

Deeply Virtual Compton Scattering off the Neutron: measurements with CLAS and CLAS12 at Jefferson Lab

Daria Sokhan

University of Glasgow, UK

on behalf of the CLAS Collaboration

NSTAR 2015 – Osaka, Japan,
25th May 2015



Wigner distributions

$$\rho(x, \vec{k}_T, \vec{b}_T)$$

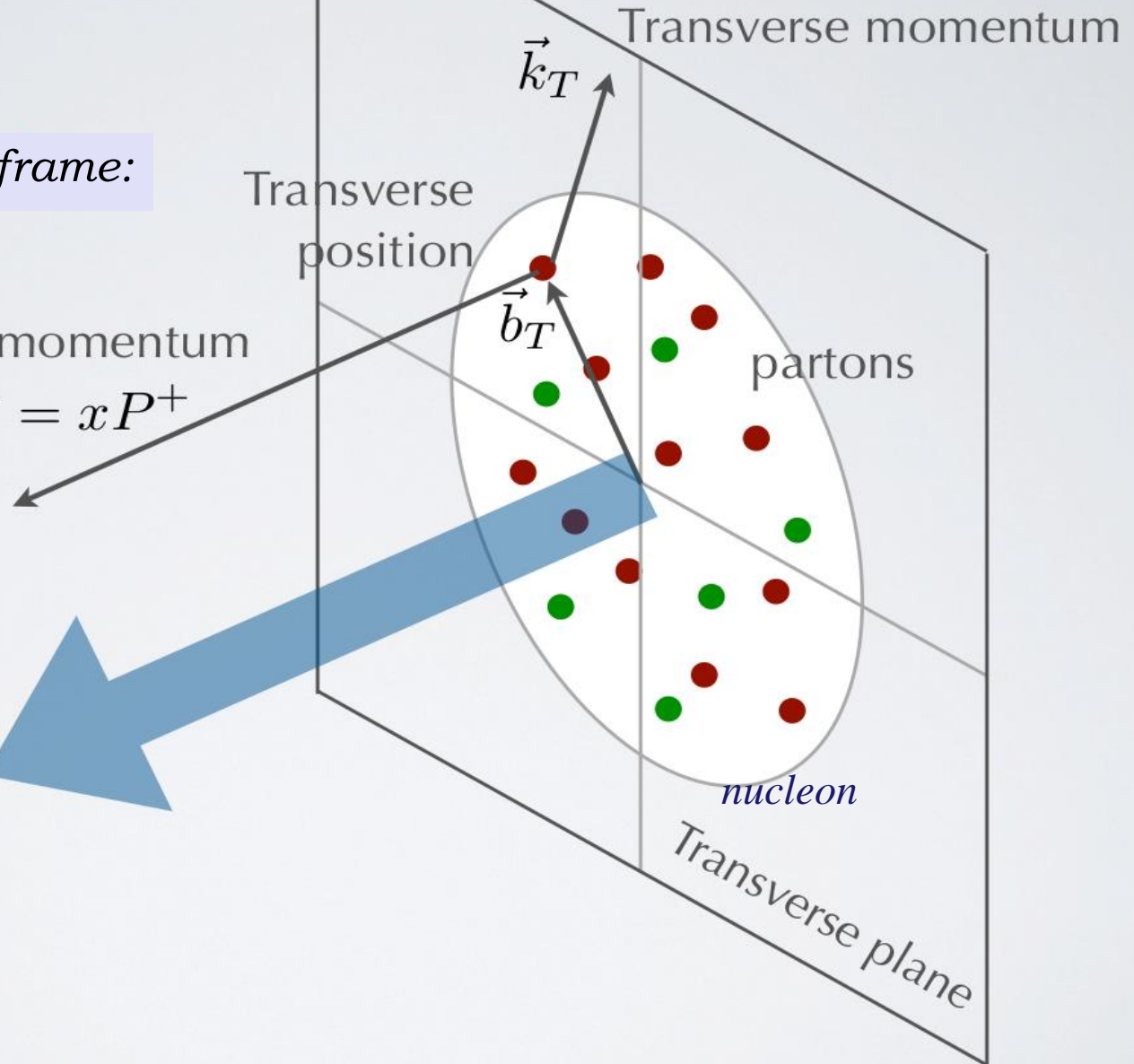
In infinite momentum frame:

Longitudinal momentum

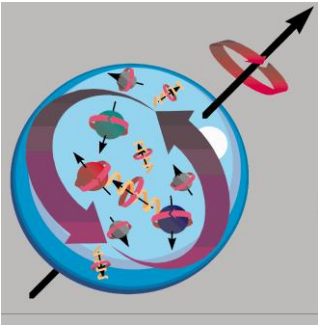
$$k^+ = xP^+$$

x : longitudinal momentum fraction carried by struck parton

Nucleon structure



Generalised Parton Distributions

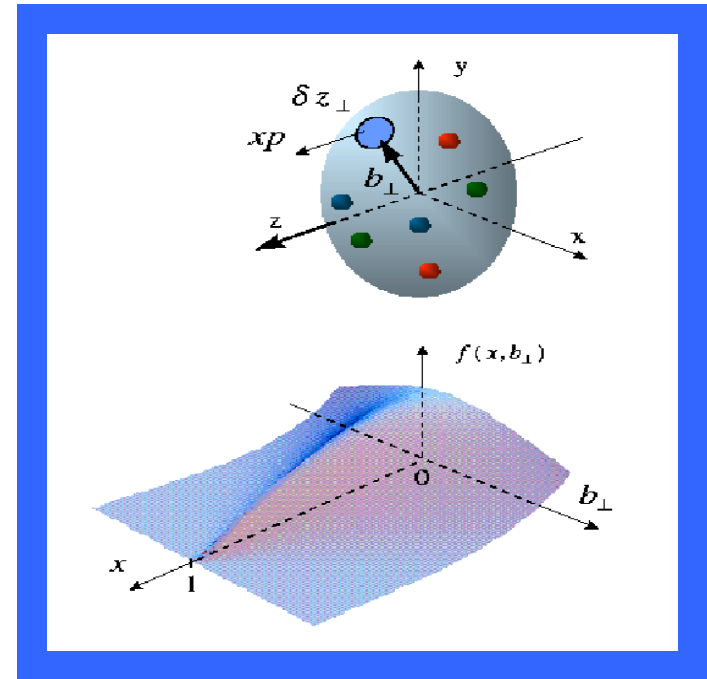


*Wigner function:
full phase space parton
distribution of the nucleon*

**Generalised
Parton
Distributions
(GPDs)**

$$\int d^2 k_T$$

- ❖ Relate transverse position of partons (b_\perp) to longitudinal momentum (x).
- ❖ Deep exclusive reactions.



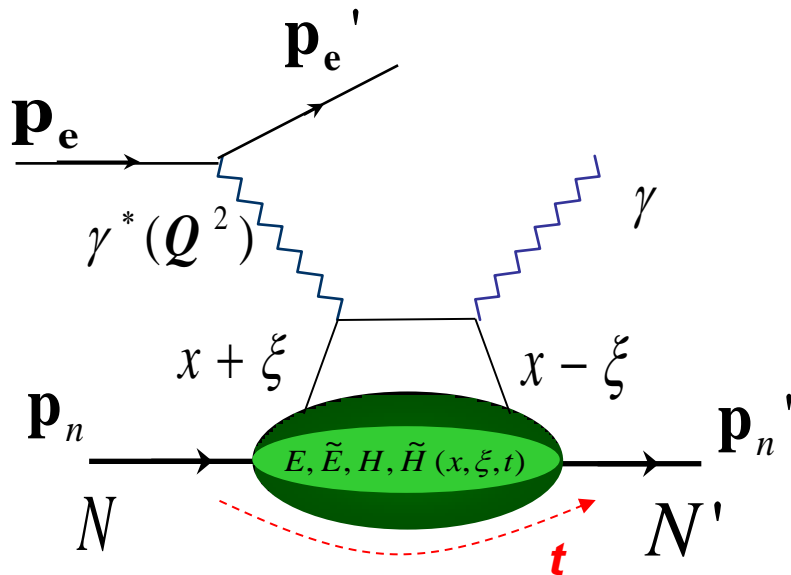
Deeply Virtual Compton Scattering

- ❖ GPDs relate transverse position of partons to longitudinal momentum.



contain information on angular momentum of quarks

- ❖ “Golden channel” for GPD extraction: **Deeply Virtual Compton Scattering (DVCS)**.



$$Q^2 = -(\mathbf{p}_e - \mathbf{p}_{e'})^2 \quad t = (\mathbf{p}_n - \mathbf{p}_{n'})^2$$

$$\text{Bjorken variable: } x_B = \frac{Q^2}{2\mathbf{p}_n \cdot \mathbf{q}}$$

$$x \pm \xi \quad \begin{array}{l} \text{longitudinal momentum} \\ \text{fractions of quarks} \end{array} \quad \xi \cong \frac{x_B}{2 - x_B}$$

- ❖ At high exchanged Q^2 , access to four GPDs: $E_q, \tilde{E}_q, H_q, \tilde{H}_q(x, \xi, t)$

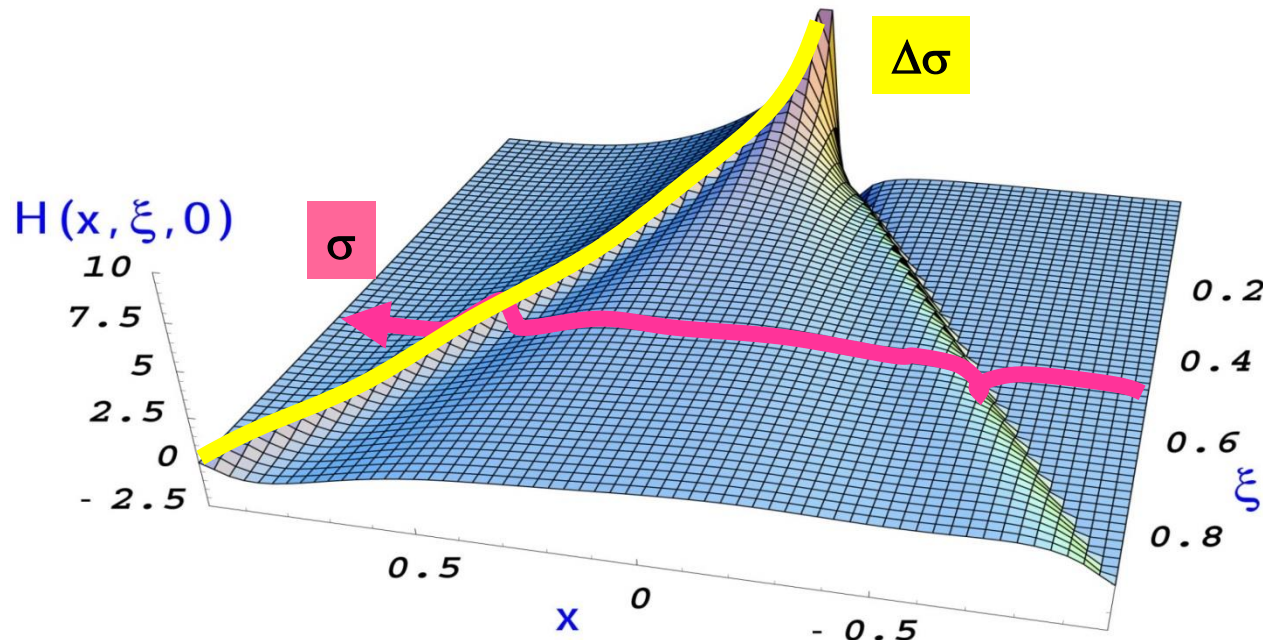
Compton Form Factors in DVCS

CFFs: complex functions directly accessible in DVCS cross-sections and spin asymmetries, eg:

$$A_{LU} = \frac{d\vec{\sigma} - d\vec{\sigma}}{d\vec{\sigma} + d\vec{\sigma}} = \frac{\Delta\sigma_{LU}}{d\vec{\sigma} + d\vec{\sigma}}$$

Related to GPDs:

$$T^{DVCS} \sim \int_{-1}^{+1} \frac{GPDs(x, \xi, t)}{x \pm \xi + i\varepsilon} dx + \dots \sim P \int_{-1}^{+1} \frac{GPDs(x, \xi, t)}{x \pm \xi} dx \pm i\pi GPDs(\pm\xi, \xi, t) + \dots$$

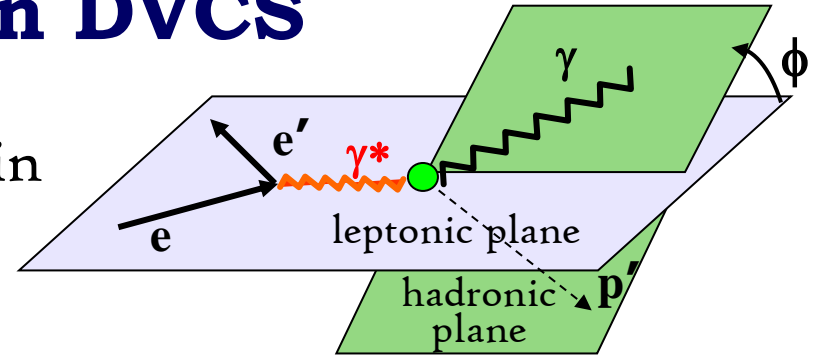


Only ξ and t are accessible experimentally!

Compton Form Factors in DVCS

❖ Experimentally accessible in DVCS spin asymmetries, eg:

$$A_{LU} = \frac{d\vec{\sigma} - d\bar{\sigma}}{d\vec{\sigma} + d\bar{\sigma}} = \frac{\Delta\sigma_{LU}}{d\vec{\sigma} + d\bar{\sigma}}$$



Beam, target polarisation

$$\xi = x_B / (2 - x_B) \quad k = t / 4M^2$$

→ e^- p/n $\Delta\sigma_{LU} \sim \sin\phi \operatorname{Im}\{F_1\mathcal{H} + \xi(F_1+F_2)\tilde{\mathcal{H}} - kF_2\mathcal{E}\}d\phi$

e^- → p/n $\Delta\sigma_{UL} \sim \sin\phi \operatorname{Im}\{F_1\tilde{\mathcal{H}} + \xi(F_1+F_2)(\mathcal{H} + x_B/2\mathcal{E}) - \xi kF_2\tilde{\mathcal{E}} + \dots\}d\phi$

e^- ↑ p/n $\Delta\sigma_{UT} \sim \cos\phi \operatorname{Im}\{k(F_2\mathcal{H} - F_1\mathcal{E}) + \dots\}d\phi$

→ e^- → p/n $\Delta\sigma_{LL} \sim (A+B\cos\phi) \operatorname{Re}\{F_1\tilde{\mathcal{H}} + \xi(F_1+F_2)(\mathcal{H} + x_B/2\mathcal{E})\dots\}d\phi$

Proton **Neutron**

$\operatorname{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p, \mathcal{E}_p\}$
 $\operatorname{Im}\{\mathcal{H}_n, \tilde{\mathcal{H}}_n, \mathcal{E}_n\}$

$\operatorname{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\}$
 $\operatorname{Im}\{\mathcal{H}_n, \mathcal{E}_n, \tilde{\mathcal{E}}_n\}$

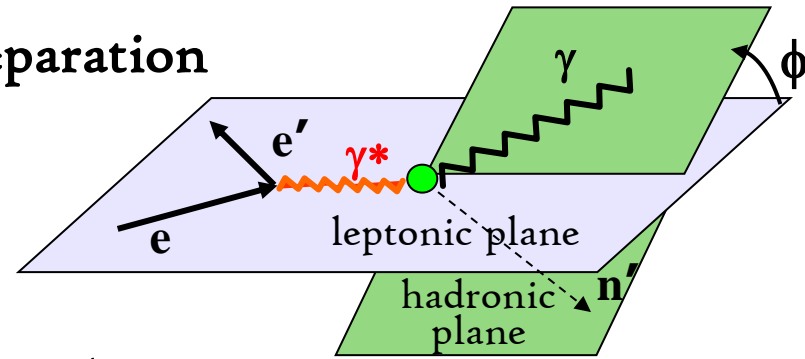
$\operatorname{Im}\{\mathcal{H}_p, \mathcal{E}_p\}$
 $\operatorname{Im}\{\mathcal{H}_n\}$

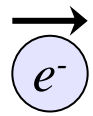
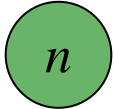
$\operatorname{Re}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\}$
 $\operatorname{Re}\{\mathcal{H}_n, \mathcal{E}_n, \tilde{\mathcal{E}}_n\}$

Neutron DVCS

❖ GPDs from proton and neutron: **flavour separation**

❖ **Neutron DVCS** extremely sensitive to E ,
least-known and least-constrained GPD



  Polarized beam, unpolarized **neutron** target:

$$\Delta\sigma_{LU} \sim \sin\phi \operatorname{Im}\{F_1\mathbf{H} + \xi(F_1+F_2)\tilde{\mathbf{H}} - kF_2\mathbf{E}\}d\phi \longrightarrow H_n, \tilde{H}_n, E_n$$

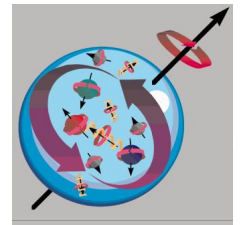
Suppressed because $F_1(t)$ is small

Suppressed because of **cancellation**
between PDF's of **u** and **d** quarks

❖ Ji's "Sum Rule": $J^q = \frac{1}{2} - J^g = \frac{1}{2} \int_{-1}^1 x dx \{ H^q(x, \xi, 0) + E^q(x, \xi, 0) \}$

$$J_N = \frac{1}{2} = \frac{1}{2} \sum_q L_q + J_g$$

Important missing link in
the **nucleon spin puzzle!**



Neutron DVCS: eg1-dvcs experiment @ Jefferson Lab (Hall B)

Data taken: Feb – Sept 2009

Longitudinally polarised targets:

Beam: polarised electrons

NH₃ (95 days)

$E_e = 4.7$ to 6 GeV

ND₃ (33 days)

polarisation $\sim 85\%$

Proton / neutron pol. $\sim 80 / 40 \%$

$$\vec{e} + \vec{d} \rightarrow e' + \gamma + n + (p_s)$$

CLAS



Exclusive reconstruction of e' , N , and γ .
Spectator proton identified
via missing mass.

plus

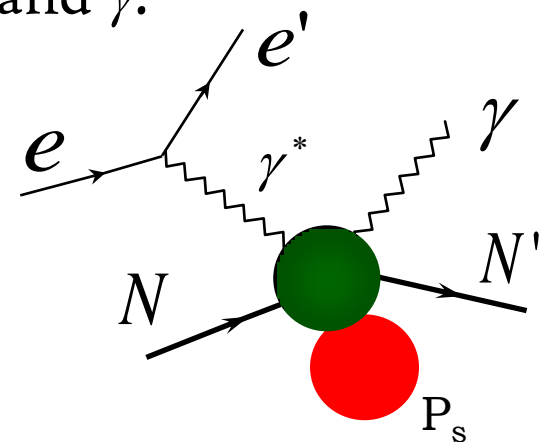
Inner

Calorimeter



(IC)

high-energy forward
photon detection

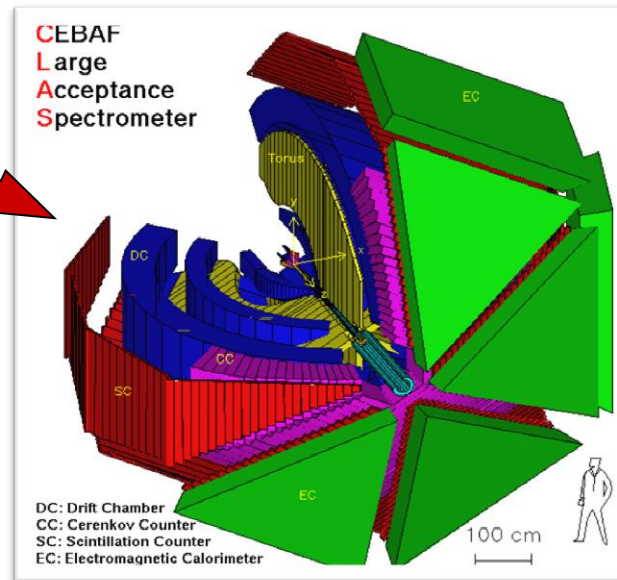


CLAS @ Jefferson Lab (Virginia, USA)



CEBAF: Continuous Electron Beam Accelerator Facility:

- ❖ Duty cycle: $\sim 100\%$
- ❖ Energy up to ~ 6 GeV
- ❖ Electron polarisation up to $\sim 85\%$



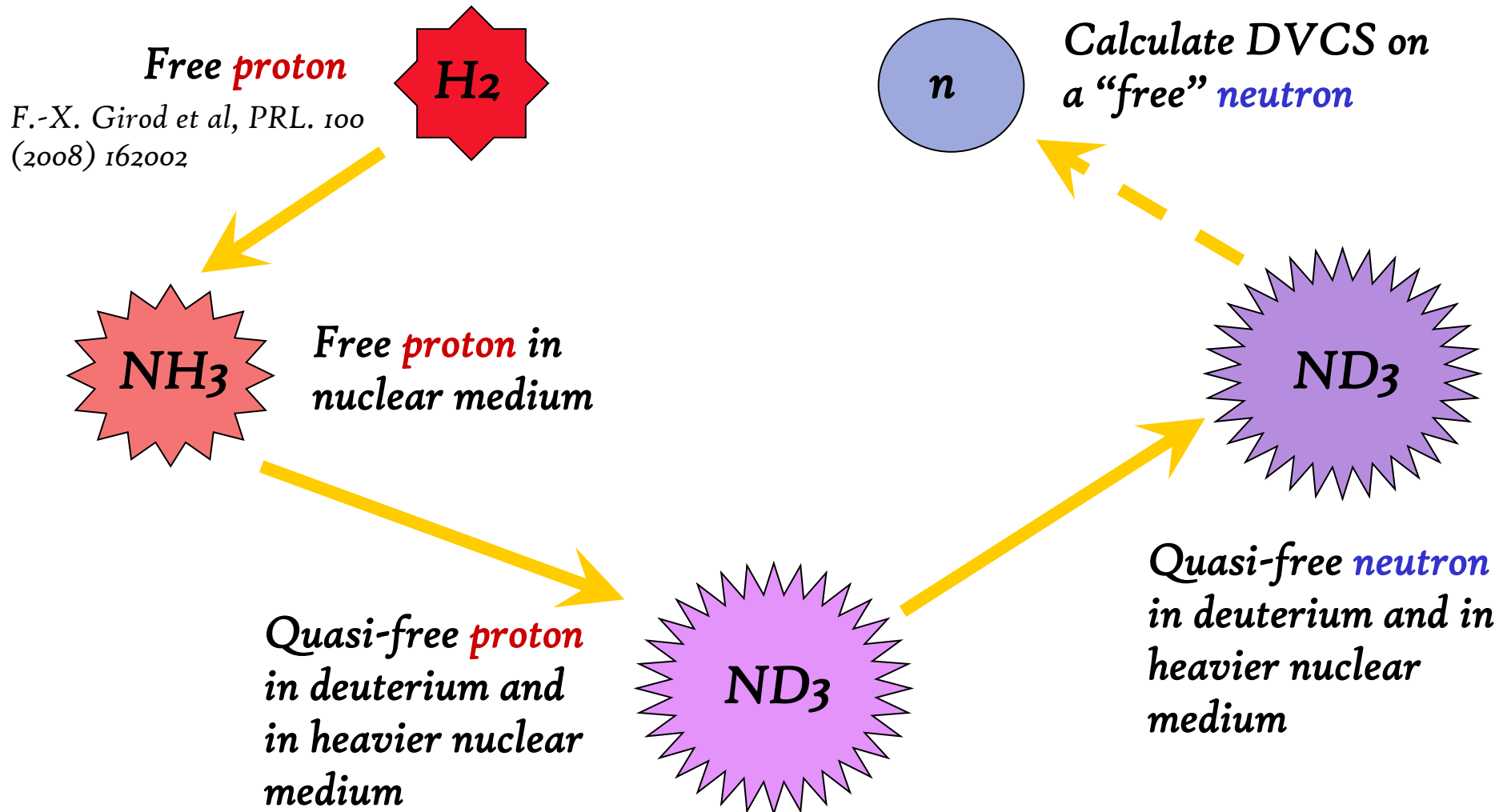
CLAS in Hall B:

- ❖ Drift chambers
- ❖ Toroidal magnetic field
- ❖ Cerenkov Counters
- ❖ Scintillator Time of Flight
- ❖ Electromagnetic Calorimeters

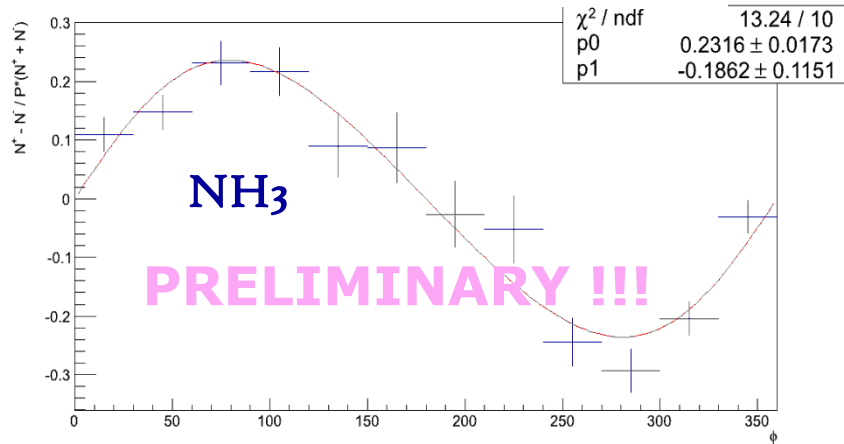


Extremely large angular coverage

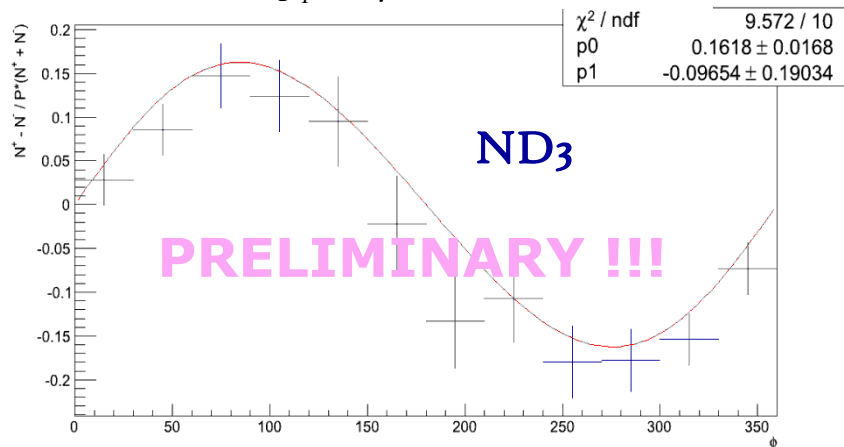
DVCS on different targets



A_{LU} – check on proton DVCS in NH_3 and ND_3



Fit: $A = \frac{p_0 \sin \varphi}{1 + p_1 \cos \varphi}$



Previously measured result on H_2 is in range 0.2 - 0.3.

F.-X. Girod et al, PRL. 100 (2008) 162002

$$\frac{N^+ - N^-}{P(N^+ + N^-)} \approx 0.23 \pm 0.02$$

Uncorrected for π^0 contamination

\longrightarrow actual A_{LU} larger!

Deuterium target – smearing due to Fermi motion requires wider data cuts.

$$\frac{N^+ - N^-}{P(N^+ + N^-)} \approx 0.16 \pm 0.02$$

π^0 contamination more significant

\longrightarrow measured A_{LU} lower than on NH_3 .

Neutron DVCS in ND_3 – identifying reaction

“Deep Inelastic Scattering” cuts:

❖ $W > 2 \text{ GeV}/c^2$ where W is the missing mass of $(eN \rightarrow e' X)$,
isolate resonance region of remaining γN

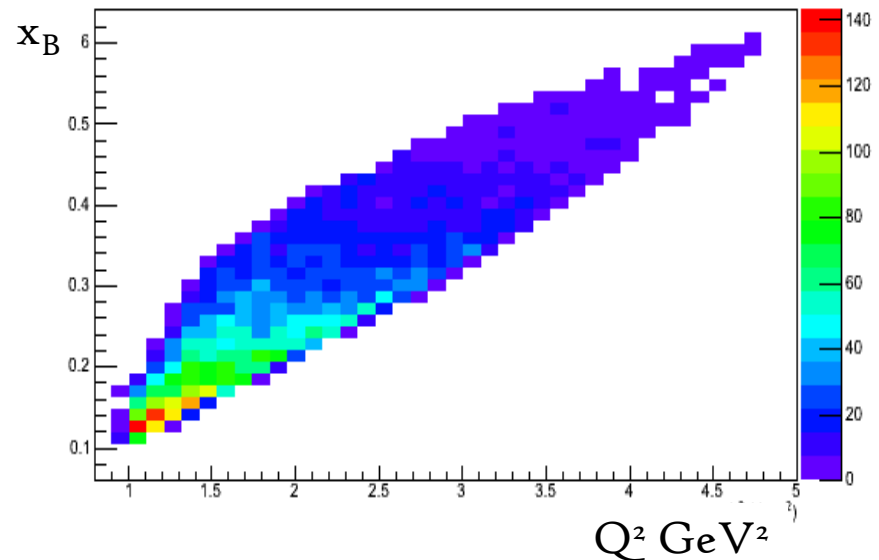
❖ $Q^2 > 1 \text{ GeV}^2$ ❖ $Q^2 > -t$ Region where factorisation applies

Additional DVCS cuts:

❖ $E_\gamma > 1 \text{ GeV}$

❖ $p_n > 0.4 \text{ GeV}/c$

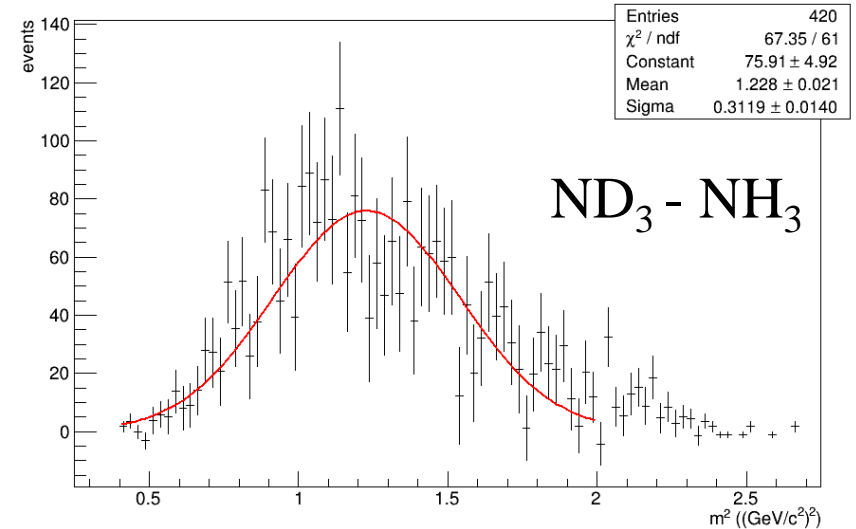
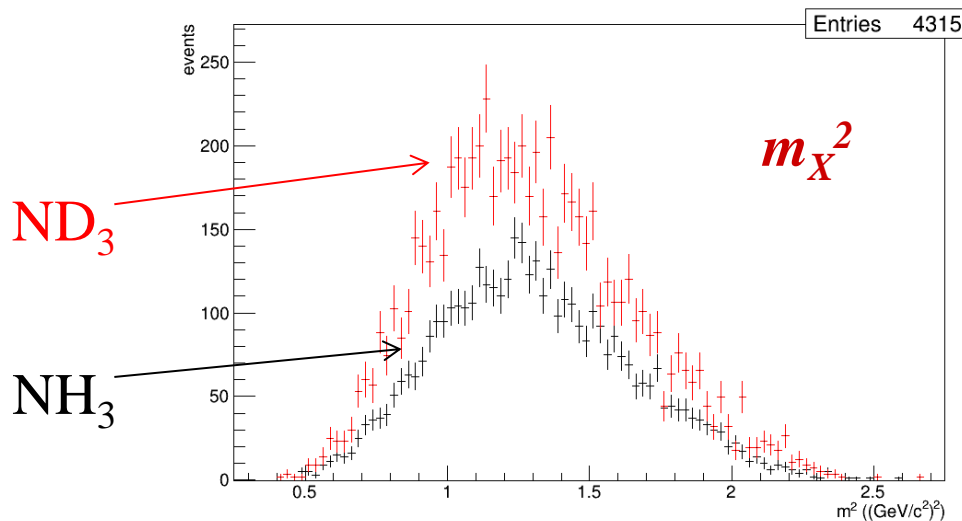
Recoiling nucleon should not
have a low p



Exclusivity cuts: spectator

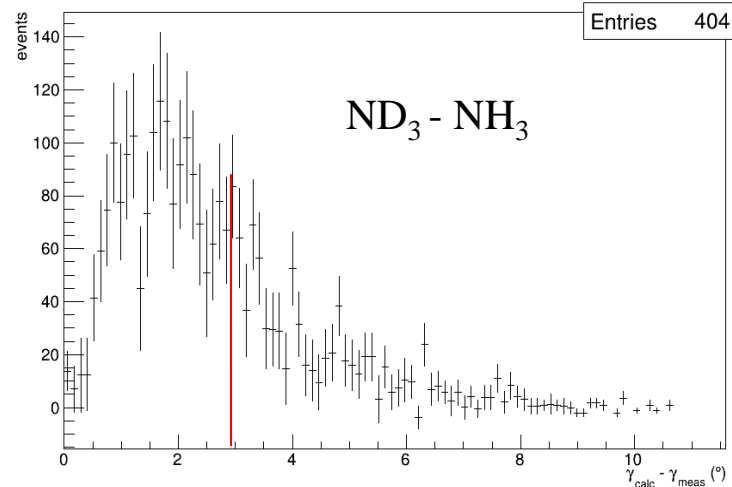
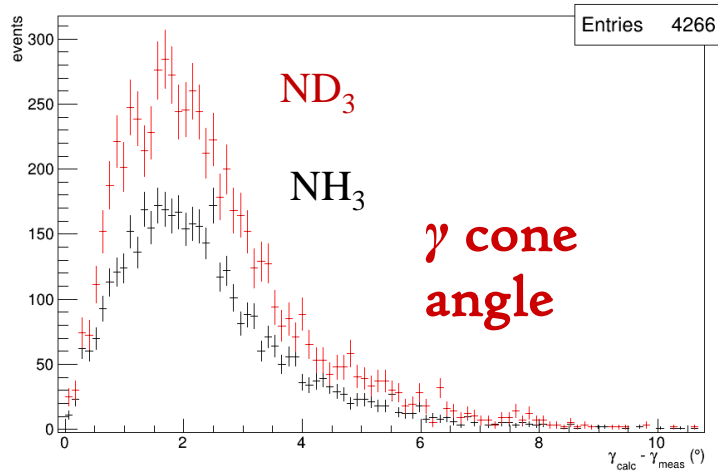
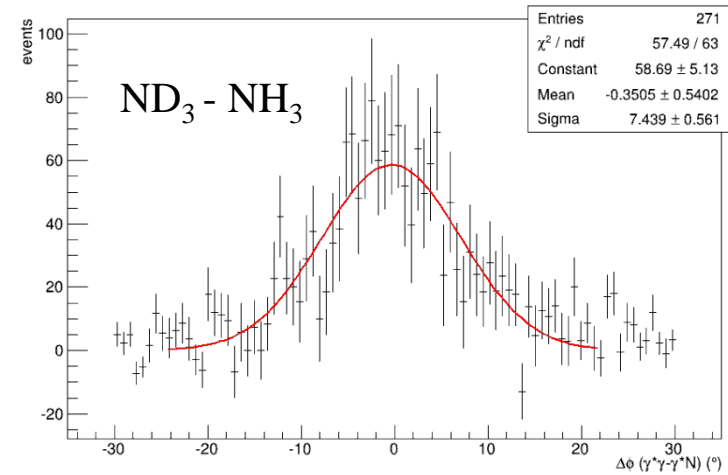
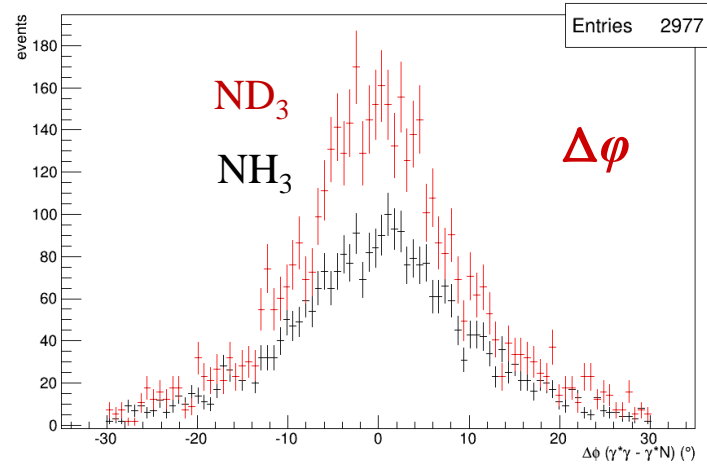
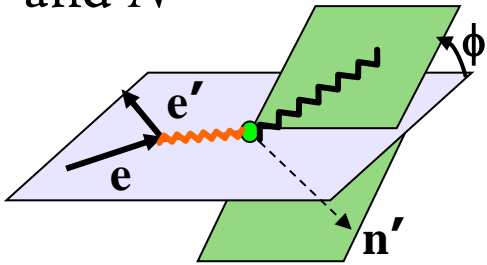
Use NH_3 data to subtract the nuclear background from the ND_3 distributions:

- ❖ **Missing momentum** from $ed \rightarrow e' N' \gamma X$ should be low for spectator nucleon in quasi-free reaction: $p_X < 0.2 \text{ GeV}/c$
- ❖ **Missing mass** of spectator from $ed \rightarrow e' N' \gamma X$: $0.5 < |m_X^2| < 2 \text{ GeV}^2/c^4$



Exclusivity cuts: angular distributions

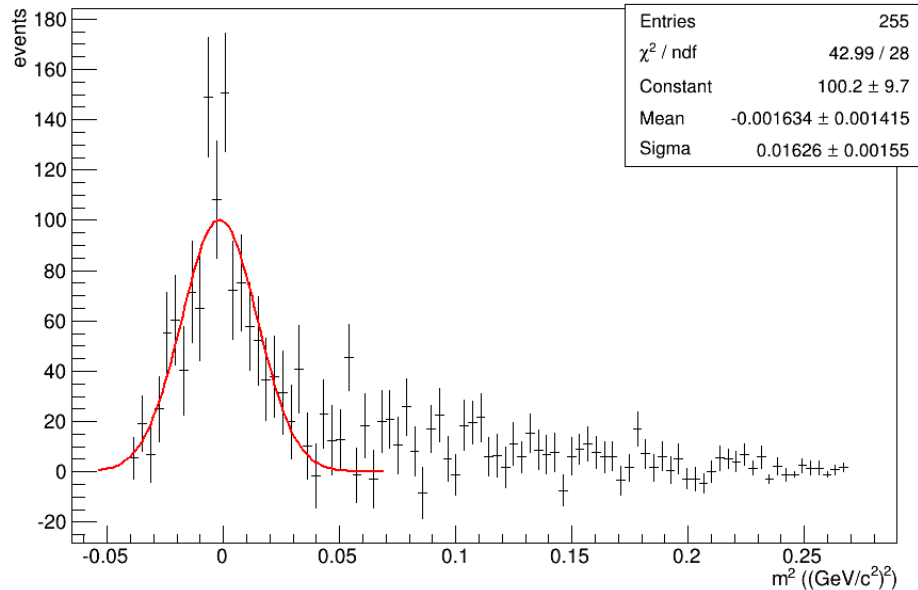
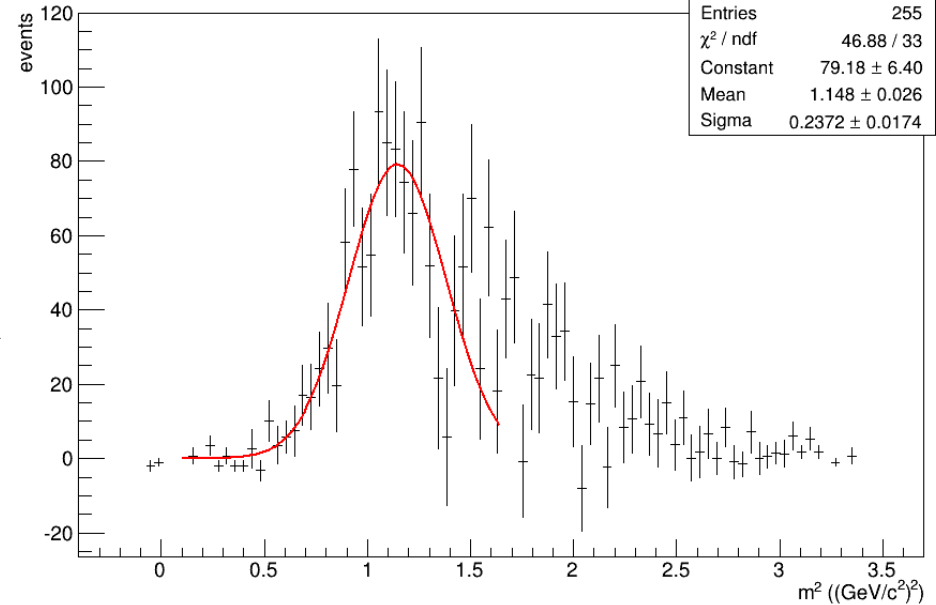
❖ $|\Delta\phi| < 5^\circ$
Coplanarity between γ and N



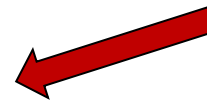
❖ γ cone angle $< 3^\circ$
Difference between calculated and measured γ direction

After exclusivity cuts

❖ **Missing mass** from $eN \rightarrow e'\gamma X$
(should correspond to recoil **neutron**)



❖ **Missing mass** from $eN \rightarrow e' N \gamma X$
(should correspond to nothing)



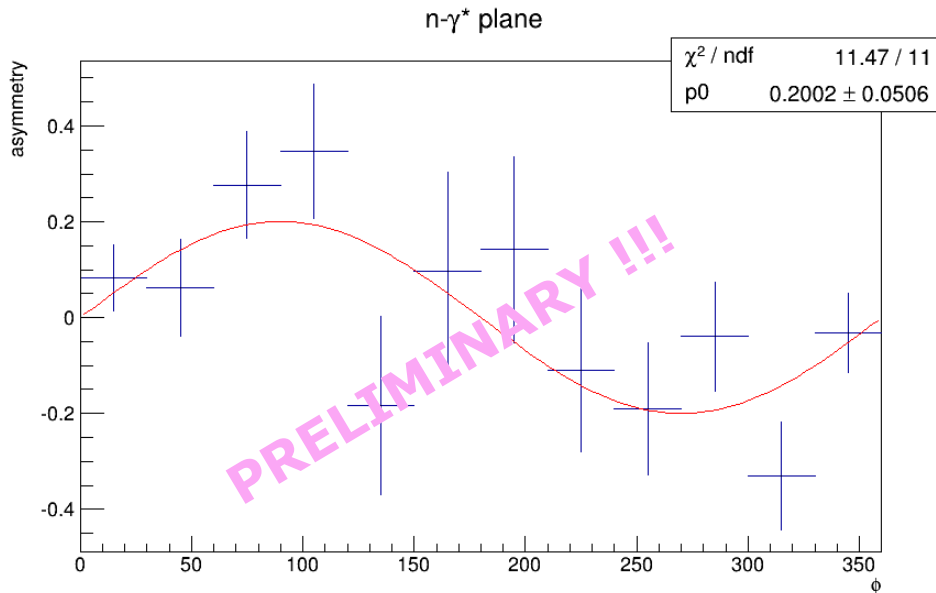
*Distributions shown
after nuclear background
subtraction.*

A_{LU} in neutron DVCS on ND_3

❖ Beam-spin asymmetry (A_{LU}):

One previous measurement from Hall A @ JLab, $A_{LU} \sim 0$. Big statistical and systematic uncertainties, slightly different kinematic region.

(M. Mazouz et al, PRL 99 (2007) 242501)



$$\text{Fit: } A_{LU} = p_0 \sin \varphi$$

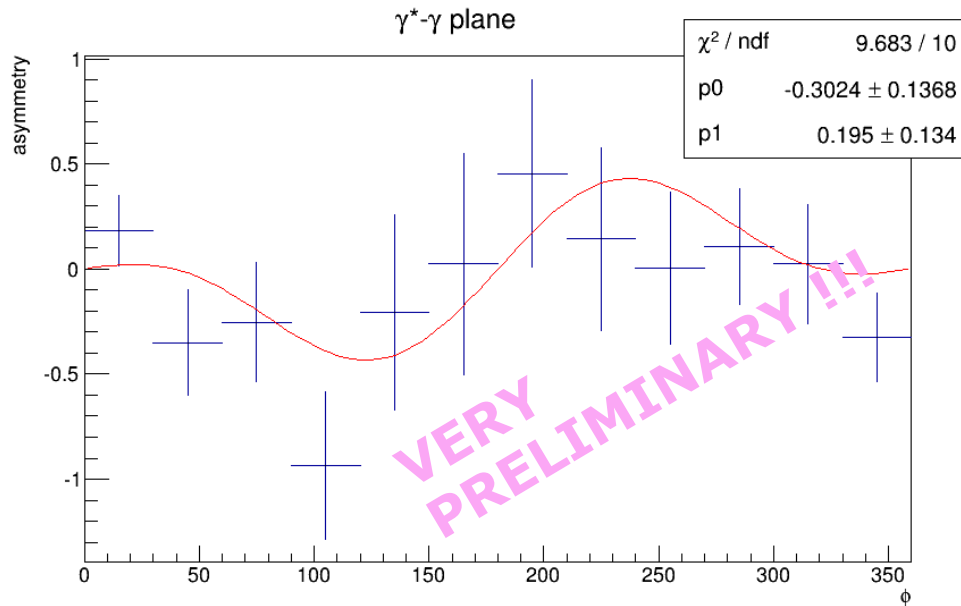
$$\frac{N^+ - N^-}{P(N^+ + N^-)} \approx 0.20 \pm 0.05$$

*Uncorrected for π^0 contamination,
which has an asymmetry of its own!*

A_{UL} in neutron DVCS on ND_3

First measurement.

❖ Target-spin asymmetry (A_{UL}):



Fit:
$$A_{UL} = \frac{p_0 \sin \varphi}{1 + p_1 \cos \varphi}$$

$$p_0 < 0$$

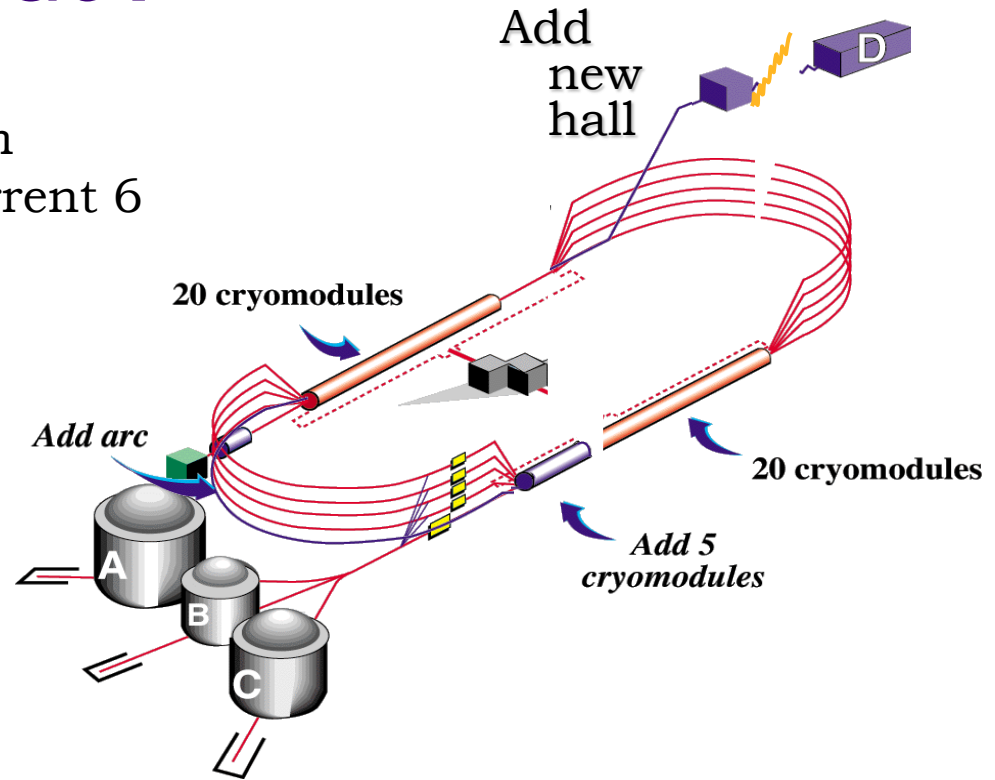
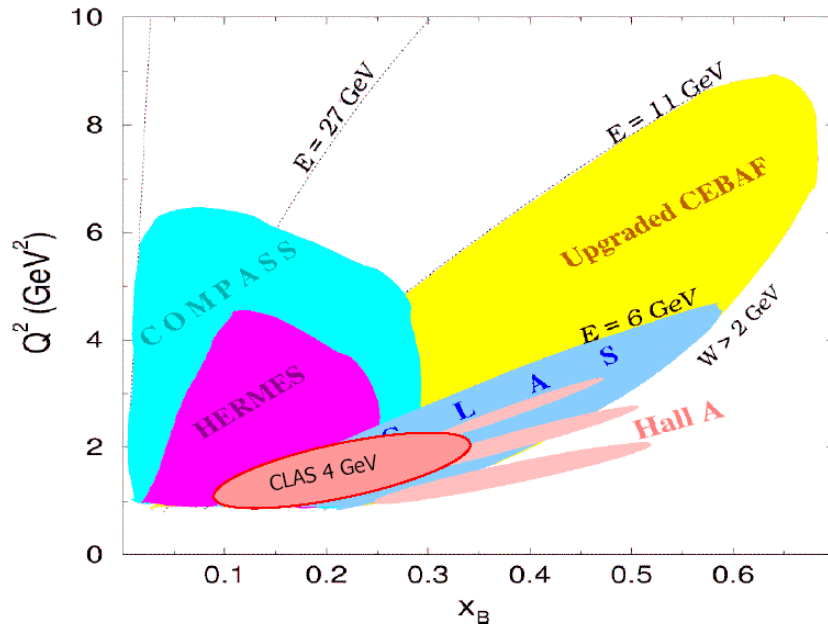
$$p_1 \text{ small}$$

*Uncorrected by the dilution factor
due to the nuclear background!*

Jefferson Lab @ 12 GeV

❖ **CEBAF**: Continuous Electron Beam Accelerator Facility, upgrade from current 6 GeV to **12 GeV** underway.

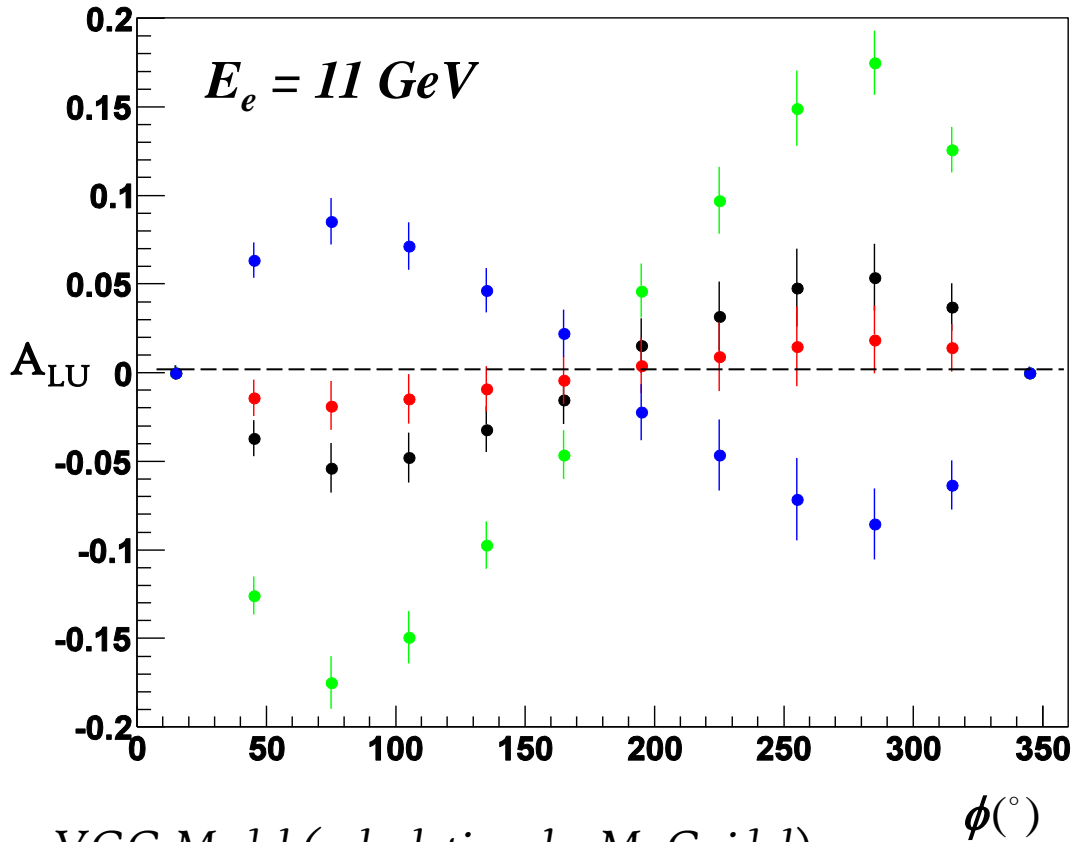
❖ Open up much larger phase space in **Q^2** and **x_B**



❖ Hall B – 11 GeV to the upgraded detector system **CLAS12**

❖ CLAS12 experiments: expected 2016

A_{LU} in Neutron DVCS @ 11 GeV



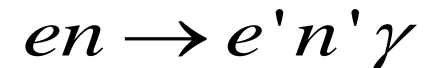
Fixed kinematics: $x_B = 0.17$ $Q^2 = 2 \text{ GeV}^2$
 $t = -0.4 \text{ GeV}^2$

$J_u = 0.3, J_d = -0.1$ $J_u = 0.3, J_d = 0.1$
 $J_u = 0.1, J_d = 0.1$ $J_u = 0.3, J_d = 0.3$

❖ At 11 GeV, beam spin asymmetry (A_{LU}) in neutron DVCS is **very** sensitive to J_u, J_d

❖ Wide coverage needed!

❖ Exclusive reconstruction of the DVCS process



requires detection and measurement of all three final state particles.

CLAS12: base detectors

Design luminosity

$$L \sim 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$$

Acceptance for
charged particles:

- Central (CD), $40^\circ < \theta < 135^\circ$
- Forward (FD), $5^\circ < \theta < 40^\circ$

Acceptance for photons:

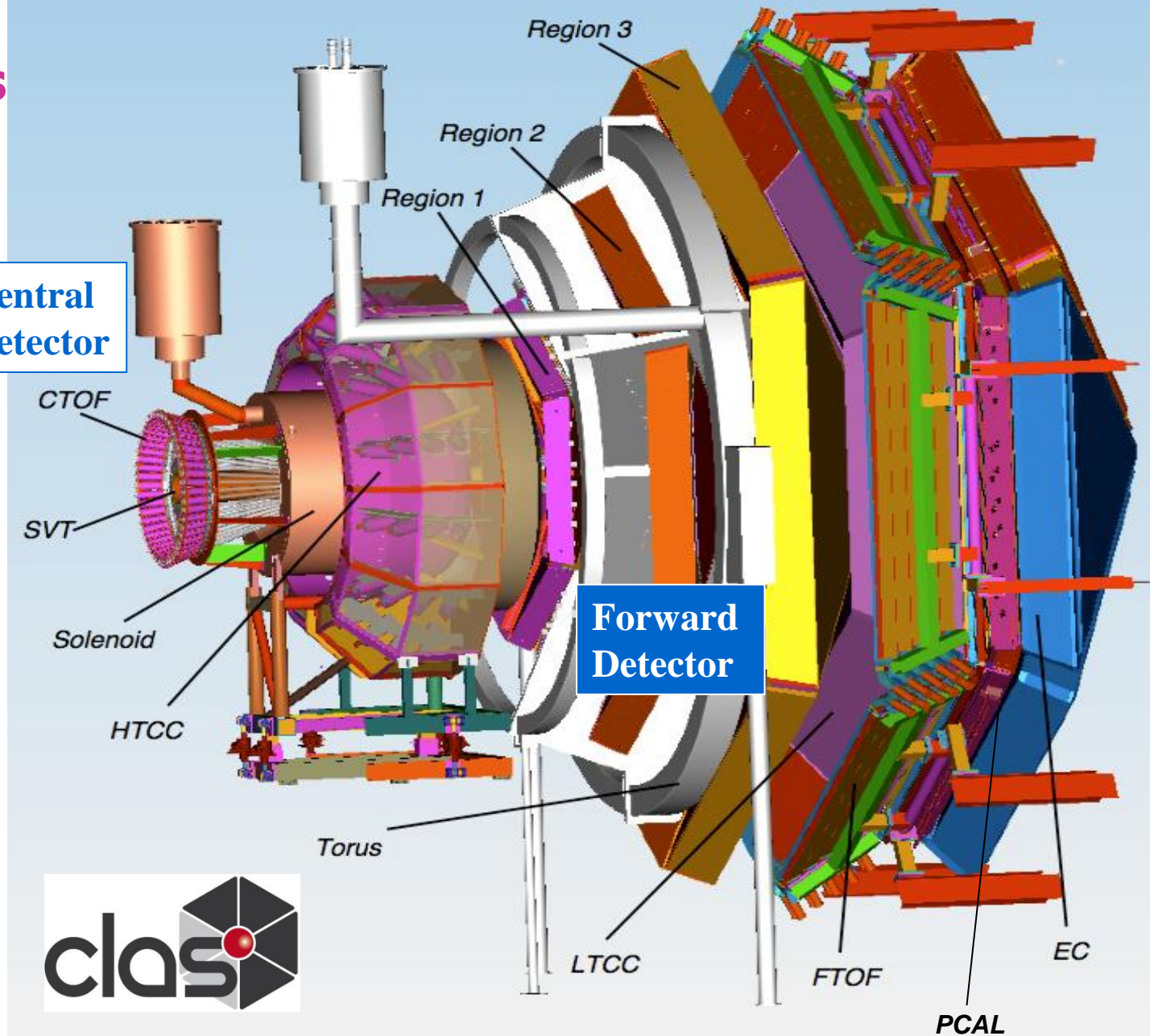
- IC $2^\circ < \theta < 5^\circ$
- EC, $5^\circ < \theta < 40^\circ$

High luminosity & large
acceptance:

Concurrent measurement
of deeply virtual **exclusive**,
semi-inclusive,
and **inclusive** processes

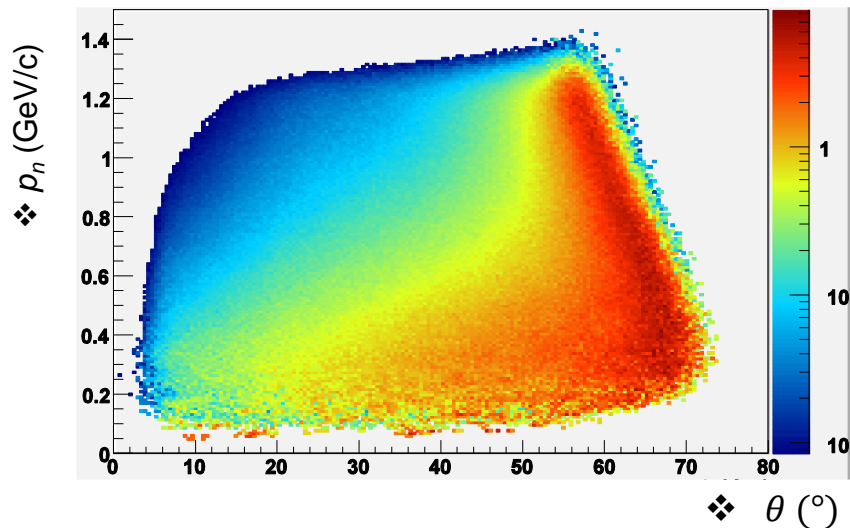
Central
Detector

Forward
Detector



Recoil DVCS neutrons in CLAS12

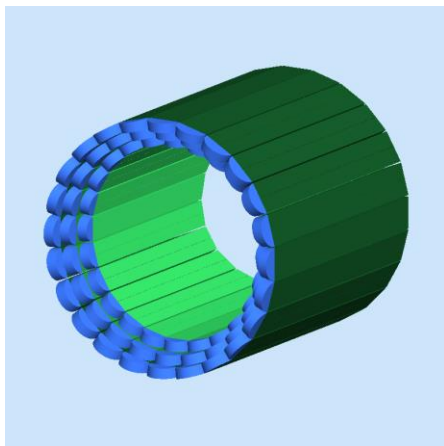
Simulation of n DVCS neutrons at $E_e = 11$ GeV.



Over **80%** of neutrons recoil at $\theta_{lab} > 40^\circ$ with peak momentum at ~ 0.4 GeV/c.

Neutron detector for CLAS12: 3-layer scintillator barrel, 48 paddles/layer.

❖ Neutron efficiency $\sim 8-9\%$

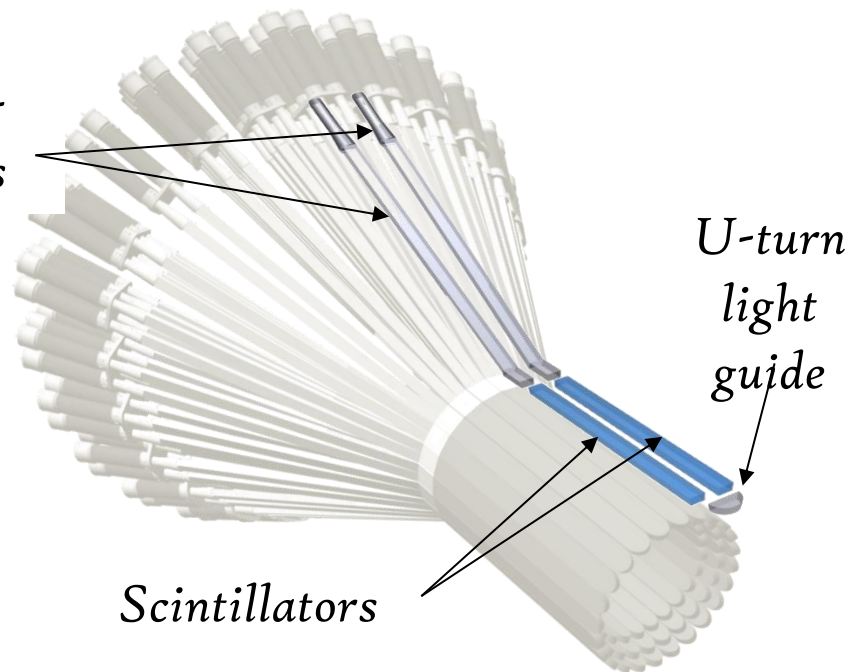


❖ $\frac{\sigma_p}{P} \approx 5-12\%$

❖ $\sigma_\theta \approx 2-3^\circ$

❖ $\sigma_\phi = 3.75^\circ$

Light-guides



Scintillators

U-turn light guide

Limitations of space and high magnetic field (5T) in central region necessitate a u-turn geometry.

A_{LU} in Neutron DVCS with CLAS12

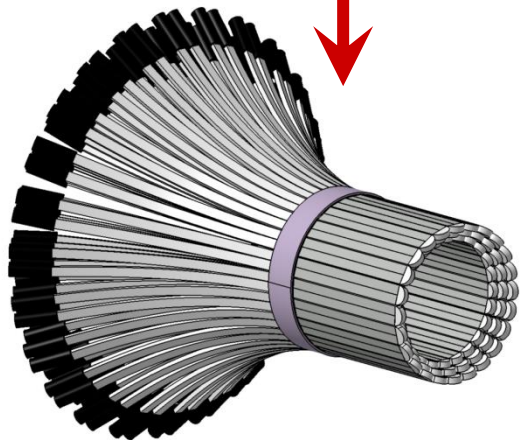
$$\vec{e} + d \rightarrow e' + n + \gamma + (p_s)$$

The **most sensitive** observable to the GPD **E**

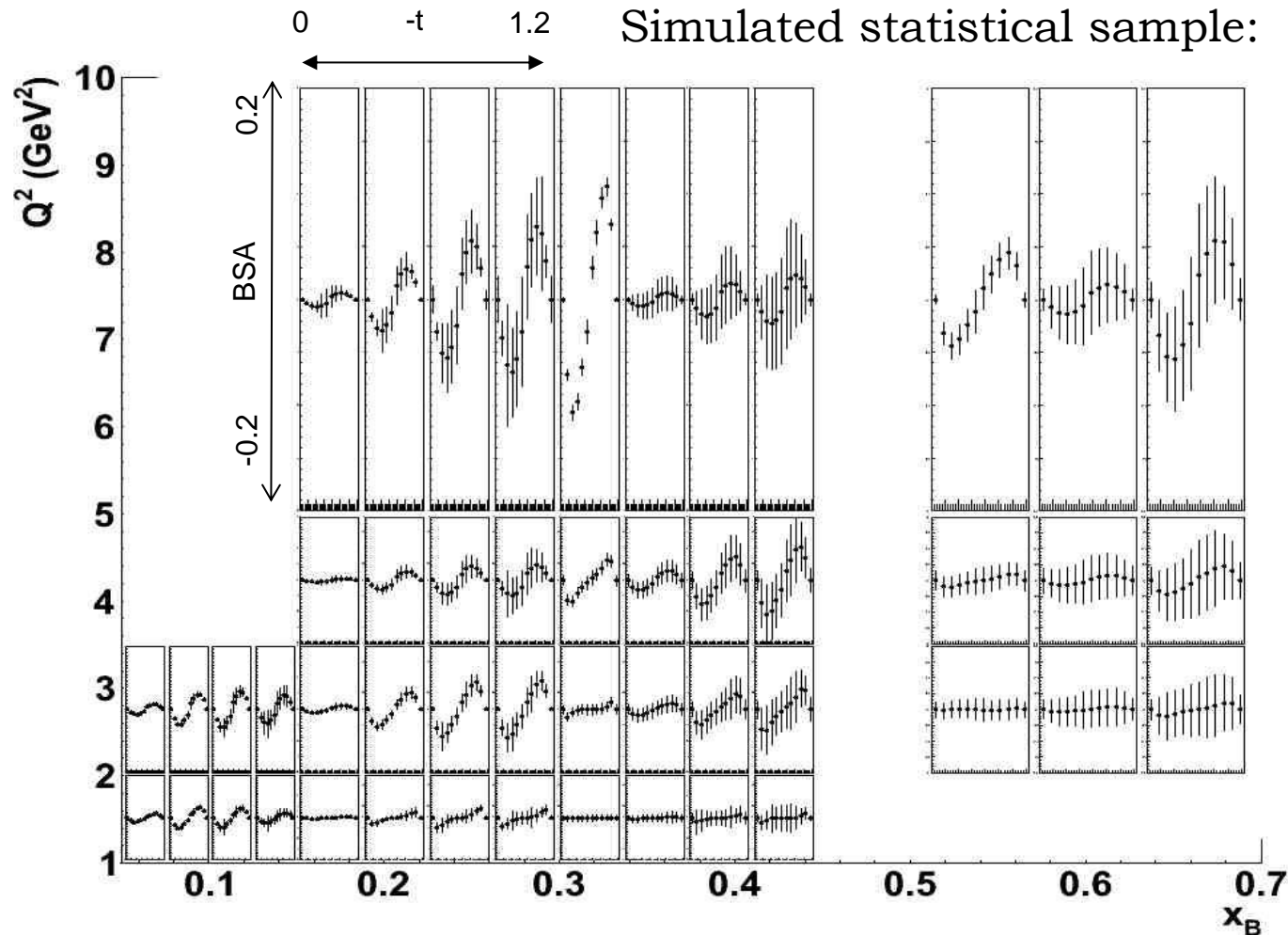
80 days of data taking
 $L = 10^{35} \text{ cm}^{-2}\text{s}^{-1}/\text{nucleon}$

CLAS12 +
Forward Calorimeter
+

Neutron Detector

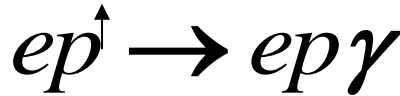


*Experiment approved,
detector constructed.*



Neutron DVCS with polarised targets

- ❖ DVCS transverse target spin asymmetry



Transverse-target spin asymmetry is **highly sensitive** to the **u-quark contributions** to proton spin.

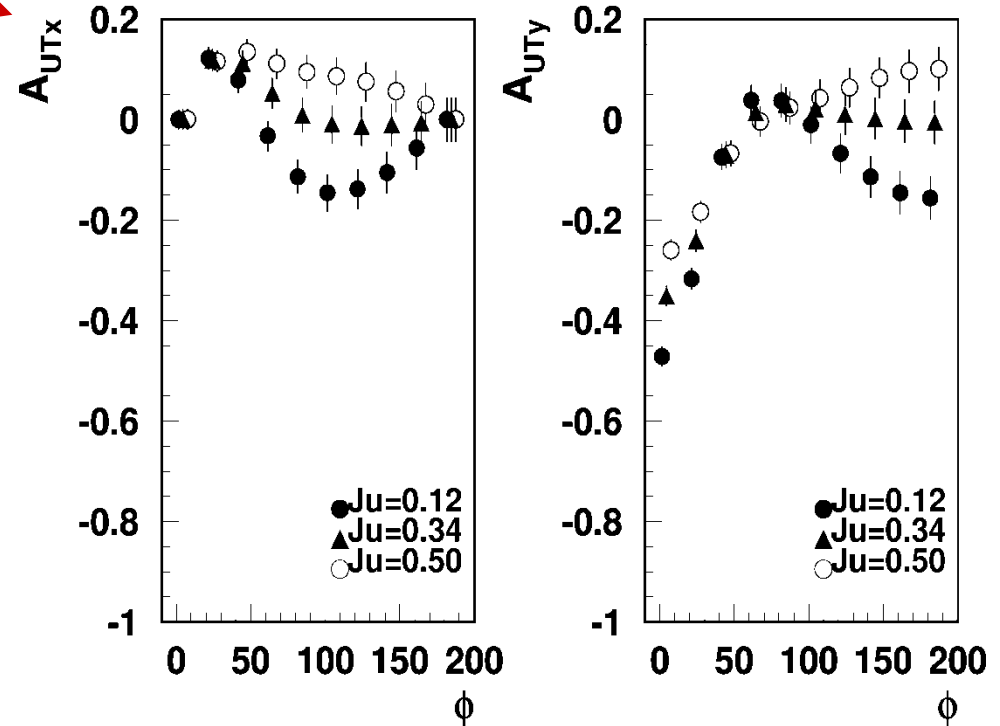
LOI approved at PAC38, target R&D under way.

- ❖ DVCS on longitudinally polarised (ND₃) target: proposal submitted to PAC40 (July 2015).

Crucial for flavour-decomposition of the GPDs.

Projected results

$Q^2=2.2 \text{ GeV}^2, x_B = 0.25, -t = 0.5 \text{ GeV}^2$



A_{UTx} Target polarization in scattering plane

A_{UTy} Target polarization perpendicular to scattering plane

Summary

- ❖ GPDs provide a 3D image of the internal dynamics of the nucleon and carry information on the composition of nucleon spin. They are experimentally accessible in exclusive reactions such as DVCS.
- ❖ Exclusive measurements of the **beam- and target-spin asymmetries in DVCS** on the **neutron** in the kinematic range opening up with CLAS12, in conjunction with those on the proton, will provide flavour decomposition of the GPDs and yield insight on the total angular momentum contribution of u, d quarks.
- ❖ An extraction of DVCS on **deuterium** @ 6GeV is underway – indications of a very low but measurable beam-spin asymmetry from the neutron and a first measurement of a target-spin asymmetry in nDVCS.

Thank you!



Back-up slides

*Wigner function:
full phase space parton
distribution of the nucleon*




$$\int d^2k_T$$

**Generalised
Parton
Distributions
(GPDs)**



*contain information on angular
momentum of quarks*


$$\int d^2b_T$$


**Transverse
Momentum
Distributions
(TMDs)**


$$\int dx$$


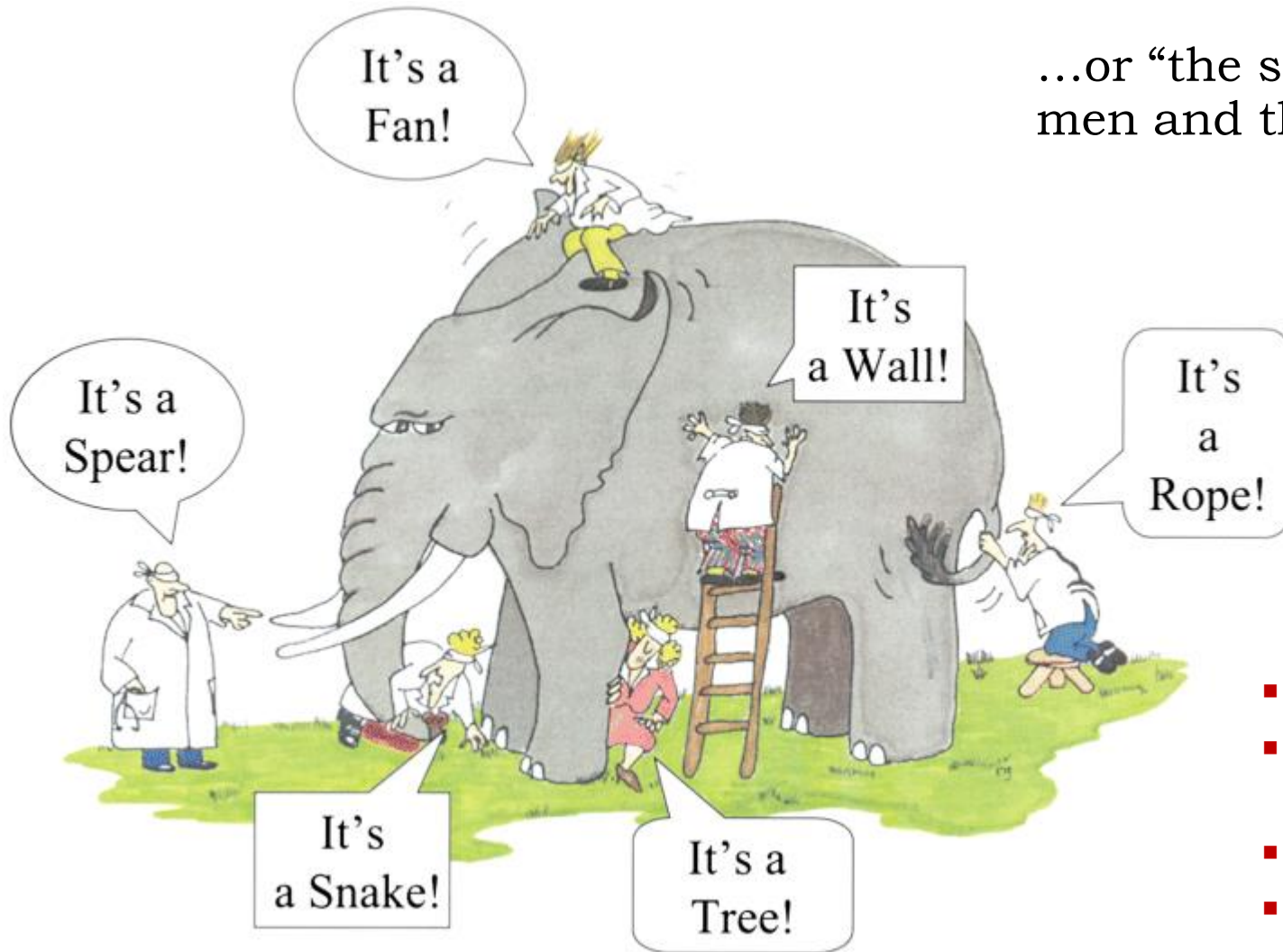
Form Factors


$$\int d^2k_T$$


**Parton Distribution
Functions (PDFs)**

A 100 views of the nucleon...

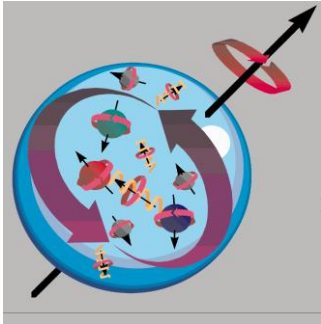
...or “the story of the blind men and the elephant”



- Elastic scattering
- Deep Inelastic Scattering (DIS)
- Semi-inclusive DIS
- Deep exclusive reactions

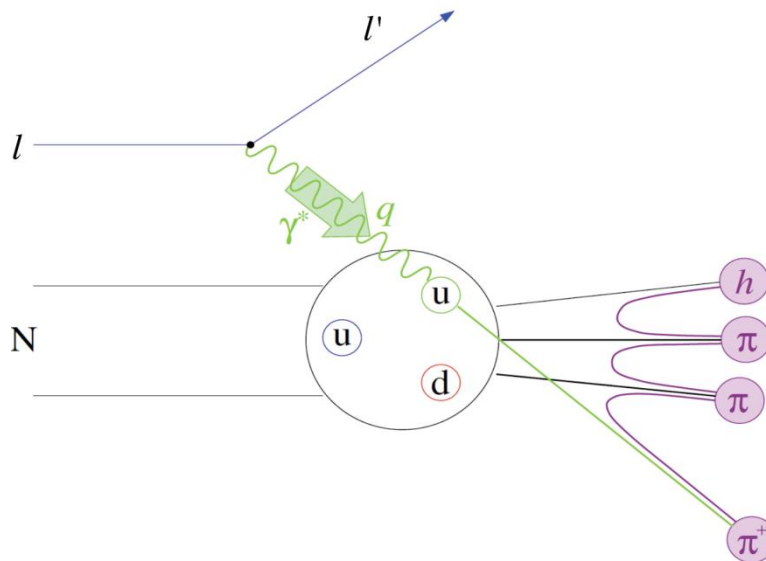
G. Renee Guzlas, artist.

Views of a nucleon: I



*Wigner function:
full phase space parton
distribution of the nucleon*

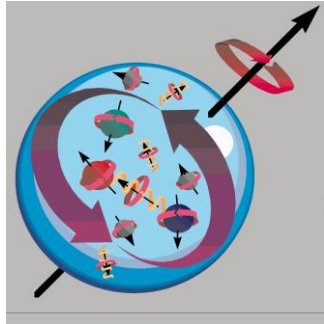
❖ Semi-inclusive Deep Inelastic Scattering:



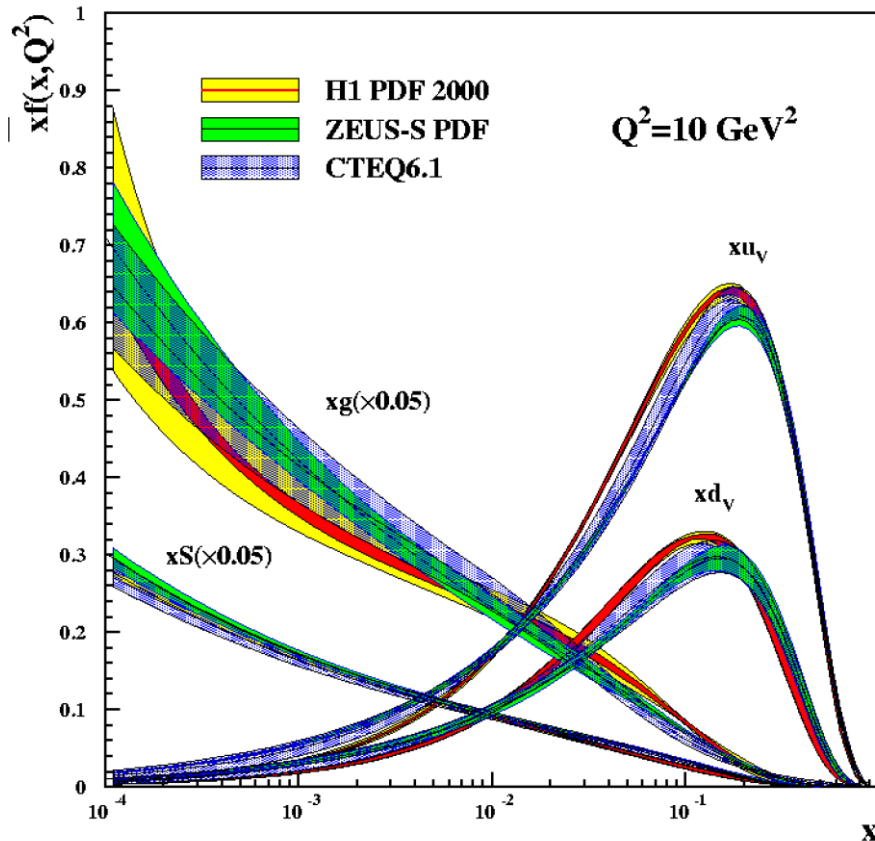
$$\int d^2b_T$$

**Transverse
Momentum
Distributions
(TMDs)**

Views of a nucleon: I



*Wigner function:
full phase space parton
distribution of the nucleon*



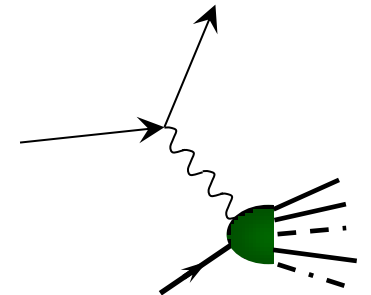
$$\int d^2 b_T$$

**Transverse
Momentum
Distributions
(TMDs)**

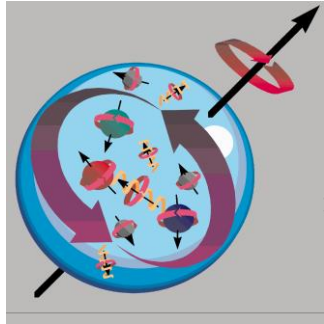
◆ Deep
Inelastic
Scattering

$$\int d^2 k_T$$

**Parton Distribution
Functions (PDFs)**



Views of a nucleon: II



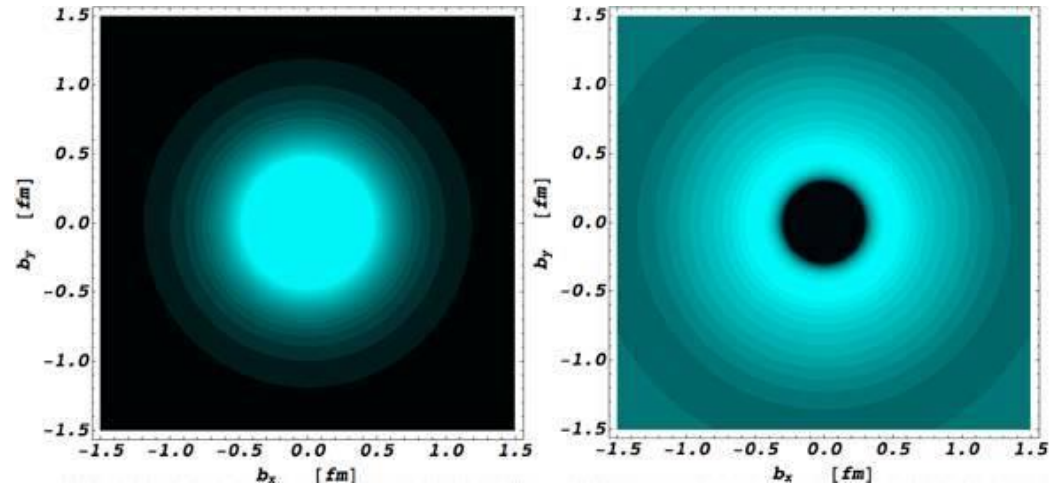
*Wigner function:
full phase space parton
distribution of the nucleon*



Fourier Transform of electric Form Factor:
transverse charge density of a nucleon

$$\int d^2 k_T$$

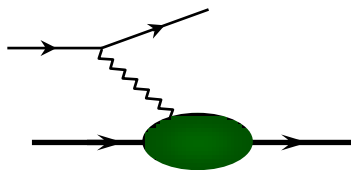
Generalised
Parton
Distributions
(GPDs)



proton

neutron

❖ Elastic
scattering



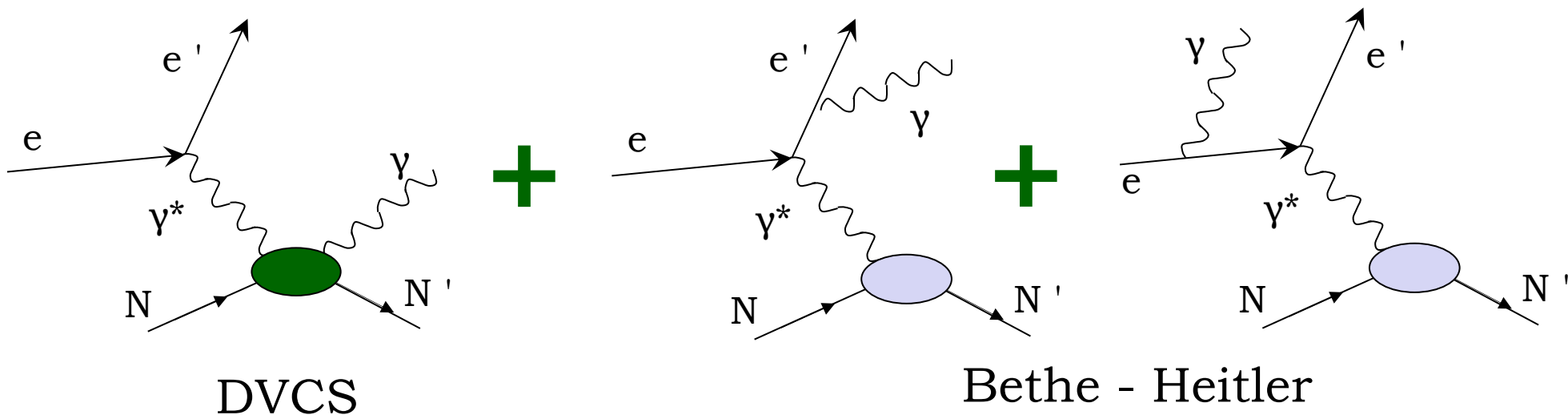
$$\int dx$$

Form Factors

G_E, G_M

Measuring DVCS

★ Process measured in experiment:



$$d\sigma \propto |T_{DVCS}|^2 + |T_{BH}|^2 + T_{BH} T_{DVCS}^* + T_{DVCS} T_{BH}^*$$

Amplitude
parameterised in
terms of Compton
Form Factors

Amplitude calculable
from elastic Form
Factors and QED!

Interference term

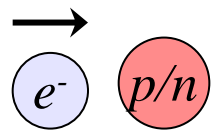
$$|T_{DVCS}|^2 \ll |T_{BH}|^2$$

Which DVCS experiment?

$H(x, \xi, t)$: Independent of quark helicity,
 $E(x, \xi, t)$: unpolarised GPDs

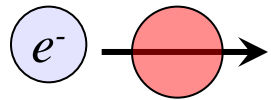
$\tilde{H}(x, \xi, t)$: Helicity-dependent,
 $\tilde{E}(x, \xi, t)$: polarised GPDs.

Beam, target polarisation



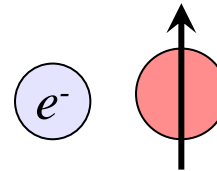
Proton Neutron

$Im\{\mathcal{H}_p, \tilde{\mathcal{H}}_p, \mathcal{E}_p\}$
 $Im\{\mathcal{H}_n, \tilde{\mathcal{H}}_n, \mathcal{E}_n\}$



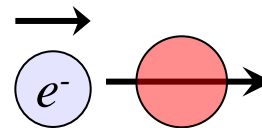
$Im\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\}$
 $Im\{\mathcal{H}_n, \mathcal{E}_n, \tilde{\mathcal{E}}_n\}$

Beam, target polarisation



Proton Neutron

$Im\{\mathcal{H}_p, \mathcal{E}_p\}$
 $Im\{\mathcal{H}_n\}$

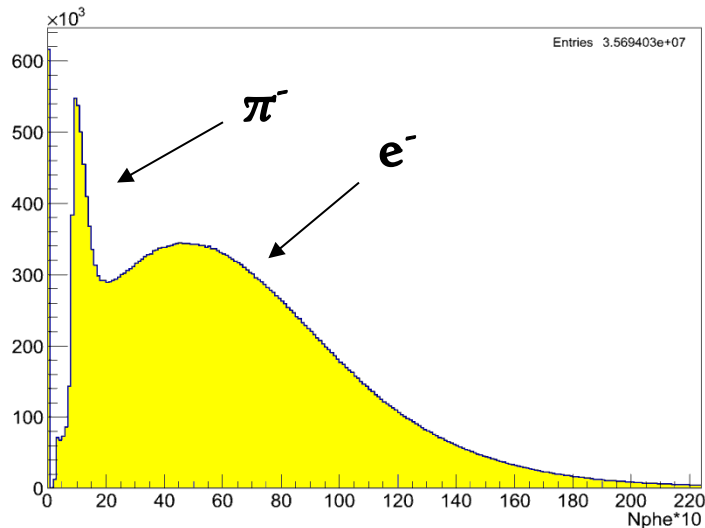


$Re\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\}$
 $Re\{\mathcal{H}_n, \mathcal{E}_n, \tilde{\mathcal{E}}_n\}$

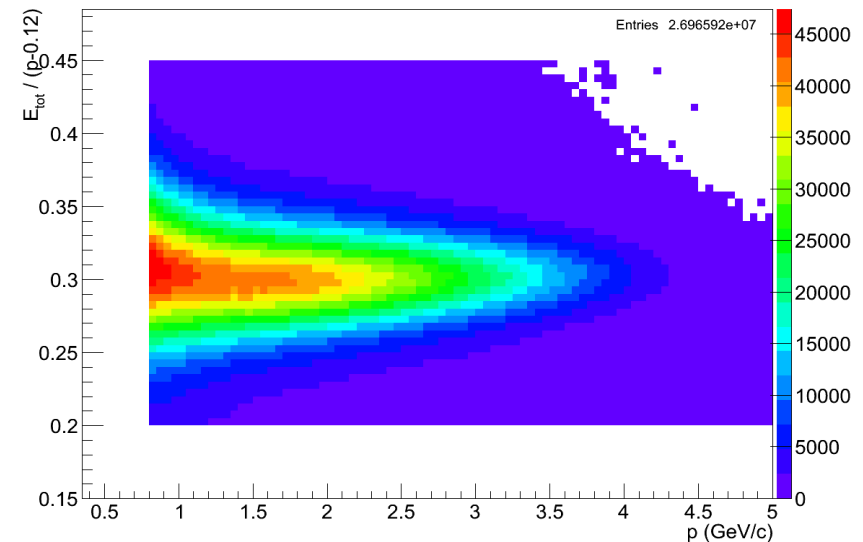
Particle ID – Electrons

- ❖ q and p from track-curvature through drift chambers in magnetic field
- ❖ **Separation from π^-** : on basis of energy deposit in electromagnetic calorimeter (EC) and number of photoelectrons produced in Cerenkov counters (CC).

of photoelectrons (x10) in CC



E deposit in EC / p vs. p



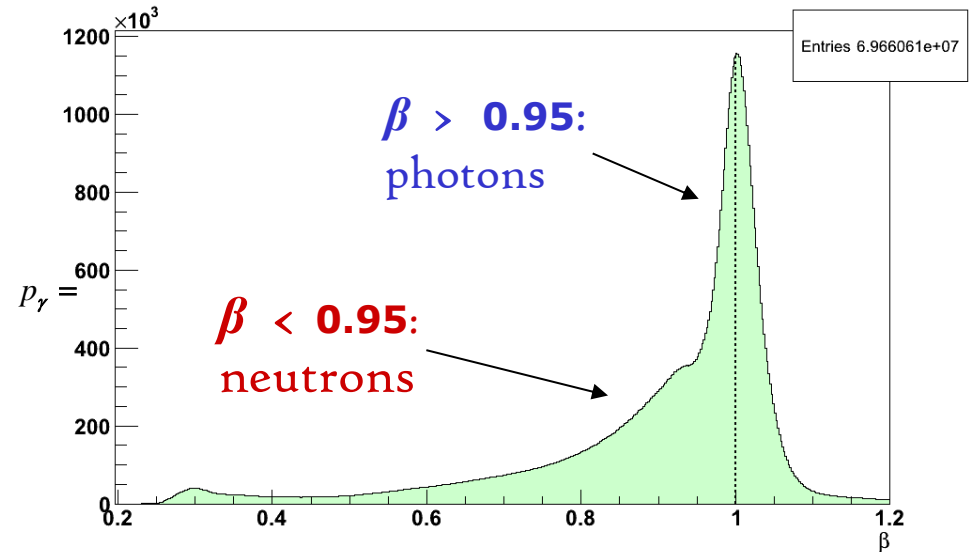
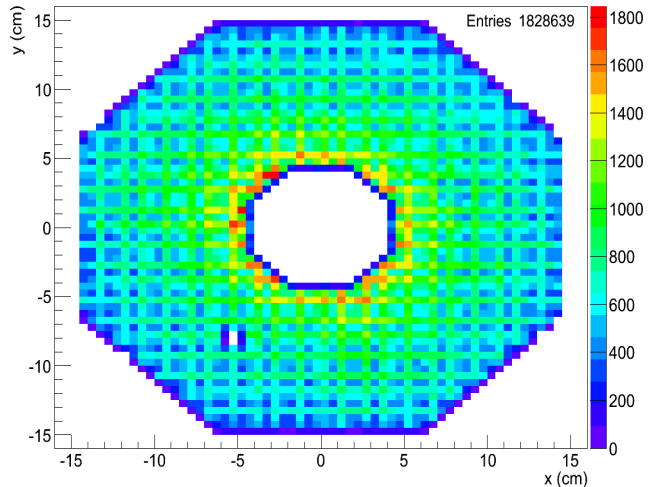
Particle ID – Photons and Neutrons

❖ β from neutral particles' time of flight to EC →

❖ Forward, low-angle photons in additional Inner Calorimeter



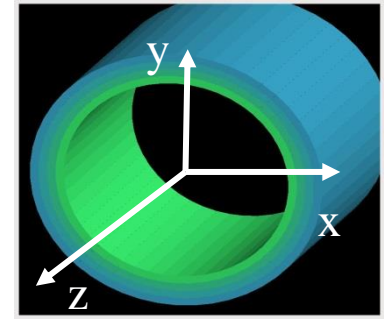
Hits in IC with E deposit > 1 GeV



❖ Neutrons:
$$p_n = \frac{\beta m_n}{\sqrt{1 - \beta^2}} \quad E_n = \sqrt{m_n^2 + p_n^2}$$

❖ Photons: E deposited in calorimeter

Neutron Detector for CLAS12

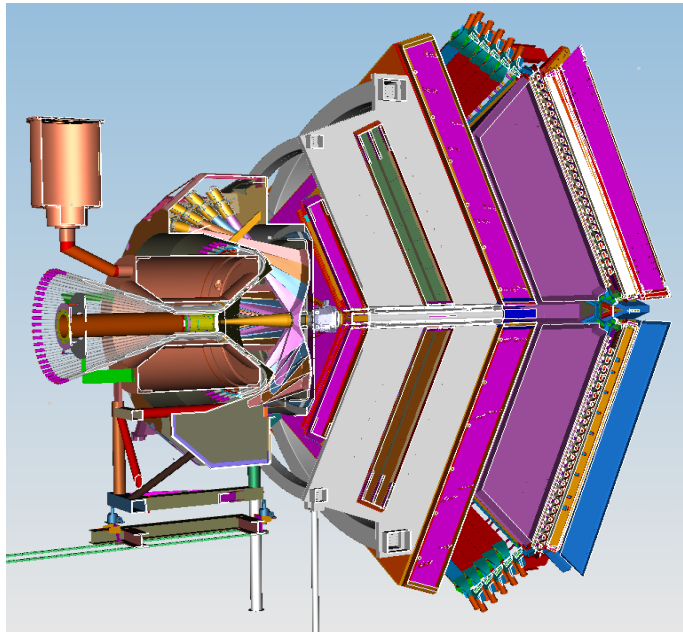


Available:

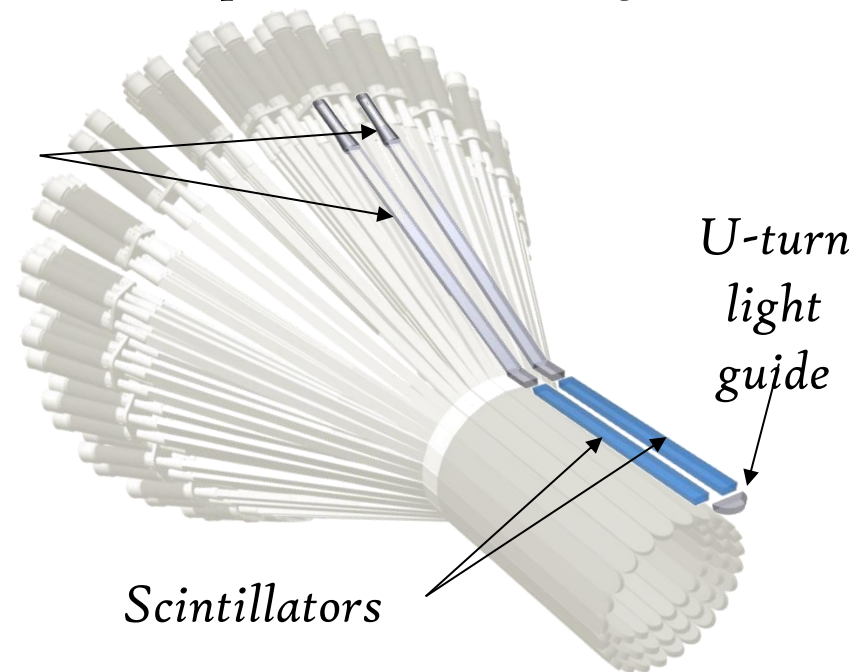
- ❖ 10 cm of radial space
- ❖ in a high magnetic field ($\sim 5\text{T}$)

Detector proposal approved:

- ❖ Plastic scintillator barrel: 3 layers, 48 paddles in each
- ❖ Length $\sim 70\text{ cm}$, inner radius 29 cm
- ❖ Long ($\sim 1.5\text{ m}$) light-guides
- ❖ PMT read-out upstream, out of high \mathbf{B} field



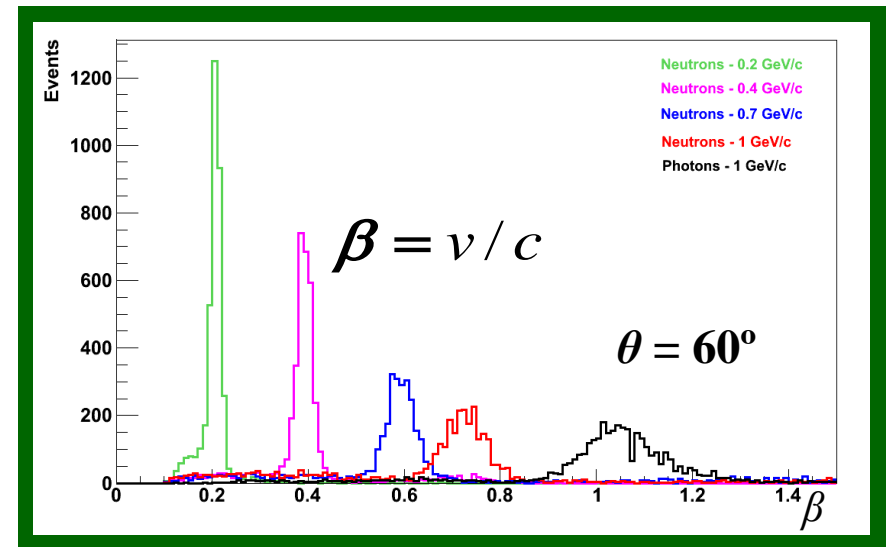
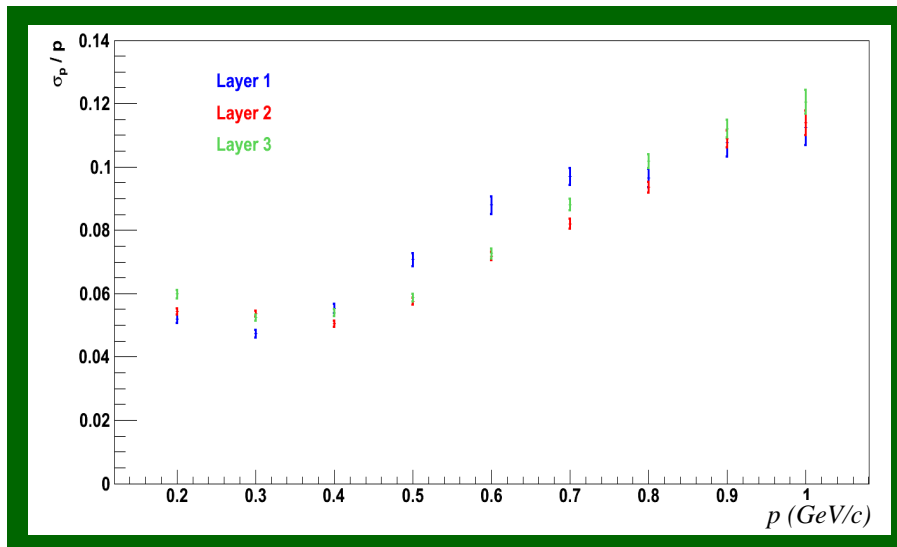
Light guides



CND Simulation (Geant 4)

- ❖ Neutron **efficiency** $\sim 8-9\%$
- ❖ Good **separation of neutrons and γ** up to $\sim 1\text{ GeV}/c$

- ❖ $\frac{\sigma_p}{p} \approx 5-12\%$ $\sigma_\theta \approx 2-3^\circ$
- ❖ 1-3% contamination from misreconstructed hits



Proposal Accepted in 2011 – detector constructed at IPN Orsay, France, by 2015. Installation in CLAS12: 2016.