

Hyperon resonance $\Lambda(1405)$ and the K^-pp three-body resonance



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1. Introduction

2. Method

- *Complex Scaling Method*
- *Feshbach projection on coupled-channel Complex Scaling Method (ccCSM+Feshbach method)*

3. Results

- *$\Lambda(1405)$ as a $K^{\text{bar}}N-\pi\Sigma$ system with ccCSM*
- *“ K^-pp ” as a $K^{\text{bar}}NN-\piYN$ system with ccCSM+Feshbach method*

4. Summary and future plan

1. Introduction

1. Introduction

The 10th International Workshop on the Physics of Excited Nucleons

NSTAR 2015

Y is also interesting!*

May 25(Mon)-28(Thu), 2015

$\Lambda(1405) 1/2^-$

$$I(J^P) = 0(\frac{1}{2}^-)$$

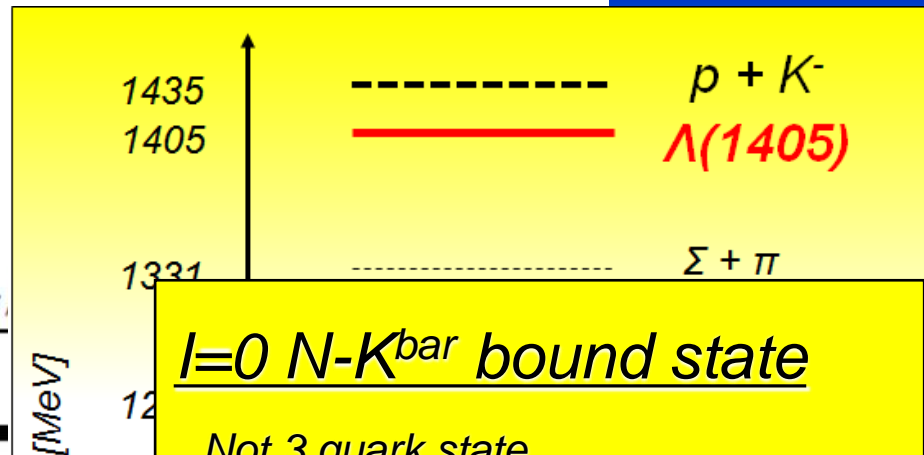
Mass $m = 1405.1^{+1.3}_{-1.0}$ MeV
Full width $\Gamma = 50 \pm 2$ MeV
Below $\bar{K}N$ threshold

$\Lambda(1405)$ DECAY MODES

$\Sigma \pi$

Fraction (Γ_i)

100 %



$I=0 N-K^{bar}$ bound state

Not 3 quark state,

← A naive quark model fails.

N. Isgar and G. Karl, Phys. Rev. D18, 4187 (1978)

But rather a molecular state

T. Hyodo and D. Jido,

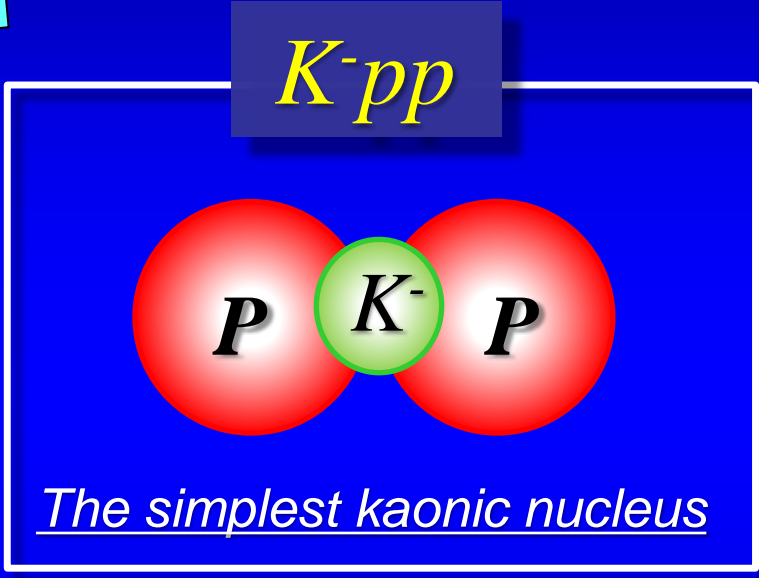
Prog. Part. Nucl. Phys. 67, 55 (2012)





$\Lambda(1405)$

... $I=0$ N - K^{bar} quasi-bound state



“Prototype of kaonic nuclei”

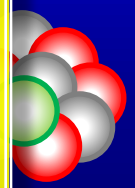


Kaonic nuclei = Exotic system??

- Doorway to dense matter
→ Chiral symmetry restoration in dense matter
- Interesting structure
- Neutron star...

Nuclear many-body system with K^-

Y. Akaishi and T. Yamazaki, PRC 52, 044005 (2002)
 A. Dote, H. Horiuchi, Y. Akaishi and T. Yamazaki, PRC70, 044313 (2004)



${}^3\text{He}K^-$, $pppK^-$,
 ${}^4\text{He}K^-$, $pppnK^-$,
 ..., ${}^8\text{Be}K^-$, ...

2. Method

- ***Complex Scaling Method***
- ***Feshbach projection on coupled-channel Complex Scaling Method***
“ccCSM+Feshbach method”

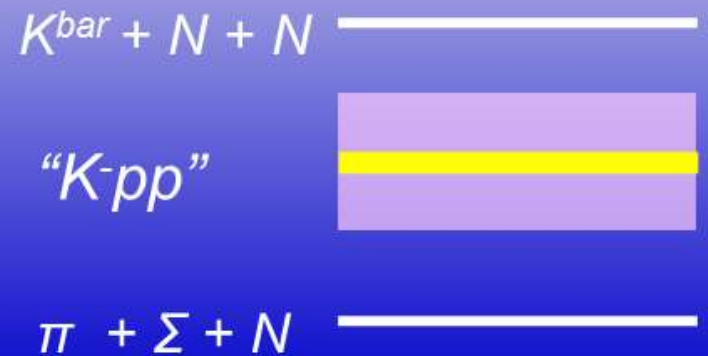
**A. D., T. Inoue, T. Myo,
PTEP 2015, 043D02 (2015)**

- $\Lambda(1405) = \text{Resonant state \& } K^{\text{bar}}N \text{ coupled with } \pi\Sigma$

- “K-pp” ... Resonant state of
 $K^{\text{bar}}NN\text{-}\pi YN$ coupled-channel system

*Doté, Hyodo, Weise, PRC79, 014003(2009). Akaishi, Yamazaki, PRC76, 045201(2007)
Ikeda, Sato, PRC76, 035203(2007). Shevchenko, Gal, Mares, PRC76, 044004(2007)
Barnea, Gal, Liverts, PLB712, 132(2012)*

- Resonant state
- Coupled-channel system



⇒ “coupled-channel
Complex Scaling Method”

Complex Scaling Method

S. Aoyama, T. Myo, K. Kato, K. Ikeda, PTP116, 1 (2006)
T. Myo, Y. Kikuchi, H. Masui, K. Kato, PPNP79, 1 (2014)

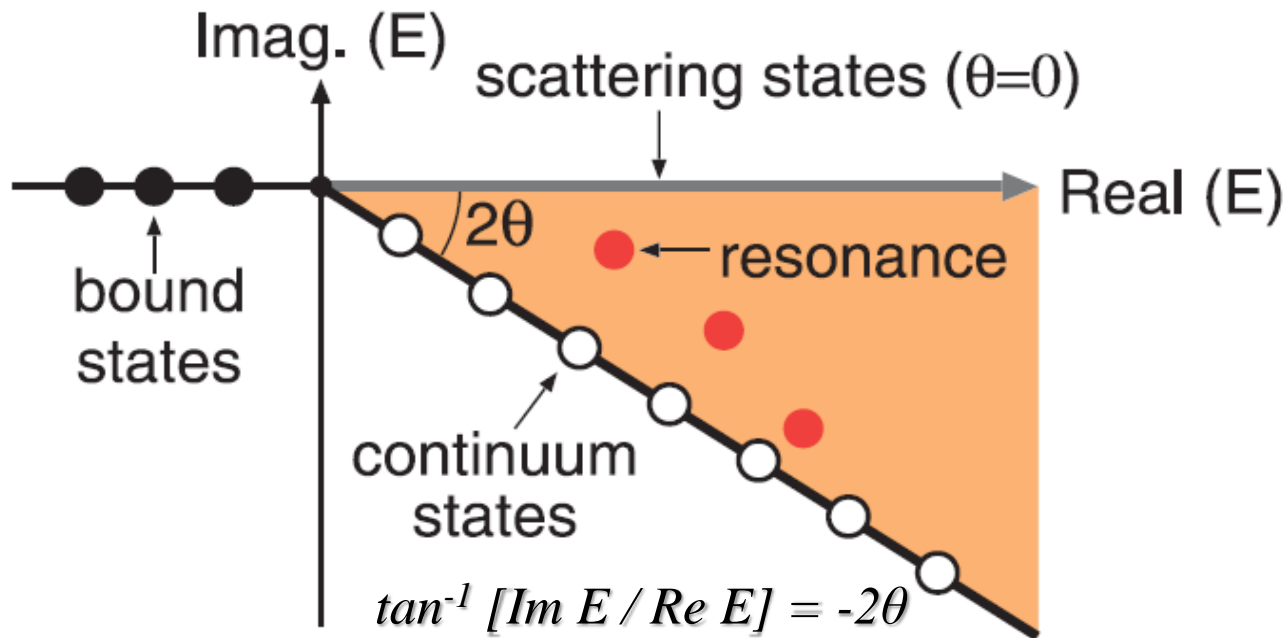
... Powerful tool for resonance study of many-body system

Complex rotation (Complex scaling) of coordinate

Resonance wave function $\rightarrow L^2$ integrable

$$U(\theta): \mathbf{r} \rightarrow \mathbf{r} e^{i\theta}, \quad \mathbf{k} \rightarrow \mathbf{k} e^{-i\theta}$$

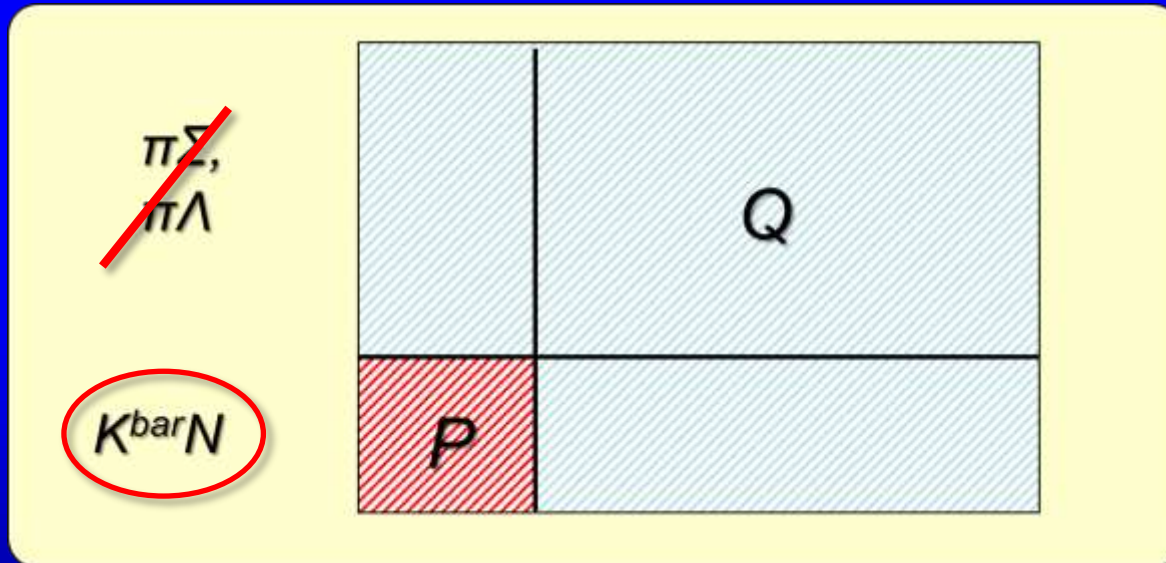
Diagonalize $H_\theta = U(\theta) H U^{-1}(\theta)$ with Gaussian base,



- Continuum state appears on 2ϑ line.
- Resonance pole is off from 2ϑ line, and independent of ϑ . (ABC theorem)

ccCSM+Feshbach method

- $\Lambda(1405) =$ *two-body* system of $K^{\text{bar}}N-\pi\Sigma$
→ Explicitly treat coupled-channel problem
- “ K^-pp ” = *three-body* system of $K^{\text{bar}}NN-\pi YN$
... High computational cost



For economical treatment of “ K^-pp ”, we construct an effective $K^{\text{bar}}N$ single-channel potential by means of *Feshbach projection* on CSM.

Formalism of ccCSM + Feshbach method

Elimination of channels by Feshbach method

Schrödinger eq.
in model space "P" and out of model space "Q"

$$\begin{pmatrix} T_P + v_P & V_{PQ} \\ V_{QP} & T_Q + v_Q \end{pmatrix} \begin{pmatrix} \Phi_P \\ \Phi_Q \end{pmatrix} = E \begin{pmatrix} \Phi_P \\ \Phi_Q \end{pmatrix}$$

Schrödinger eq. in P-space : $(T_P + U_P^{Eff}(E))\Phi_P = E\Phi_P$

Effective potential for P-space

$$U_P^{Eff}(E) = v_P + V_{PQ} G_Q(E) V_{QP}$$

Q-space Green function:

$$G_Q(E) = \frac{1}{E - H_{QQ}}$$

Extended Closure Relation in Complex Scaling Method

$$H_{QQ}^\theta |\chi_n^\theta\rangle = \varepsilon_n^\theta |\chi_n^\theta\rangle$$

$$H_{QQ}^\theta = U(\theta) H_{QQ} U^{-1}(\theta)$$

$$\int_C \sum_{R+B} |\chi_n^\theta\rangle \langle \chi_n^\theta| = 1$$



Diagonalize H_{QQ}^θ with Gaussian base,

$$\sum_n |\chi_n^\theta\rangle \langle \chi_n^\theta| \approx 1$$

Well approximated

T. Myo, A. Ohnishi and K. Kato, PTP 99, 801 (1998)
R. Suzuki, T. Myo and K. Kato, PTP 113, 1273 (2005)

Express the $G_Q(E)$ with Gaussian base using ECR

$$G_Q^\theta(E) = \frac{1}{E - H_{QQ}^\theta} \approx \sum_n |\chi_n^\theta\rangle \frac{1}{E - \varepsilon_n^\theta} \langle \chi_n^\theta|$$

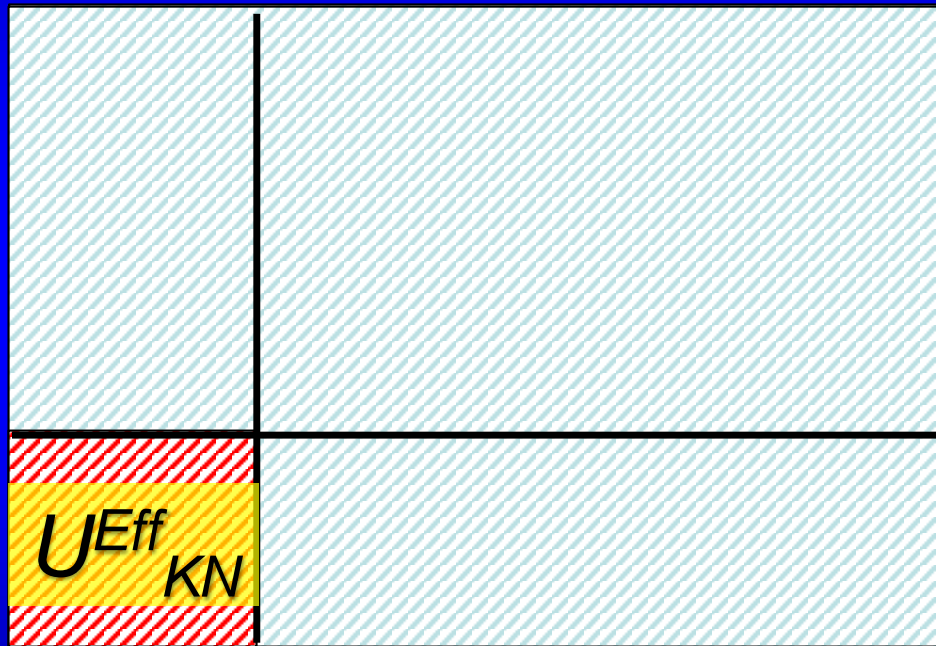


$$U_P^{Eff}(E) = v_P + V_{PQ} \underbrace{U^{-1}(\theta) G_Q^\theta(E) U(\theta)}_{G_Q(E)} V_{QP}$$

$\{ |\chi_n^\theta\rangle \}$: expanded with Gaussian base.

Applying this technique
to the **two-body** $K^{\text{bar}}N\text{-}\pi Y$ system,

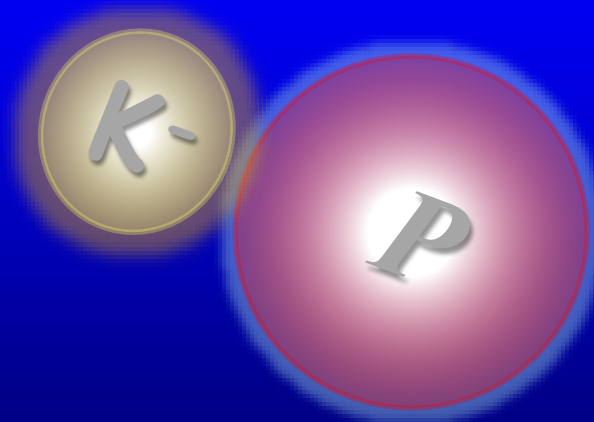
Effective single-channel $K^{\text{bar}}N$ potential
is constructed.



Using the U^{Eff}_{KN} in “ $K\text{-}pp$ ” three-body calculation,
the $K^{\text{bar}}NN\text{-}\pi YN$ coupled-channel problem is reduced
to the $K^{\text{bar}}NN$ single-channel problem.

3. Result

Hyperon resonance Λ (1405)



Chiral SU(3) potential with a Gaussian form

A. D., T. Inoue, T. Myo, Nucl. Phys. A 912, 66 (2013)

• **Anti-kaon = Nambu-Goldstone boson**

⇒ Chiral SU(3)-based $K^{\text{bar}}N$ potential

- Weinberg-Tomozawa term of effective chiral Lagrangian
- Gaussian form in r -space
- Semi-rela. / Non-rela.
- Based on Chiral SU(3) theory
→ **Energy dependence**

A non-relativistic potential (NRv2c)

$$V_{ij}^{(I=0,1)}(r) = -\frac{C_{ij}^{(I=0,1)}}{8f_{\pi}^2} (\omega_i + \omega_j) \sqrt{\frac{1}{m_i m_j}} g_{ij}(r)$$

$$g_{ij}(r) = \frac{1}{\pi^{3/2} d_{ij}^3} \exp\left[-(r/d_{ij})^2\right] : \text{Gaussian form}$$

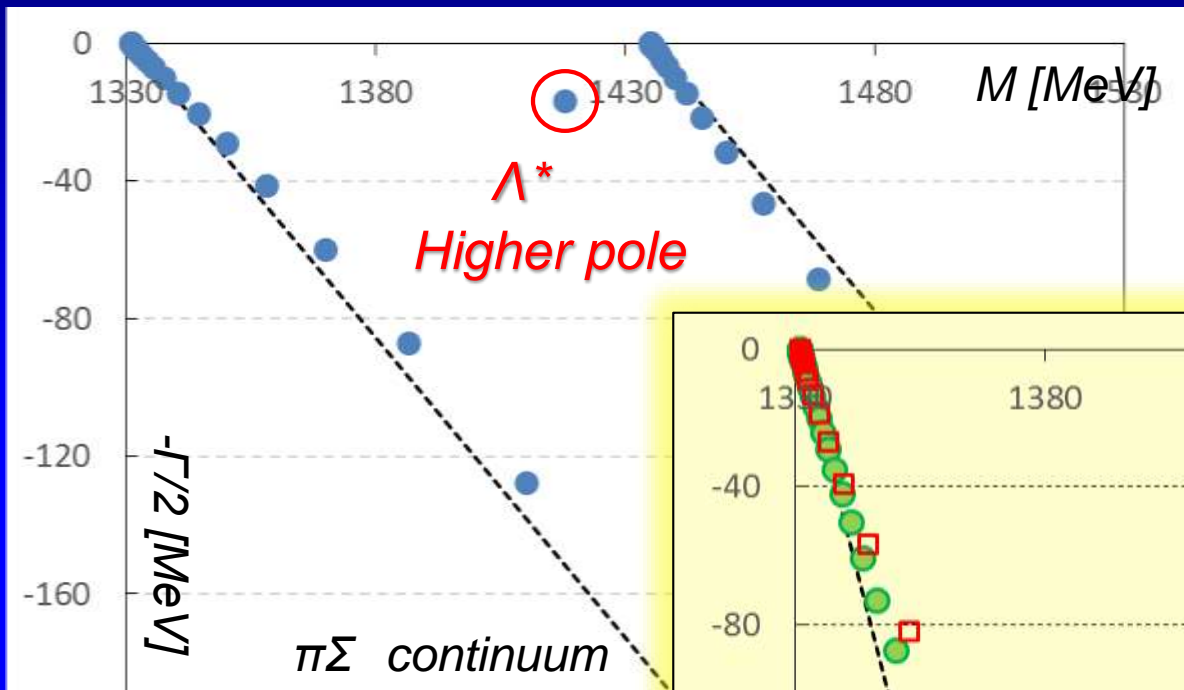
ω_i : meson energy

Constrained by $K^{\text{bar}}N$ scattering length

$$a_{KN(I=0)} = -1.70 + i0.67 \text{ fm}, \quad a_{KN(I=1)} = 0.37 + i0.60 \text{ fm}$$

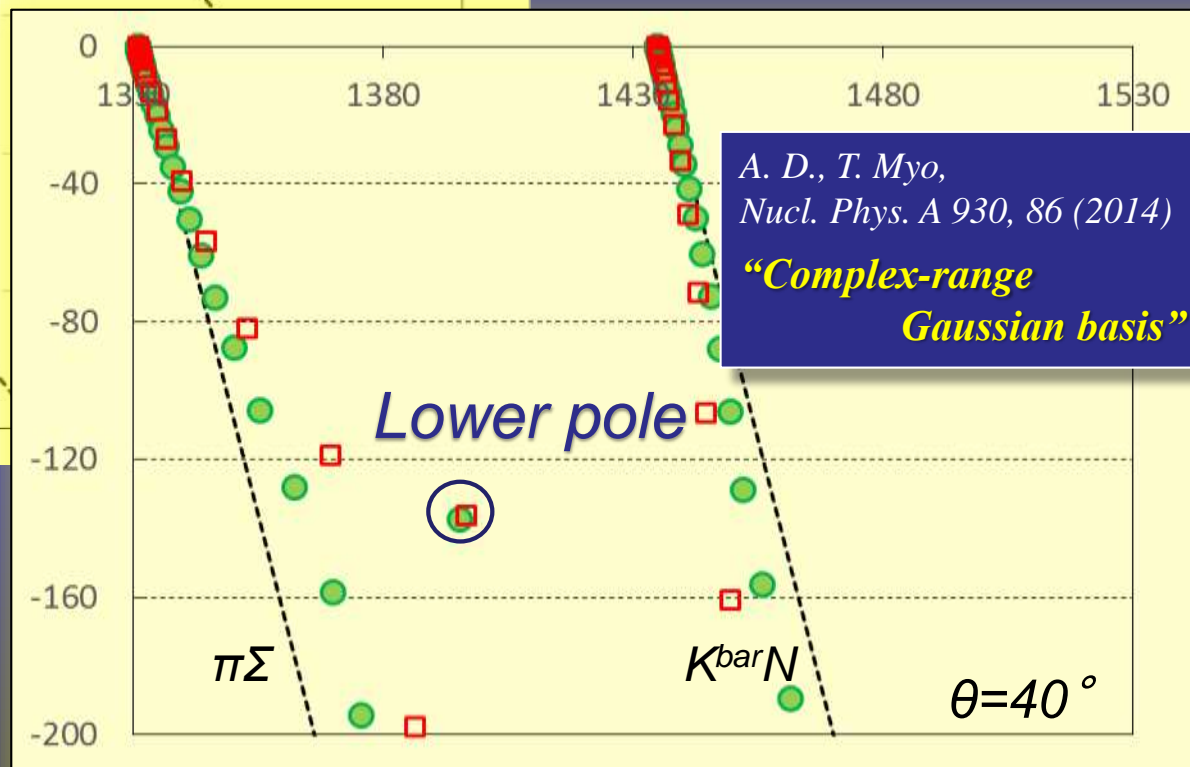
A. D. Martin, NPB179, 33(1979)

Poles of $I=0$ $K^{\text{bar}}N$ - $\pi\Sigma$ system found by ccCSM



$K^{\text{bar}}N$ potential:
a chiral $SU(3)$ potential
(NRv2, $f_\pi=110$)

A. D., T. Inoue, T. Myo,
Nucl. Phys. A 912, 66 (2013)



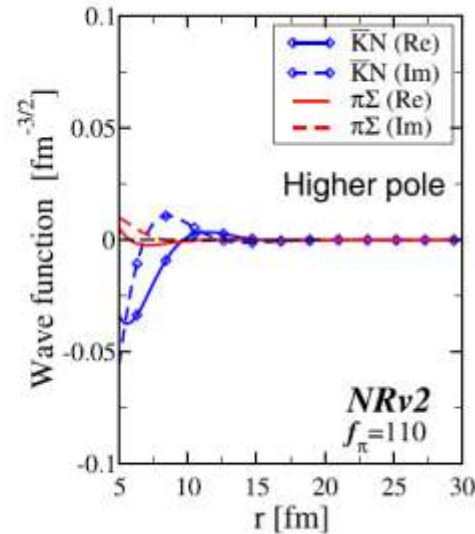
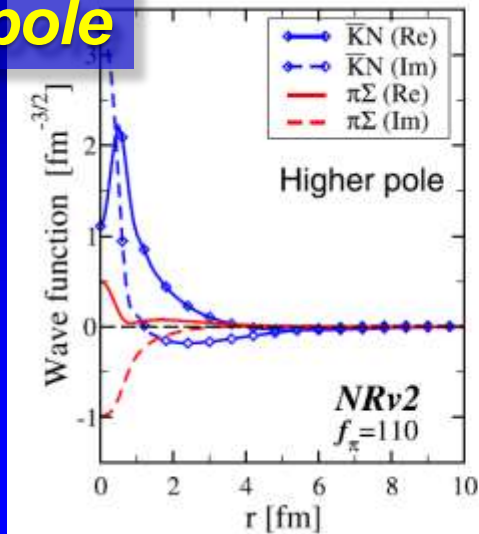
A. D., T. Myo,
Nucl. Phys. A 930, 86 (2014)
"Complex-range
Gaussian basis"

Double-pole structure of $\Lambda(1405)$ is confirmed!

ccCSM wfnc. of double pole

A. D., T. Myo,
Nucl. Phys. A 930, 86 (2014)

Higher pole

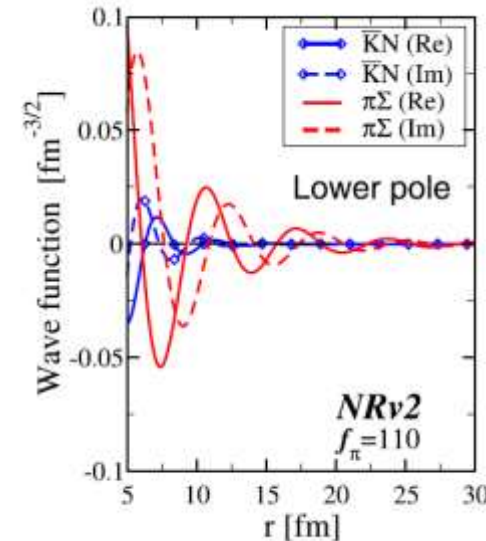
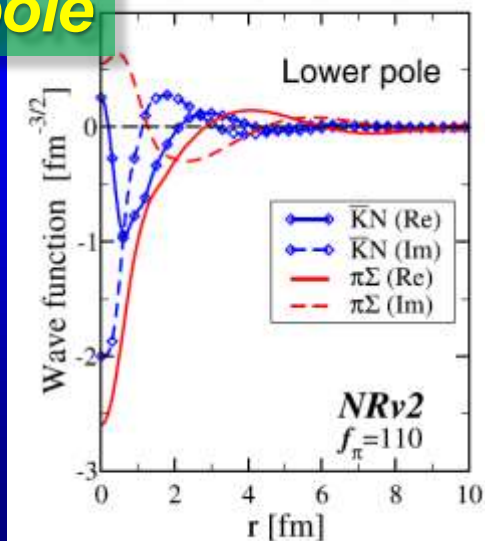


Norm ($K^{\text{bar}}N$)
 $1.115+0.098i$

Norm ($\pi\Sigma$)
 $-0.115-0.098i$

$K^{\text{bar}}N$ dominant

Lower pole



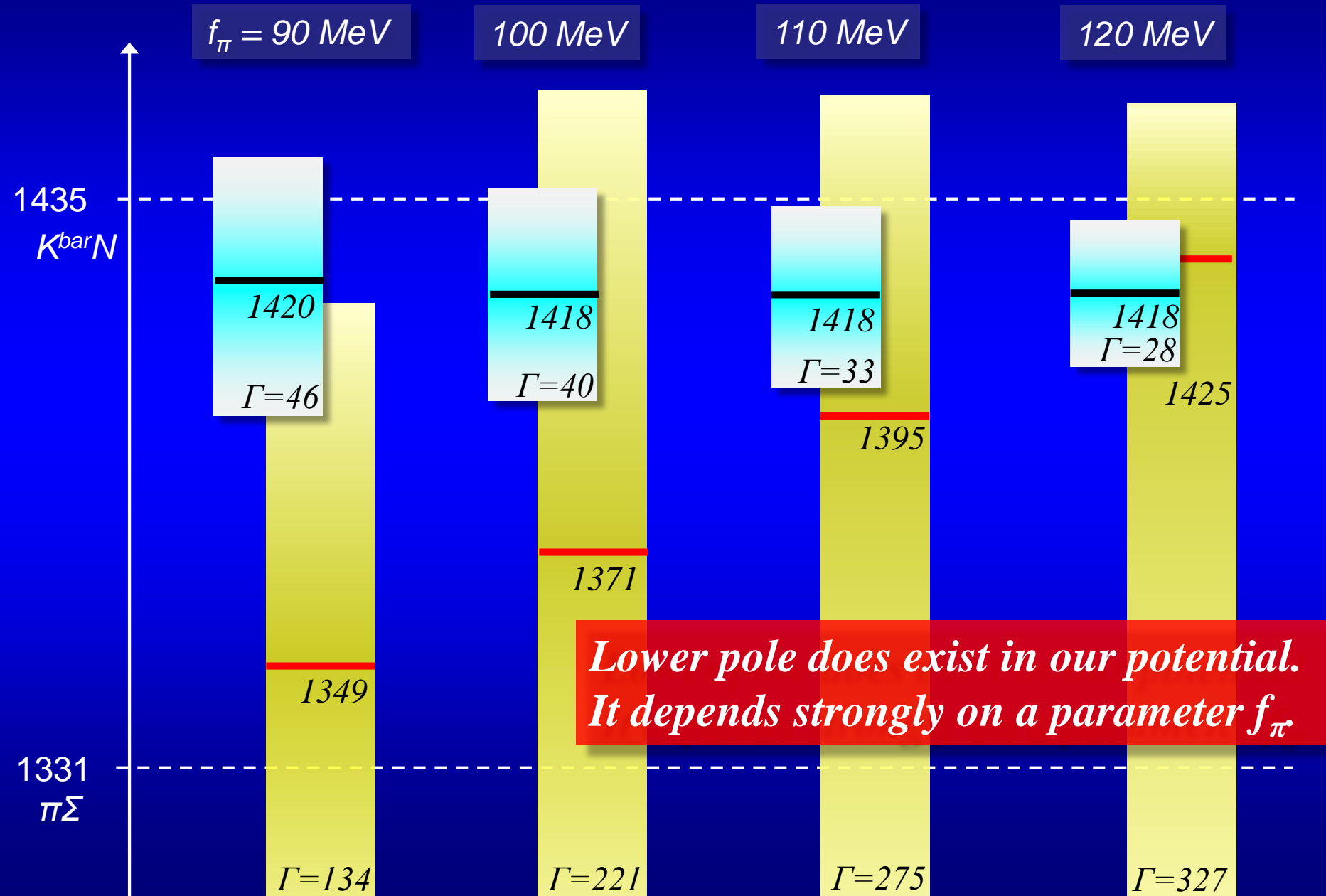
Norm ($K^{\text{bar}}N$)
 $0.097+0.154i$

Norm ($\pi\Sigma$)
 $0.903-0.154i$

$\pi\Sigma$ dominant

Double-pole structure of $\Lambda(1405)$

A. D., T. Myo,
Nucl. Phys. A 930, 86 (2014)



3. Result

Three-body “K-pp” resonance

A Venn diagram consisting of three overlapping circles. The left and right circles are purple, and the middle circle is yellow. Each circle contains a letter: 'P' in the left circle, 'K' in the middle circle, and 'P' in the right circle. The circles overlap in the center, and each pair of circles also overlaps.

$$\text{“K-pp”} = K^{\text{bar}}NN - \pi\Sigma N - \pi\Lambda N \quad (J^\pi = 0^-, T=1/2)$$

Apply ccCSM + Feshbach method to K^-pp

“ K^-pp ” ... $K^{bar}NN - \pi\Sigma N - \pi\Lambda N$ ($J^\pi=0^-, T=1/2$)

For the two-body system, $P = K^{bar}N$, $Q = \pi Y$

$$\begin{matrix} V(K^{bar}N - \pi Y; I=0,1) \\ V(\pi Y - \pi Y'; I=0,1) \end{matrix} \xrightarrow{\text{Feshbach + ccCSM}} U_{K^{bar}N(I=0,1)}^{Eff}(E)$$

- Schrödinger eq. for $K^{bar}NN$ channel :

$$\left(T_{K^{bar}NN} + V_{NN} + \sum_{i=1,2} U_{K^{bar}N_i(I)}^{Eff}(E_{K^{bar}N}) \right) \Phi_{K^{bar}NN} = E \Phi_{K^{bar}NN}$$

- Trial wave function

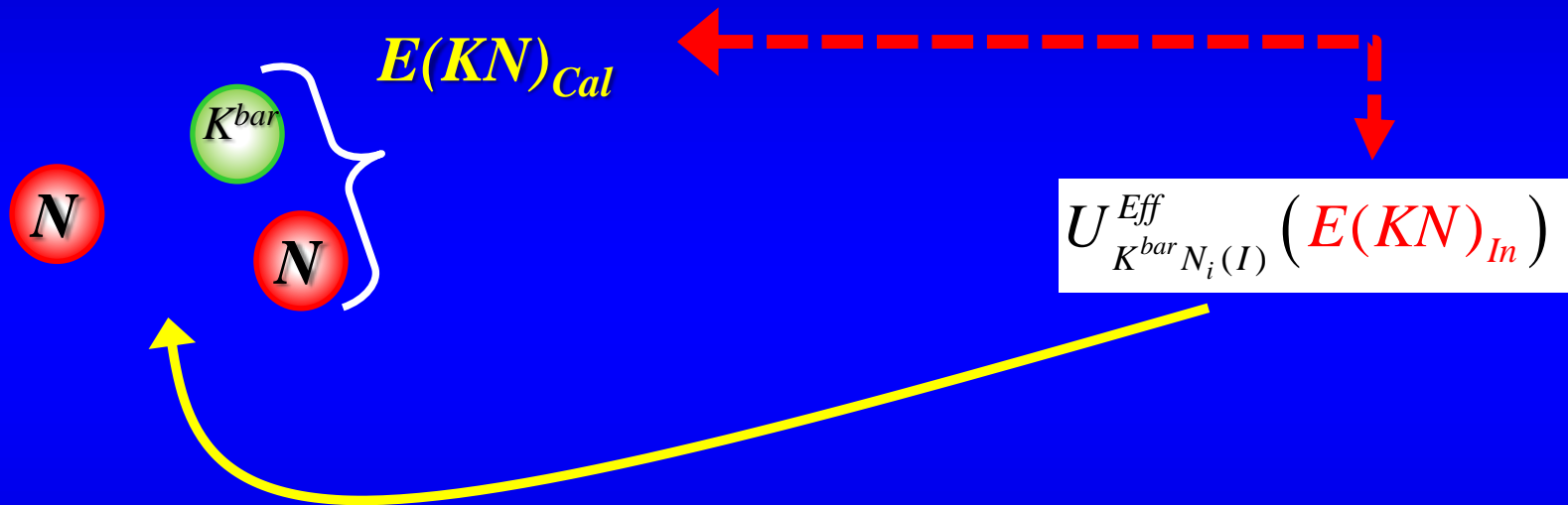
$$\begin{aligned} |"K^-pp"> &= \sum_a C_a^{(KNN,1)} \left\{ G_a^{(KNN,1)}(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) + G_a^{(KNN,1)}(-\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) \right\} |S_{NN}=0> \left[[K[NN]_1]_{T=1/2} \right] \\ &+ \sum_a C_a^{(KNN,2)} \left\{ G_a^{(KNN,2)}(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) - G_a^{(KNN,2)}(-\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) \right\} |S_{NN}=0> \left[[K[NN]_0]_{T=1/2} \right] \end{aligned} \quad \begin{matrix} \text{Ch. 1: } K^{bar}NN, \quad NN:1E \\ \text{Ch. 2: } K^{bar}NN, \quad NN:1O \end{matrix}$$

- Basis function = Correlated Gaussian
...including 3-types Jacobi-coordinates

$$G_a^{(KNN,i)}(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) = N_a^{(KNN,i)} \exp \left[-(\mathbf{x}_1^{(3)}, \mathbf{x}_2^{(3)}) A_a^{(KNN,i)} \begin{pmatrix} \mathbf{x}_1^{(3)} \\ \mathbf{x}_2^{(3)} \end{pmatrix} \right]$$

Self-consistency for **complex** $K^{\text{bar}}N$ energy

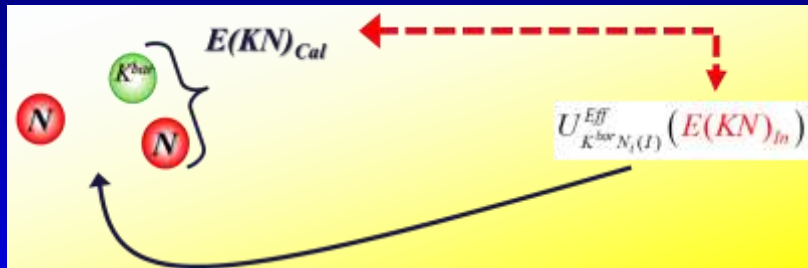
Effective $K^{\text{bar}}N$ potential has energy dependence...



- $E(KN)_{\text{In}}$: assumed in the $K^{\text{bar}}N$ potential
- $E(KN)_{\text{Cal}}$: calculated with the obtained K - pp

When $E(KN)_{\text{In}} = E(KN)_{\text{Cal}}$,
a self-consistent solution is obtained.

Self-consistency for **complex** K^{bar} N energy



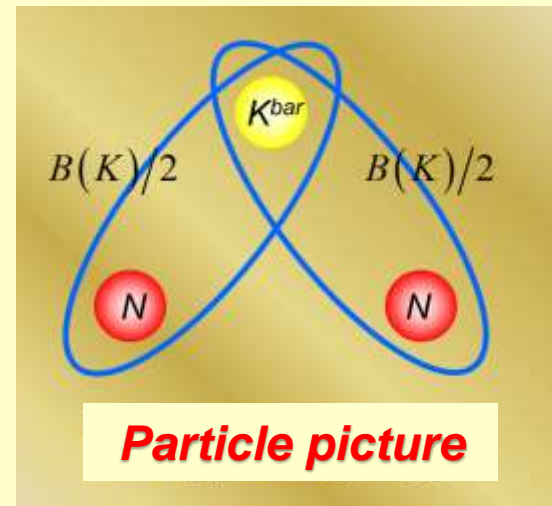
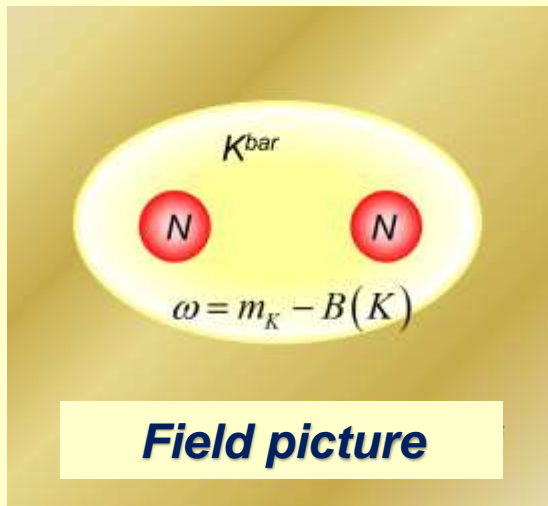
How to determine the two-body energy in the three-body system?

A. D., T. Hyodo, W. Weise,
PRC79, 014003 (2009)

1. Kaon's binding energy: $B(K) \equiv -\left\{ \langle H \rangle - \langle H_{NN} \rangle \right\}$ H_{NN} : Hamiltonian of two nucleons

2. Define a K^{bar} N-bond energy in two ways

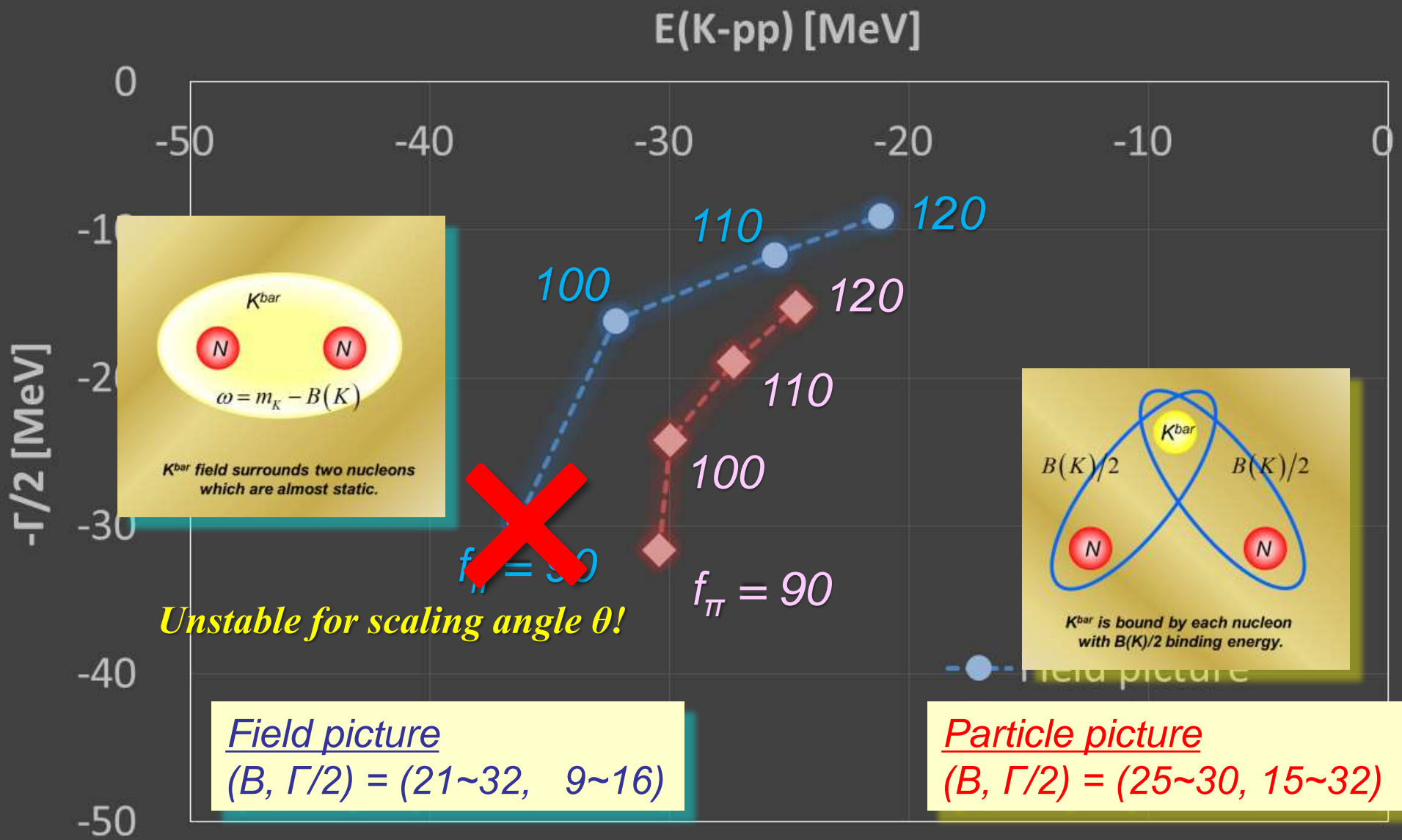
$$E_{KN} = M_N + \omega = \begin{cases} M_N + m_K - B(K) & : \text{Field picture} \\ M_N + m_K - B(K)/2 & : \text{Particle picture} \end{cases}$$



Self-consistent results

$f_\pi = 90 \sim 120 \text{ MeV}$

$NN \text{ pot.} : \text{Av18 (Central)}$
 $K^{\text{bar}}N \text{ pot.} : \text{NRv2c potential}$
 $(f_\pi = 90 - 120 \text{ MeV})$



Field picture
 $(B, \Gamma/2) = (21 \sim 32, 9 \sim 16)$

Particle picture
 $(B, \Gamma/2) = (25 \sim 30, 15 \sim 32)$

NN correlation density

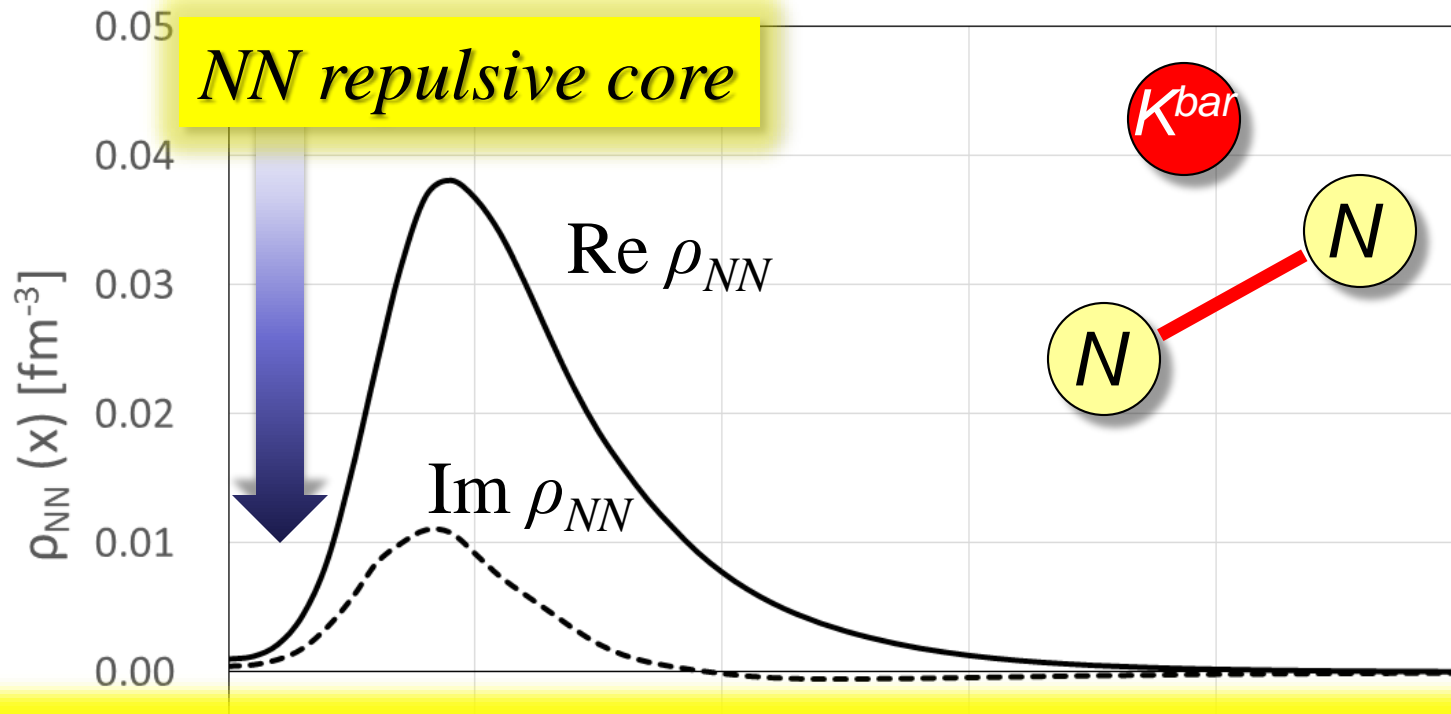
NN pot. : Av18 (Central)
 $K^{\text{bar}}N$ pot. : NRv2c potential
 $f_{\pi}=110$, Particle pict.

Correlation density in Complex Scaling Method

$$\rho_{NN,\theta}(\mathbf{x}) = \delta^3(\hat{\mathbf{r}}_{NN,\theta} - \mathbf{x})$$
$$\hat{\mathbf{r}}_{NN,\theta} = \hat{\mathbf{r}}_{NN} e^{i\theta}$$



$$\rho_{NN}(\mathbf{x}) \equiv \langle \Phi_{\theta} | \rho_{NN,\theta}(\mathbf{x}) | \Phi_{\theta} \rangle$$
$$= e^{-3i\theta} \int d^3\mathbf{R} \Phi_{\theta}^2(\mathbf{x}e^{-i\theta}, \mathbf{R})$$

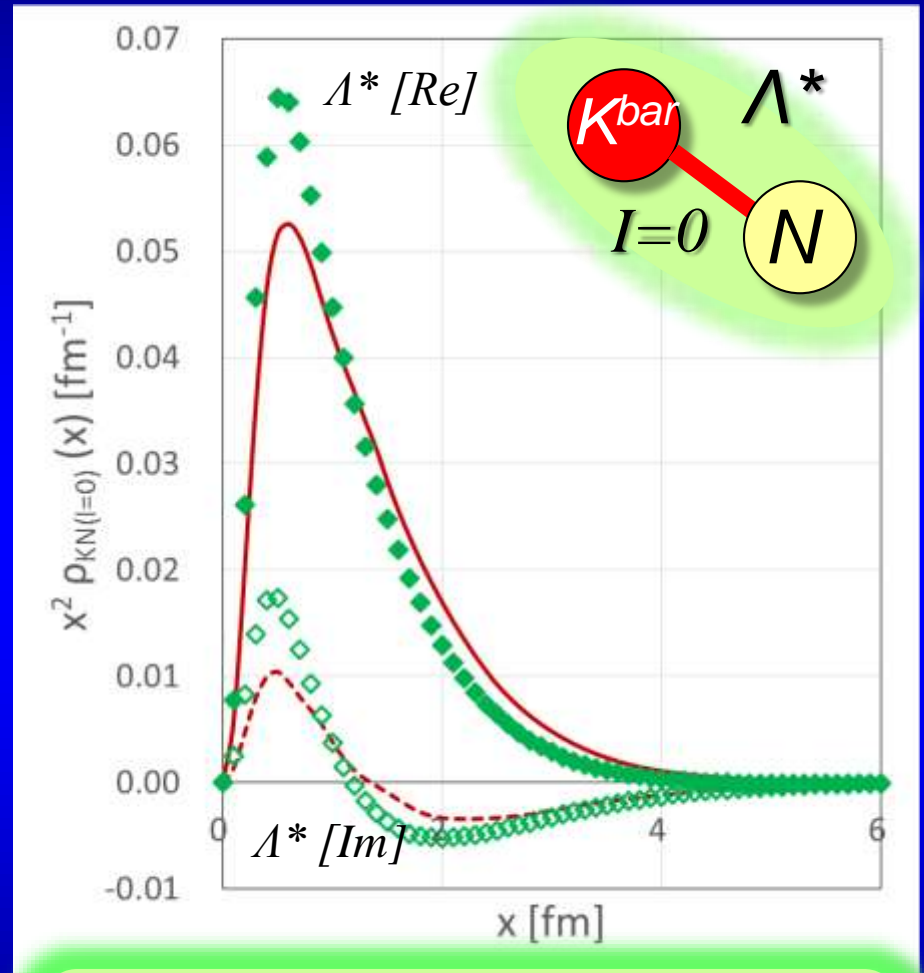
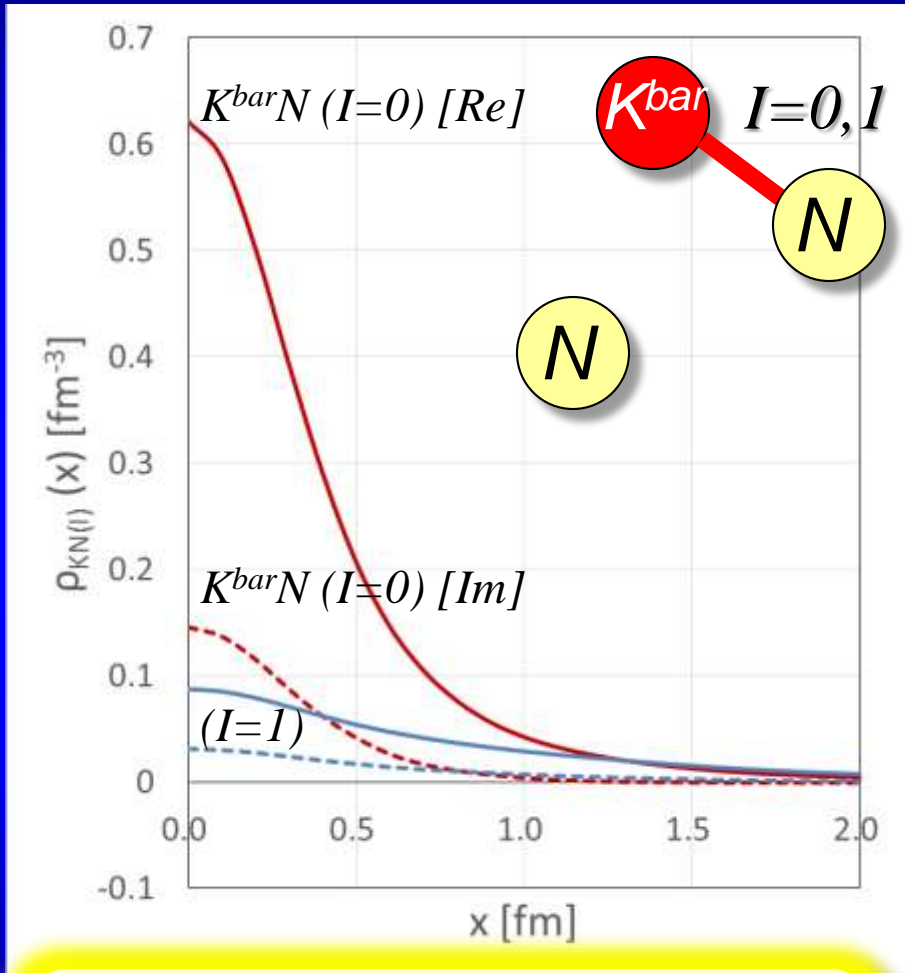


NN distance = 2.1 - i 0.3 fm

*~ Mean distance of 2N in nuclear matter at **normal density!***

$K^{\text{bar}}N$ correlation density

NN pot. : Av18 (Central)
 $K^{\text{bar}}N$ pot. : NRv2c potential
 $f_{\pi}=110$, Particle pict.



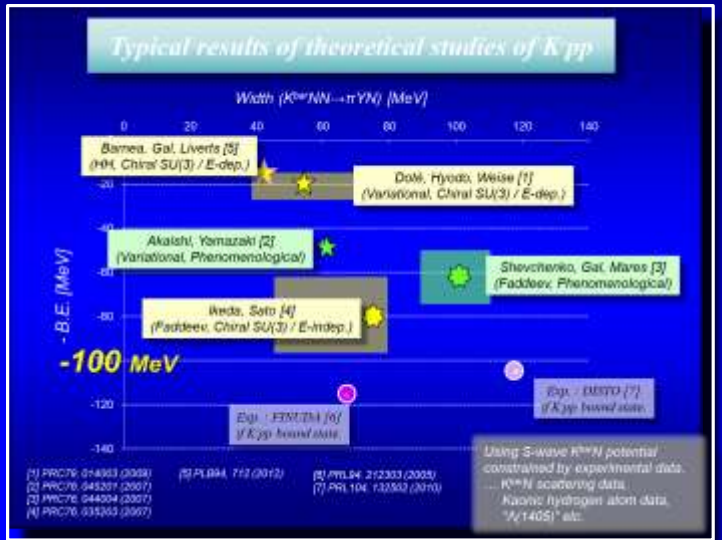
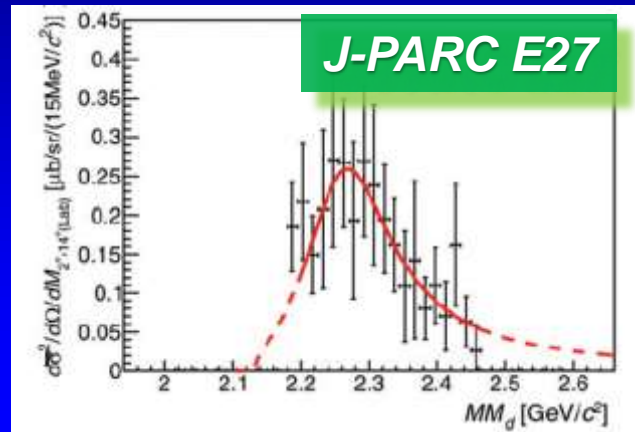
$I=0$ $K^{\text{bar}}N$ compacter than $I=1$ one
 ← Strong $K^{\text{bar}}N$ attraction in $I=0$

$I=0$ $K^{\text{bar}}N$ seems similar to Λ^*
 → Λ^* survives in K -pp

How to understand experimental results?

FINUDA

DISTO



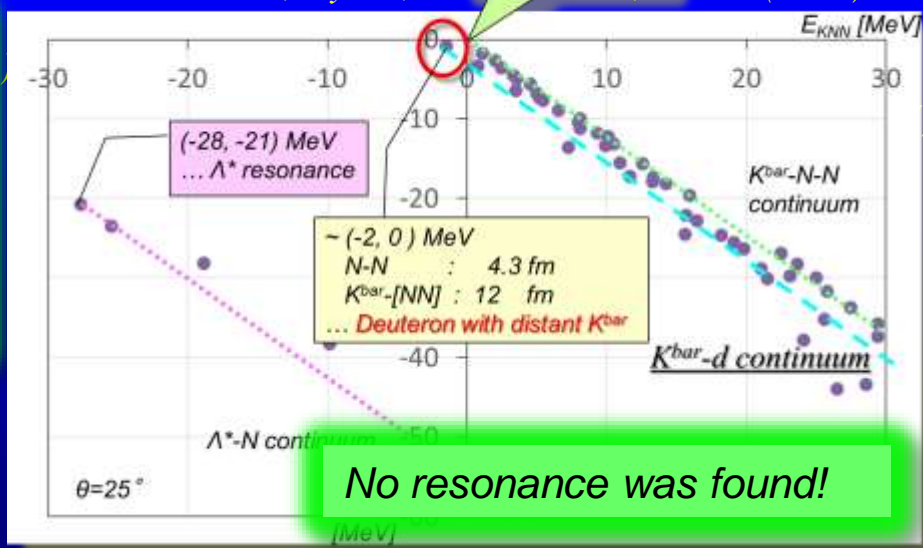
Signal at ~100 MeV below $K^{\bar{b}ar}NN$ threshold

FINUDA: PRL 94, 212303 (2005)

$B(K^-pp) < 100$ MeV

Doté, Hyodo, Weise, PRC 79, 014003(2009)

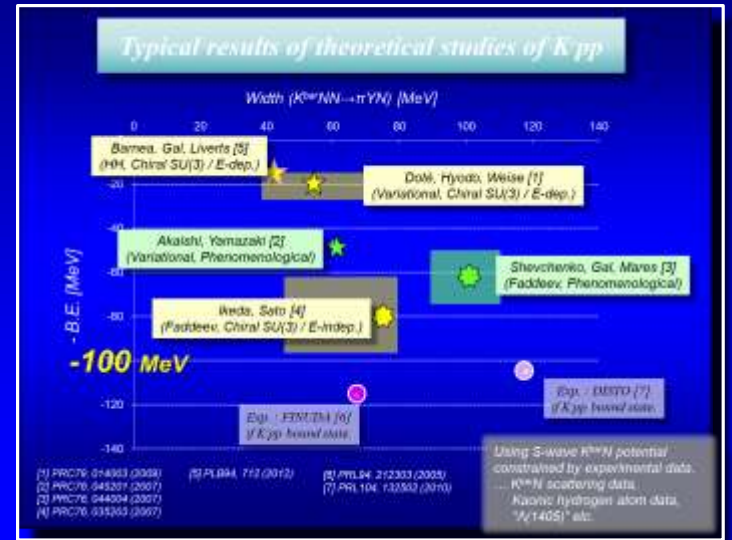
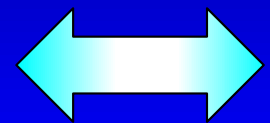
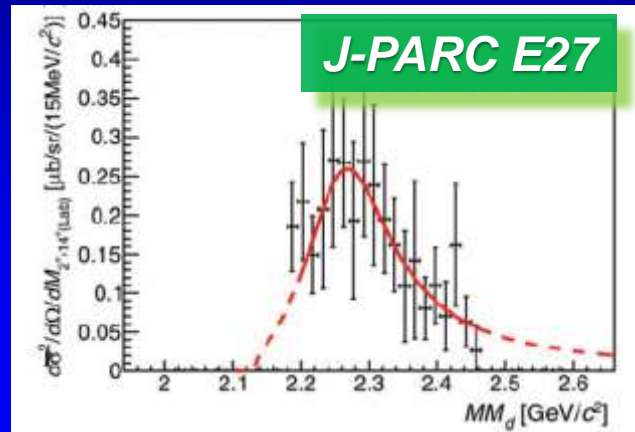
J-PARC E27: $d(\pi^+, K^+)$ reaction
to search for K^-pp
 ⇒ The observed state may be “ K^-pp ”
 with $J^{\pi} = 1^-$ and $l = 1/2$,
 because target deuteron has spin 1.



How to understand experimental results?

FINUDA

DISTO



Signal at ~100 MeV below $K^{\text{bar}}NN$ threshold

$B(Kpp) < 100$ MeV

FINUDA: PRF 94, 212303 (2005)

Dote, Inoue, Myo: PRC79, 014003(2009)

1. Partial restoration of chiral symmetry
... $K^{\text{bar}}N$ potential is enhanced by 17%.

S. Maeda, Y. Akaishi, T. Yamazaki, Proc. Jpn. Acad., Ser. B 89, 418 (2013)

2. Double-pole structure of "Kpp" ($J^{\pi}=0^{-}, l=1/2$)
... "Kpp" has two poles similarly to $\Lambda(1405)$. The lower pole appears.

Y. Ikeda, H. Kamano, T. Sato, PTP124, 533 (2010)

A. Dote, T. Inoue, T Myo, PTEP 2015, 043D02 (2015)

3. Pion assisted dibaryon " $Y = \pi\Sigma N - \pi\Lambda N$ ($J^{\pi}=2^{+}, l=3/2$)"

A. Gal, arXiv:1412.0198 (Proceeding of EXA2014)

***4. Summary
and future plans***

4. Summary and future plans

A prototype of K^{bar} nuclei “ K^-pp ” = Resonance state of $K^{\text{bar}}NN-\pi YN$ coupled system

“coupled-channel Complex Scaling Method + Feshbach projection”

... Represent the **Q-space Green function** with the **Extended Complete Set** well approximated by **Gaussian base**

⇒ Eliminate πY channels to reduce the problem to a $K^{\text{bar}}NN$ single channel problem.

K^-pp studied with ccCSM+Feshbch method

- Used a Chiral SU(3)-based potential (Gaussian form in r -space)
- Self-consistency for kaon's **complex** energy
- Correlation density in CSM shows effect of NN repulsive core and Λ^* survival in K^-pp resonance.
- $J^\pi=1^-$ state (“Deuteron+ K^- ”-like channel) seems not to exist as a resonance state.

K^-pp ($J^\pi=0^-, T=1/2$) --- NRv2c potential case

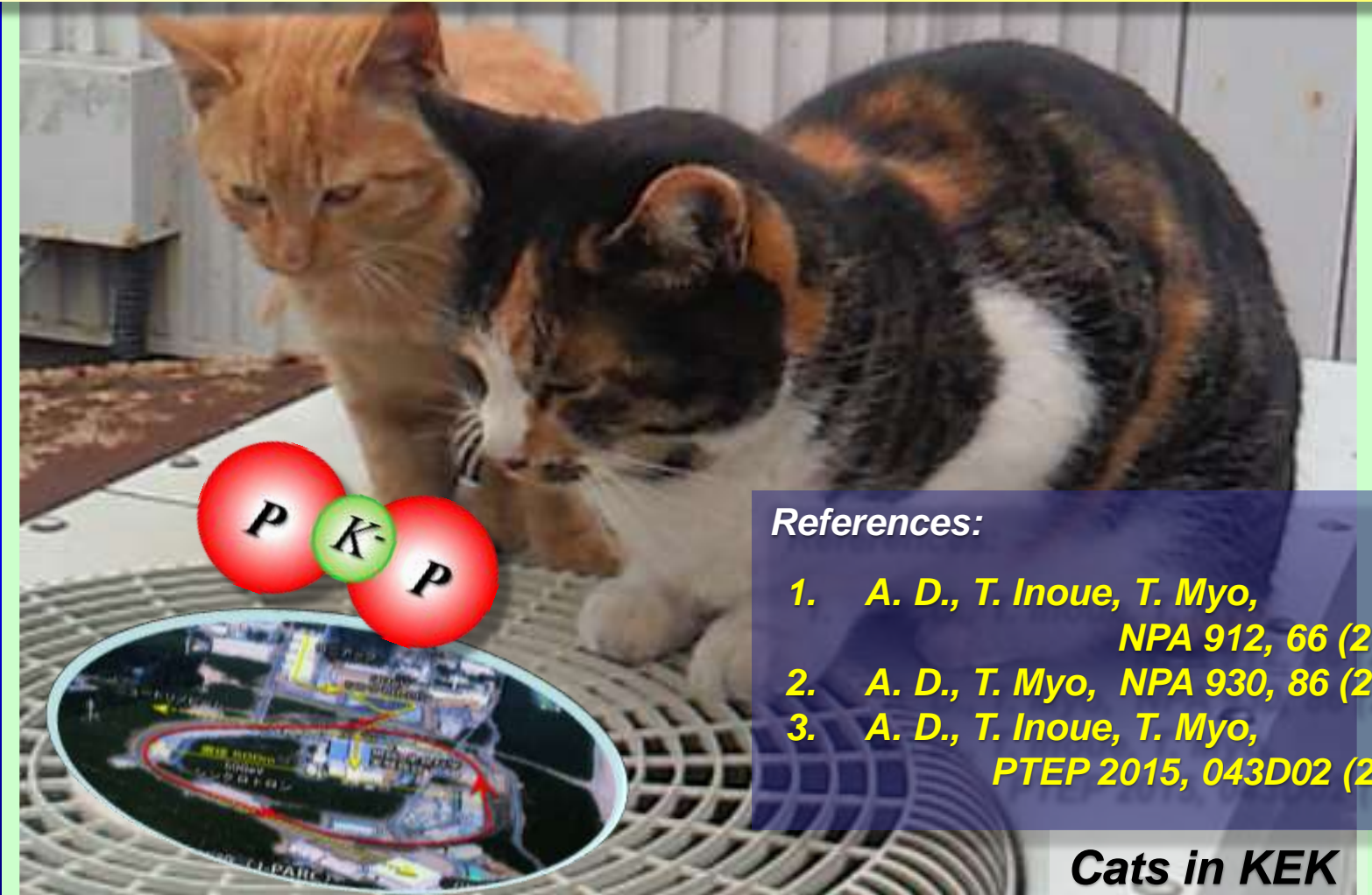
($B, \Gamma/2$) = (21~31, 9~16) MeV : “Field picture”
(25~30, 15~32) MeV : “Particle pict.”

Mean NN distance ~ 2.2 fm → Normal density

Future plans

- Full-coupled channel calculation of K^-pp
... Deailed study for the double pole structure of K^-pp
- Application to resonances of other hadronic systems

Thank you for your attention!



References:

1. *A. D., T. Inoue, T. Myo, NPA 912, 66 (2013)*
2. *A. D., T. Myo, NPA 930, 86 (2014)*
3. *A. D., T. Inoue, T. Myo, PTEP 2015, 043D02 (2015)*

Cats in KEK