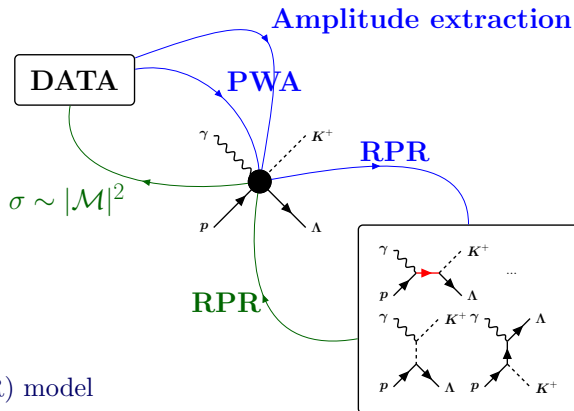


Hunting the resonances in $p(\gamma, K^+)\Lambda$: (over)complete measurements and partial-wave analyses

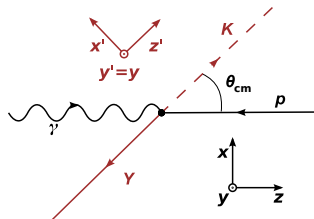
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NSTAR2015, Osaka



- 1 Introduction
- 2 Regge-plus-resonance (RPR) model
- 3 Amplitude representations
- 4 Traditional method: multipole decomposition (PWA)
- 5 Alternative (complementary) method: amplitude extraction
 - Partial amplitude extraction using real data
 - From complete to overcomplete sets
 - Amplitude comparison
- 6 Conclusions

Case study of $p(\gamma, K^+)\Lambda$ 

| | | |
|-------------------|-----------------|-------------|
| Photon: γ | 1^- | / |
| Proton: p | $\frac{1}{2}^+$ | uud |
| Kaon: K^+ | 0^- | u \bar{s} |
| Lambda: Λ | $\frac{1}{2}^+$ | uds |

Two independent kinematic variables

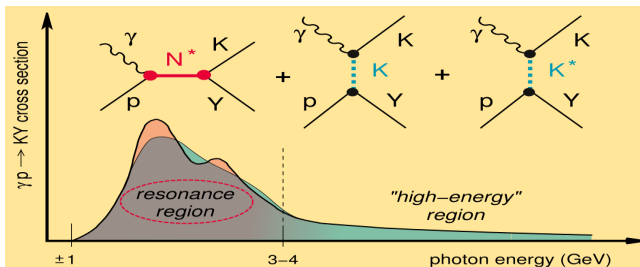
- Invariant mass W
- Kaon angle $\theta_{\text{c.m.}}$

Dynamics

- 2 spin-1/2 particles and a real photon
→ 8 combinations
- Parity conservation
- **4 independent COMPLEX**

REACTION AMPLITUDES

$$\mathcal{M}_{\lambda_p, \lambda_\Lambda}^{\lambda_\gamma} \rightarrow \mathcal{M}_{i=1,2,3,4}$$

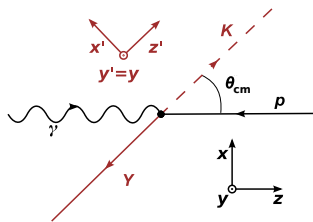


- Regge background: exchange of $K(494)$ and $K^*(892)$ Regge trajectories in t channel
- Enrich Reggeized background with N^* : $J = \frac{1}{2}, \frac{3}{2}, \frac{5}{2}$ with $M_{N^*} \leq 2$ GeV

Bayesian inference of the resonance content of $p(\gamma, K^+)\Lambda$ [PRL108 (2012) 182002]

$S_{11}(1535)$, $S_{11}(1650)$, $F_{15}(1680)$, $P_{13}(1720)$,
 $D_{13}(1875)$, $P_{13}(1900)$, $P_{11}(1900)$, and $F_{15}(2000)$

- **17 parameters**

Transversity Amplitudes (TA) $b_{i=1,\dots,4}$


$$b_1 \equiv y \langle + | J_y | + \rangle_y$$

$$b_2 \equiv y \langle - | J_y | - \rangle_y$$

$$b_3 \equiv y \langle + | J_x | - \rangle_y$$

$$b_4 \equiv y \langle - | J_x | + \rangle_y$$

 Normalized TA $a_{i=1,\dots,4}$

$$a_i = \frac{b_i}{\sqrt{|b_1|^2 + |b_2|^2 + |b_3|^2 + |b_4|^2}} = r_i e^{i\alpha_i}$$

CGLN amplitudes and multipole decomposition

$$\mathcal{M} = \langle m_{s_\Lambda} | -iF_1 \sigma \cdot \mathbf{e}_p \gamma - F_2 (\sigma \cdot \mathbf{e}_p) [\sigma \cdot (\mathbf{e}_k \times \mathbf{e}_p \gamma)] - iF_3 (\sigma \cdot \mathbf{e}_k) (\mathbf{e}_p \cdot \mathbf{e}_p \gamma) - iF_4 (\sigma \cdot \mathbf{e}_p) (\mathbf{e}_p \cdot \mathbf{e}_p \gamma) | m_{s_p} \rangle$$

$$F_1 = \sum_l P'_{l+1}(\cos \theta_{c.m.}) [E_{l+} + lM_{l+}] + P'_{l-1}(\cos \theta_{c.m.}) [E_{l-} + (l+1)M_{l-}]$$

$$F_2 = \sum_l P'_l(\cos \theta_{c.m.}) [(l+1)M_{l+} + lM_{l-}]$$

$$F_3 = \sum_l P''_{l+1}(\cos \theta_{c.m.}) [E_{l+} - M_{l+}] + P''_{l-1}(\cos \theta_{c.m.}) [E_{l-} + M_{l-}]$$

$$F_4 = \sum_l P''_l(\cos \theta_{c.m.}) [-E_{l-} - M_{l-} - E_{l+} + M_{l+}]$$

Multipoles (RPR-2011): **BACKGROUND DOMINANCE**

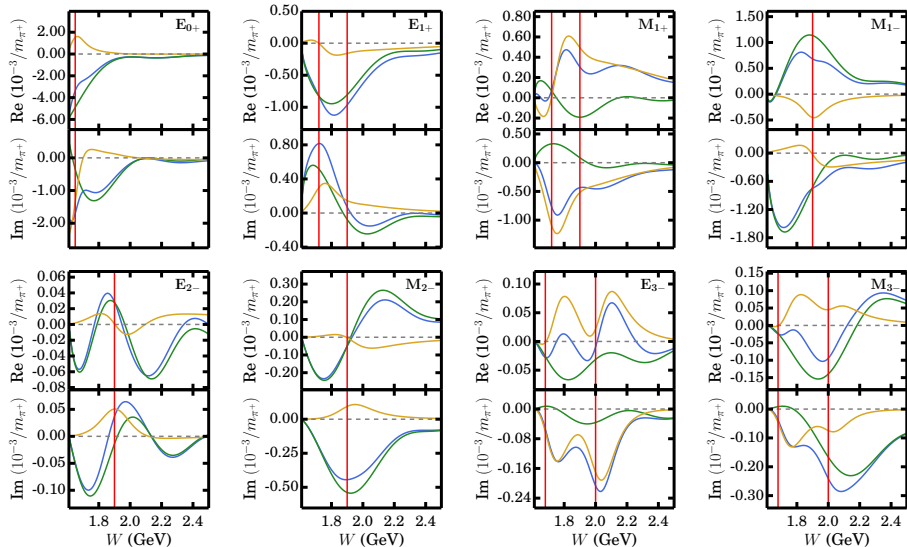


Figure : RPR-2011, RPR-2011 \Resonances and RPR-2011 \Regge.

Polarization observables in pseudoscalar-meson photoproduction

| | $(\mathcal{B}_1, \mathcal{T}_1, \mathcal{R}_1)$ | $(\mathcal{B}_2, \mathcal{T}_2, \mathcal{R}_2)$ | Transversity expression |
|----------|---|---|---|
| Σ | $(y, 0, 0)$ | $(x, 0, 0)$ | $r_1^2 + r_2^2 - r_3^2 - r_4^2$ |
| T | $(0, +y, 0)$ | $(0, -y, 0)$ | $r_1^2 - r_2^2 - r_3^2 + r_4^2$ |
| P | $(0, 0, +y)$ | $(0, 0, -y)$ | $r_1^2 - r_2^2 + r_3^2 - r_4^2$ |
| C_x | $(+, 0, +x)$ | $(+, 0, -x)$ | $-2 \operatorname{Im}(a_1 a_4^* + a_2 a_3^*)$ |
| C_z | $(+, 0, +z)$ | $(+, 0, -z)$ | $+2 \operatorname{Re}(a_1 a_4^* - a_2 a_3^*)$ |
| O_x | $(+\frac{\pi}{4}, 0, +x)$ | $(+\frac{\pi}{4}, 0, -x)$ | $+2 \operatorname{Re}(a_1 a_4^* + a_2 a_3^*)$ |
| O_z | $(+\frac{\pi}{4}, 0, +z)$ | $(+\frac{\pi}{4}, 0, -z)$ | $+2 \operatorname{Im}(a_1 a_4^* - a_2 a_3^*)$ |
| E | $(+, -z, 0)$ | $(+, +z, 0)$ | $+2 \operatorname{Re}(a_1 a_3^* - a_2 a_4^*)$ |
| F | $(+, +x, 0)$ | $(+, -x, 0)$ | $-2 \operatorname{Im}(a_1 a_3^* + a_2 a_4^*)$ |
| G | $(+\frac{\pi}{4}, +z, 0)$ | $(+\frac{\pi}{4}, -z, 0)$ | $-2 \operatorname{Im}(a_1 a_3^* - a_2 a_4^*)$ |
| H | $(+\frac{\pi}{4}, +x, 0)$ | $(+\frac{\pi}{4}, -x, 0)$ | $+2 \operatorname{Re}(a_1 a_3^* + a_2 a_4^*)$ |
| T_x | $(0, +x, +x)$ | $(0, +x, -x)$ | $+2 \operatorname{Re}(a_1 a_2^* + a_3 a_4^*)$ |
| T_z | $(0, +x, +z)$ | $(0, +x, -z)$ | $+2 \operatorname{Im}(a_1 a_2^* + a_3 a_4^*)$ |
| L_x | $(0, +z, +x)$ | $(0, +z, -x)$ | $-2 \operatorname{Im}(a_1 a_2^* - a_3 a_4^*)$ |
| L_z | $(0, +z, +z)$ | $(0, +z, -z)$ | $+2 \operatorname{Re}(a_1 a_2^* - a_3 a_4^*)$ |

■ $\frac{d\sigma}{d\Omega}(\mathcal{B}, \mathcal{T}, \mathcal{R})$: cross section for given beam (\mathcal{B}), target (\mathcal{T}), recoil (\mathcal{R}) polarization

■ Asymmetries

$$\mathcal{A} = \frac{\frac{d\sigma}{d\Omega}(\mathcal{B}_1, \mathcal{T}_1, \mathcal{R}_1) - \frac{d\sigma}{d\Omega}(\mathcal{B}_2, \mathcal{T}_2, \mathcal{R}_2)}{\frac{d\sigma}{d\Omega}(\mathcal{B}_1, \mathcal{T}_1, \mathcal{R}_1) + \frac{d\sigma}{d\Omega}(\mathcal{B}_2, \mathcal{T}_2, \mathcal{R}_2)}$$

■ $\frac{d\sigma}{d\Omega}(0,0,0) = \frac{\rho}{4} \sum_{i=1}^4 |b_i|^2$

SINGLE asymmetries: MODULI

DOUBLE asymmetries: PHASES

4 complex amplitudes, or 8 real variables

- There is one arbitrary global phase

$$\delta_i^{\alpha_4} = \alpha_i - \alpha_4.$$

- Take $\alpha_4 = 0$ and use normalized transversity amplitudes

$$1 = |a_1|^2 + |a_2|^2 + |a_3|^2 + |a_4|^2$$

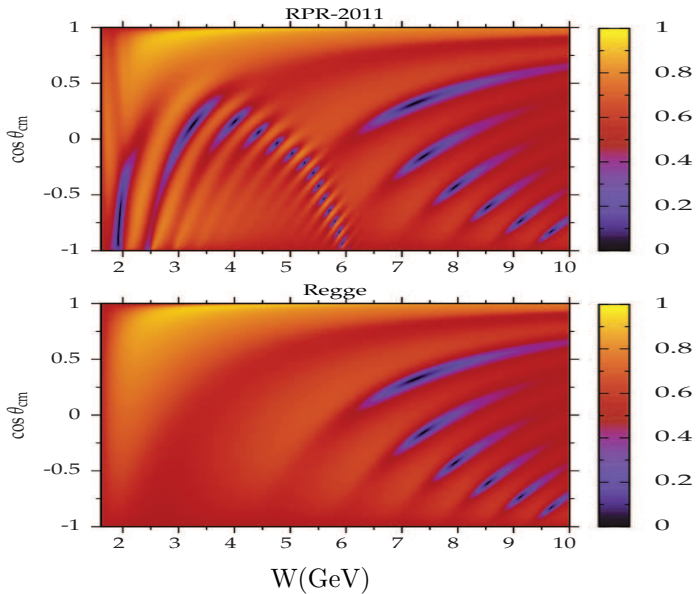
We need **6** real variables and an independent scaling factor

Definition **COMPLETE SET**

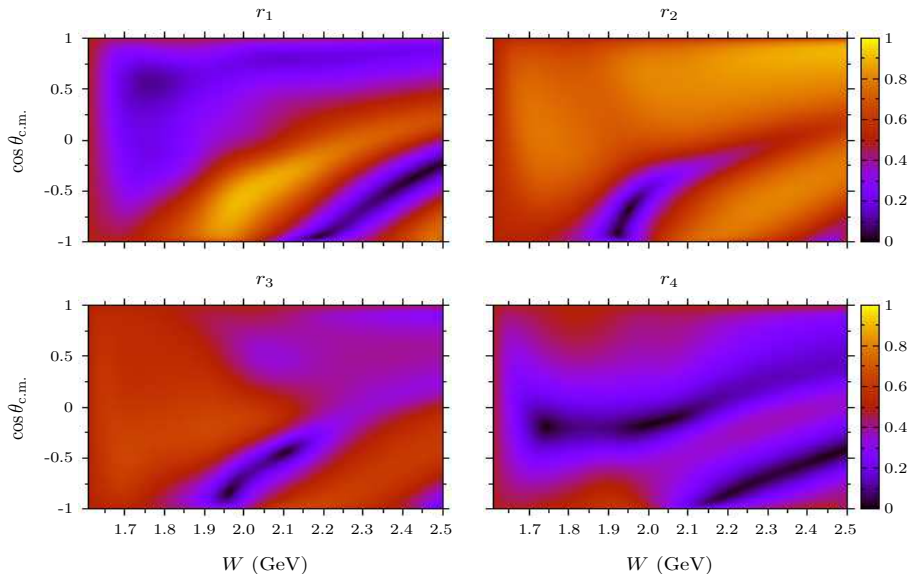
A complete set is a minimum set of observables from which one can determine the underlying reaction amplitudes **unambiguously**.

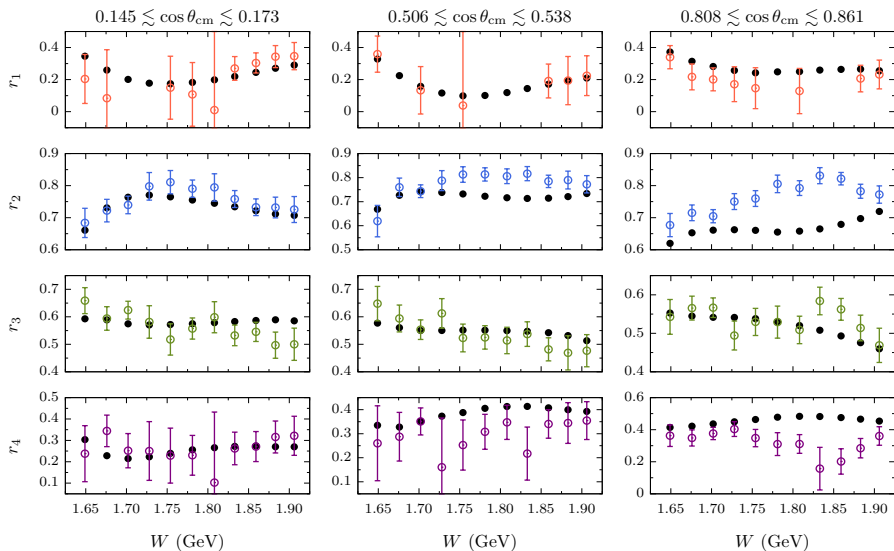
[Chiang & Tabakin PRC55 (1997) 2054]: 8 observables

Role of resonances for the NTA moduli (r_2)

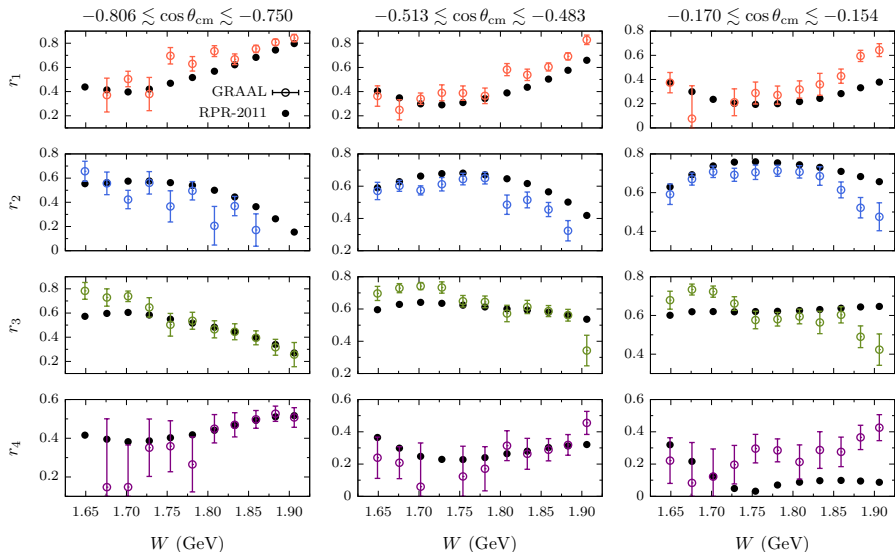


RPR-2011 predictions for $(W, \cos \theta_{c.m.})$ dependence of NTA moduli for $p(\gamma, K^+)\Lambda$



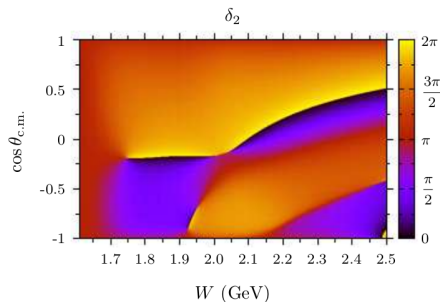
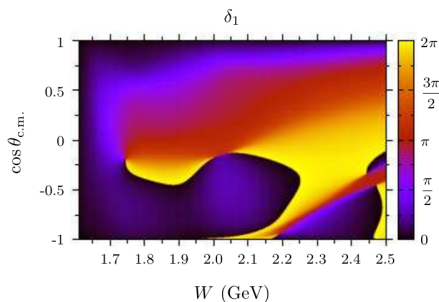


$$r_2 : b_2 = {}_y \langle -|J_y| - \rangle_y$$

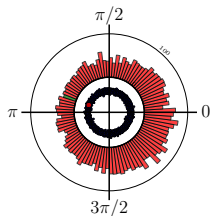


$$r_2 : b_2 = {}_y \langle -|J_y| - \rangle_y$$

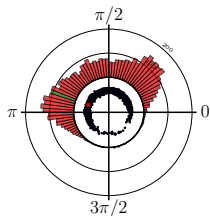
RPR-2011 predictions for $(W, \cos \theta_{c.m.})$ dependence of NTA relative phases
 $\delta_i = \alpha_i - \alpha_4$ for $p(\gamma, K^+)\Lambda$



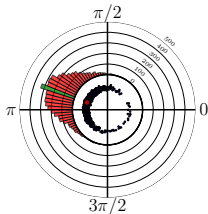
- at forward angles the background dominates and the W -dependence of δ_i is mild
- at backward angles large N^* contributions and the W -dependence of δ_i is wild



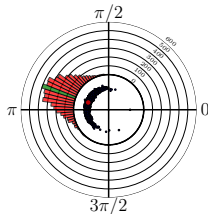
(a) **“Complete”**



(b) **Discrete amb.**



(c) **Unimodal**



(d) **Improved**

$$\mathcal{M}_a(s, t) \equiv \begin{pmatrix} a_1(s, t) \\ a_2(s, t) \\ a_3(s, t) \\ a_4(s, t) \end{pmatrix}$$

$$\mathcal{M}_a^\dagger \mathcal{M}_a = 1$$

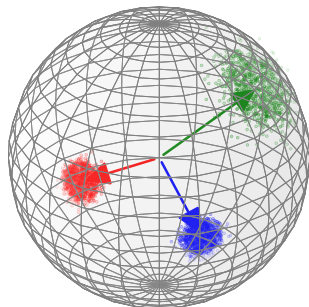
Q: What is the distance between \mathcal{M}_1 and \mathcal{M}_2 ?

A: $\mathcal{D}[\mathcal{M}_1, \mathcal{M}_2] = \arccos \operatorname{Re} \mathcal{M}_1^\dagger \mathcal{M}_2$

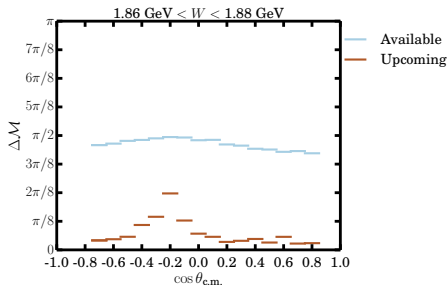
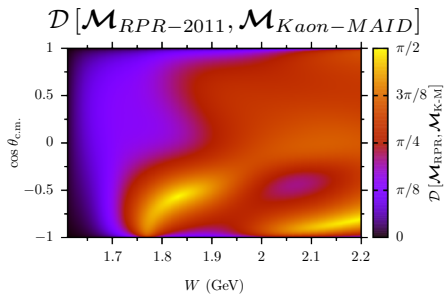
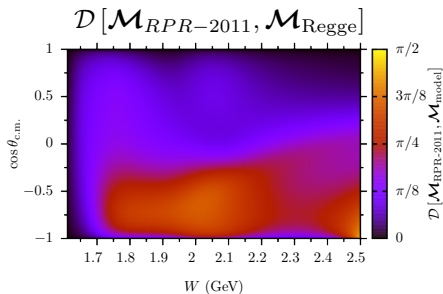
Both $\mathcal{M}_{i=1,2}$ have an unknown α_4 .

Q: How to calculate $\mathcal{D}[\mathcal{M}_1, \mathcal{M}_2]$ independent of choice α_4 ?

A: $\alpha_4 = \operatorname{argmin}_{\alpha_4} (\mathcal{D}[\mathcal{M}_1(\alpha_4), \mathcal{M}_2(\alpha_4' = 0)])$



Model comparison in amplitude space

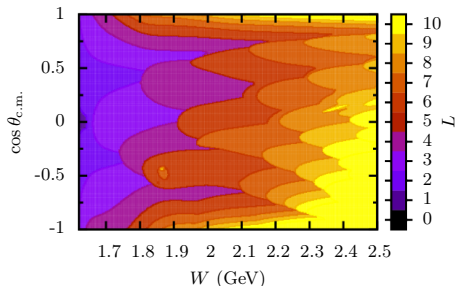


Resolution of the data

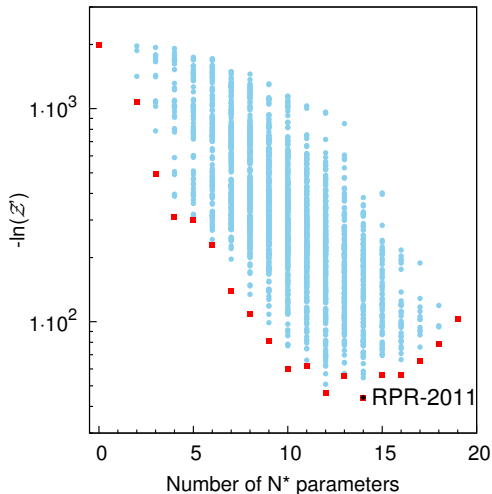
- Obtaining resonance information in background-dominated reactions requires background-subtraction schemes, such as RPR-2011.
- Hierarchy in the quality/quantity of the data!
- Quadratic equations connect $\{\Sigma, P, T\}$ to the **moduli** $\{r_1, r_2, r_3, r_4\}$ of the normalized transversity amplitudes
 - 1 Analysis of $\gamma p \rightarrow K^+ \Lambda$ with $\{\Sigma, T, P\}$ from GRAAL ($1.65 \lesssim W \lesssim 1.91$ GeV) allowed to extract $\{r_1, r_2, r_3, r_4\}$ in $\approx 95\%$ of considered ($W, \cos \theta_{c.m.}$)
 - 2 RPR-2011 is in reasonable agreement with the extracted r_i
- Extracting the NTA independent **phases** $\{\delta_1, \delta_2, \delta_3\}$ is far more challenging (connected to asymmetries by means of non-linear equations)
- **Mathematical Completeness does not imply Practical Completeness!!**
- **Overcomplete sets** provide a solution!

- J. Nys, T. Vrancx and J. Ryckebusch
Amplitude extraction in pseudoscalar-meson photoproduction: towards a situation of complete information
J. Phys. G **42** (2015) 3, 034016
- D. G. Ireland
Information Content of Polarization Measurements
Phys. Rev. C **82** (2010) 025204
- T. Vrancx, J. Ryckebusch, T. Van Cuyck T, P. Vancraeyveld
Incompleteness of complete pseudoscalar-meson photoproduction
Phys. Rev. C **87** (2013) 055205.
- L. De Cruz, J. Ryckebusch, T. Vrancx, P. Vancraeyveld
A Bayesian analysis of kaon photoproduction with the Regge-plus-resonance model
Phys. Rev. C **86** (2012) 015212
- L. De Cruz, T. Vrancx, P. Vancraeyveld, J. Ryckebusch
Bayesian inference of the resonance content of $p(\gamma, K^+)\Lambda$
Phys. Rev. Lett. **108** (2012) 182002

Backup slides



Bayesian evidence map for the 2^{11} model variants



RPR-2011 (PDG-2010)

- $S_{11}(1535)$ ****
- $S_{11}(1650)$ ****
- $D_{15}(1675)$ ****
- $F_{15}(1680)$ ****
- $D_{13}(1700)$ ***
- $P_{11}(1710)$ **
- $P_{13}(1720)$ ****
- $D_{13}(1875)$ *m*
- $P_{13}(1900)$ **
- $P_{11}(1900)$ *m*
- $F_{15}(2000)$ **

PRL108 (2012) 182002

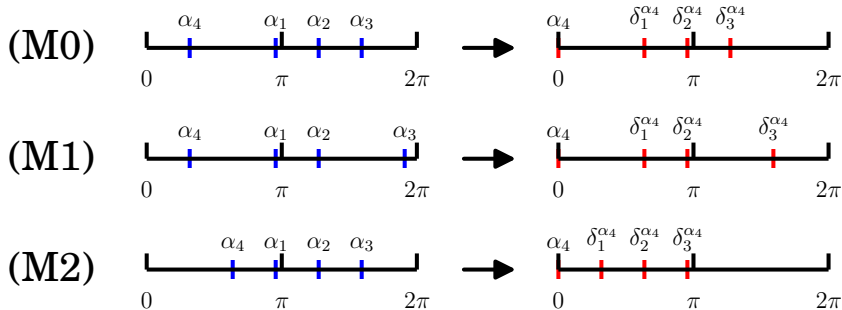


Figure : Example of a situation where a global phase transformation, followed by $\alpha_4 = 0$ can give a distorted picture of the degree of compatibility of two models.

Observables for particular experimental setups

| Configuration | | | $\frac{d\sigma}{d\Omega}(\text{conf.}) / \frac{d\sigma}{d\Omega}(0,0,0)$ |
|---------------|---------------|---------------|---|
| \mathcal{B} | \mathcal{T} | \mathcal{R} | |
| 0 | 0 | N | 1 |
| 0 | 0 | Y | $1 + PP_y^R$ |
| 0 | L | N | 1 |
| 0 | L | Y | $1 + PP_y^R + P_z^T(L_{x'}P_{x'}^R + L_{z'}P_{z'}^R)$ |
| 0 | T | N | $1 + TP_y^T$ |
| 0 | T | Y | $1 + \Sigma P_y^T P_y^R + TP_y^T + PP_y^R + P_x^T(T_{x'}P_{x'}^R + T_{z'}P_{z'}^R)$ |
| c | 0 | N | 1 |
| c | 0 | Y | $1 + PP_y^R + P_c^\gamma(C_{x'}P_{x'}^R + C_{z'}P_{z'}^R)$ |
| c | L | N | $1 - EP_c^\gamma P_z^T$ |
| c | L | Y | $1 + PP_y^R - EP_c^\gamma P_z^T - HP_c^\gamma P_z^T P_y^R + P_c^\gamma(C_{x'}P_{x'}^R + C_{z'}P_{z'}^R) + P_z^T(L_{x'}P_{x'}^R + L_{z'}P_{z'}^R)$ |
| c | T | N | $1 + TP_y^T + FP_c^\gamma P_x^T$ |
| c | T | Y | $1 + \Sigma P_y^T P_y^R + TP_y^T + PP_y^R + GP_c^\gamma P_y^R P_x^T + FP_c^\gamma P_x^T + P_c^\gamma P_y^T(C_{x'}P_{x'}^R + C_{z'}P_{z'}^R) + P_c^\gamma P_y^T(O_{x'}P_{x'}^R + O_{z'}P_{z'}^R) + P_x^T(T_{x'}P_{x'}^R + T_{z'}P_{z'}^R)$ |
| l | 0 | N | $1 - \Sigma P_l^\gamma \cos(2\phi_\gamma)$ |
| l | 0 | Y | $1 - \Sigma P_l^\gamma \cos(2\phi_\gamma) - TP_l^\gamma P_y^R \cos(2\phi_\gamma) + PP_y^R + P_l^\gamma \sin(2\phi_\gamma)(O_{x'}P_{x'}^R + O_{z'}P_{z'}^R)$ |
| l | L | N | $1 - \Sigma P_l^\gamma \cos(2\phi_\gamma) + GP_l^\gamma P_z^T \sin(2\phi_\gamma)$ |
| l | L | Y | $1 + PP_y^R - P_l^\gamma \cos(2\phi_\gamma) \left(TP_y^R + \Sigma + P_x^T(T_{x'}P_{x'}^R - T_{z'}P_{z'}^R) \right) + P_l^\gamma \sin(2\phi_\gamma) \left(GP_z^T + FP_y^R P_z^T + O_{x'}P_{x'}^R + O_{z'}P_{z'}^R \right) + P_z^T(L_{x'}P_{x'}^R + L_{z'}P_{z'}^R)$ |
| l | T | N | $1 + TP_y^T - P_l^\gamma \cos(2\phi_\gamma)(PP_y^T + \Sigma) + HP_l^\gamma P_x^T \sin(2\phi_\gamma)$ |
| l | T | Y | $1 - P_l^\gamma P_y^T P_y^R \cos(2\phi_\gamma) + \Sigma P_y^T P_y^R + TP_y^T + PP_y^R + P_x^T(T_{x'}P_{x'}^R + T_{z'}P_{z'}^R) - P_l^\gamma \cos(2\phi_\gamma) \left(-P_x^T(L_{x'}P_{x'}^R - L_{z'}P_{z'}^R) + PP_y^T + \Sigma + TP_y^R \right) + P_l^\gamma \sin(2\phi_\gamma) \left((O_{x'}P_{x'}^R + O_{z'}P_{z'}^R) + HP_x^T + EP_y^R P_x^T - P_y^T(C_{x'}P_{x'}^R - C_{z'}P_{z'}^R) \right)$ |

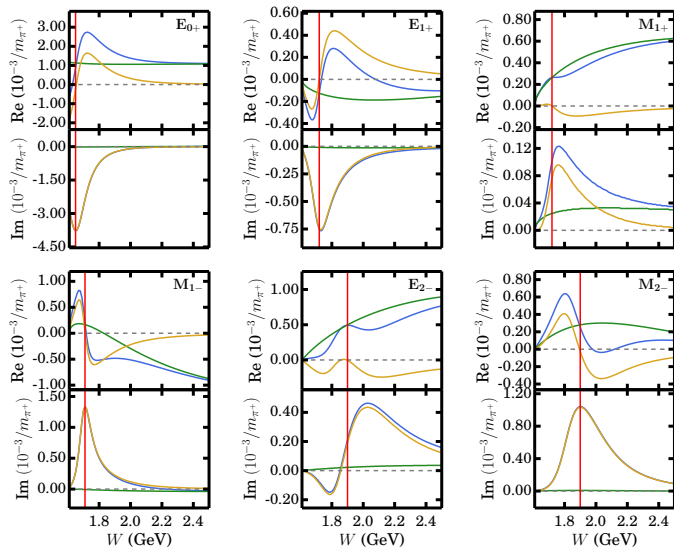
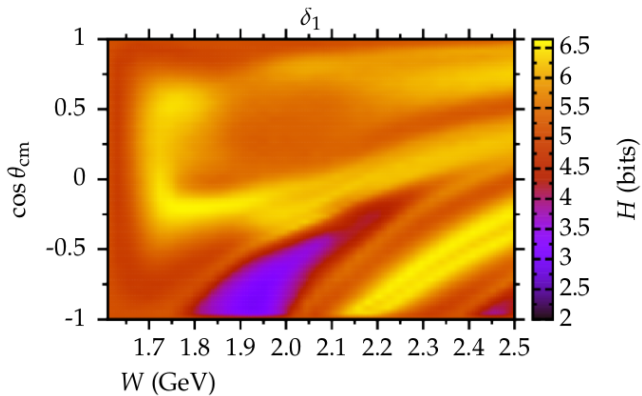
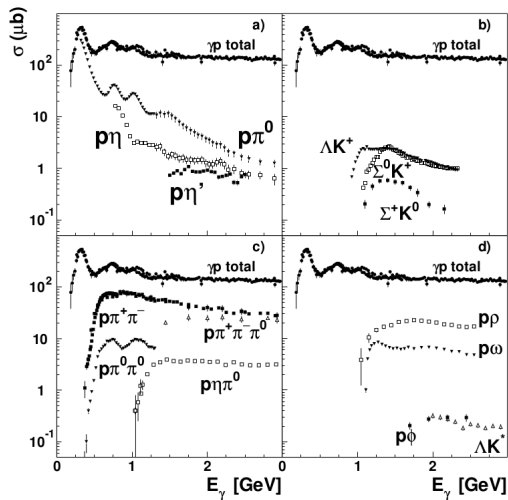


Figure : Kaon-MAID, Kaon-MAID \Resonances and Kaon-MAID \Bg.

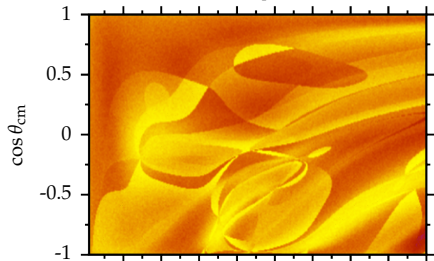
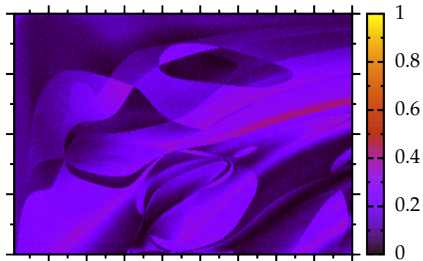
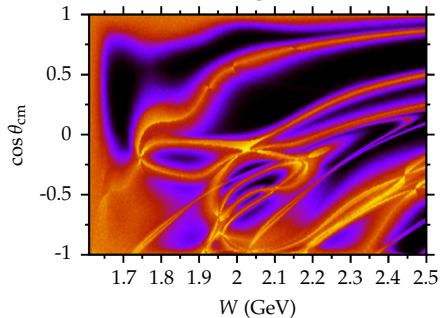
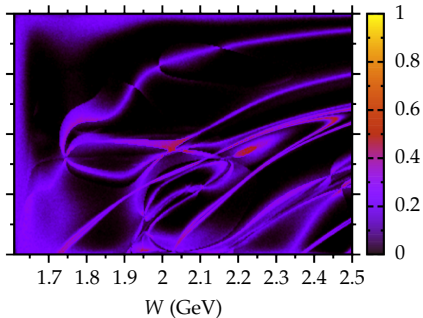


In the following, we study the effect of additional observables on the precision of the extracted amplitudes.

| Set number | Observables |
|------------|--|
| 1 | $\{C_x, O_x, E, F\}$ |
| 2 | $\{C_x, O_x, E, F, C_z\}$ |
| 3 | $\{C_x, O_x, E, F, C_z, O_z\}$ |
| 4 | $\{C_x, O_x, E, F, C_z, O_z, G\}$ |
| 5 | $\{C_x, O_x, E, F, C_z, O_z, G, H\}$ |
| 6 | $\{C_x, O_x, E, F, C_z, O_z, G, H, T_x\}$ |
| 7 | $\{C_x, O_x, E, F, C_z, O_z, G, H, T_x, L_x\}$ |
| 8 | $\{C_x, O_x, E, F, C_z, O_z, G, H, T_x, L_x, T_z\}$ |
| 9 | $\{C_x, O_x, E, F, C_z, O_z, G, H, T_x, L_x, T_z, L_z\}$ |



| | |
|----------|---|
| Σ | $(R_1^2 + R_2^2 - R_3^2 - R_4^2)/\mathcal{N}$ |
| T | $(R_1^2 - R_2^2 - R_3^2 + R_4^2)/\mathcal{N}$ |
| P | $(R_1^2 - R_2^2 + R_3^2 - R_4^2)/\mathcal{N}$ |
| C_x | $-2(R_1 R_4 \sin \delta_1 + R_2 R_3 \sin(\delta_2 - \delta_3))/\mathcal{N}$ |
| C_z | $+2(R_1 R_4 \cos \delta_1 - R_2 R_3 \cos(\delta_2 - \delta_3))/\mathcal{N}$ |
| O_x | $+2(R_1 R_4 \cos \delta_1 + R_2 R_3 \cos(\delta_2 - \delta_3))/\mathcal{N}$ |
| O_z | $+2(R_1 R_4 \sin \delta_1 - R_2 R_3 \sin(\delta_2 - \delta_3))/\mathcal{N}$ |
| E | $+2(R_1 R_3 \cos(\delta_1 - \delta_3) - R_2 R_4 \cos \delta_2)/\mathcal{N}$ |
| F | $-2(R_1 R_3 \sin(\delta_1 - \delta_3) + R_2 R_4 \sin \delta_2)/\mathcal{N}$ |
| G | $-2(R_1 R_3 \sin(\delta_1 - \delta_3) - R_2 R_4 \sin \delta_2)/\mathcal{N}$ |
| H | $+2(R_1 R_3 \cos(\delta_1 - \delta_3) + R_2 R_4 \cos \delta_2)/\mathcal{N}$ |
| T_x | $+2(R_1 R_2 \cos(\delta_1 - \delta_2) + R_3 R_4 \cos \delta_3)/\mathcal{N}$ |
| T_z | $+2(R_1 R_2 \sin(\delta_1 - \delta_2) + R_3 R_4 \sin \delta_3)/\mathcal{N}$ |
| L_x | $-2(R_1 R_2 \sin(\delta_1 - \delta_2) - R_3 R_4 \sin \delta_3)/\mathcal{N}$ |
| L_z | $+2(R_1 R_2 \cos(\delta_1 - \delta_2) - R_3 R_4 \cos \delta_3)/\mathcal{N}$ |

$\eta_{\text{total}} (\sigma_{\text{exp}} = 0.1)$  $\eta_{\text{incorrect}} (\sigma_{\text{exp}} = 0.1)$  $\eta_{\text{total}} (\sigma_{\text{exp}} = 0.01)$  $\eta_{\text{incorrect}} (\sigma_{\text{exp}} = 0.01)$ 

| | | Kinematics nr. | | | | |
|--------|------|----------------|--------------|--------------|-------------|-------------|
| | | 1 | 2 | 3 | 4 | |
| Single | S | 0.21 | 0.43 | 0.21 | 0.47 | |
| | T | -0.89 | -0.57 | -0.52 | 0 | |
| | P | -0.15 | -0.54 | 0.25 | 0.03 | |
| Double | BR | C_x | -0.28 | -0.51 | -0.32 | -0.16 |
| | | C_z | 0.84 | 0.28 | 0.62 | 0.85 |
| | | O_x | -0.92 | -0.64 | -0.74 | 0.02 |
| | | O_z | -0.33 | -0.31 | -0.37 | -0.19 |
| | BT | E | 0.03 | 0.02 | 0.44 | 0.22 |
| | | F | -0.09 | 0.56 | -0.27 | 0.83 |
| | | G | -0.30 | -0.55 | -0.37 | 0.08 |
| | | H | 0.29 | 0.41 | 0.58 | -0.17 |
| | TR | T_x | -0.24 | -0.59 | -0.39 | -0.30 |
| | | T_z | -0.24 | -0.16 | -0.49 | 0.93 |
| | | L_x | 0.33 | 0.43 | 0.52 | -0.40 |
| | | L_z | 0.02 | -0.09 | -0.21 | -0.34 |

