

Baryon resonances in a combined analysis of pion- and photon-induced reactions

Recent results from the Juelich PWA

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In collaboration with:

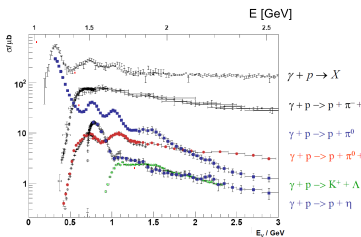
M. Döring, H. Haberzettl, J. Haidenbauer, C. Hanhart, F. Huang, S. Krewald, U.-G. Meißner,
and K. Nakayama

NSTAR2015

May 25, 2015, Osaka, Japan

The excited hadron spectrum: Connection between experiment and QCD in the non-perturbative regime

Experimental study of hadronic reactions



Major progress in recent years:

- enlarged data base with high quality for different final states
→ alternative source of information besides $\pi N \rightarrow X$
- measurement of several (double) polarization observables in $\gamma N \rightarrow X$
→ towards a *complete experiment*: unambiguous determination of the amplitude (up to an overall phase)

Extract information from experimental data:

e.g. unitarized ChPT, “classical” PWA, K-Matrix, unitary isobar models ...

Dynamical coupled channel (DCC) models:

- combined analysis of different reactions
- wide energy range
- theoretical constraints of the S -matrix are met (or approximated)

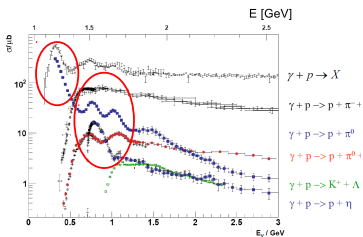
This talk: N and Δ
resonances



The excited hadron spectrum:

Connection between experiment and QCD in the non-perturbative regime

Experimental study of hadronic reactions



source: ELSA; data: ELSA, JLab, MAMI

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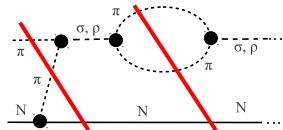
This talk: N and Δ resonances



Theoretical constraints of the S -matrix

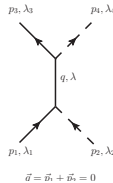
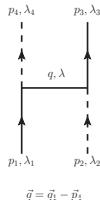
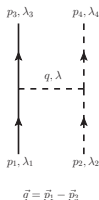
Unitarity: probability conservation

- 2-body unitarity
- 3-body unitarity:
 - discontinuities from t -channel exchanges
 - Meson exchange from requirements of the S -matrix [Aaron, Almado, Young, Phys. Rev. 174, 2022 (1968)]



Analyticity: from unitarity and causality

- correct structure of branch point, right-hand cut (real, dispersive parts)
- to approximate left-hand cut → Baryon u -channel exchange



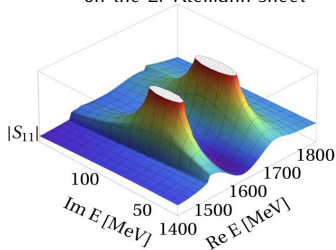
→ Resonances



Analytic structure of the amplitude

important information for a reliable determination of the resonance spectrum

Resonances: poles in the T -matrix
on the 2. Riemann sheet

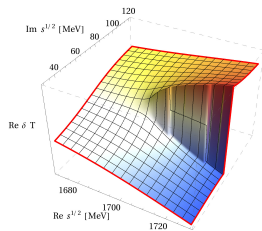


$\text{Re}(E_0)$ = "mass", $-2\text{Im}(E_0)$ = "width"

- pole position E_0 is the same in all channels
- residues \rightarrow branching ratios

Opening of inelastic channels:

\Rightarrow branch point and new Riemann sheet



Example: ρN branch point at

$$M_N + m_{\rho} = 1700 \pm i75 \text{ MeV}$$

Inclusion of branch points important to avoid
false resonance signal!

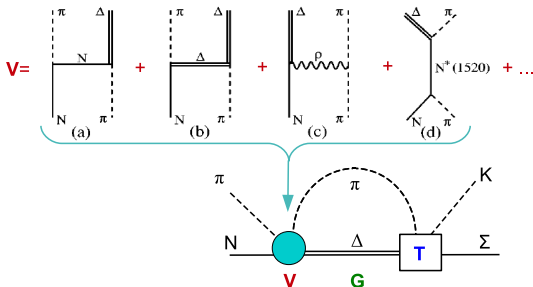


A dynamical coupled-channel approach: the Jülich model

Dynamical coupled-channels (DCC): simultaneous analysis of different reactions

The scattering equation in partial-wave basis

$$\langle L'S'p' | T_{\mu\nu}^I | LSp \rangle = \langle L'S'p' | V_{\mu\nu}^I | LSp \rangle + \sum_{\gamma, L''S''} \int_0^\infty dq \, q^2 \langle L'S'p' | V_{\mu\gamma}^I | L''S''q \rangle \frac{1}{E - E_\gamma(q) + i\epsilon} \langle L''S''q | T_{\gamma\nu}^I | LSp \rangle$$



- potentials V constructed from effective \mathcal{L}
- s -channel diagrams: T^P
genuine resonance states
- t - and u -channel: T^{NP}
dynamical generation of poles
partial waves strongly correlated

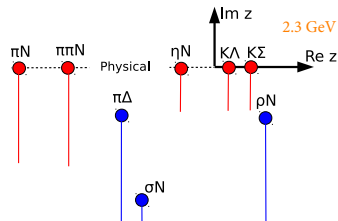
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$$\langle L'S'p' | T_{\mu\nu}^J | LSp \rangle = \langle L'S'p' | V_{\mu\nu}^J | LSp \rangle + \sum_{\gamma, L''S''} \int_0^\infty dq \, q^2 \langle L'S'p' | V_{\mu\gamma}^J | L''S''q \rangle \frac{1}{E - E_\gamma(q) + i\epsilon} \langle L''S''q | T_{\gamma\nu}^J | LSp \rangle$$

- (2-body) unitarity and analyticity respected
 - 3-body $\pi\pi N$ channel:
 - parameterized effectively as $\pi\Delta$, σN , ρN
 - $\pi N/\pi\pi$ subsystems fit the respective phase shifts
- ↳ branch points move into complex plane

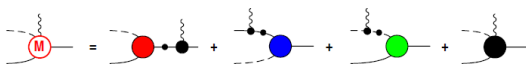


Photoproduction

Different approaches

- **Field theoretical** approaches : DMT, ANL-Osaka, **Jülich-Athens-Washington**, ...

Example: Gauge invariant formulation by Haberzettl, Huang and Nakayama
[Phys. Rev. C56 \(1997\)](#), [Phys. Rev. C74 \(2006\)](#), [Phys. Rev. C85 \(2012\)](#)



complicated, involved construction/calculation

Focus of the present analysis:

extraction of resonance parameters

⇒ flexible, **phenomenological** parameterization of photo excitation

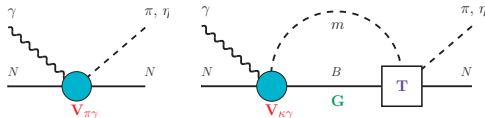
- Advantage: easy to implement, analyze large amounts of data
- Disadvantage: no information on microscopic reaction dynamics

Photoproduction in a semi-phenomenological approach

Multipole amplitude

$$M_{\mu\gamma}^{IJ} = V_{\mu\gamma}^{IJ} + \sum_{\kappa} T_{\mu\kappa}^{IJ} G_{\kappa} V_{\kappa\gamma}^{IJ}$$

(partial wave basis)



$$m = \pi, \eta, B = N, \Delta$$

$T_{\mu\kappa}$: Jülich hadronic T -matrix → Watson's theorem fulfilled by construction
→ **analyticity of T**: extraction of resonance parameters

Phenomenological potential:

$$V_{\mu\gamma}(E, q) = \begin{array}{c} \gamma \\ \text{wavy line} \\ \text{N} \text{---} \text{B} \\ \text{P}_{\mu}^{NP} \end{array} + \begin{array}{c} \gamma \\ \text{wavy line} \\ \text{N} \text{---} \text{N}^*, \Delta^* \\ \text{P}_i^P \end{array} \begin{array}{c} m \\ \text{dashed line} \\ \text{B} \\ \gamma_{\mu}^a \end{array} = \frac{\tilde{\gamma}_{\mu}^a(q)}{m_N} P_{\mu}^{NP}(E) + \sum_i \frac{\gamma_{\mu;i}^a(q) P_i^P(E)}{E - m_i^b}$$

$\tilde{\gamma}_{\mu}^a, \gamma_{\mu;i}^a$: hadronic vertices → correct threshold behaviour, cancellation of singularity at $E = m_i^b$
→ $\gamma_{\mu;i}^a$ affects **pion-** and **photon-**induced production of final state mB

i : resonance number per multipole; μ : channels $\pi N, \eta N, \pi \Delta$

Polynomials



Data analysis and fit results



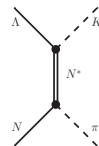
Combined analysis of pion- and photon-induced reactions

Simultaneous fit

Fit parameter:

- $\pi N \rightarrow \pi N$
 $\pi^- p \rightarrow \eta n, K^0 \Lambda, K^0 \Sigma^0, K^+ \Sigma^-$
 $\pi^+ p \rightarrow K^+ \Sigma^+$

s-channel: resonances (T^P)

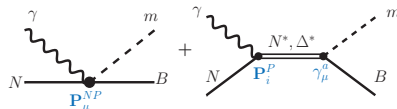


$$m_{bare} + f_{\pi NN^*}$$

⇒ 128 free parameters

11 N^* resonances \times (1 m_{bare} + couplings to $\pi N, \rho N, \eta N, \pi \Delta, K \Lambda, K \Sigma$)
 + 10 Δ resonances \times (1 m_{bare} + couplings to $\pi N, \rho N, \pi \Delta, K \Sigma$)

- $\gamma p \rightarrow \pi^0 p, \pi^+ n, \eta p$
 ⇒ up to 456 free parameters
 couplings of the polynomials





↳ calculations on the JUROPA supercomputer: parallelization in energy ($\sim 300 - 400$ processes)



Data base

simultaneous fit to π^- - and γ -induced reactions

	Fit A	Fit B
		
$\pi N \rightarrow \pi N$	PWA GW-SAID WI08 [Arndt <i>et al.</i> , PRC 86 (2012)]	
$\pi^- p \rightarrow \eta n$	$d\sigma/d\Omega, P$	
$\pi^- p \rightarrow K^0 \Lambda$	$d\sigma/d\Omega, P, \beta$	
$\pi^- p \rightarrow K^0 \Sigma^0$	$d\sigma/d\Omega, P$	
$\pi^- p \rightarrow K^+ \Sigma^-$	$d\sigma/d\Omega$	
$\pi^+ p \rightarrow K^+ \Sigma^+$	$d\sigma/d\Omega, P, \beta$	
	~ 6000 data points	
$\gamma p \rightarrow \pi^0 p$	$d\sigma/d\Omega, \Sigma, P, T, \Delta\sigma_{31}, G, H$	
$\gamma p \rightarrow \pi^+ n$	$d\sigma/d\Omega, \Sigma, P, T, \Delta\sigma_{31}, G, H$	
$\gamma p \rightarrow \eta p$	$d\sigma/d\Omega, P, \Sigma$	$d\sigma/d\Omega, P, \Sigma, T, F$
	29,392 data points	29,680 data points

- More single/double polarization:
 $E, C_{x'L}, C_{z'L}, T, P, H$ (ELSA 2014)
 $(\gamma p \rightarrow \pi^0 p)$
 \Rightarrow predictions

$\rightarrow T, F$: Akondi *et al.*
(A2 at MAMI)

PRL 113 10, 102001 (2014)



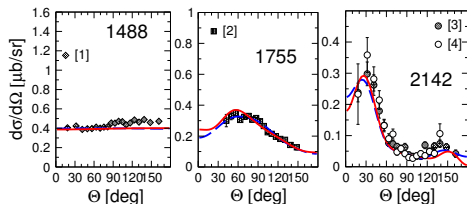
Fit results $\gamma p \rightarrow \eta p$

selected results, arXiv:1504.01643 [nucl-th]

--- T, F not included

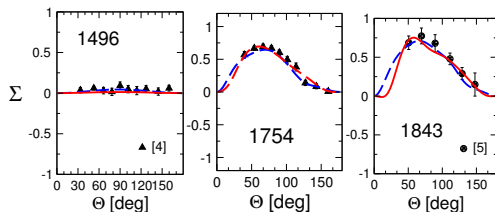
— T, F included

• Differential cross section



[1] McNicoll 2010 (MAMI), [2] Williams 2009 (JLab), [3] Credé 2009 (ELSA), [4] Credé 2005 (ELSA)

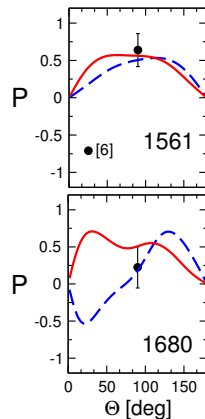
• Beam asymmetry



[4] Bartalini 2007 (GRAAL), [5] Elsner 2007 (ELSA)

• Recoil polarization

- only 7 data points in total -

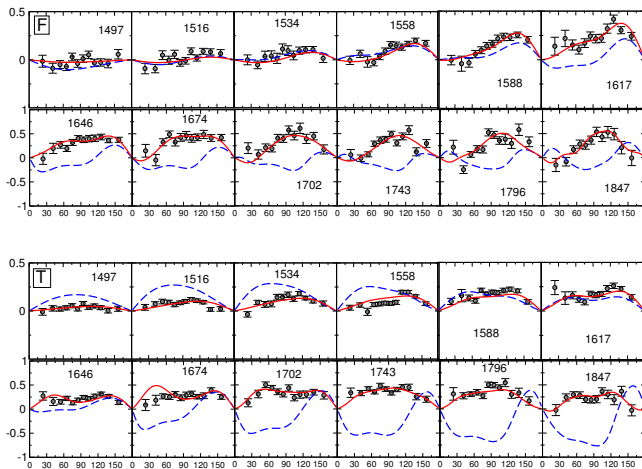


[6] Heusch (Caltech),
PRL25, 1381 (1970)



F and T in $\gamma p \rightarrow \eta p$

arXiv:1504.01643 [nucl-th]

 Data: Akondi *et al.* (A2 at MAMI) PRL 113 10, 102001 (2014)


--- prediction
— fit

Polarization:

Beam	Target	Recoil
+1	+x	0
-1	+x	0

Polarization:

Beam	Target	Recoil
0	+y	0
0	-y	0



Resonance content: $I=1/2$

arXiv:1504.01643 [nucl-th]

Pole search on the 2^{nd} sheet of the scattering matrix $T_{\mu\nu}$

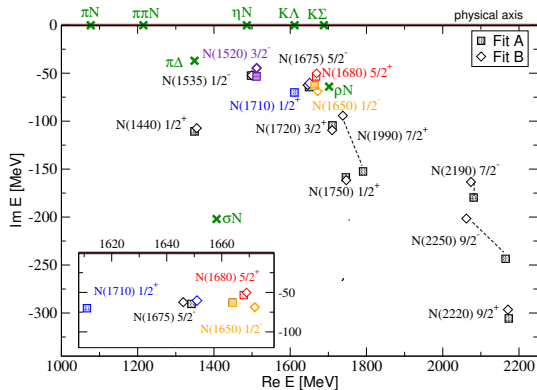
Resonance parameter:

- "mass" = $Re(E_0)$
- "width" = $-2Im(E_0)$
- Residues \rightarrow branching ratios

E_0 : pole position

\rightarrow no new states compared to Jülich2012
(Jülich2012: only pion-induced data)

\rightarrow no narrow structure at 1.68 GeV
(seen in eta photoproduction on the neutron)



x: branch points Notation: $N(\text{"name"}) J^{\text{parity}}$



Resonance parameters

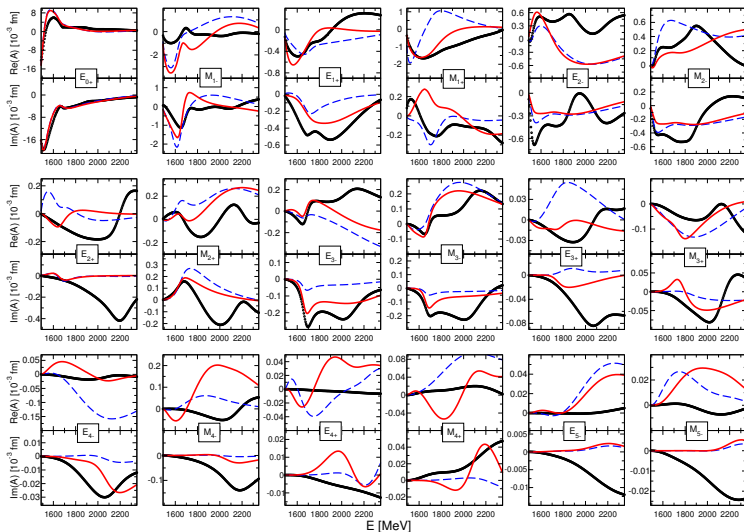
selected results, arXiv:1504.01643 [nucl-th]

fit	Re E_0	$-2\text{Im } E_0$	$ r_{\pi N} $	$\theta_{\pi N \rightarrow \pi N}$	$\frac{\Gamma_{\pi N}^{1/2} \Gamma_{\eta N}^{1/2}}{\Gamma_{\text{tot}}}$	$\theta_{\pi N \rightarrow \eta N}$	
	[MeV]	[MeV]	[MeV]	[deg]	[%]	[deg]	
$N(1535) 1/2^-$	A	1497	105	23	-48	51	110
	B	1499	104	22	-46	51	112
	A_{had}	1498	74	17	-37	51	120
$N(1710) 1/2^+$	A	1611	140	2.7	-40	6.1	175
	B	1651	121	3.2	55	16	-180
	A_{had}	1637	97	4	-30	24	130

fit A: T, F not includedfit B: T, F includedfit A_{had} : Jülich2012, only pion-induced data

Multipoles for $\gamma p \rightarrow \eta p$

Comparison with the Bonn-Gatchina 2014-02 solution



(Black) solid line: BG 2014-02. Dashed (blue) line: fit A; solid (red) line: fit B.

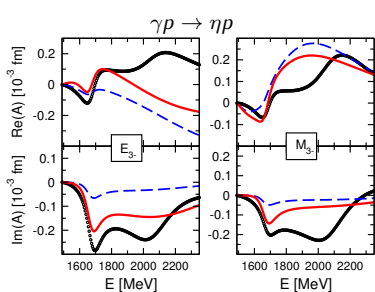
$\gamma p \rightarrow \pi^0 p$



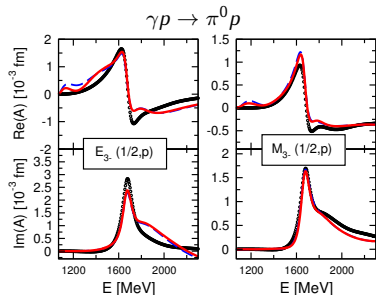
Multipoles: $\gamma p \rightarrow \eta p$ vs $\gamma p \rightarrow \pi^0 p$

Comparison with the Bonn-Gatchina 2014-02 solution

- Example: E_{3-} and M_{3-} multipoles (F_{15} , F_{35})



(Black) solid line: BG 2014-02. Dashed (blue) line: fit A; solid (red) line: fit B.



(Black) solid line: BG 2014-02. Dashed (blue) line: fit A; solid (red) line: fit B.

- ⇒ Multipole content of $\gamma p \rightarrow \eta p$ seems less well established than for $\gamma p \rightarrow \pi N$
- ⇒ Convergence with larger data base of $\gamma p \rightarrow \eta p$?



Summary

Extraction of the N^* and Δ resonance spectrum

from a simultaneous analysis of pion- and photon-induced reactions

- DCC analysis of $\pi N \rightarrow \pi N, \eta N, K\Lambda$ and $K\Sigma$
- π and η photoproduction in a semi-phenomenological approach

Comparison of 3 different fits:

- simultaneous fit of $\pi N, \gamma N \rightarrow X$ **without** recent MAMI T and F data
- simultaneous fit of $\pi N, \gamma N \rightarrow X$ **with** recent MAMI T and F data
- earlier fit (Jülich2012), only $\pi N \rightarrow X$

- ⇒ noticeable influence of photoproduction data in general / new polarization observables
- on pole positions and photocouplings
 - on hadronic couplings



Thank you for your attention!



Photocouplings at the pole

selected results, arXiv:1504.01643 [nucl-th]

$$\tilde{A}_{pole}^h = A_{pole}^h e^{i\vartheta^h}$$

$$h = 1/2, 3/2$$

$$\tilde{A}_{pole}^h = I_F \sqrt{\frac{q_p}{k_p} \frac{2\pi(2J+1)E_0}{m_N r_{\pi N}}} \text{Res } A_{L\pm}^h$$

I_F : isospin factor
 q_p (k_p): meson (photon) momentum at the pole
 $J = L \pm 1/2$ total angular momentum
 E_0 : pole position
 $r_{\pi N}$: elastic πN residue

	fit	$A_{pole}^{1/2}$	$\vartheta^{1/2}$
		[$10^{-3} \text{ GeV}^{-\frac{1}{2}}$]	[deg]
N(1535) $1/2^-$	A	107	4.6
	B	106	5.2
	2	50_{-4}^{+4}	-14_{-10}^{+12}
N(1710) $1/2^+$	A	7.1	-177
	B	20	-83
	2	28_{-2}^{+9}	103_{-6}^{+20}

fit A: T, F not included

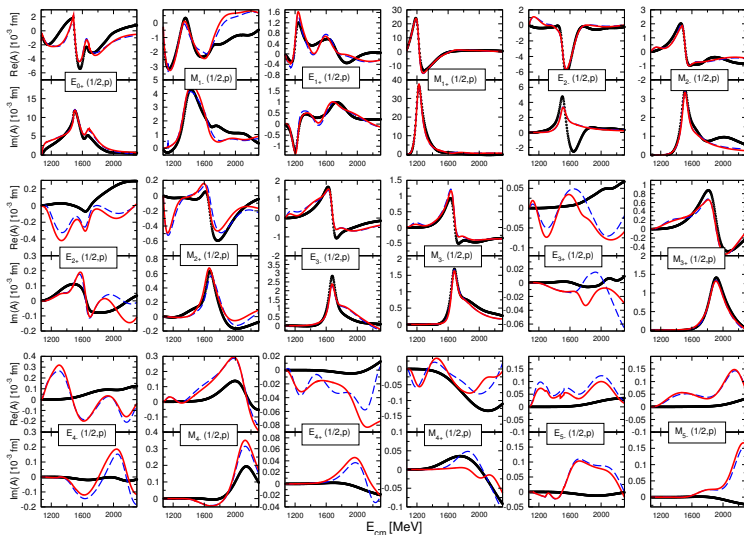
fit B: T, F included

fit 2: Jülich2013, only
 pion photoproduction
 (same pole positions as fit A_{had})



Multipoles: $\gamma p \rightarrow \pi^0 p$

Comparison with the Bonn-Gatchina 2014-02 solution

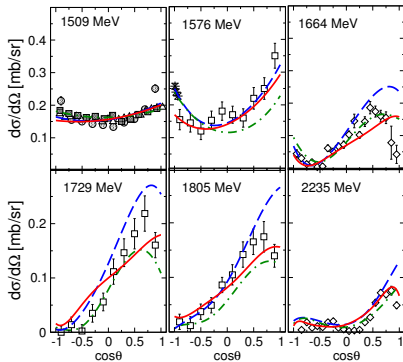
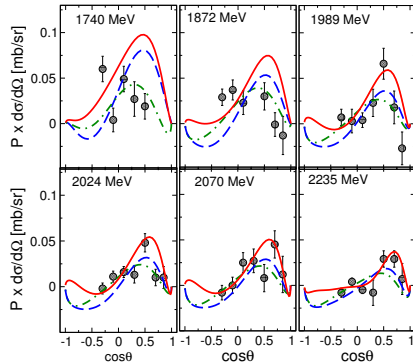


(Black) solid line: BG 2014-02. Dashed (blue) line: fit A; solid (red) line: fit B.

[← back](#)

Pion-induced eta production: $\pi^- p \rightarrow \eta n$

arXiv:1504.01643 [nucl-th]

fit A: T, F not includedfit B: T, F includedfit A_{had}: Jülich2012, only pion-induced data

Resonance content: $I=3/2$

arXiv:1504.01643 [nucl-th]

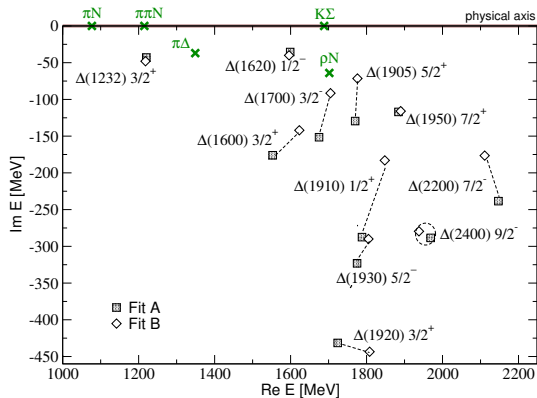
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- Residues \rightarrow branching ratios

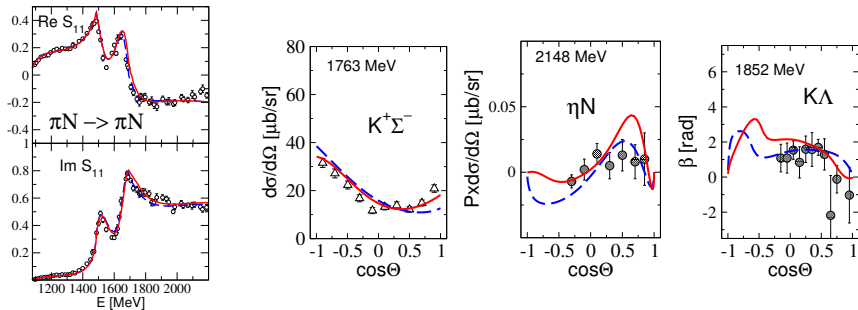
E_0 : pole position

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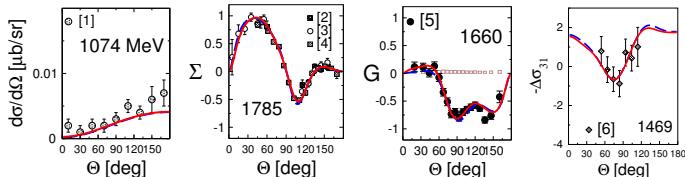
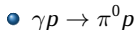
x: branch points Notation: $N(\text{"name"}) J^{\text{parity}}$



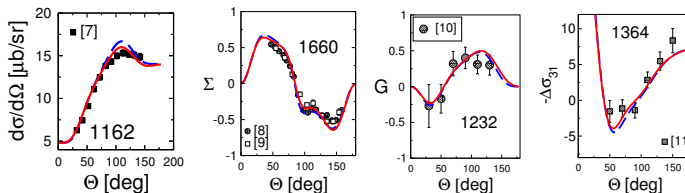
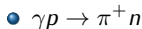
Selected results: **Fit A** and **Fit B**

Pion photoproduction: selected fit results

arXiv:1504.01643 [nucl-th]



- [1] Schmidt 2001 (MAMI)
 [2] Elsner 2009 (ELSA)
 [3] Sparks 2010 (ELSA)
 [4] Bartalini 2005 (GRAAL)
 [5] Thiel 2012 (ELSA)
 [6] Ahrens 2002 (MAMI)



- [7] Ahrens 2004 (MAMI)
 [8] Bartalini 2002 (GRAAL)
 [9] Ajaka 2000 (GRAAL)
 [10] Ahrens 2005 (MAMI)
 [11] Ahrens 2006 (MAMI)



Double polarization in $\gamma p \rightarrow \pi^0 p$

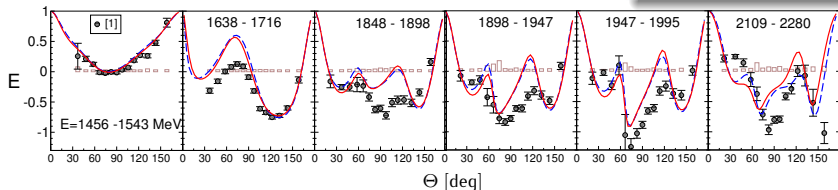
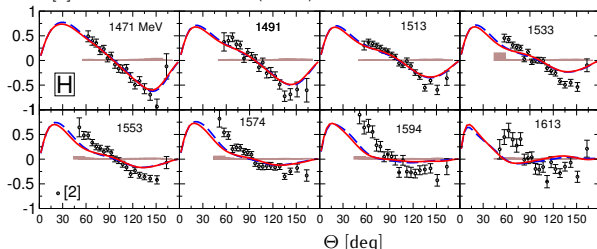
Data NOT included in fit

selected results, arXiv:1504.01643 [nucl-th]

 [1] Gottschall *et al.* 2013 (ELSA) PRL 112 1, 012003

Polarization

Beam	Target	Recoil
+1	-z	0
-1	-z	0

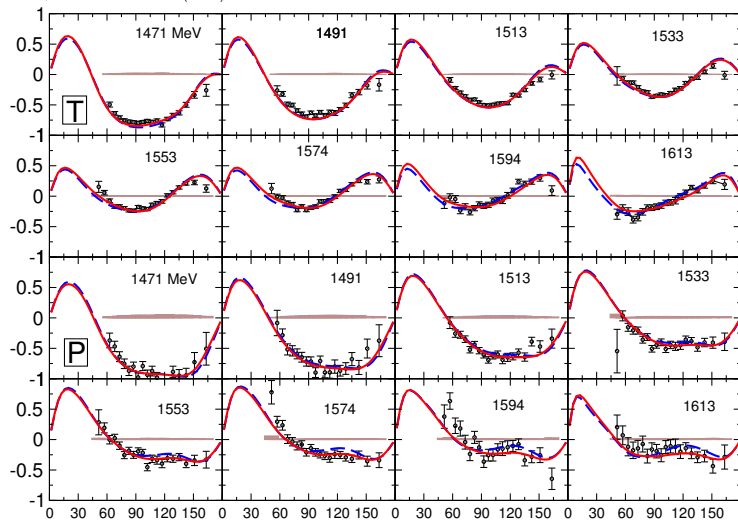

 [2] Hartmann *et al.* 2014 (ELSA) PRL 113, 062001


Polarization

Beam	Target	Recoil
\perp'	x	0
\parallel'	x	0

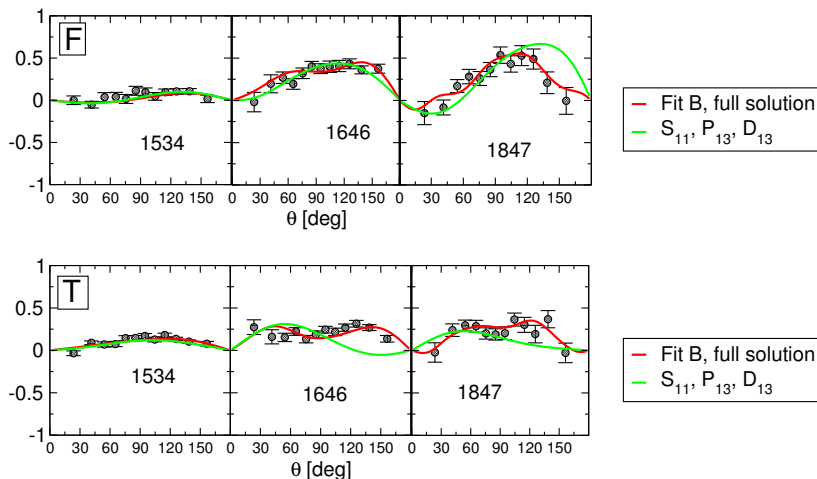
new ELSA T,P data



T, P in $\gamma p \rightarrow \pi^0 p$
 Data NOT included in the Fit
Data: J. Hartmann *et al.* 2014 (ELSA)[nucl-th] [← back](#)

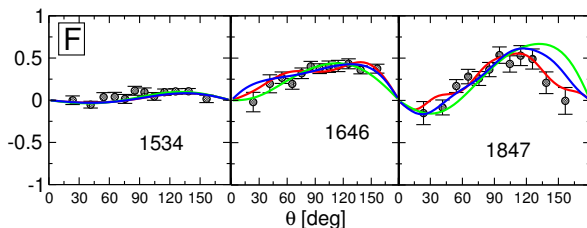
Partial wave contribution in F and T in $\gamma p \rightarrow \eta p$

arXiv:1504.01643 [nucl-th]

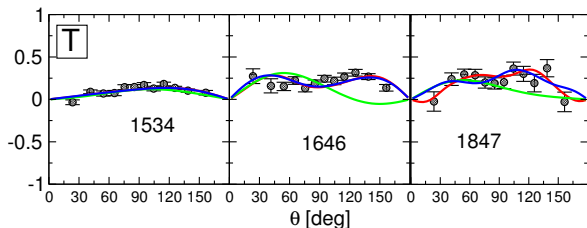


Partial wave contribution in F and T in $\gamma p \rightarrow \eta p$

arXiv:1504.01643 [nucl-th]



— Fit B, full solution
— S_{11}, P_{13}, D_{13}
— all s-, p-, d-, f-waves



— Fit B, full solution
— S_{11}, P_{13}, D_{13}
— all s-, p-, d-, f-waves



Details of the formalism

Polynomials:

$$P_i^P(E) = \sum_{j=1}^n g_{i,j}^P \left(\frac{E - E_0}{m_N} \right)^j e^{-g_{i,n+1}^P (E - E_0)}$$

$$P_\mu^{\text{NP}}(E) = \sum_{j=0}^n g_{\mu,j}^{\text{NP}} \left(\frac{E - E_0}{m_N} \right)^j e^{-g_{\mu,n+1}^{\text{NP}} (E - E_0)}$$

◀ back

- $E_0 = 1077$ MeV
- $g_{i,j}^P, g_{\mu,j}^{\text{NP}}$: fit parameter
- $e^{-g(E-E_0)}$: appropriate high energy behavior
- $n = 3$



Data base

simultaneous fit to $\pi N \rightarrow \pi N, \eta N, K\Lambda, K\Sigma$ World data base on $\eta N, K\Lambda, K\Sigma$

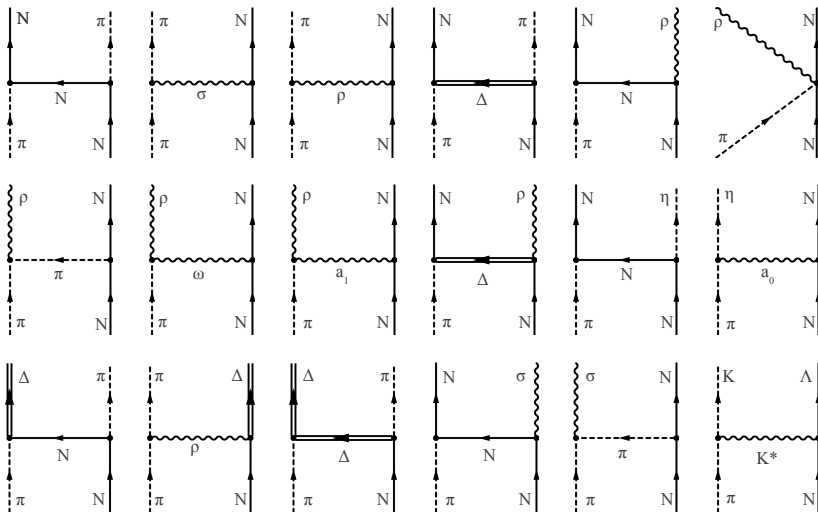
	PWA	σ_{tot}	$\frac{d\sigma}{d\Omega}$	P	β
$\pi N \rightarrow \pi N$	GWU/SAID 2006 up to J=9/2				
$\pi^- p \rightarrow \eta n$		62 data points	38 energy points z=1489 to 2235 MeV	12 energy points 1740 to 2235 MeV	
$\pi^- p \rightarrow K^0 \Lambda$		66 data points	46 energy points 1626 to 1405 MeV	27 energy points 1633 to 2208 MeV	7 energy points 1852 to 2262 MeV
$\pi^- p \rightarrow K^0 \Sigma^0$		16 data points	29 energy points 1694 to 2405 MeV	19 energy points 1694 to 2316 MeV	
$\pi^- p \rightarrow K^+ \Sigma^-$		14 data points	15 energy points 1739 to 2405 MeV		
$\pi^+ p \rightarrow K^+ \Sigma^+$		18 data points	32 energy points 1729 to 2318 MeV	32 energy points 1729 to 2318 MeV	2 energy points 2021 and 2107 MeV

~ 6000 data points

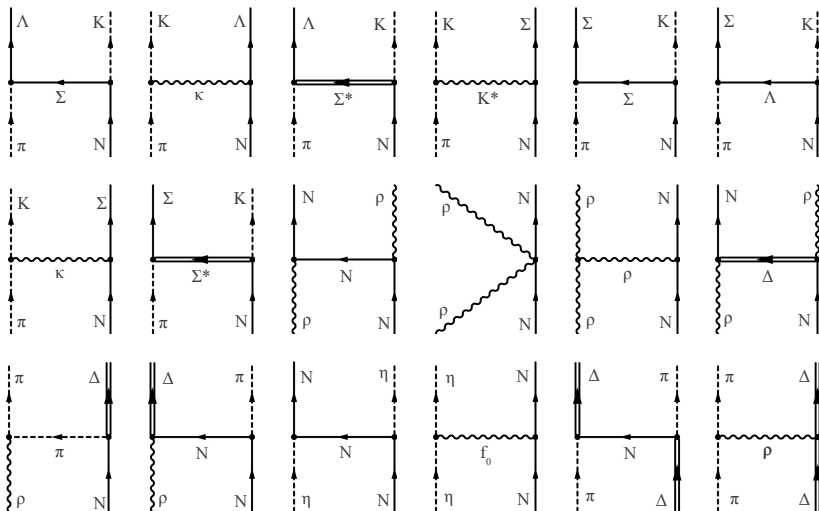
◀ back



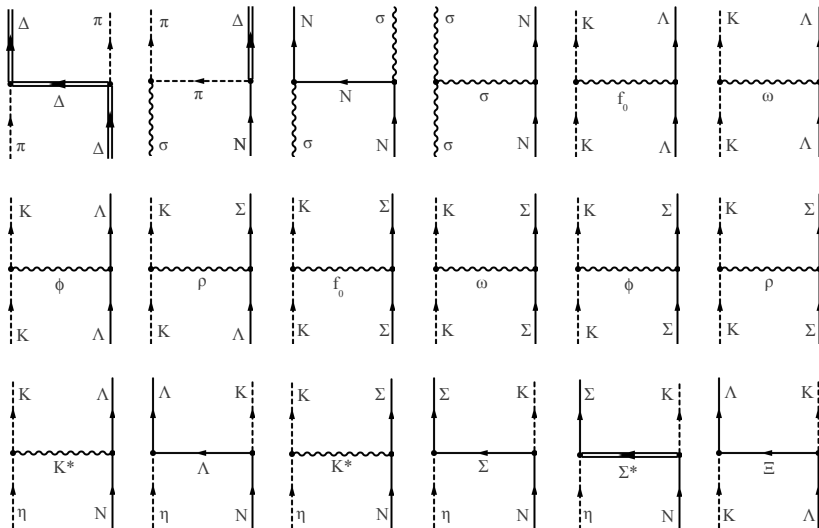
◀ back



◀ back

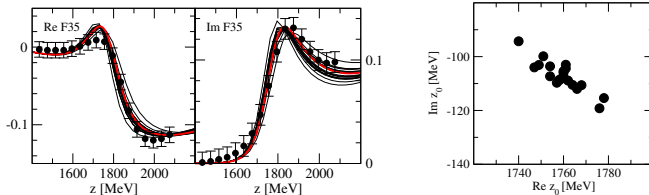


◀ back



Error analysis

- $\chi^2 + 1$ criterion: determination of the non-linear parameter error
 - error of parameter p_i determined by range of p_i such that χ_{min}^2 rises by less than 1
- ⇒ error on pole positions and residues.



NPBA 851, 58 (2011)

BUT: numerically not possible with ≥ 500 free parameters

Work in progress: Developing of techniques to apply Monte-Carlo error propagation using bootstrap method (M. Döring et al.)



Matching to lattice

Prediction & analysis of lattice data

[M. Döring et al., EPJ A47, 163 (2011)]

Scattering equation:

$$T(q'', q') = V(q'', q') + \int_0^\infty dq q^2 V(q'', q) \frac{1}{z - E_1(q) - E_2(q) + i\epsilon} T(q, q')$$

Discretization of momenta in the scattering equation:

$$\int \frac{\vec{d}^3 q}{(2\pi)^3} f(|\vec{q}|^2) \rightarrow \frac{1}{L^3} \sum_{\vec{n}_i} f(|\vec{q}_i|^2), \quad \vec{q}_i = \frac{2\pi}{L} \vec{n}_i, \quad \vec{n}_i \in \mathbb{Z}^3$$

$$T(q'', q') = V(q'', q') + \frac{2\pi^2}{L^3} \sum_{i=0}^{\infty} \vartheta(i) V(q'', q_i) \frac{1}{z - E_1(q_i) - E_2(q_i)} T(q_i, q')$$

$\vartheta^{(P)}(i)$ series

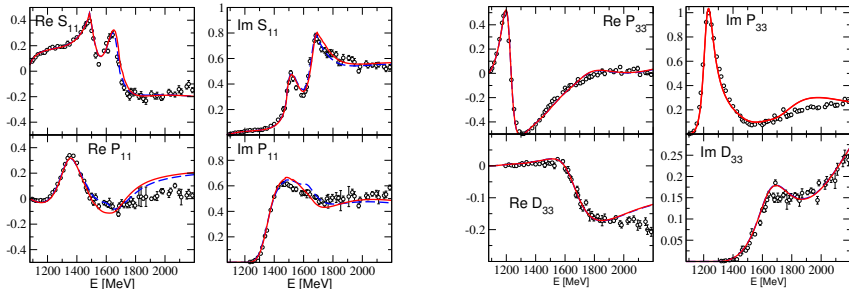
- Study finite-volume effects
- Predict lattice spectra



$\pi N \rightarrow \pi N$ partial wave amplitudes

selected results, arXiv:1504.01643 [nucl-th]

Fit A and Fit B

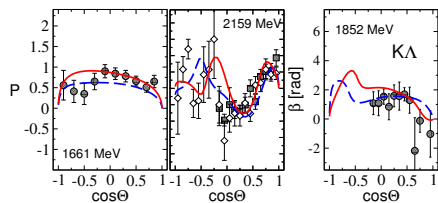
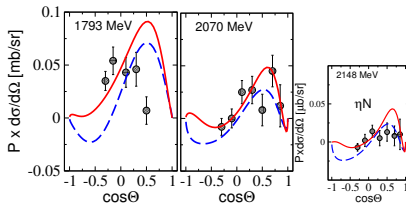
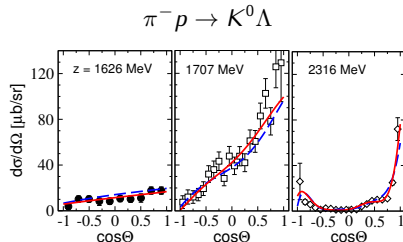
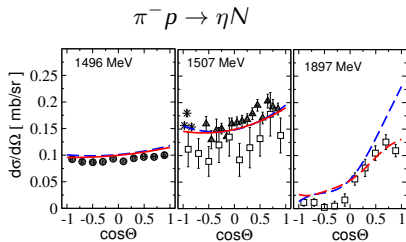


- Notation: L_{2I2J}
- Input to fit: energy-dependent partial wave analysis, GWU/SAID 2006 up to $J = 9/2$ ($\sim H_{39}$)



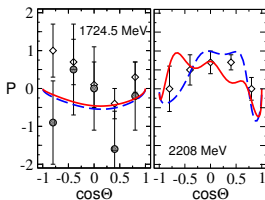
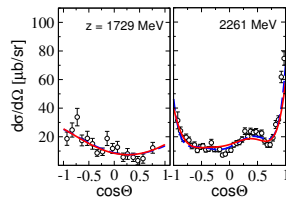
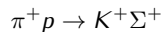
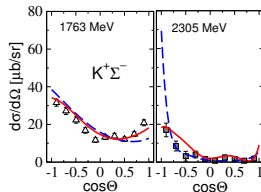
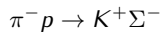
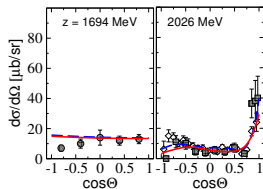
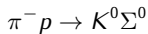
$\pi N \rightarrow \eta N, K\Lambda$

selected results, arXiv:1504.01643 [nucl-th]

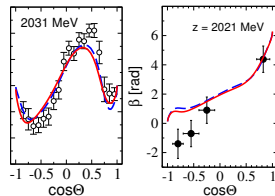


$\pi N \rightarrow K \Sigma$

selected results, arXiv:1504.01643 [nucl-th]

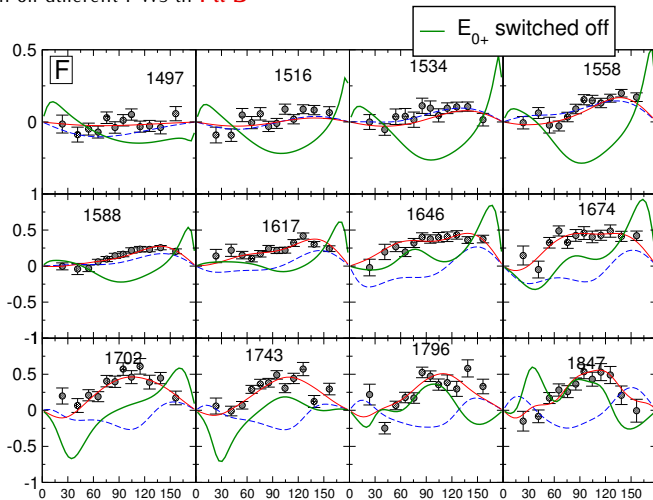


No polarization data!



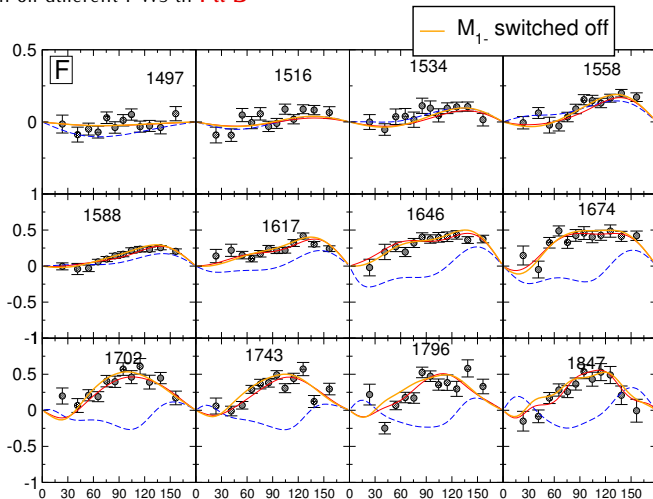
Partial wave contribution to F in $\gamma p \rightarrow \eta p$

arXiv:1504.01643 [nucl-th]

Switch off different PWs in **Fit B**

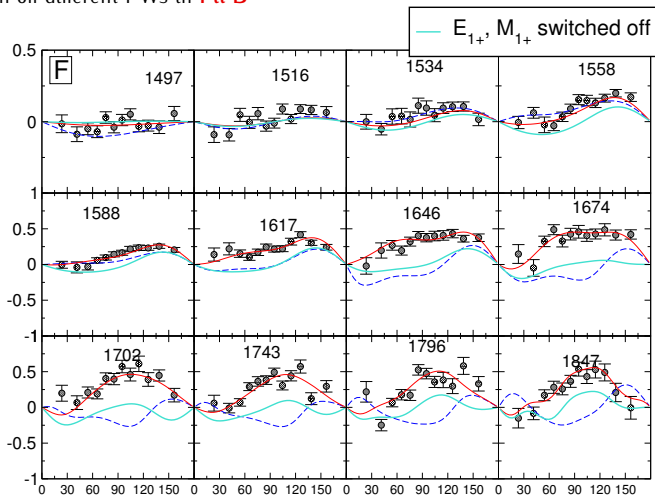
Partial wave contribution to F in $\gamma p \rightarrow \eta p$

arXiv:1504.01643 [nucl-th]

Switch off different PWs in **Fit B**

Partial wave contribution to F in $\gamma p \rightarrow \eta p$

arXiv:1504.01643 [nucl-th]

Switch off different PWs in **Fit B**

Photoproduction of pseudoscalar meson

- Photocouplings of resonances
- high precision data from ELSA, MAMI, JLab... → resolve questionable/find new states

Photoproduction amplitude of pseudoscalar mesons:

[Chew, Goldberger, Low, and Nambu, Phys. Rev. 106, 1345 \(1957\)](#)

$$\hat{\mathcal{M}} = F_1 \vec{\sigma} \cdot \vec{\epsilon} + i F_2 \vec{\epsilon} \cdot (\hat{k} \times \hat{q}) + F_3 \vec{\sigma} \cdot \hat{k} \hat{q} \cdot \vec{\epsilon} + F_4 \vec{\sigma} \cdot \hat{q} \hat{q} \cdot \vec{\epsilon}$$

\vec{q} : meson momentum
 \vec{k} ($\vec{\epsilon}$): photon momentum
 (polarization)

F_i : complex functions of the scattering angle, constructed from multipole amplitudes $M_{\mu\gamma}^J$

⇒ 16 polarization observables:
 asymmetries composed of **beam**, **target** and/or **recoil** polarization measurements

⇒ **Complete Experiment**: unambiguous determination of the amplitude

8 carefully selected observables, including

[Chiang and Tabakin, PRC 55, 2054 \(1997\)](#)

- **single** and **double** polarization observables
- measurement of **beam**, **target** and **recoil** polarization

↳ easier to realize in K than in π or η photoproduction

↔ Caveat: in reality more observables needed (data uncertainties)

