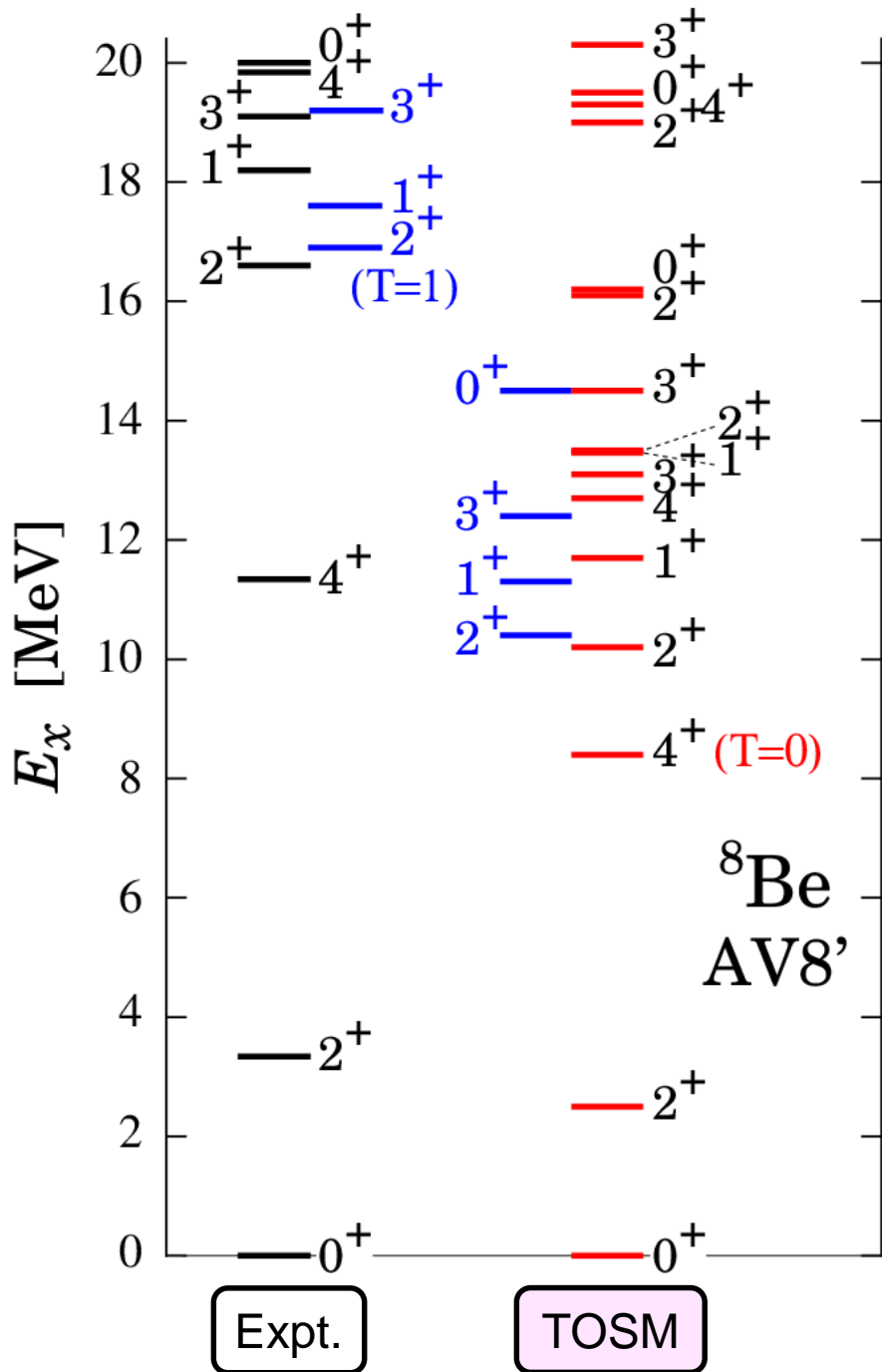
${}^8\text{Be}$ spectrum

- Green's function Monte Carlo
C. Pieper, R.B. Wiringa,
Annu.Rev.Nucl.Part.Sci.51 (2001)



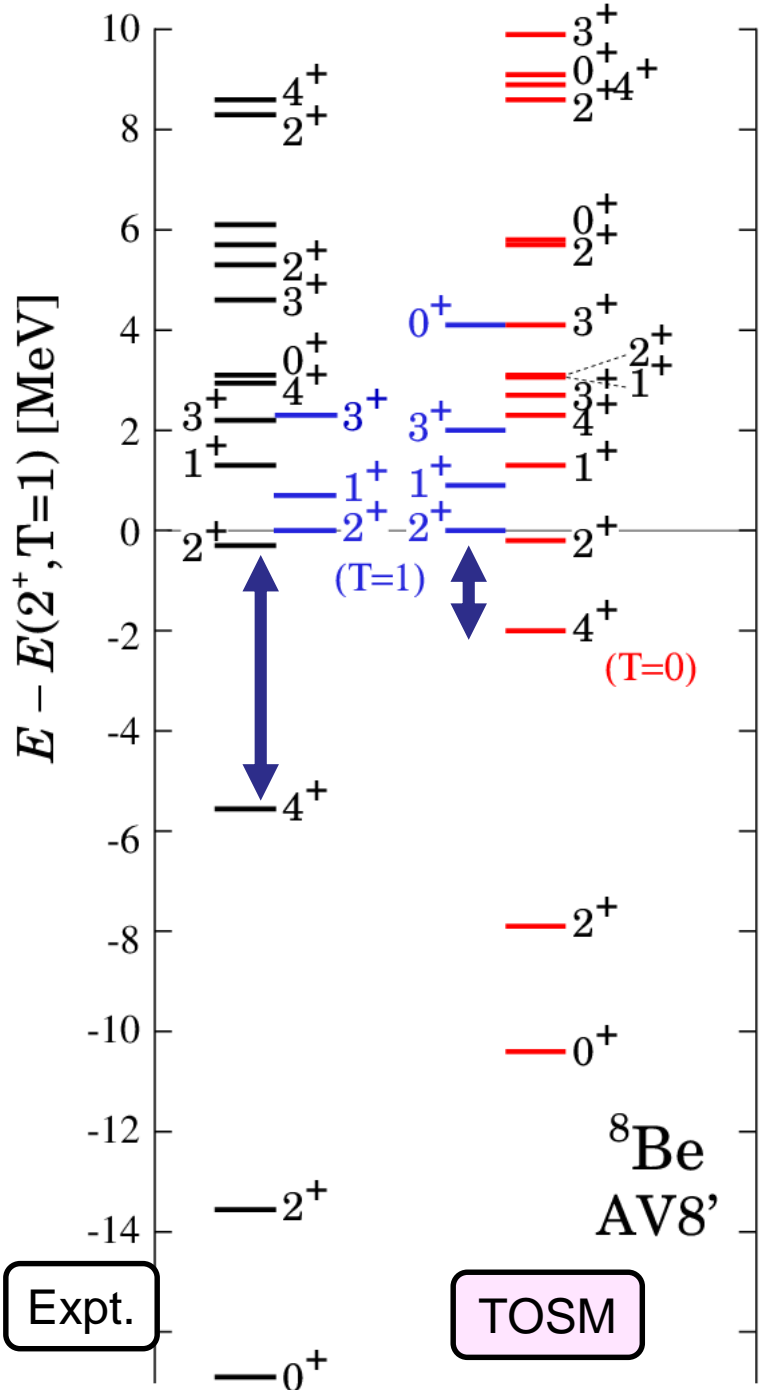
α - α structure

^8Be in TOSM – AV8' –

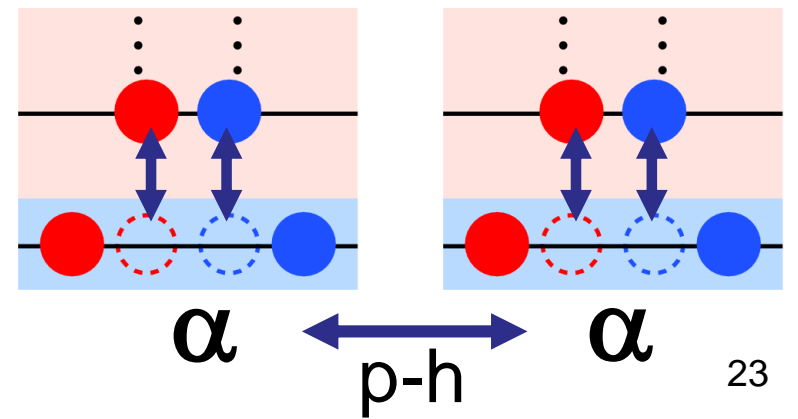


- $V_T \times 1.1, V_{LS} \times 1.4$
 - simulate ^4He benchmark (Kamada et al., PRC64)
- ground band
- highly excited states
 - small E_x
 - correct level order (T=0,1)
- $R_m(^8\text{Be})=2.21$ fm
 - 2α THSR : 2.8 fm
 - ^4He : 1.52 fm
 - ^{12}C : 2.35 fm $\sim R_m(\text{exp})$

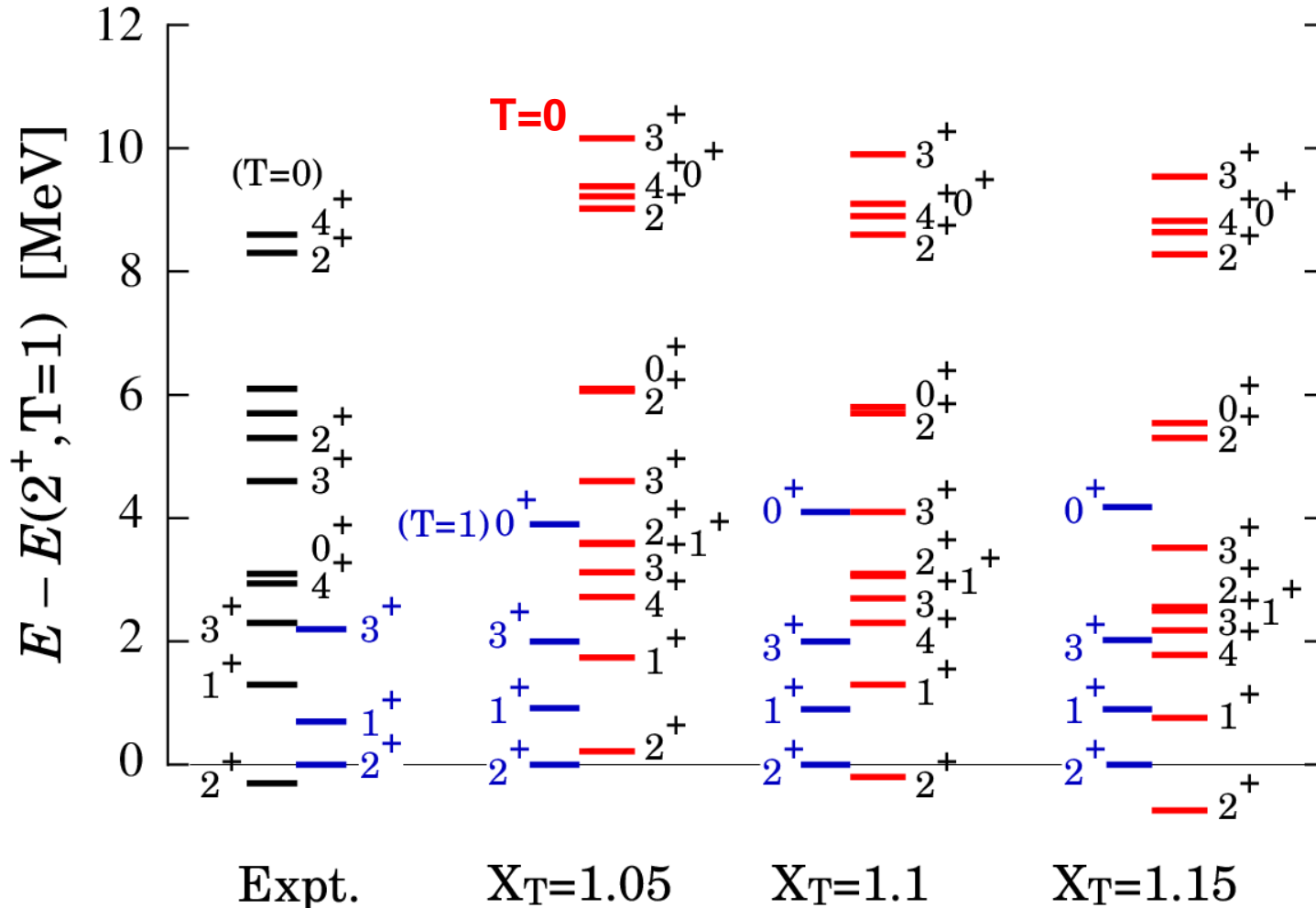
^8Be in TOSM – AV8' –



- $V_T \times 1.1, V_{LS} \times 1.4$
 - simulate ^4He benchmark (Kamada et al., PRC64)
- correct level order ($T=0, 1$)
- α : $0p0h+2p2h$ with high- k
 - 2α needs $4p4h$.
 - spatial asymptotic form of 2α



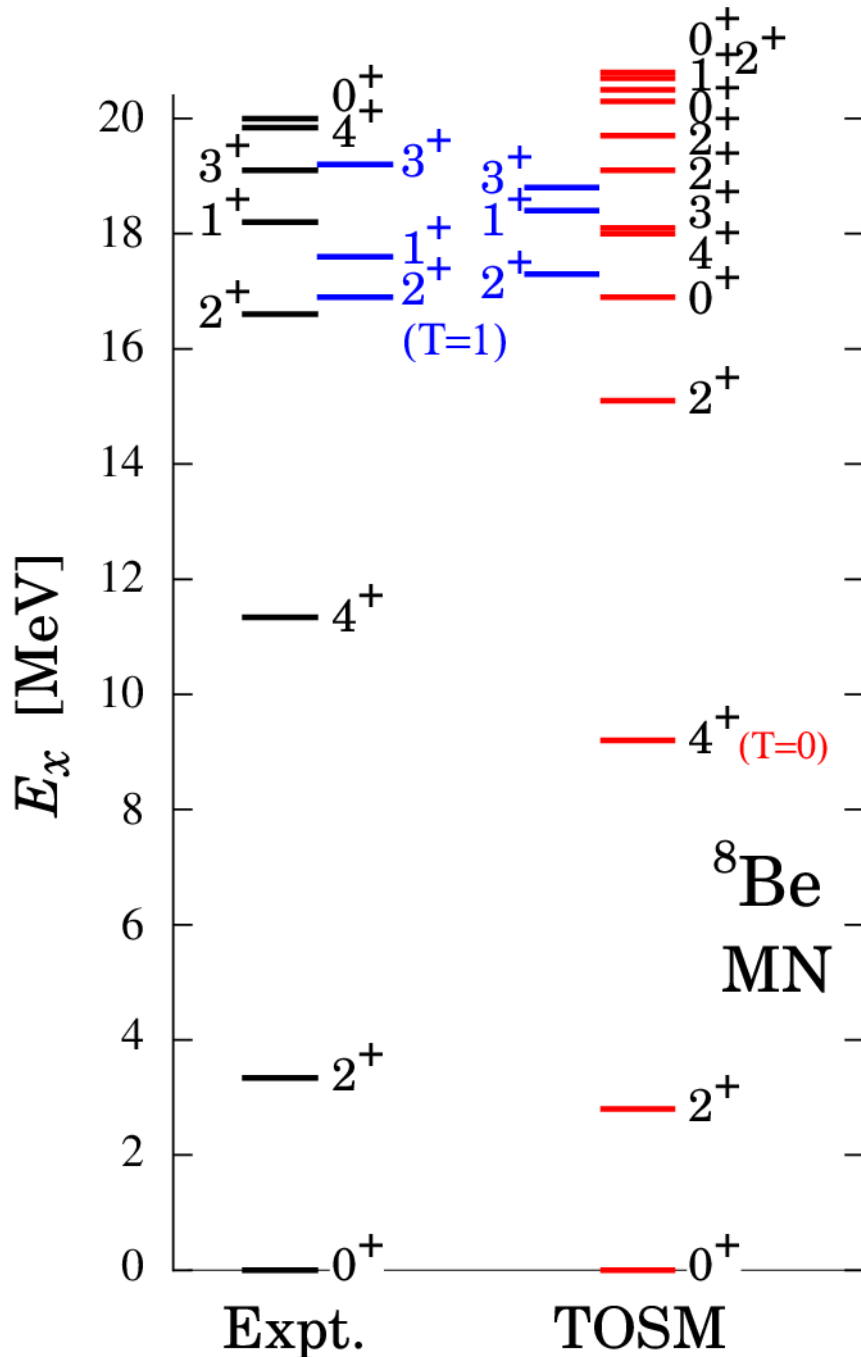
V_{tensor} dependence of ${}^8\text{Be}$



- S-wave UCOM can be simulated with $X_T \sim 1.1$ (PTP121)
- Stronger tensor correlation in **T=0 states** than T=1 states.

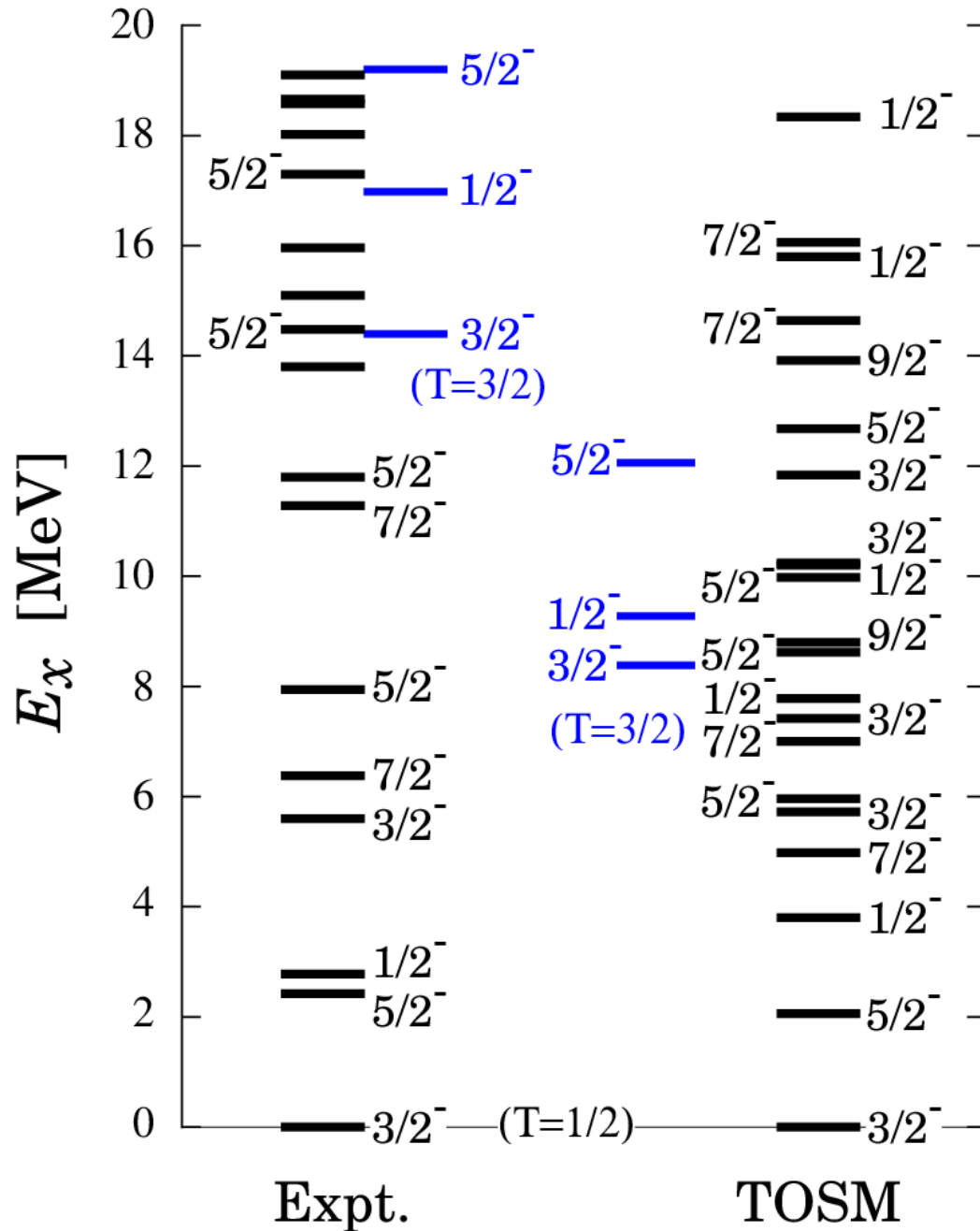
^8Be in TOSM

– Minnesota –



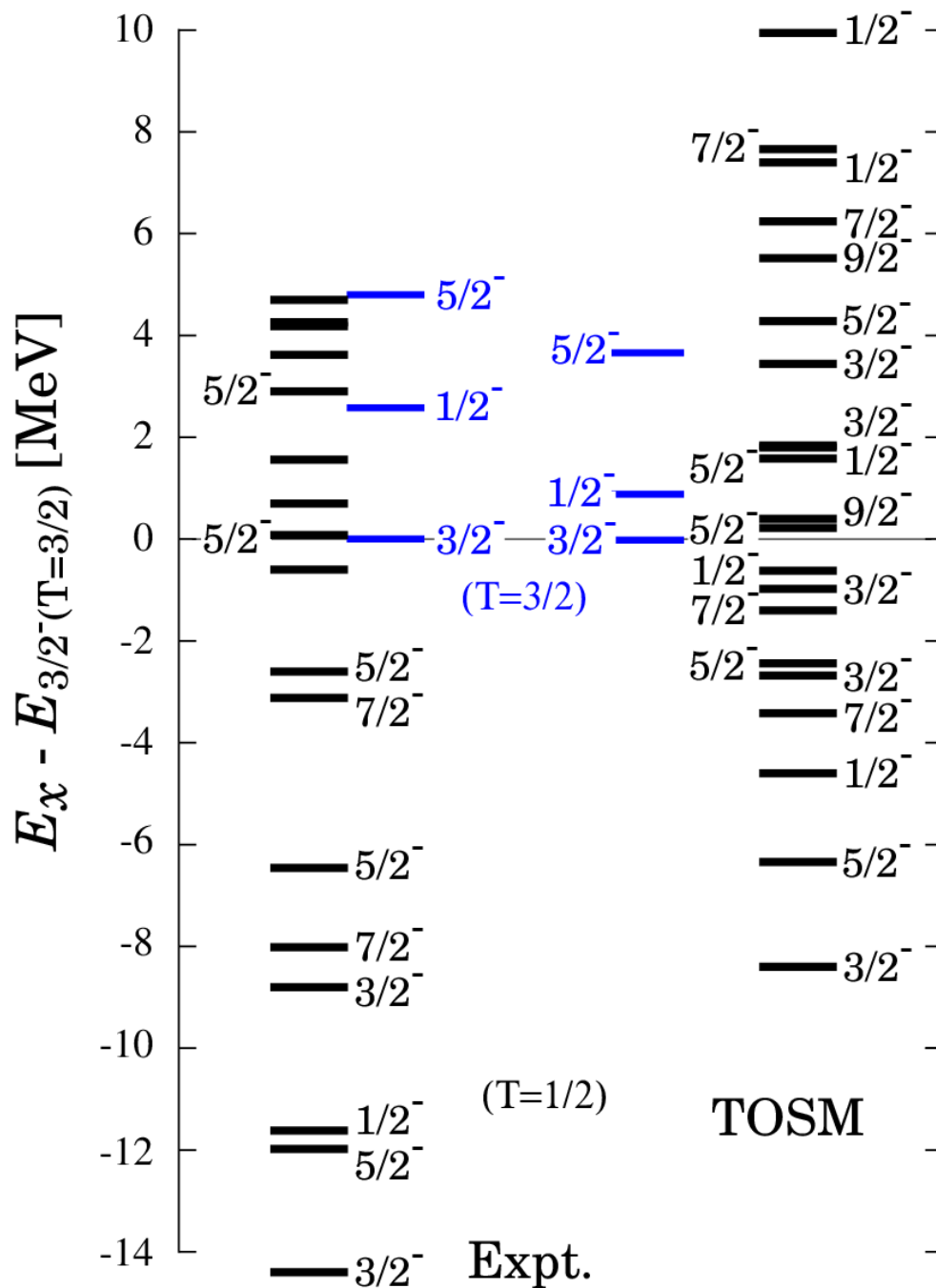
- ground band & highly excited states
 - good E_x
 - incorrect level order (0^+ , 1^+ , 3^+)
 - $E_x(T=0) < E_x(T=1)$
- Radius (R_m) is small
 - ^4He 1.39 fm
 - ^8Be 1.89 fm
 - ^{12}C 1.85 fm

${}^9\text{Be}$ in TOSM



- low-lying states
 - correct level order
- highly excited states
 - small E_x with T=1/2,3/2
- $R_m({}^9\text{Be})=2.32$ fm
 - exp: 2.45(1) fm
 - ${}^4\text{He}$: 1.52 fm
 - ${}^8\text{Be}$: 2.21 fm < 2.8 fm
(2α THSR, Funaki)
 - ${}^{12}\text{C}$: 2.35 fm $\sim R_m(\text{exp})$

${}^9\text{Be}$ in TOSM



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 - correct level order
- highly excited states
 - small E_x with $T=1/2, 3/2$
- $R_m({}^9\text{Be})=2.32$ fm
 - exp: 2.45(1) fm
 - ${}^4\text{He}$: 1.52 fm
 - ${}^8\text{Be}$: 2.21 fm < 2.8 fm
(2α THSR, Funaki)
 - ${}^{12}\text{C}$: 2.35 fm $\sim R_m(\text{exp})$

Summary

- **TOSM+UCOM** using V_{bare} .
 - Strong tensor correlation from $0p0h-2p2h$.
- He, Li, & Be isotopes
 - Energy spectra, Radius
 - ${}^4\text{He}$ contains “ **pn -pair of $p_{1/2}$** ” owing to S_{12} of V_T , which affect the spectrum of n -rich nuclei.
- ${}^8\text{Be}$, ${}^9\text{Be}$
 - Two aspects : grand state region & highly excited states.
 - For Be isotopes, indication of more configurations to describe 2α structure states.