

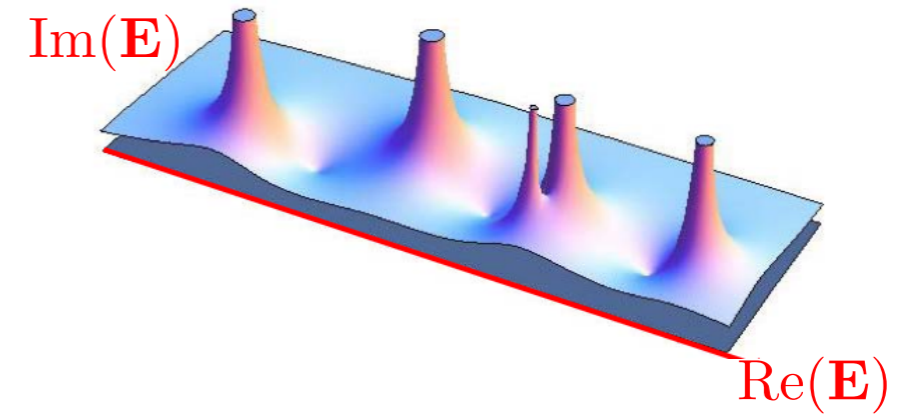
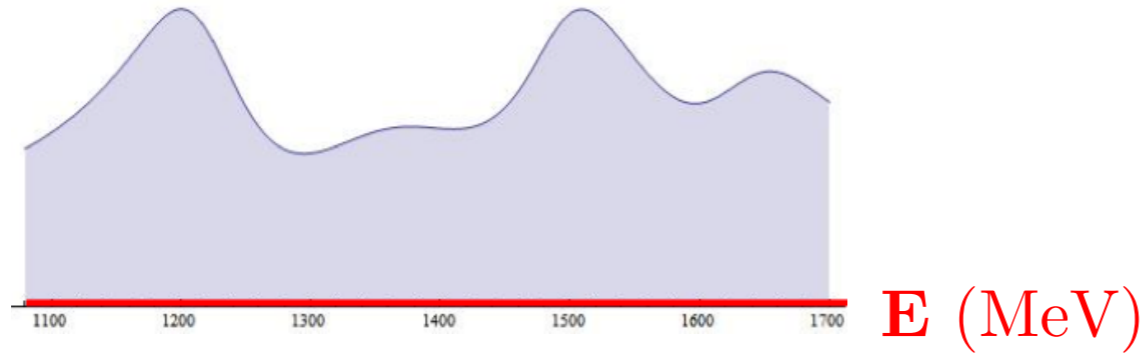
Hadron Spectroscopy and Interactions from Lattice QCD

Robert Edwards



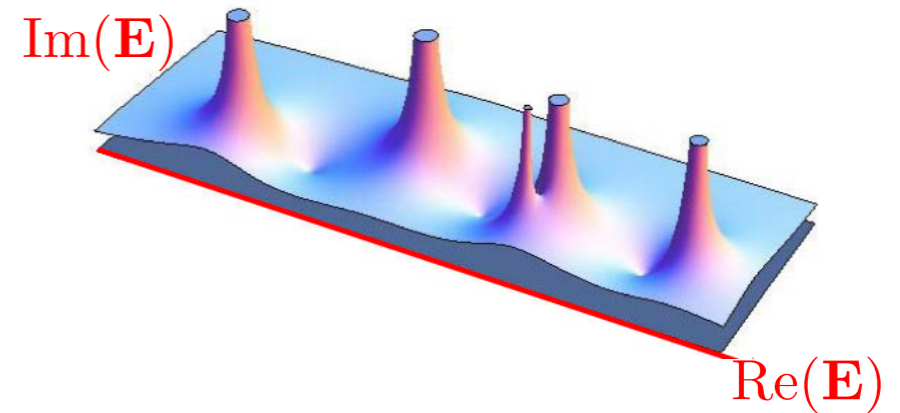
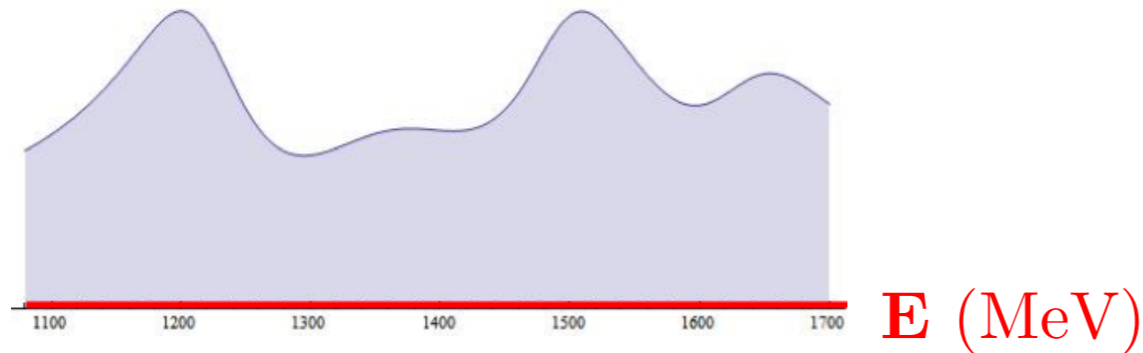
Resonances

- Most hadrons are resonances
 - E.g., $\pi N \pi N$



Resonances

- Most hadrons are resonances
 - E.g., $\pi N \pi N$



- Formally defined as a pole in a partial-wave scattering amplitude

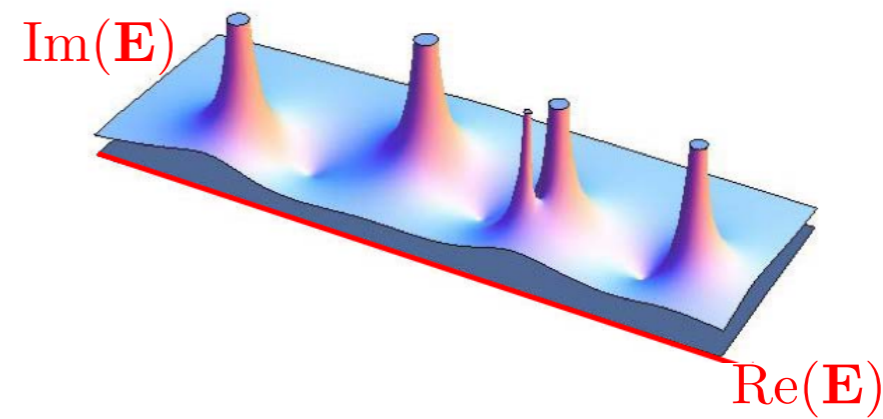
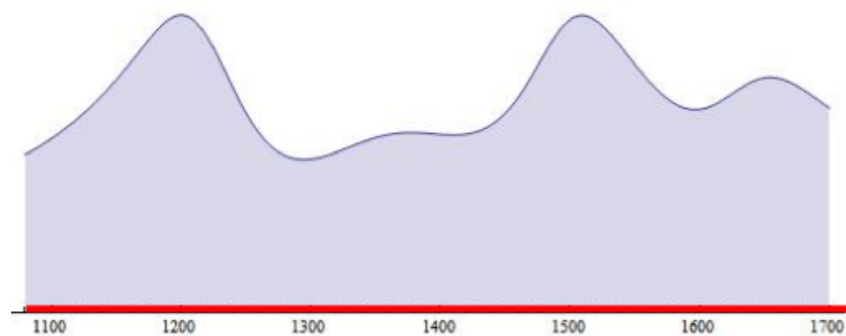
$$t_l(s) \sim \frac{R}{s_0 - s} + \dots$$



- Different channels should have same pole location
- Pole structure gives decay information

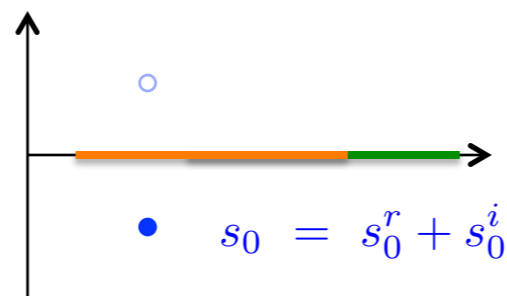
Resonances

- Most hadrons are resonances
 - E.g., $\pi N \pi N$



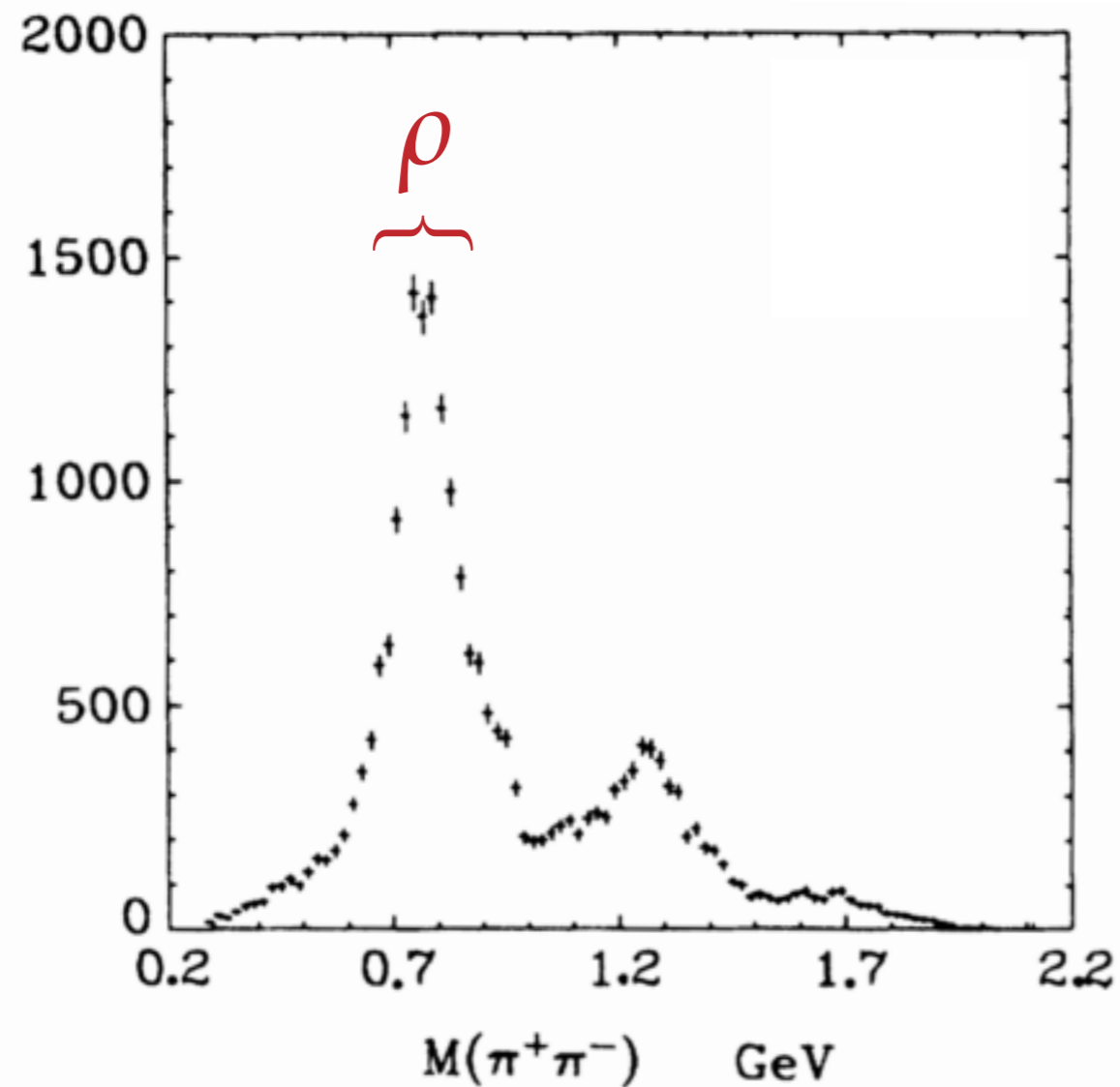
- Formally defined as a pole in a partial-wave scattering amplitude

$$t_l(s) \sim \frac{R}{s_0 - s} + \dots$$



- Different channels should have same pole location
 - Pole structure gives decay information
- Can we predict hadron properties from first principles?

Isospin=1 $\pi\pi$ P-wave



expand angular dependence
in *partial waves*

PARTIAL WAVE AMPLITUDE

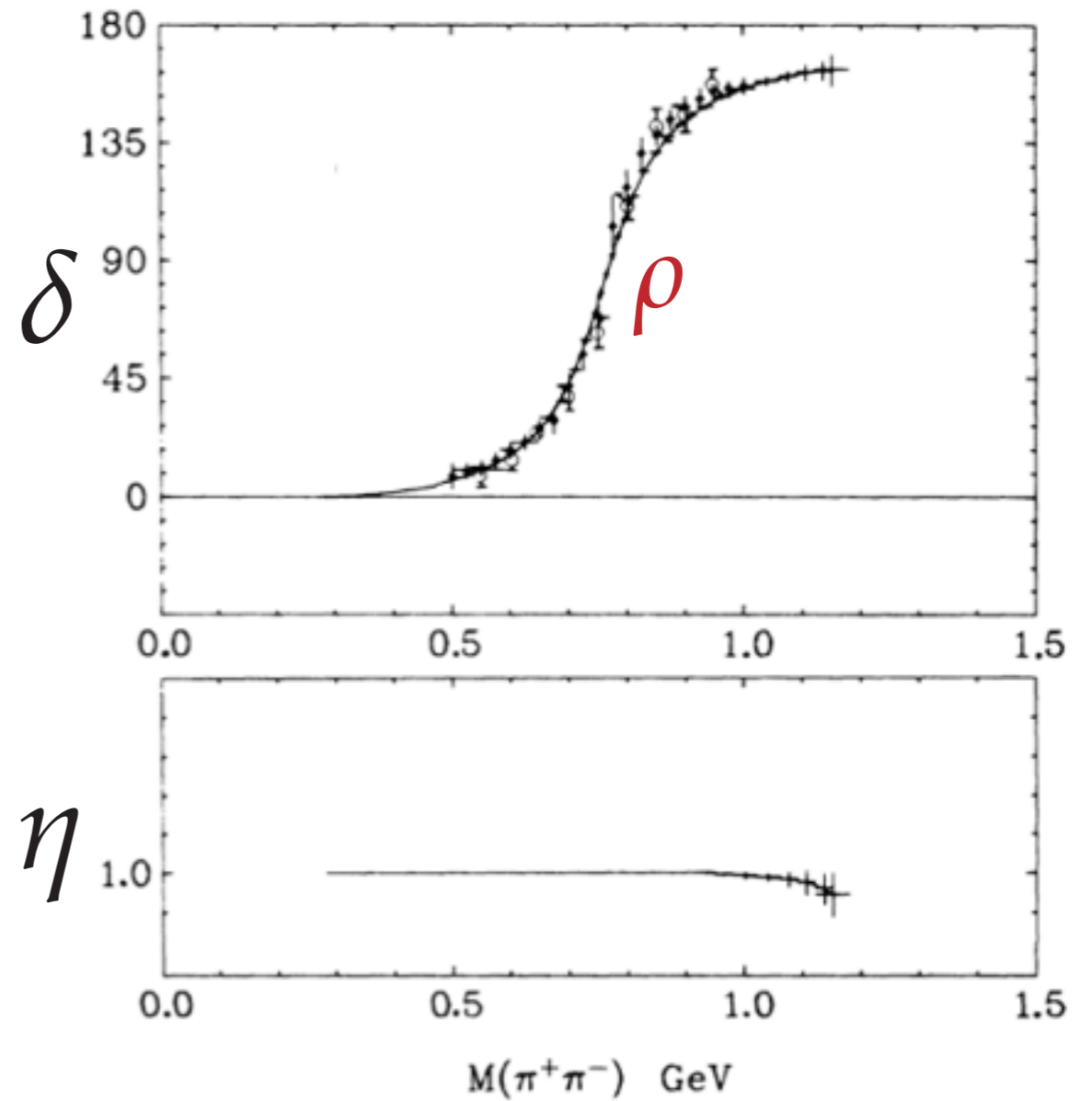
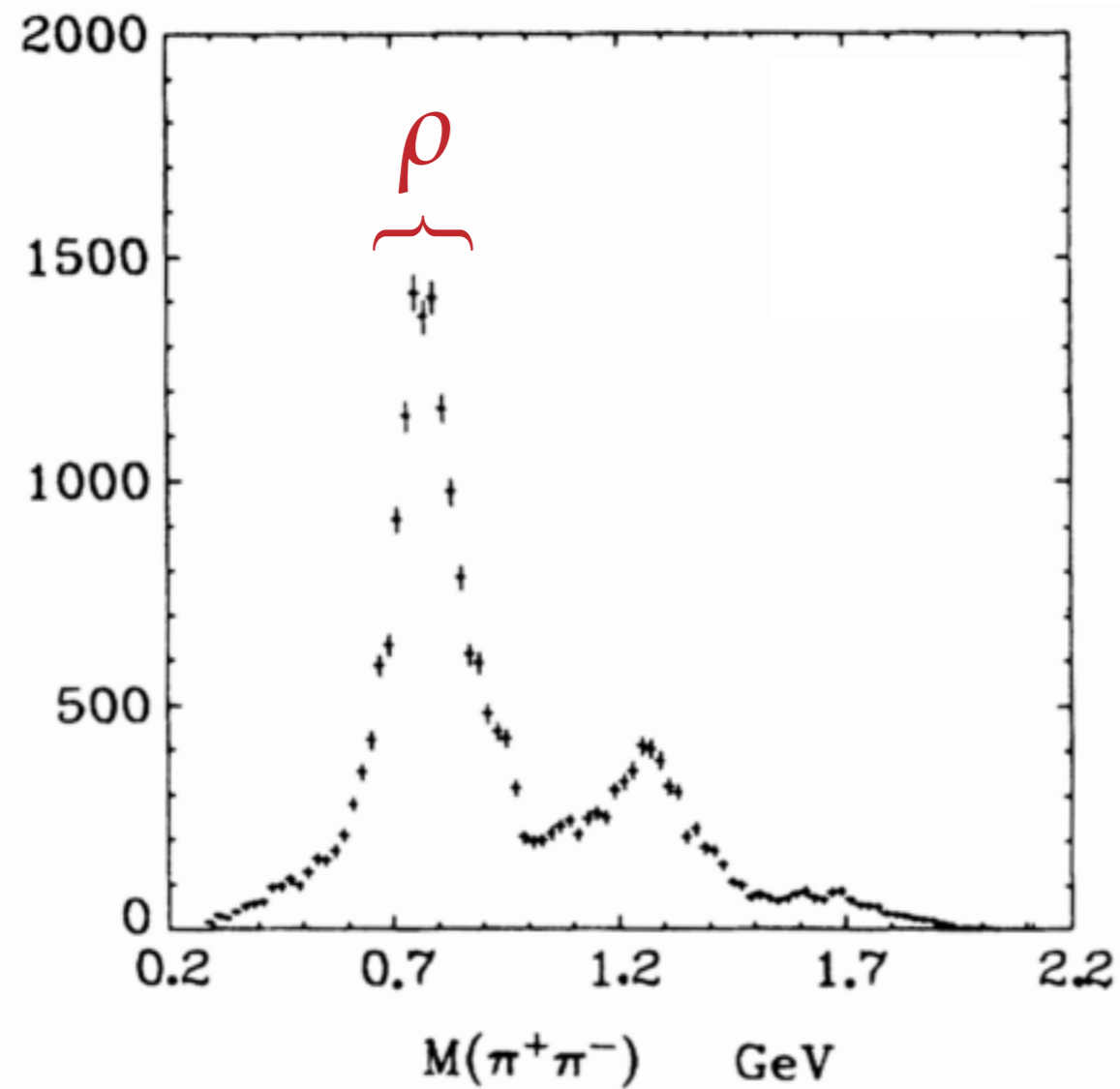
$$f_\ell = \frac{1}{2i} \left(\eta_\ell e^{2i\delta_\ell} - 1 \right)$$

$\eta = 1$ elastic

$\eta \leq 1$ inelastic

Isospin=1 $\pi\pi$ P-wave

RESONANT PHASE SHIFT



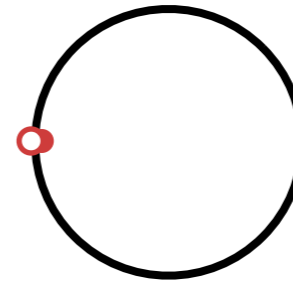
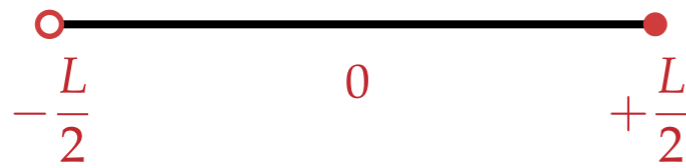
Finite-volume

- Where's the meson-meson continuum ?
 - there isn't one !

Finite-volume

- Where's the meson-meson continuum ?
 - there isn't one !
 - in a finite-volume the spectrum is discrete

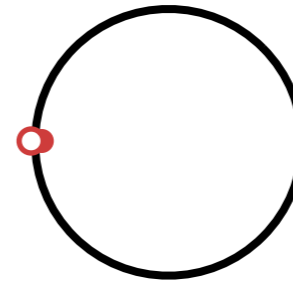
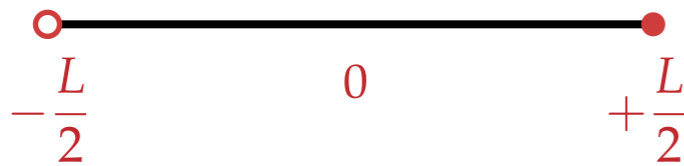
one-dim :



Finite-volume

- Where's the meson-meson continuum ?
 - there isn't one !
 - in a finite-volume the spectrum is discrete

one-dim :



e.g. a free particle

$$\psi(x) \sim e^{ipx}$$

» periodic boundary condition

$$\psi(x) = \psi(x + L)$$

$$e^{ipx} = e^{ip(x+L)}$$

$$e^{ipL} = 1$$

$$p = \frac{2\pi}{L}n$$

discrete
energy
spectrum

Interacting particles in a finite-volume

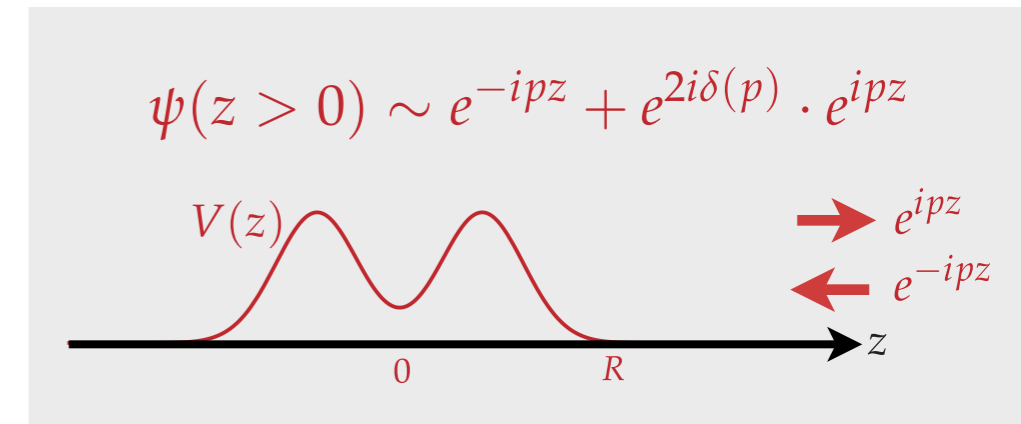
- Two identical bosons **interacting** through a finite-range potential

$$\psi(z) \sim \cos [p|z| + \delta(p)] \quad \text{outside the range of the potential, } |z| > R$$

Interacting particles in a finite-volume

- Two identical bosons **interacting** through a finite-range potential

$$\psi(z) \sim \cos [p|z| + \delta(p)] \quad \text{outside the range of the potential, } |z| > R$$



Interacting particles in a finite-volume

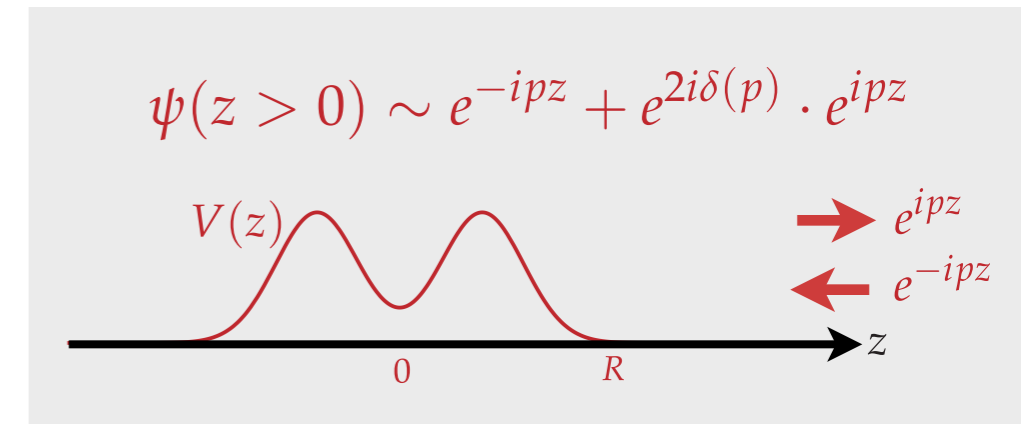
- Two identical bosons **interacting** through a finite-range potential

$$\psi(z) \sim \cos [p|z| + \delta(p)] \quad \text{outside the range of the potential, } |z| > R$$

- periodic boundary conditions

$$\left. \begin{aligned} \psi\left(-\frac{L}{2}\right) &= \psi\left(\frac{L}{2}\right) \\ \frac{d\psi}{dz}\left(-\frac{L}{2}\right) &= \frac{d\psi}{dz}\left(\frac{L}{2}\right) \end{aligned} \right\} \implies 0 = \sin\left[\frac{pL}{2} + \delta(p)\right]$$

$$\frac{pL}{2} + \delta(p) = n\pi$$



Interacting particles in a finite-volume

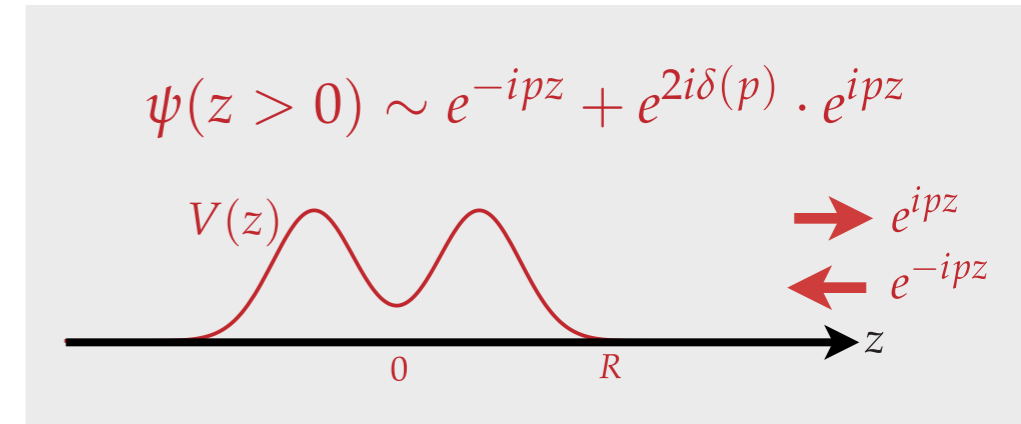
- Two identical bosons **interacting** through a finite-range potential

$$\psi(z) \sim \cos [p|z| + \delta(p)] \quad \text{outside the range of the potential, } |z| > R$$

- periodic boundary conditions

$$\left. \begin{aligned} \psi\left(-\frac{L}{2}\right) &= \psi\left(\frac{L}{2}\right) \\ \frac{d\psi}{dz}\left(-\frac{L}{2}\right) &= \frac{d\psi}{dz}\left(\frac{L}{2}\right) \end{aligned} \right\} \implies 0 = \sin\left[\frac{pL}{2} + \delta(p)\right]$$

$$\frac{pL}{2} + \delta(p) = n\pi$$



$$p = \frac{2\pi}{L}n - \frac{2}{L}\delta(p)$$

**discrete
energy
spectrum**

Interacting particles in a finite-volume

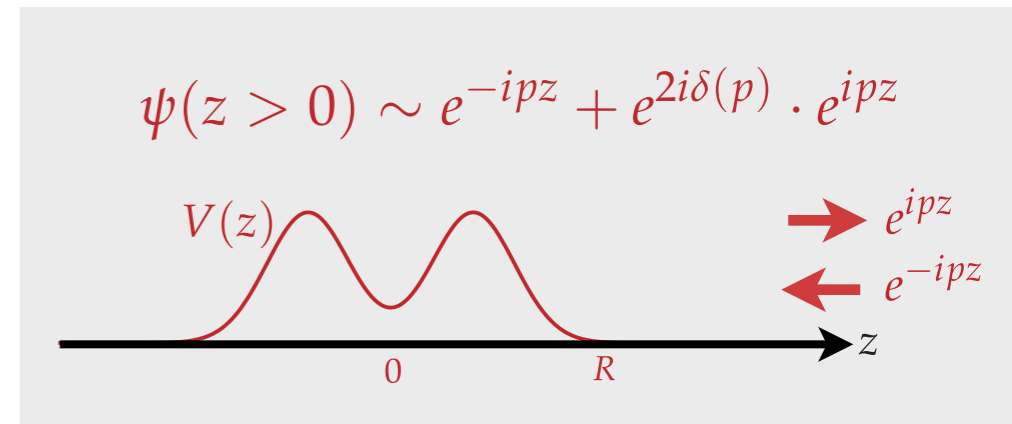
- Two identical bosons **interacting** through a finite-range potential

$$\psi(z) \sim \cos [p|z| + \delta(p)] \quad \text{outside the range of the potential, } |z| > R$$

- periodic boundary conditions

$$\left. \begin{aligned} \psi\left(-\frac{L}{2}\right) &= \psi\left(\frac{L}{2}\right) \\ \frac{d\psi}{dz}\left(-\frac{L}{2}\right) &= \frac{d\psi}{dz}\left(\frac{L}{2}\right) \end{aligned} \right\} \implies 0 = \sin\left[\frac{pL}{2} + \delta(p)\right]$$

$$\frac{pL}{2} + \delta(p) = n\pi$$



$$p = \frac{2\pi}{L}n - \frac{2}{L}\delta(p)$$

**discrete
energy
spectrum**

discrete energy spectrum is determined by the scattering amplitude

(or vice-versa)

Scattering in a finite cubic volume

- Expect a discrete spectrum in a finite periodic volume

$$\psi(x + L) = \psi(x)$$

e.g. free particle $e^{ip(x+L)} = e^{ipx}$

quantized momentum $p = \frac{2\pi}{L}n$

- For an interacting theory

$$\cot \delta_\ell(E) = \mathcal{M}_\ell(E, L)$$

LÜSCHER ...

elastic scattering
phase-shift

known
function

Discrete energies
in a finite-volume

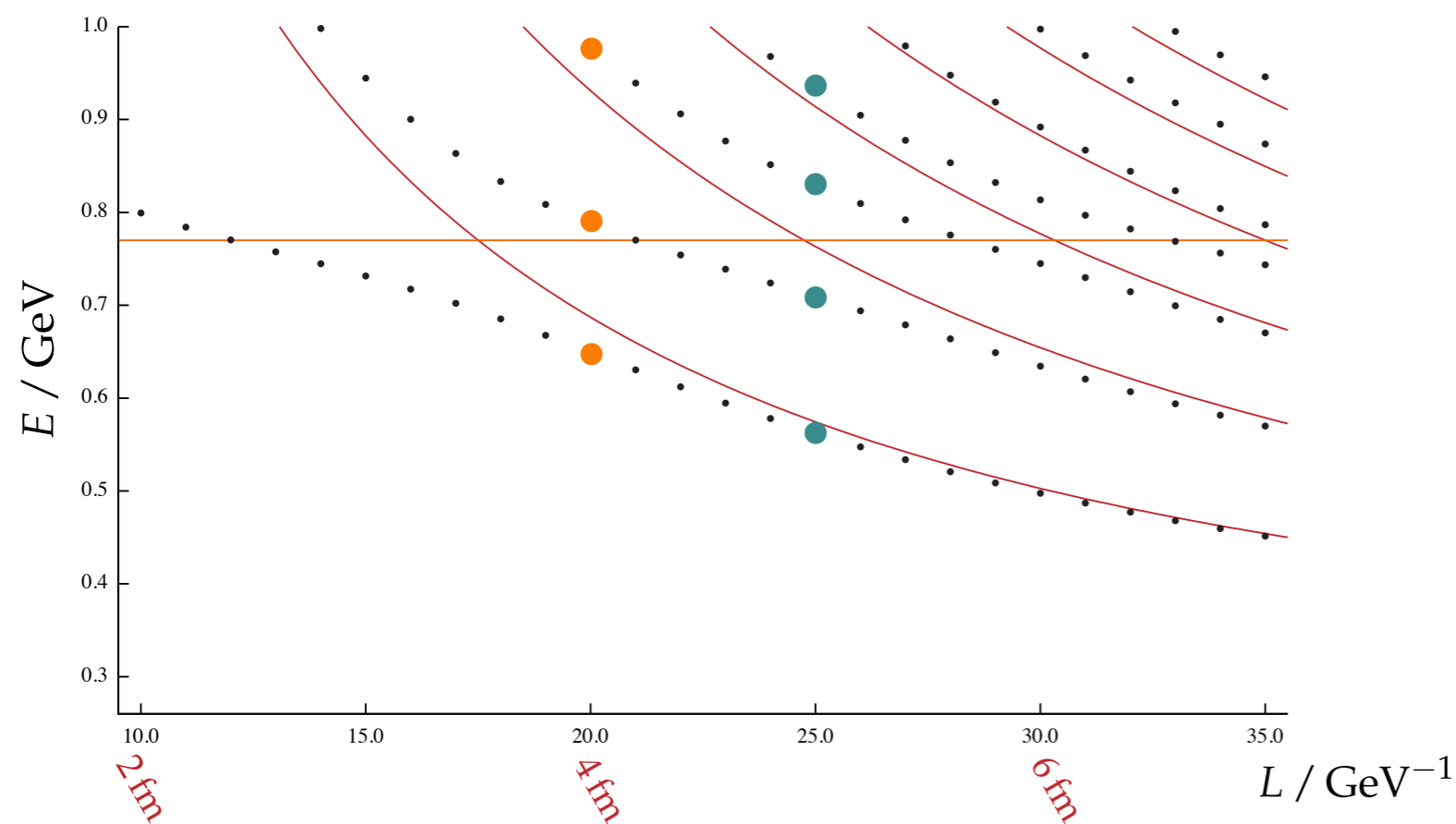


Discrete values
of the phase-shift

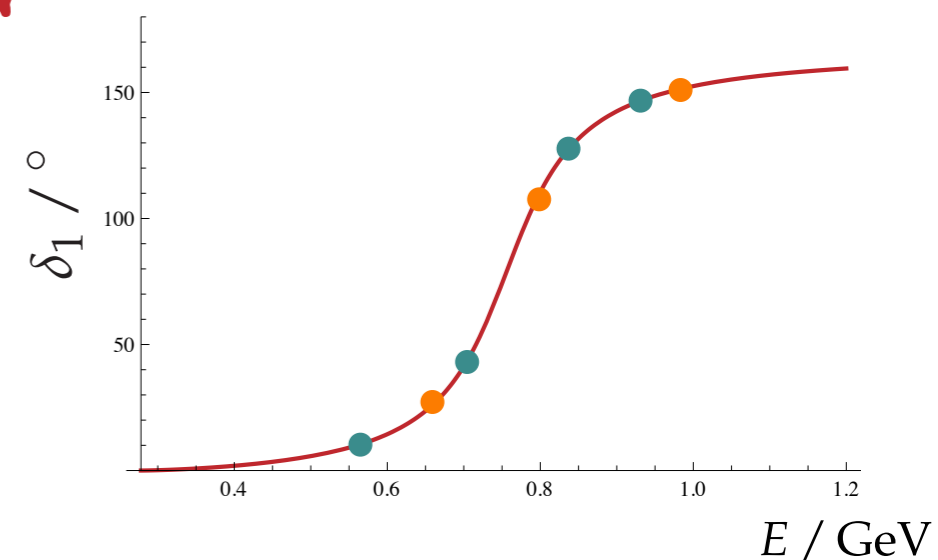
Scattering in a finite cubic volume

- Experimental $\pi\pi$ $I=1$ P -wave scattering amplitude

CUBIC BOX SPECTRUM



P -WAVE PHASE SHIFT



Coupled-channel scattering

HE, JHEP 0507 011
HANSEN, PRD86 016007
BRICENO, PRD88 094507
GUO, PRD88 014051

- Finite-volume formalism recently derived (multiple methods)

$$\det \left[\left([t^{(\ell)}(E)]_{ij}^{-1} + i\rho_i(E) \delta_{ij} \right) - \delta_{ij} \mathcal{M}_\ell(p_i(E)L) \right] = 0$$

scattering matrix phase space known functions *matrices in partial-wave space ..*

- However, this is one equation for multiple unknowns (per energy level) $\frac{1}{2}N(N+1)$ for N channels
 - parameterize the energy dependence of t
 - try to describe a spectrum globally

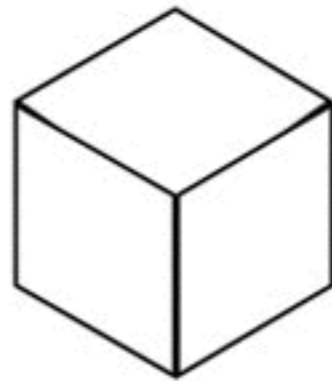
“Energy-dependent” analysis

Finite volume QCD & the hadron spectrum

- Compute correlation functions as an average over field configurations

$$\text{e.g. } \int \mathcal{D}\psi \mathcal{D}\bar{\psi} \mathcal{D}A_\mu \bar{\psi} \Gamma \psi(t) \bar{\psi} \Gamma \psi(0) e^{-\int d^4x \mathcal{L}_{\text{QCD}}(\psi, \bar{\psi}, A_\mu)}$$

‘sum’ ‘field correlation’ ‘probability weight’



*Field integration within a finite, but continuous, hypercube
Need some kind of ultraviolet regulator...*

- Spectrum from two-point correlation functions

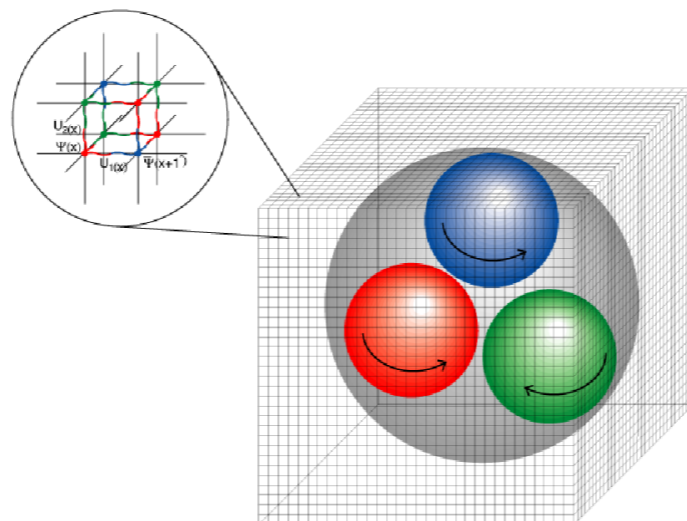
$$\begin{aligned} C(t) &= \langle 0 | \mathcal{O}(t) \mathcal{O}^\dagger(0) | 0 \rangle \\ &= \sum_{\mathbf{n}} e^{-E(\mathbf{n})t} \langle 0 | \mathcal{O}(0) | \mathbf{n} \rangle \langle \mathbf{n} | \mathcal{O}^\dagger(0) | 0 \rangle \end{aligned}$$

Lattice QCD & the hadron spectrum

- Compute correlation functions as a Monte Carlo average over field configurations

$$\text{e.g. } \int \mathcal{D}\psi \mathcal{D}\bar{\psi} \mathcal{D}A_\mu \bar{\psi} \Gamma \psi(t) \bar{\psi} \Gamma \psi(0) e^{-\int d^4x \mathcal{L}_{\text{QCD}}(\psi, \bar{\psi}, A_\mu)}$$

‘sum’ ‘field correlation’ ‘probability weight’



Discretize the action over sites

Serves as an ultraviolet regulator

- Spectrum from two-point correlation functions

$$\begin{aligned} C(t) &= \langle 0 | \mathcal{O}(t) \mathcal{O}^\dagger(0) | 0 \rangle \\ &= \sum_{\mathbf{n}} e^{-E(\mathbf{n})t} \langle 0 | \mathcal{O}(0) | \mathbf{n} \rangle \langle \mathbf{n} | \mathcal{O}^\dagger(0) | 0 \rangle \end{aligned}$$

Excited states from correlators

- how to get at excited QCD eigenstates ?

- optimal operator for state $|\mathbf{n}\rangle$: $\Omega_{\mathbf{n}}^{\dagger} \sim \sum_i v_i^{(\mathbf{n})} \mathcal{O}_i^{\dagger}$

for a basis of meson operators $\{\mathcal{O}_i\}$

- can be obtained (in a variational sense) from the matrix of correlators

$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^{\dagger}(0) | 0 \rangle$$

- by solving a generalized eigenvalue problem

$$C(t)v^{(\mathbf{n})} = C(t_0)v^{(\mathbf{n})} \lambda_{\mathbf{n}}(t)$$

eigenvalues

$$\lambda_{\mathbf{n}}(t) \sim e^{-E_{\mathbf{n}}(t-t_0)}$$

‘diagonalize the correlation matrix’

- a large basis can be constructed using covariant derivatives :

$$\mathcal{O} \sim \bar{\psi} \Gamma \overleftrightarrow{D} \dots \overleftrightarrow{D} \psi$$

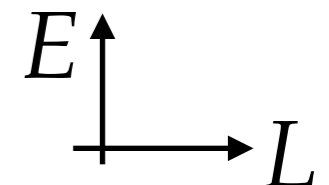
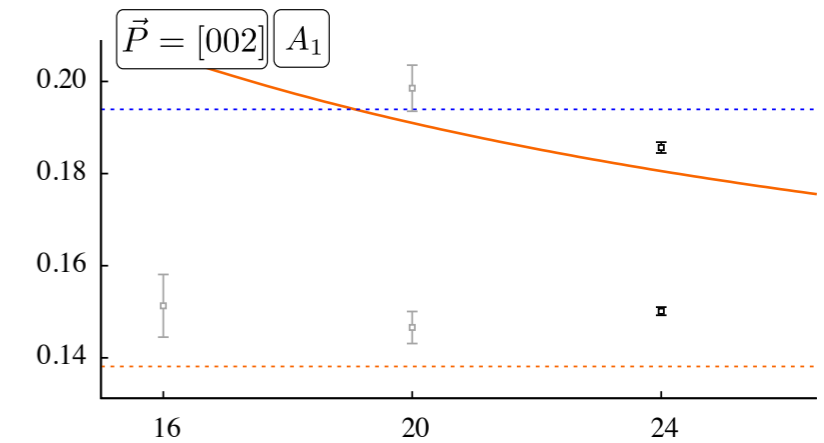
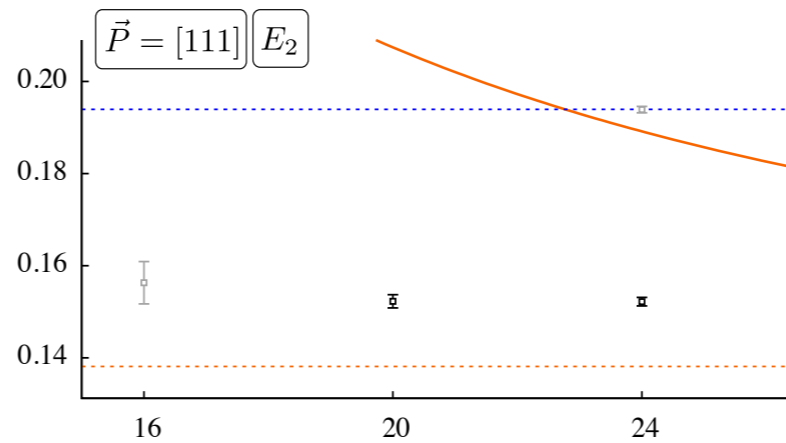
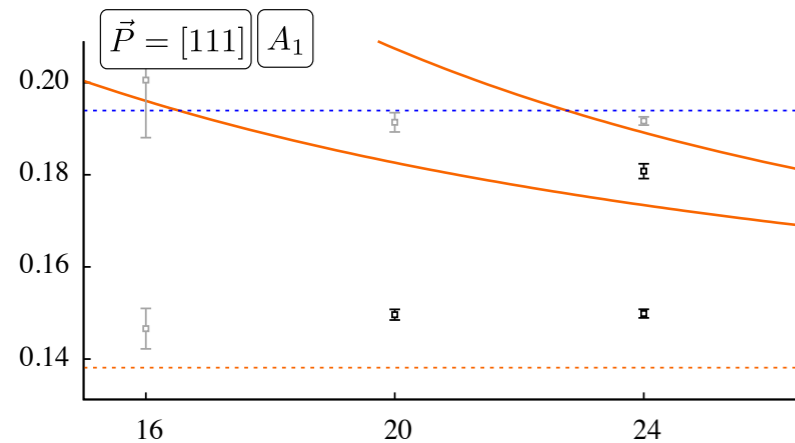
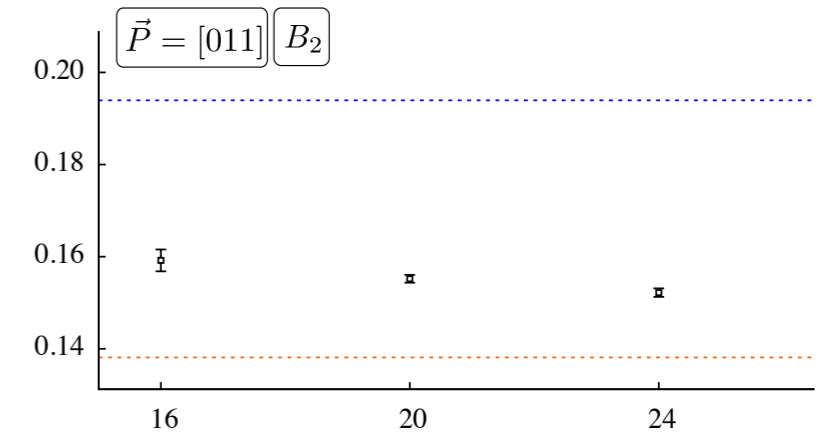
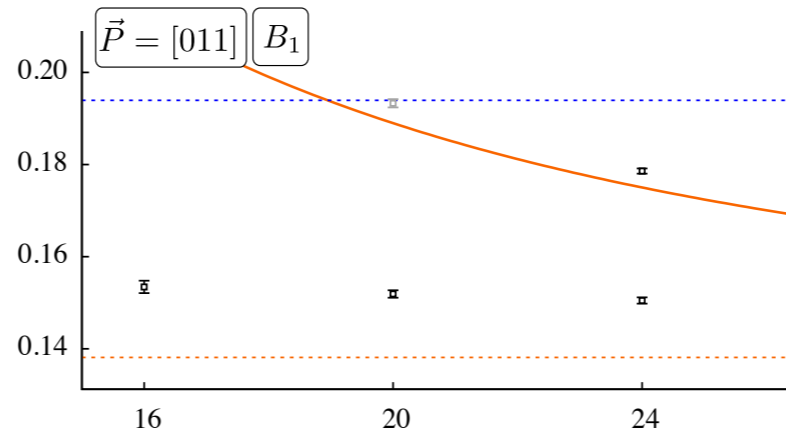
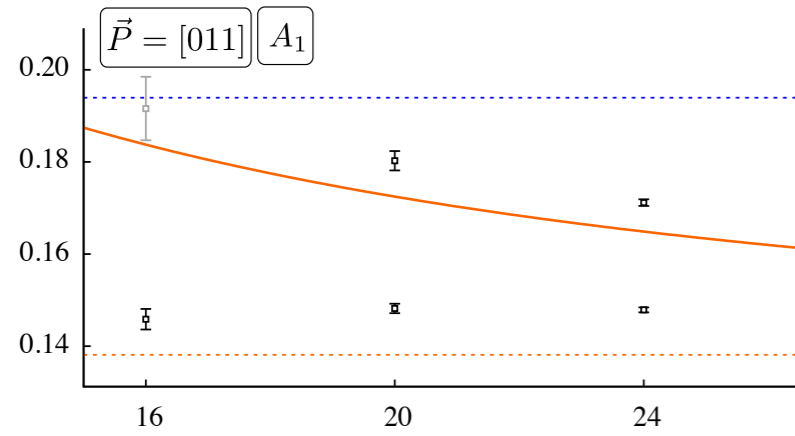
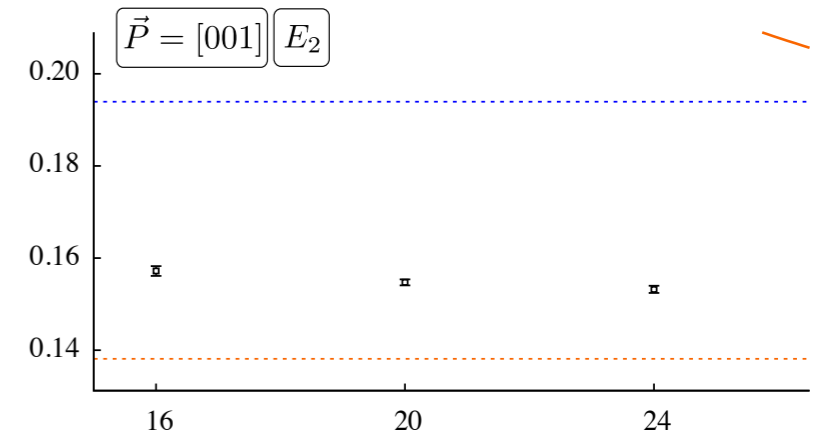
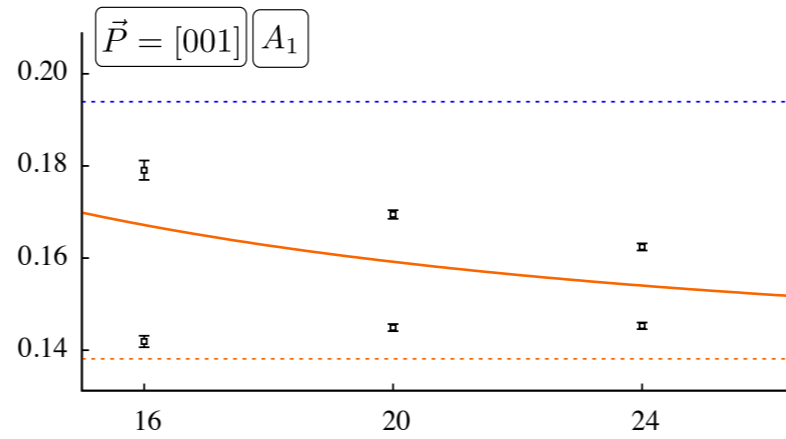
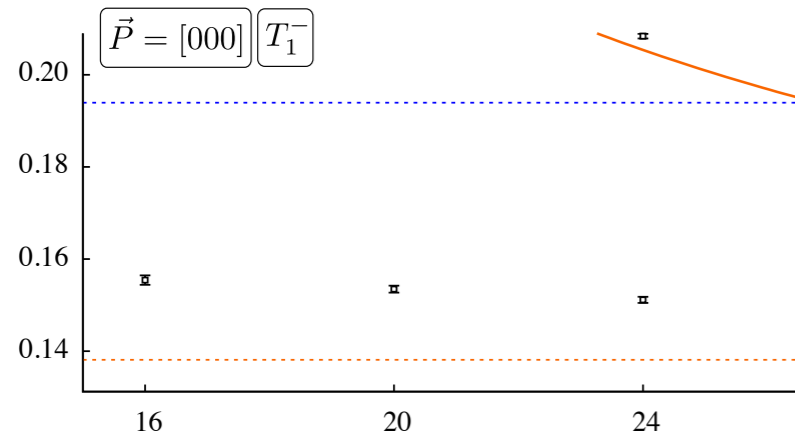
Including multi-meson operators

- form correlator matrix with both $\bar{\psi}\Gamma\psi$ and $\pi\pi$ -like

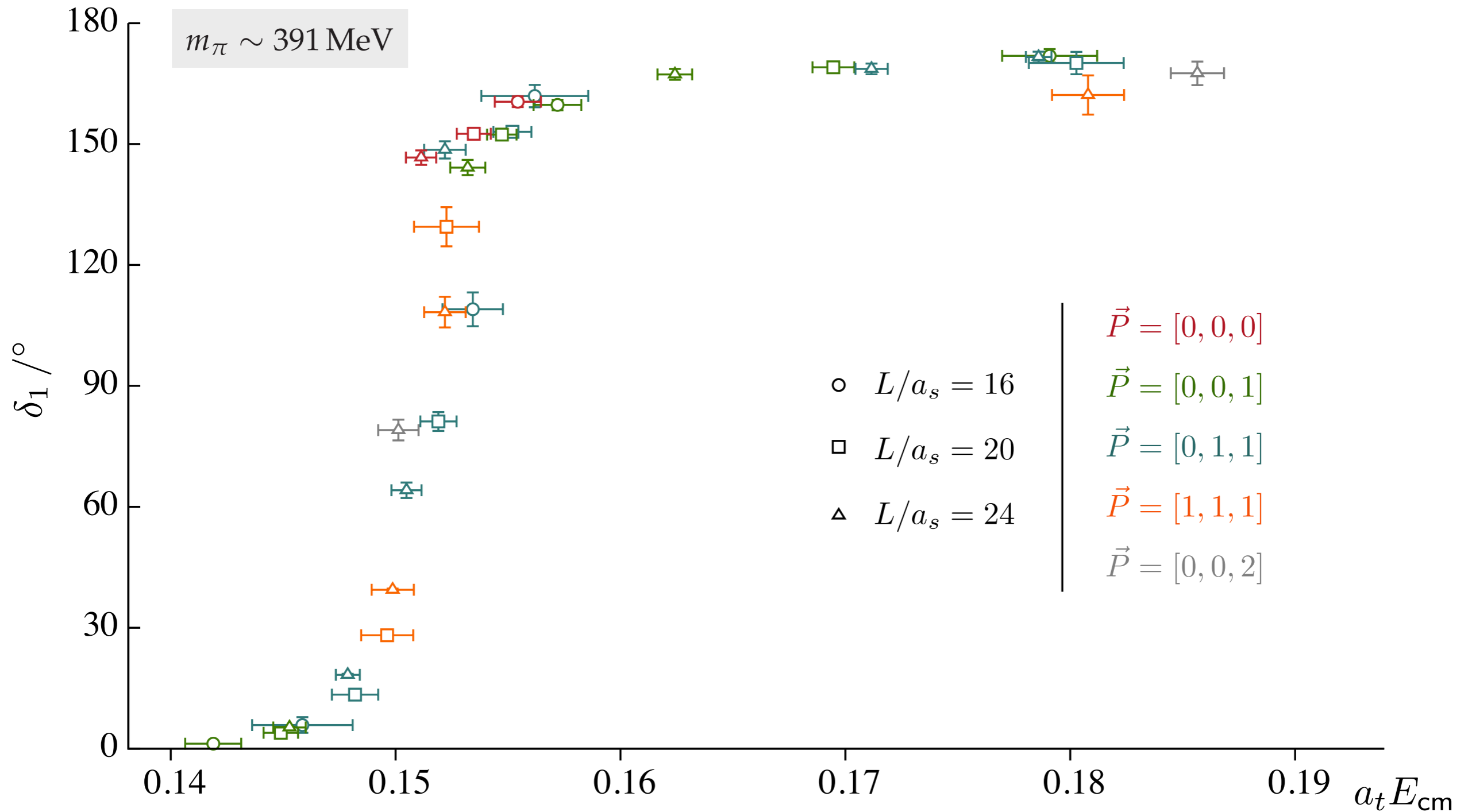
- include operators which resemble a pair of pions $\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}; \vec{k}_1, \vec{k}_2) \pi^\dagger(\vec{k}_1) \pi^\dagger(\vec{k}_2)$
 $\pi^\dagger \sim \bar{\psi}\Gamma\psi$

Finite-volume spectrum - moving frames

$m_\pi \sim 391 \text{ MeV}$

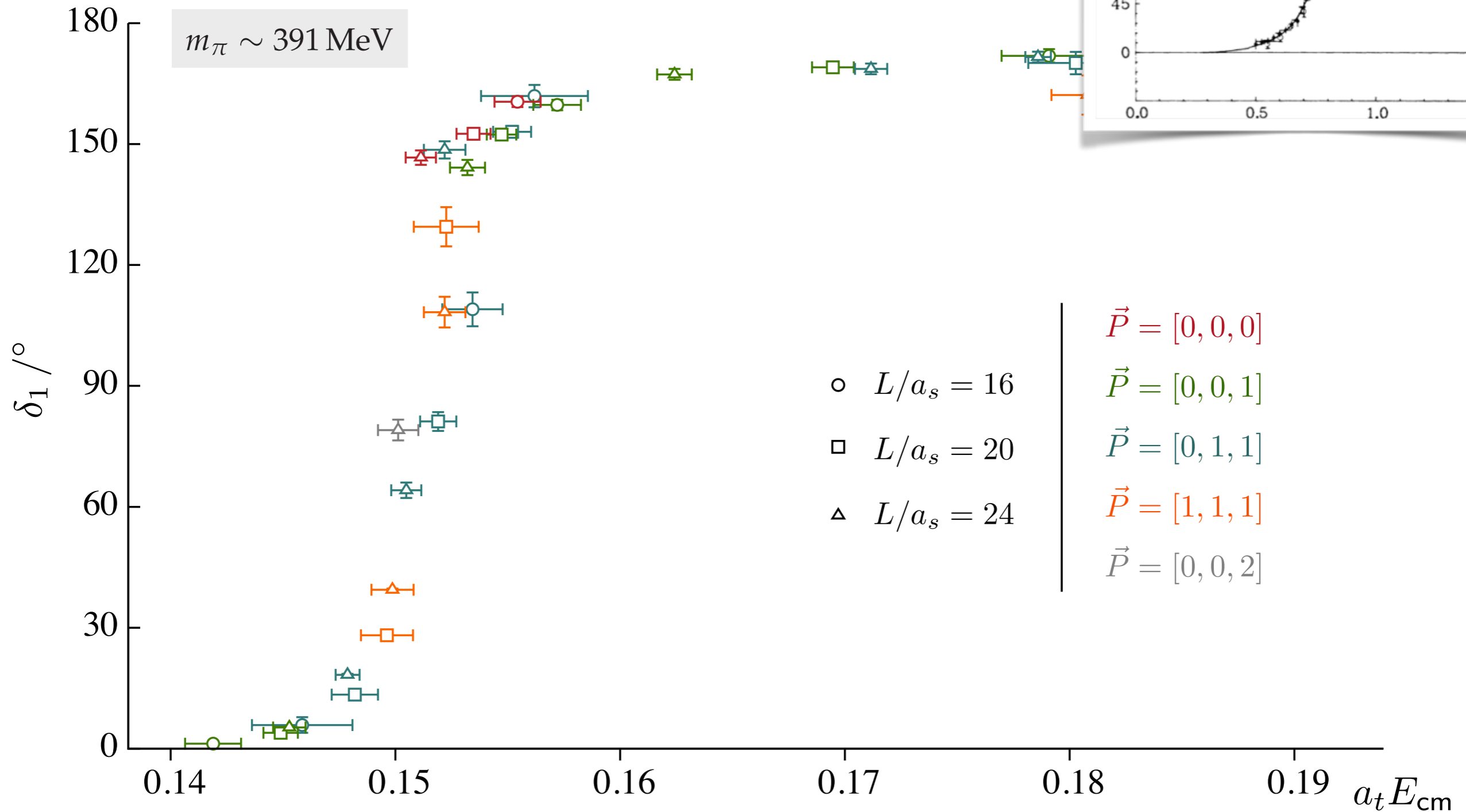


$\pi\pi$ P -wave phase-shift



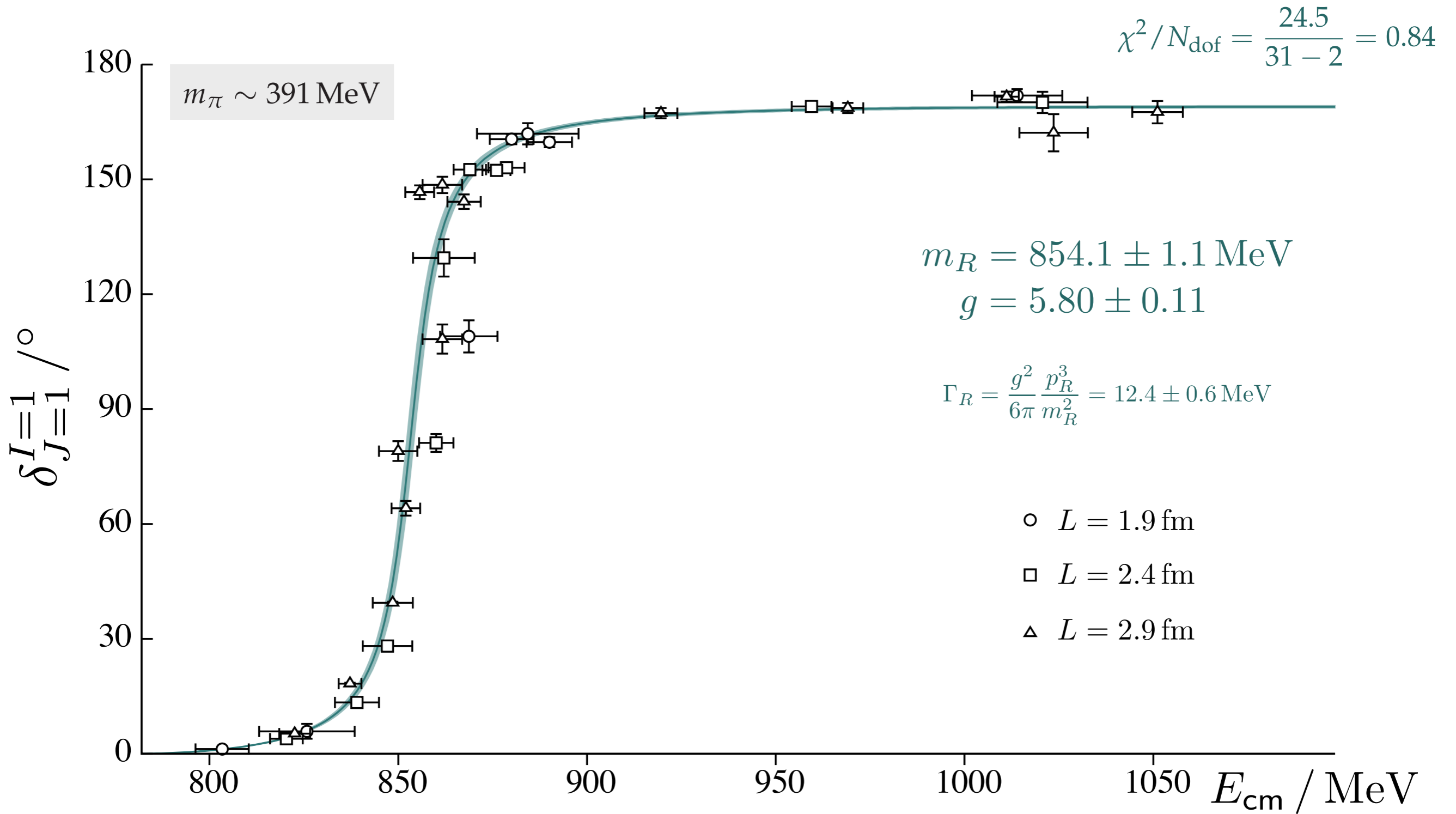
PRD87 034505 (2013)

$\pi\pi$ P-wave phase-shift



PRD87 034505 (2013)

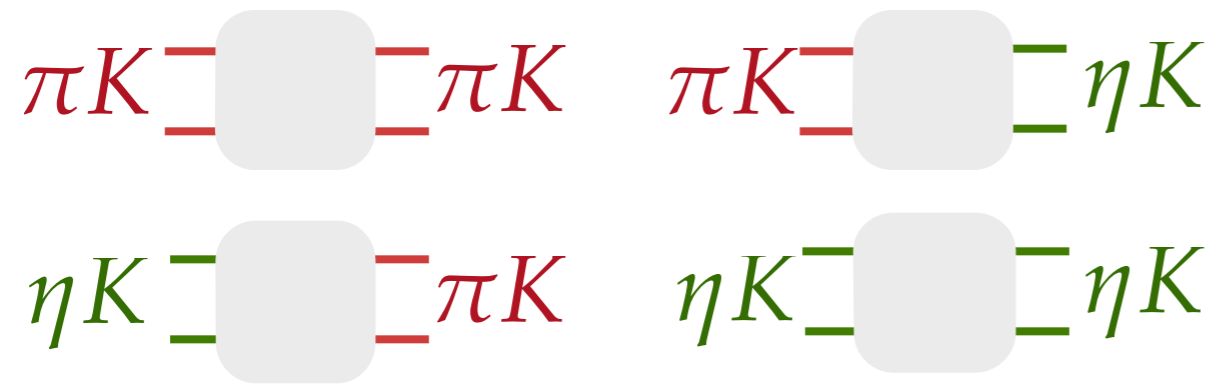
ρ resonance



$g_{\text{phys.}} = 5.95(2)$ *PDG*

$\pi K/\eta K$ scattering & kaon resonances

- Example of coupled-channel scattering



- Compute finite-volume spectrum

$\bar{u}\Gamma s$

$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}; \vec{k}_1, \vec{k}_2) \pi^\dagger(\vec{k}_1) K^\dagger(\vec{k}_2)$$

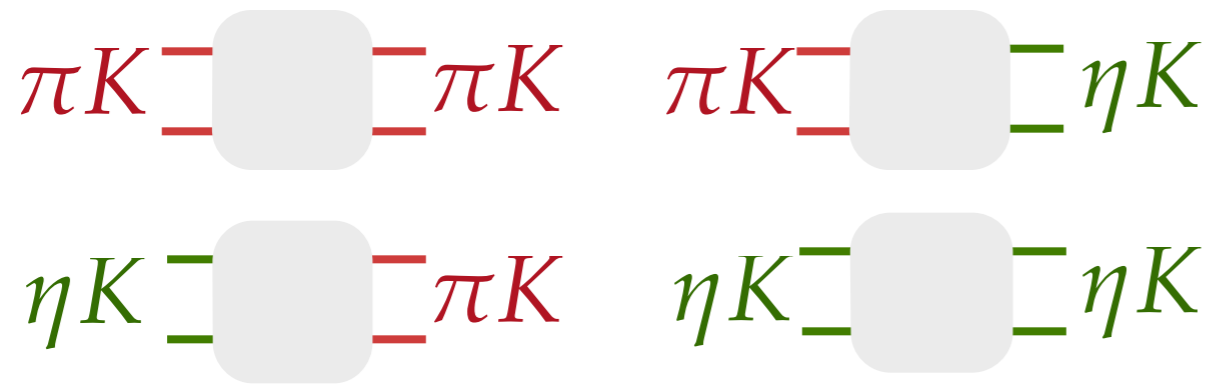
$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}; \vec{k}_1, \vec{k}_2) \eta^\dagger(\vec{k}_1) K^\dagger(\vec{k}_2)$$

PRL 113 182001

PRD 91 054008

$\pi K/\eta K$ scattering & kaon resonances

- Example of coupled-channel scattering



- Compute finite-volume spectrum

$\bar{u}\Gamma s$

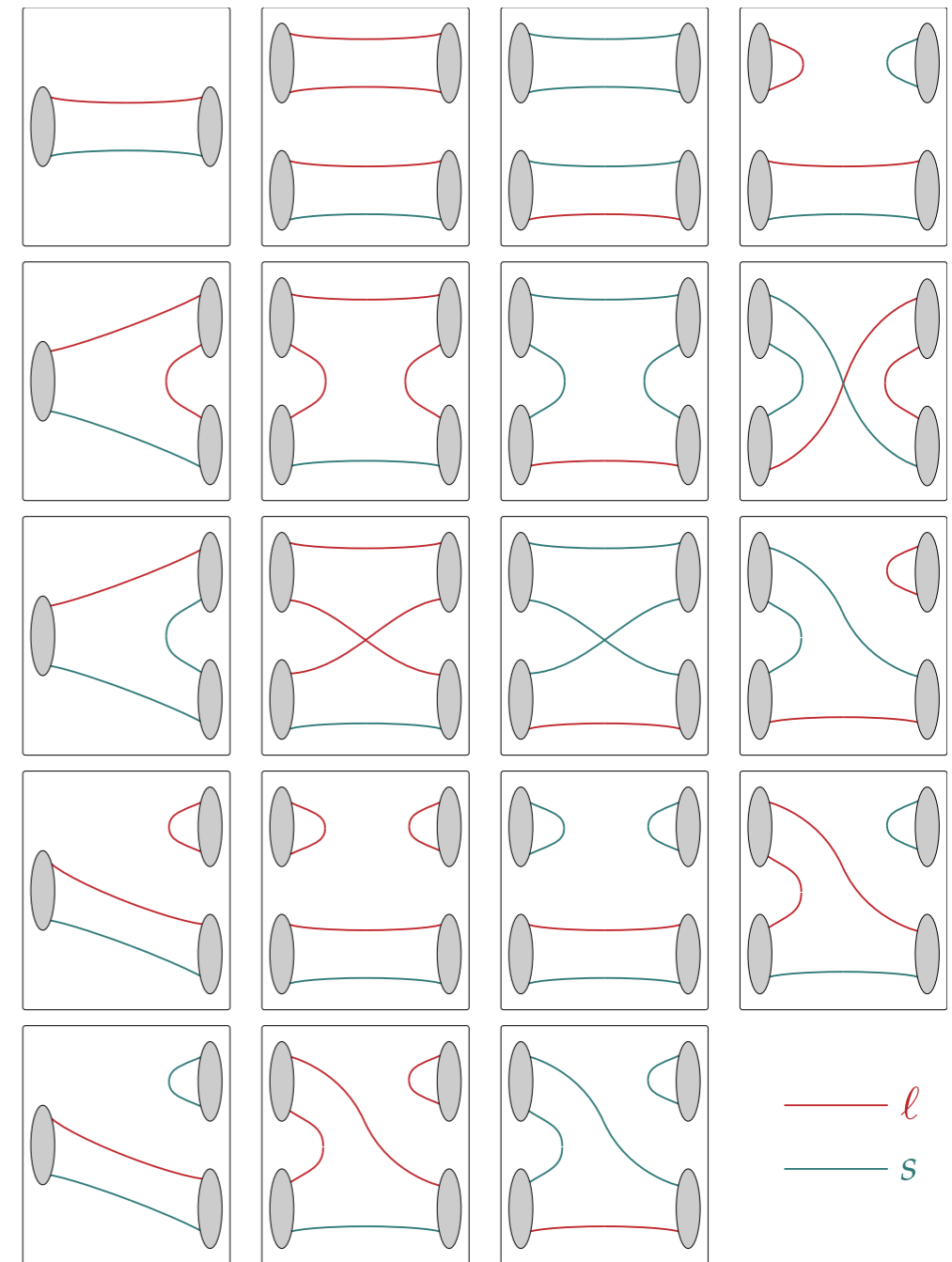
$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}; \vec{k}_1, \vec{k}_2) \pi^+(\vec{k}_1) K^+(\vec{k}_2)$$

$$\sum_{\hat{k}_1, \hat{k}_2} C(\Lambda, \vec{P}; \vec{k}_1, \vec{k}_2) \eta^+(\vec{k}_1) K^+(\vec{k}_2)$$

PRL 113 182001

PRD 91 054008

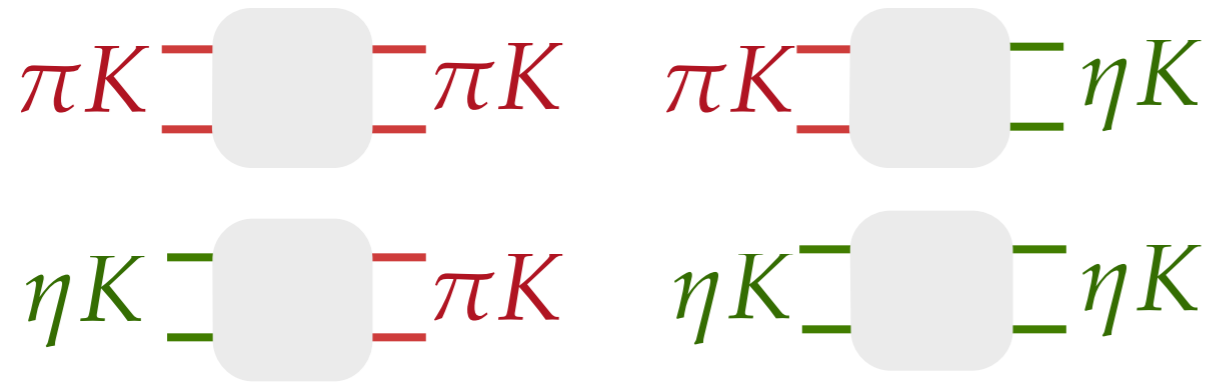
WICK CONTRACTIONS



— l
— s

$\pi K/\eta K$ scattering & kaon resonances

- Parameterize the t -matrix in a unitarity conserving way



$$t_{ij}^{-1}(E) = K_{ij}^{-1}(E) + \delta_{ij} I_i(E)$$

$$K_{ij}(E) = \frac{g_i g_j}{m^2 - E^2} + \gamma_{ij}$$

- Vary the parameters, solving

$$\det \left[([t^{(\ell)}(E)]_{ij}^{-1} + i\rho_i(E) \delta_{ij}) - \delta_{ij} \mathcal{M}_\ell(E, L) \right] = 0$$

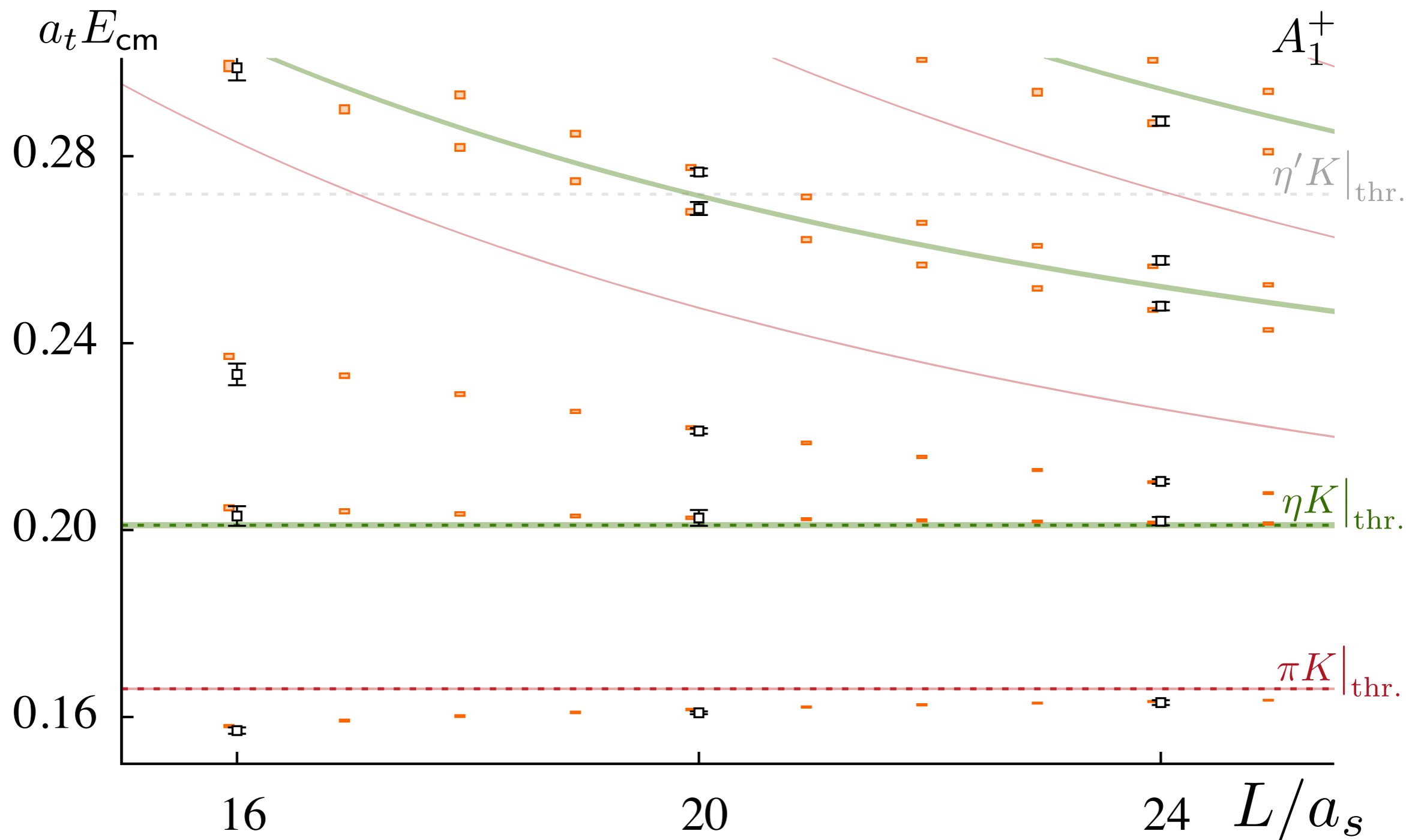
for the spectrum in each irreducible representation & momentum

Want pole mass and couplings of t -matrix

$\pi K/\eta K$ scattering

Model description $\chi^2/N_{\text{dof}} = \frac{6.40}{15-6} = 0.71$

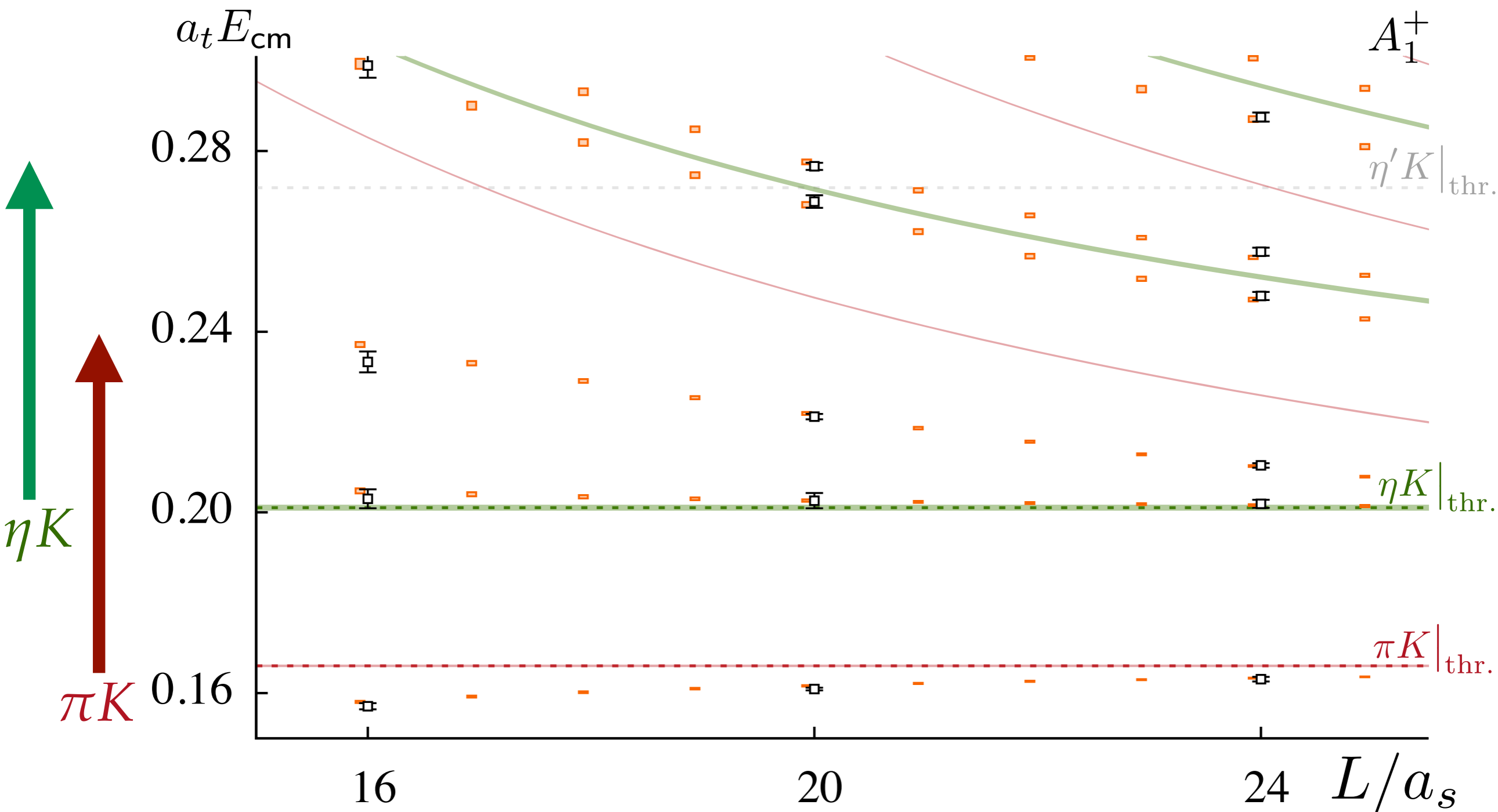
$m_\pi \sim 391 \text{ MeV}$



$\pi K/\eta K$ scattering

Model description $\chi^2/N_{\text{dof}} = \frac{6.40}{15-6} = 0.71$

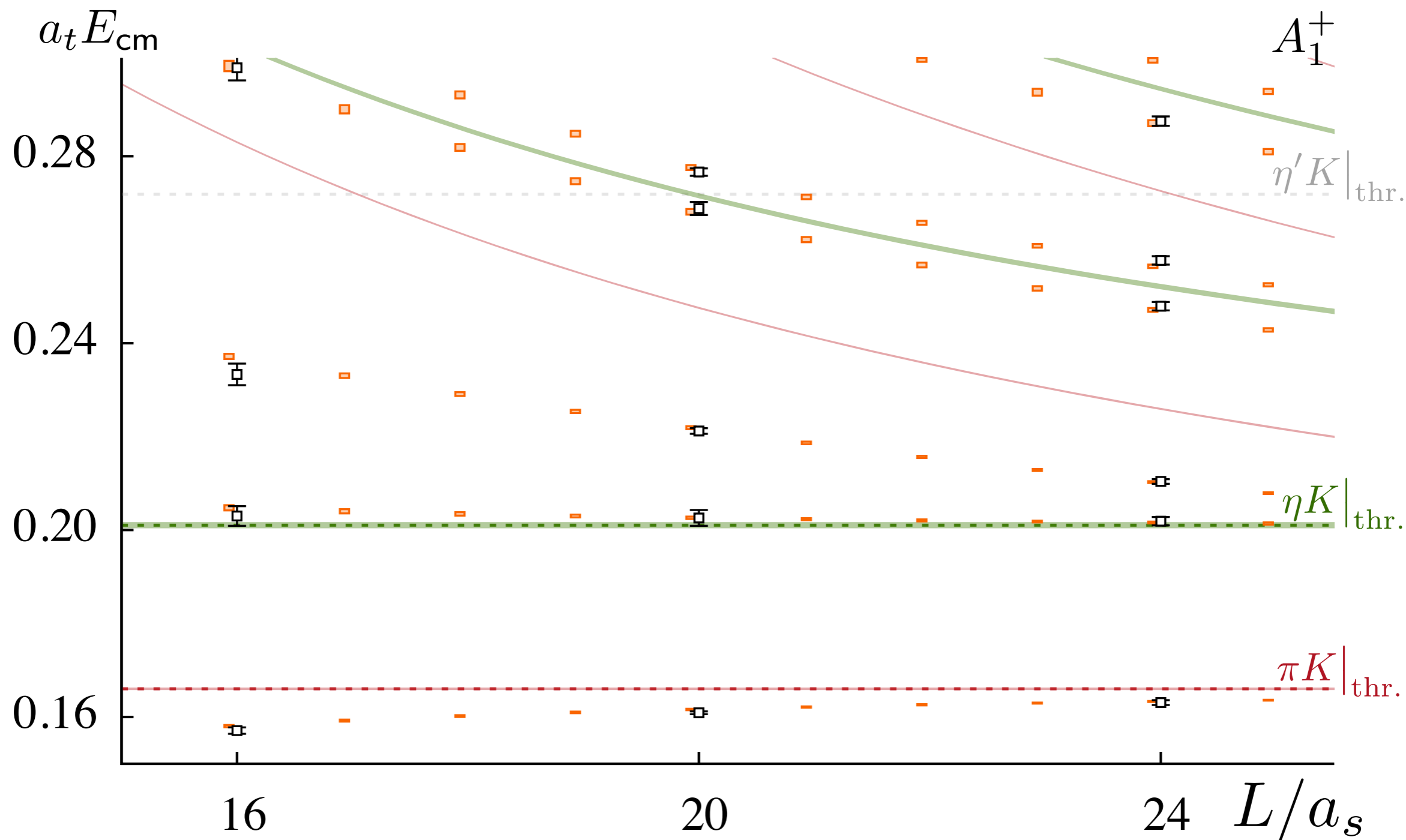
$m_\pi \sim 391$ MeV



$\pi K/\eta K$ scattering

Model description $\chi^2/N_{\text{dof}} = \frac{6.40}{15-6} = 0.71$

$m_\pi \sim 391 \text{ MeV}$



$\pi K/\eta K$ scattering

$m_\pi \sim 391 \text{ MeV}$

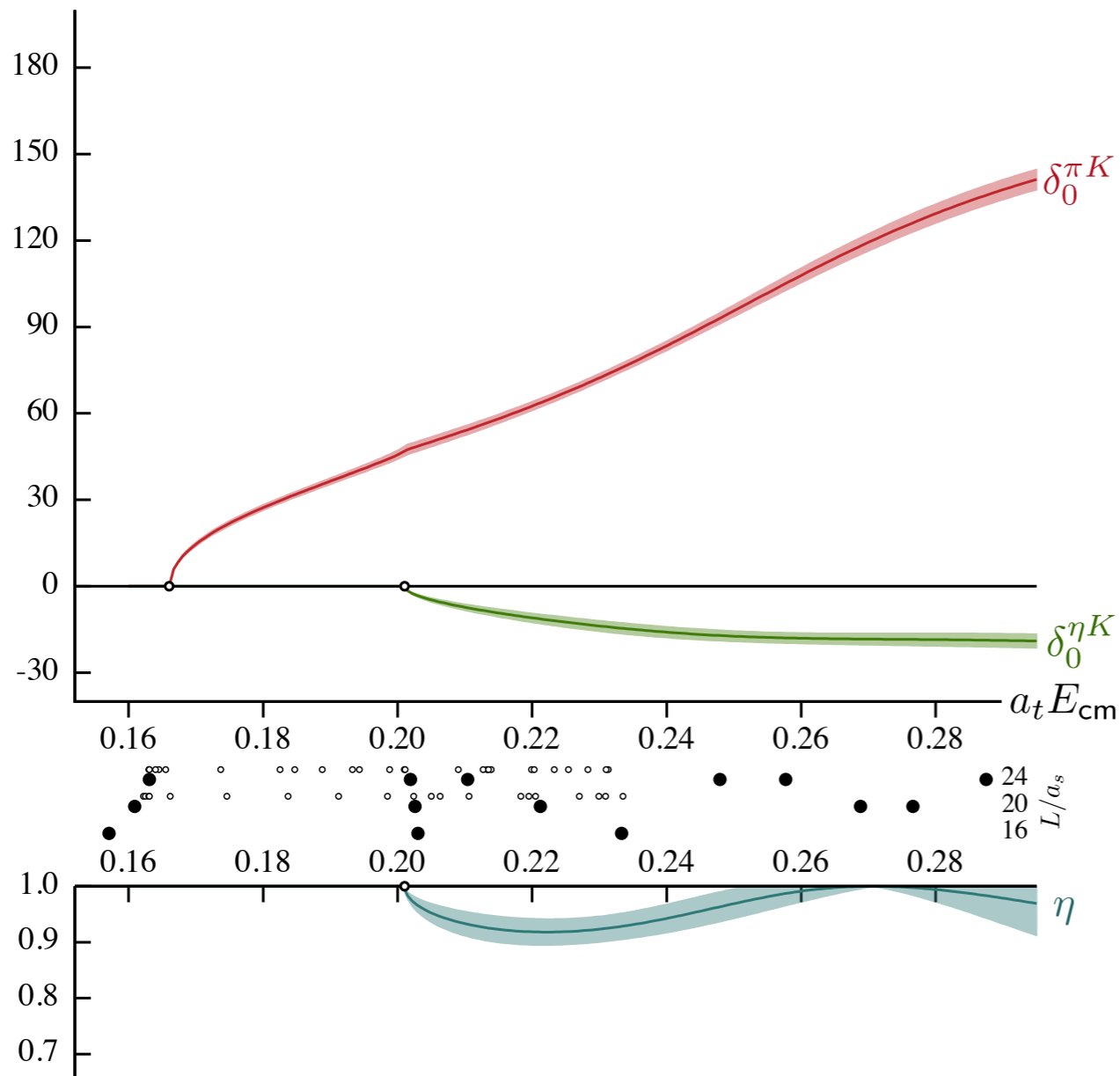
- Describe all the finite-volume spectra

$$\chi^2/N_{\text{dof}} = \frac{49.1}{61 - 6} = 0.89$$

$$S_{\pi K, \pi K} = \eta e^{2i\delta^{\pi K}}$$

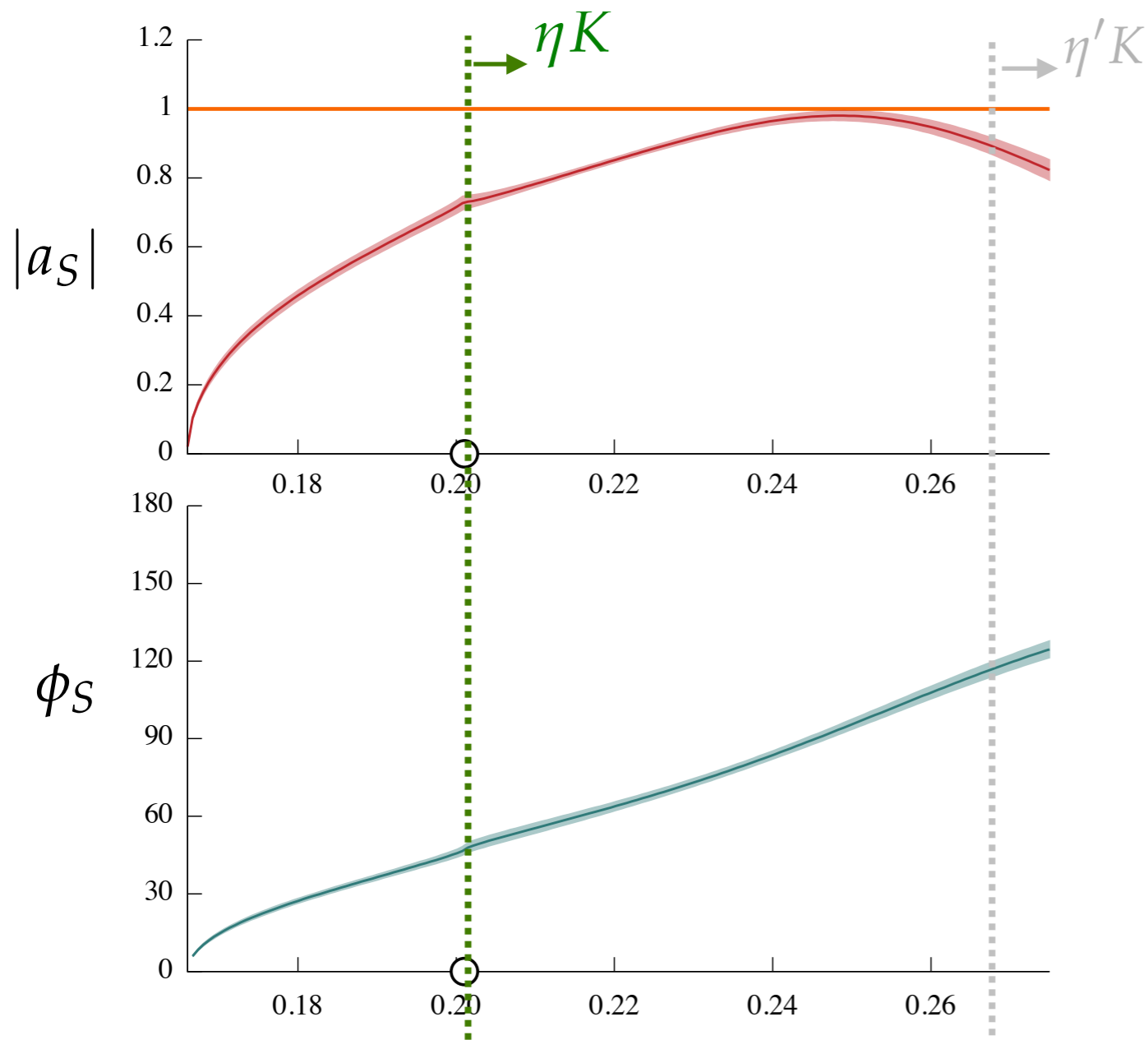
$$S_{\eta K, \eta K} = \eta e^{2i\delta^{\eta K}}$$

S-WAVE $\pi K/\eta K$ SCATTERING

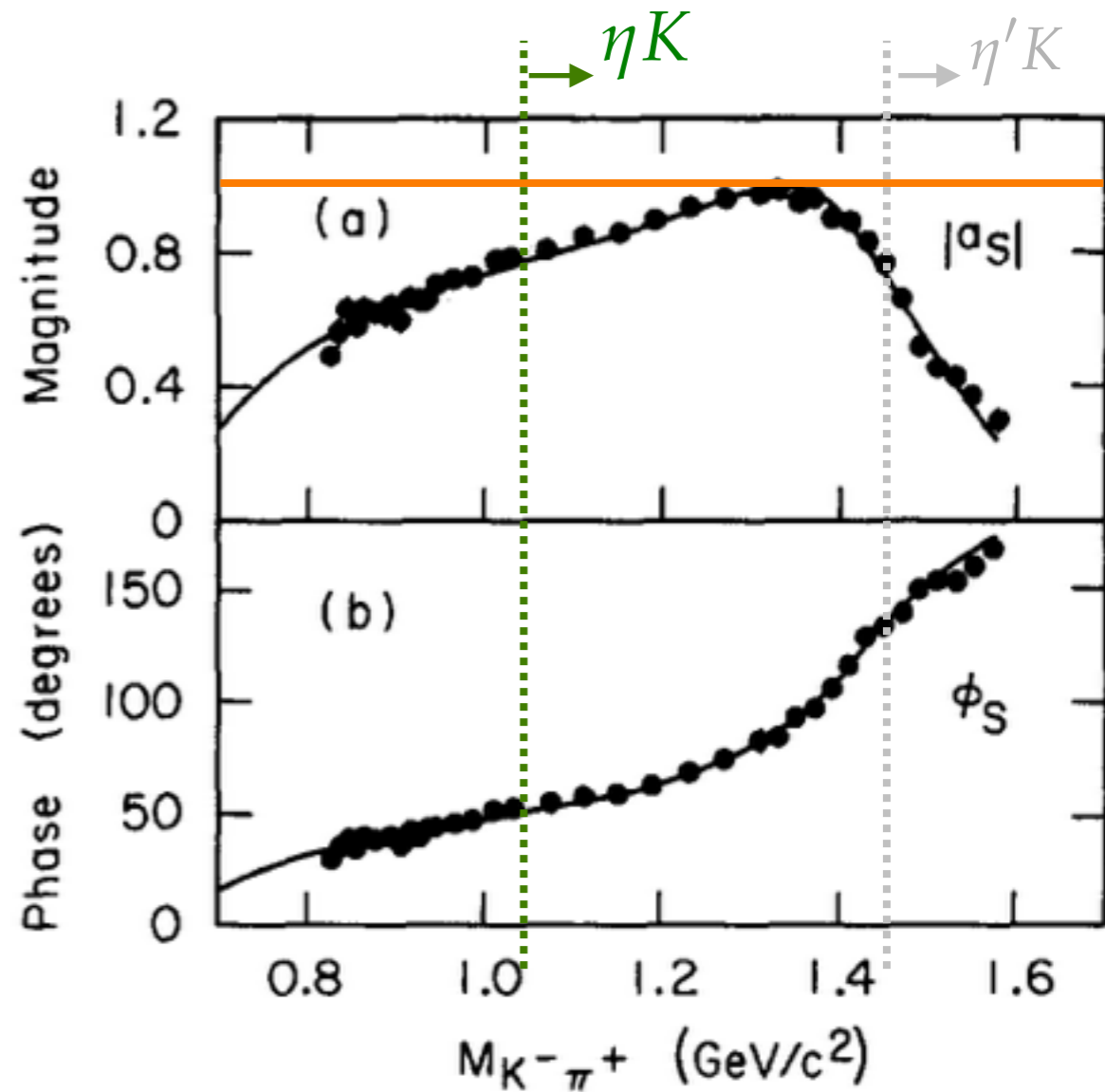


Versus experimental scattering

S-WAVE $\pi K \rightarrow \pi K$ AMPLITUDE



$m_\pi \sim 391$ MeV

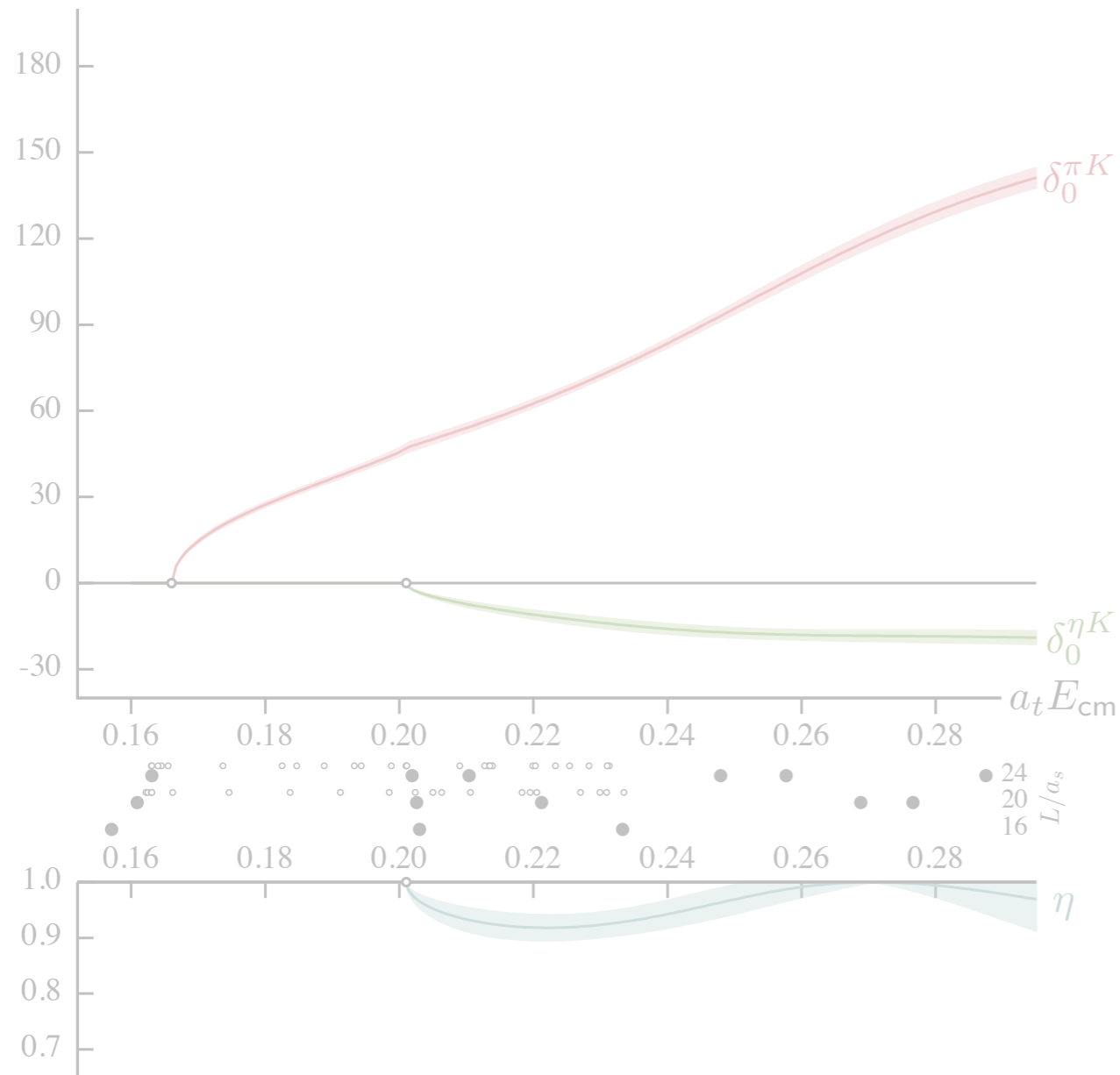


LASS, NPB296 493

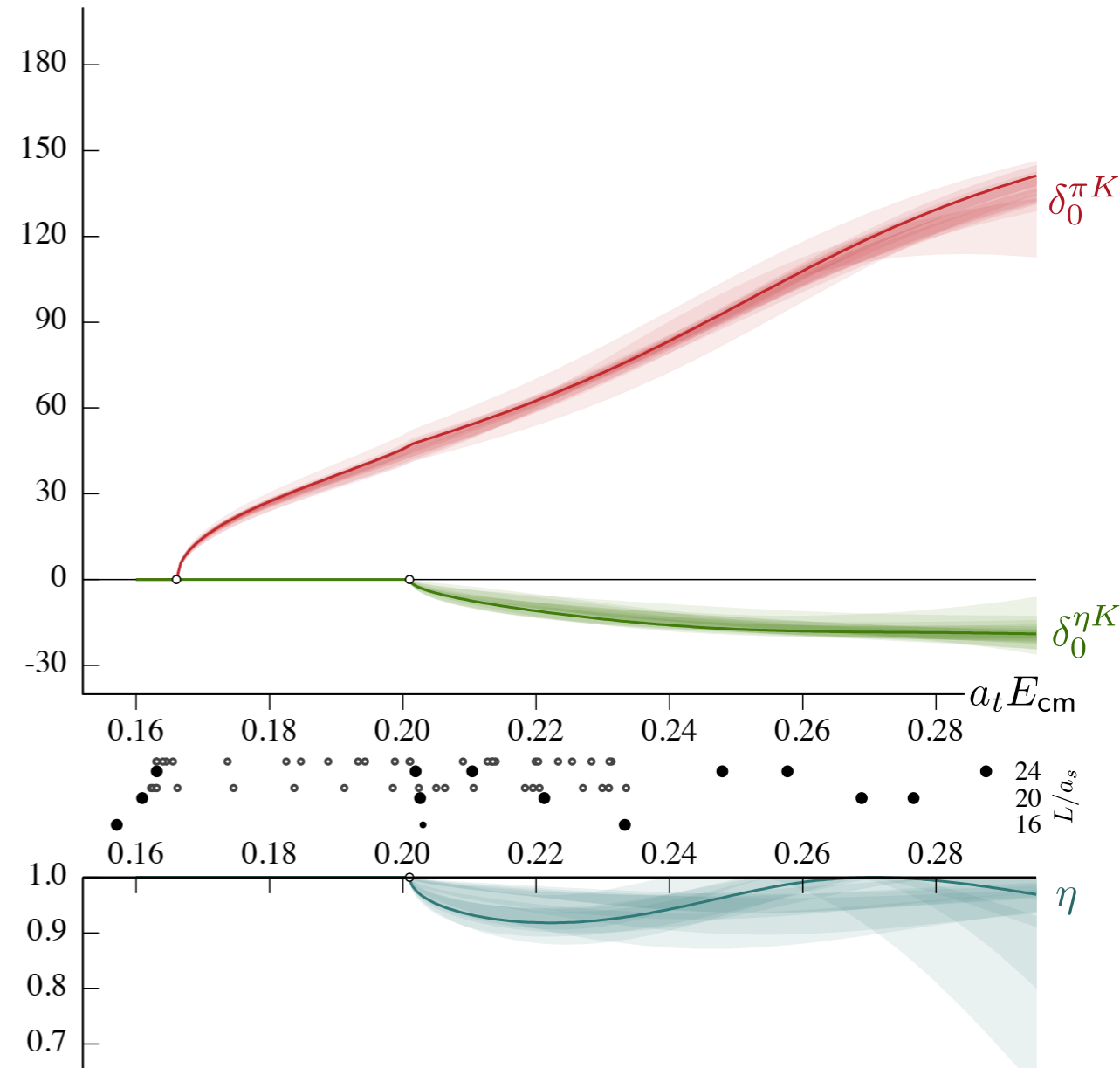
$\pi K/\eta K$ scattering

$$m_\pi \sim 391 \text{ MeV}$$

- Are the result parameterization dependent ?
 - Try a range of parameterizations ...



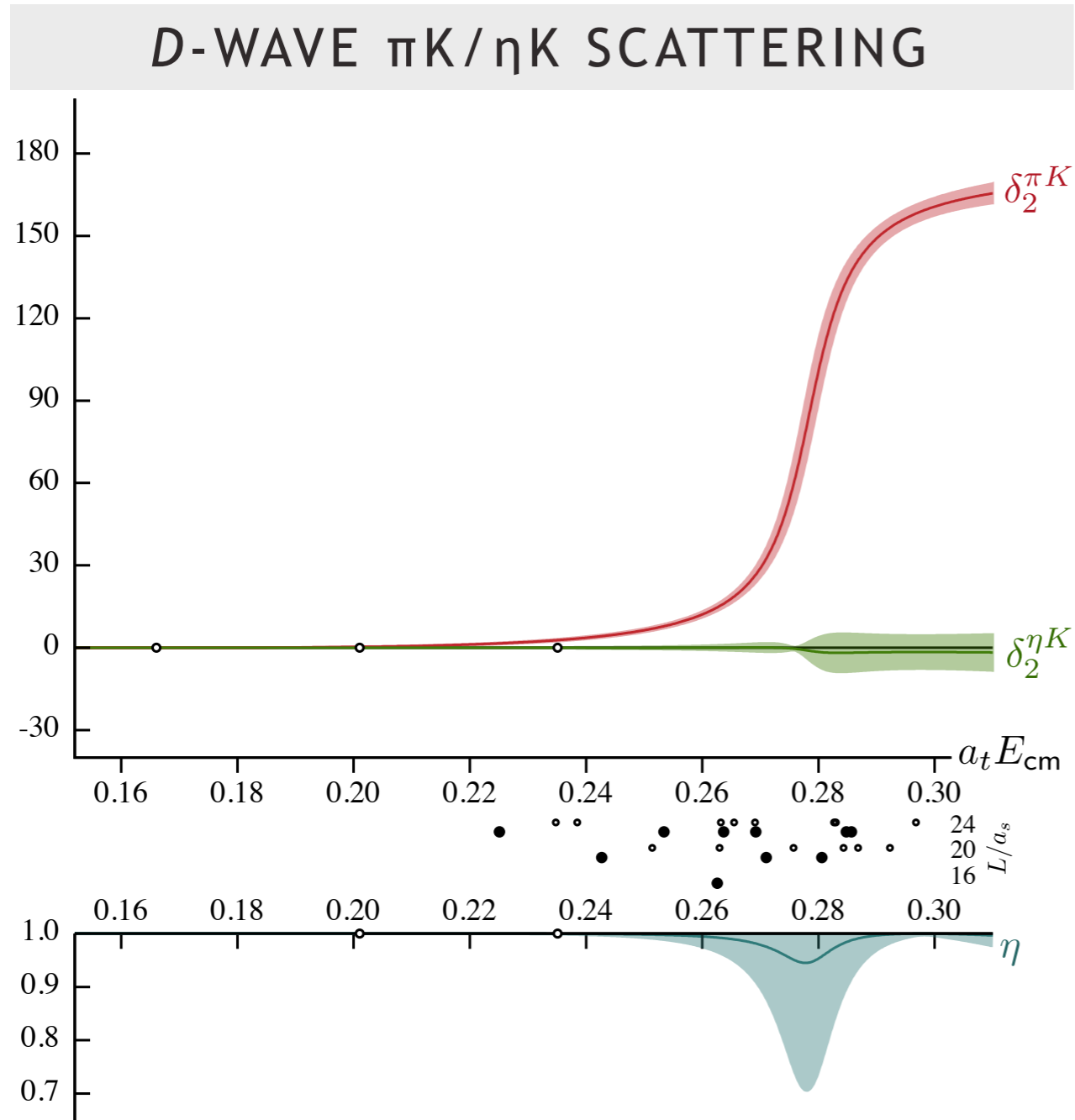
S-WAVE $\pi K/\eta K$ SCATTERING



- gross features are robust

$\pi K/\eta K$ scattering

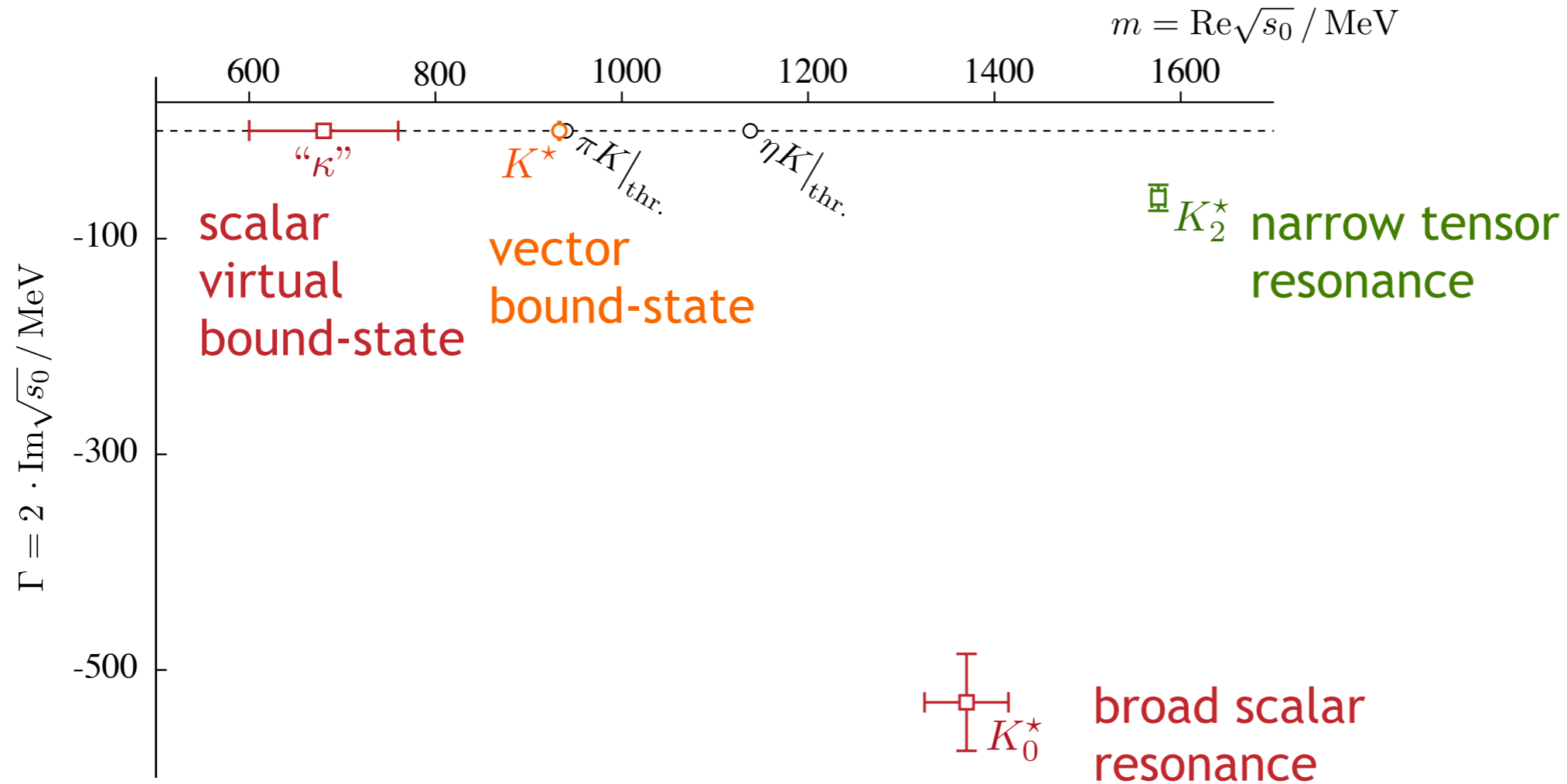
- Clear narrow resonance in D -wave scattering



$m_\pi \sim 391 \text{ MeV}$

Singularity content

- t -matrix poles as least model-dependent characterization of resonances



$m_\pi \sim 391 \text{ MeV}$

PRL 113 182001
PRD 91 054008

Scattering with external currents ?

- E.g. $\pi\gamma \rightarrow \pi\pi$ in P -wave : the ρ appears as a resonance

- The observables are the amplitudes $A_\ell(E_{\pi\pi}, Q^2)$

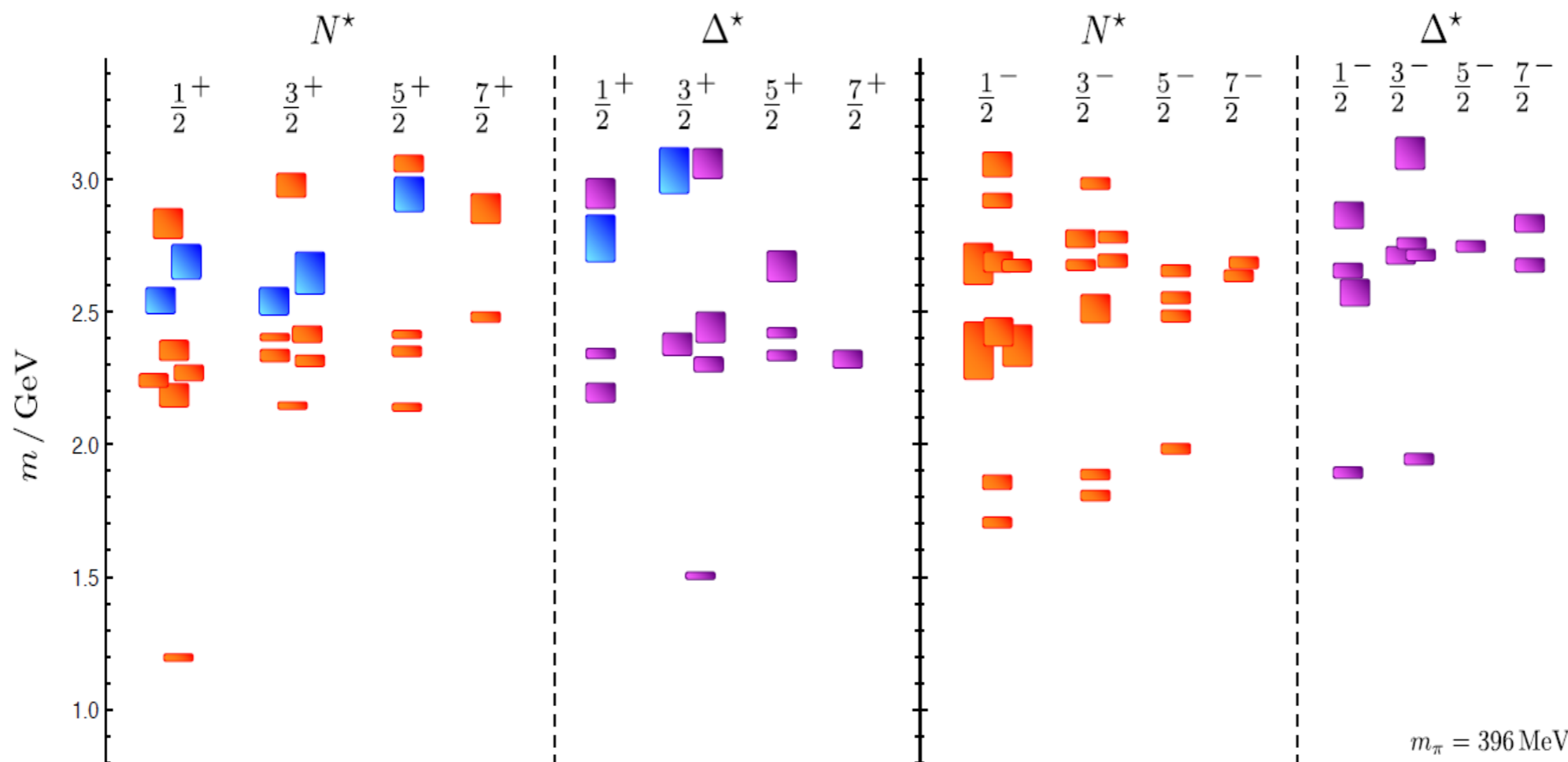
$\rho \rightarrow \pi\gamma$ form-factor
defined at the ρ -pole

$$A_1(s \rightarrow s_\rho) \sim \frac{c_{\rho\pi\pi} c_{\rho\pi\gamma}(Q^2)}{s - s_\rho}$$

Hey, this is NSTAR - where are the baryons?

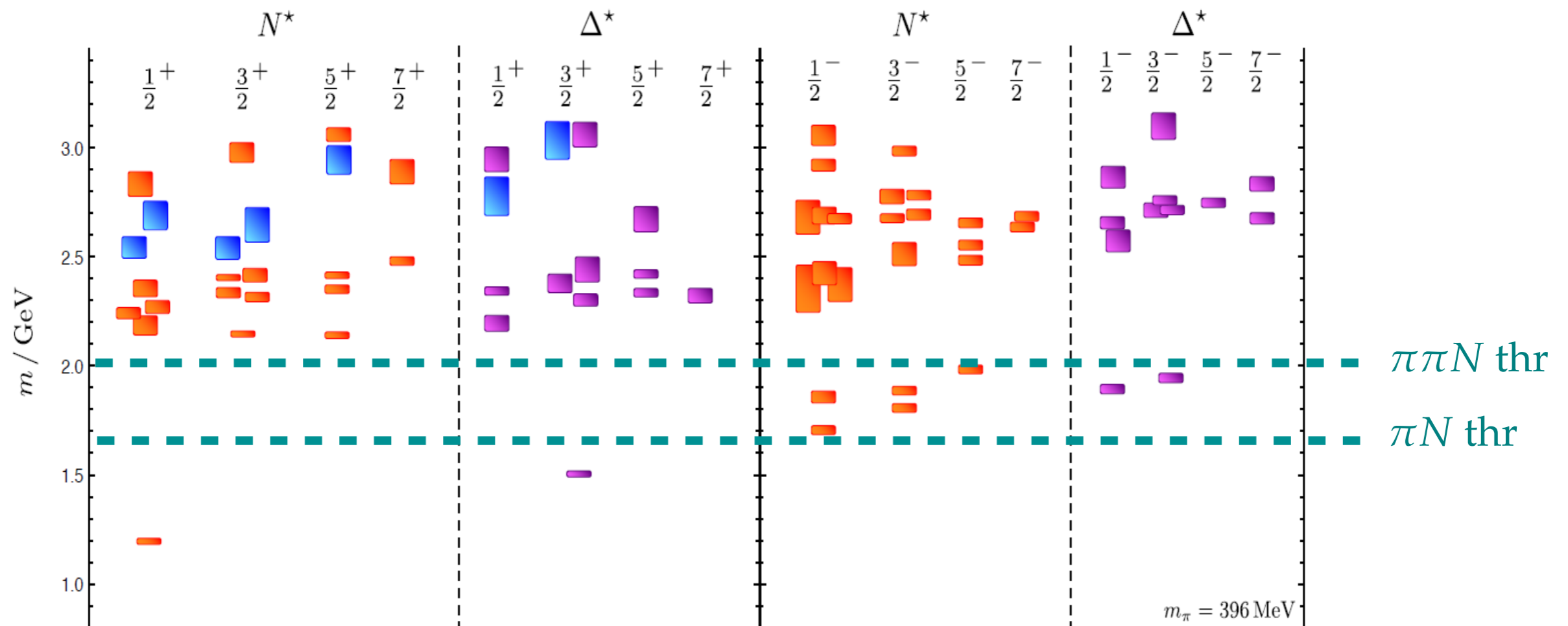
- Initial determination of spectrum with only qqq style operators
 - See rich spectrum, including hybrid-like states

PRD 84 & 85



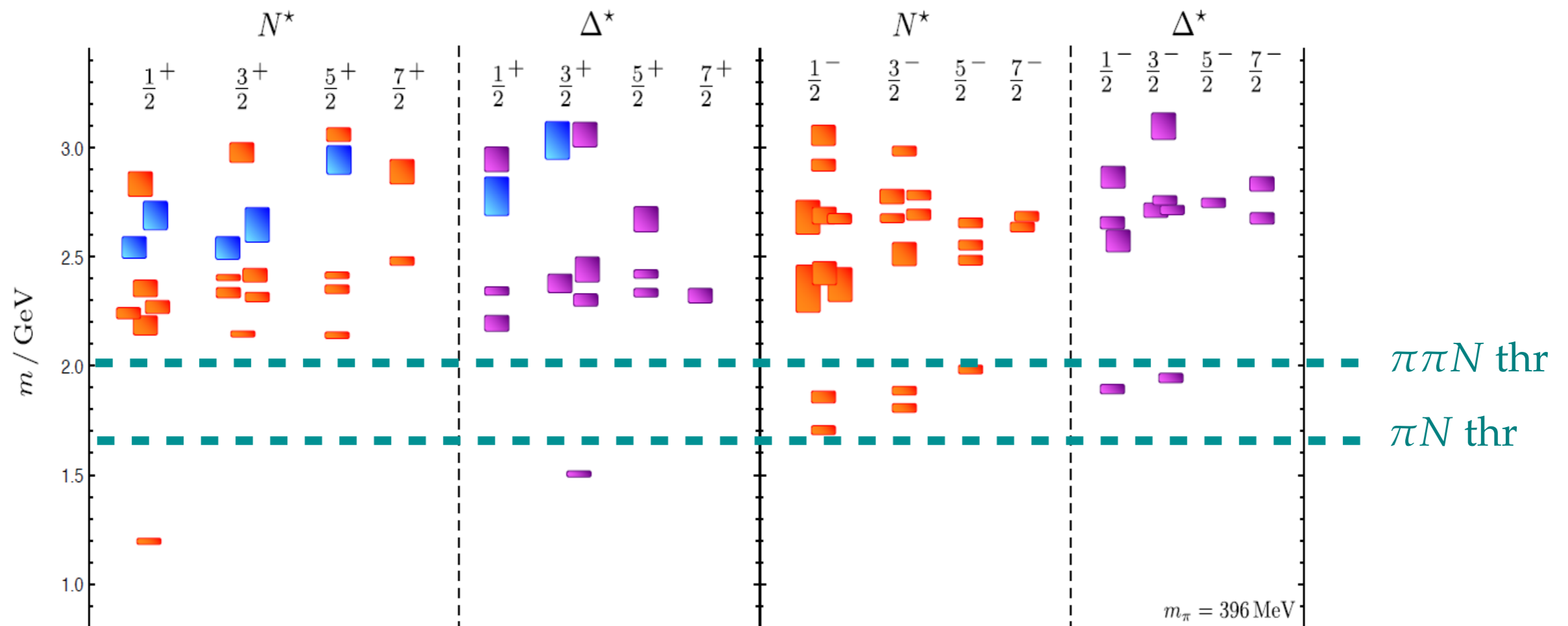
Hey, this is NSTAR - where are the baryons?

- Initial determination of spectrum with only qqq style operators PRD 84 & 85
 - See rich spectrum, including hybrid-like states
 - However, no operators that look like πN or $\pi\pi N$ - missing scattering states



Hey, this is NSTAR - where are the baryons?

- Initial determination of spectrum with only qqq style operators PRD 84 & 85
 - See rich spectrum, including hybrid-like states
 - However, no operators that look like πN or $\pi\pi N$ - missing scattering states
 - Some initial results in S11 have appeared GRAZ GROUP

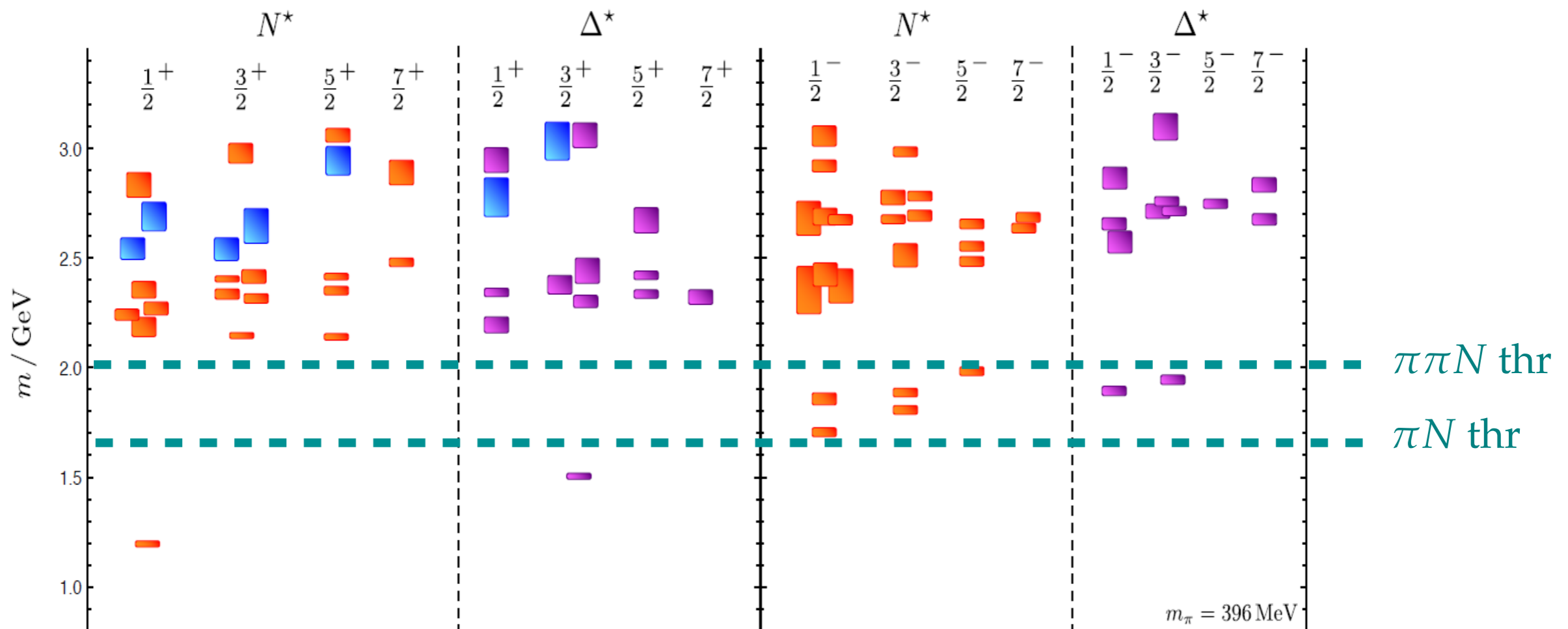


Hey, this is NSTAR - where are the baryons?

- Initial determination of spectrum with only qqq style operators *PRD 84 & 85*
 - See rich spectrum, including hybrid-like states
 - However, no operators that look like πN or $\pi\pi N$ - missing scattering states
 - Some initial results in S_{11} have appeared *GRAZ GROUP*

- Development of three-body formalism required

HANSEN & SHARPE - MUCH PROGRESS



Hadron Spectrum Collaboration

JEFFERSON LAB

Jozef Dudek
Robert Edwards
Balint Joo
David Richards
Frank Winter

TRINITY COLLEGE, DUBLIN

Mike Peardon
Sinead Ryan

TATA, MUMBAI

Nilmani Mathur

CAMBRIDGE UNIVERSITY

Christopher Thomas

U. OF MARYLAND

Steve Wallace

& postdocs, students

MESON SPECTRUM

PRL103 262001 (2009) $I = 1$
PRD82 034508 (2010) $I = 1, K^*$
PRD83 111502 (2011) $I = 0$
JHEP07 126 (2011) $c\bar{c}$
PRD88 094505 (2013) $I = 0$
JHEP05 021 (2013) D, D_s

BARYON SPECTRUM

PRD84 074508 (2011) $(N, \Delta)^*$
PRD85 054016 (2012) $(N, \Delta)_{\text{hyb}}$
PRD87 054506 (2013) $(N \dots \Xi)^*$
PRD90 074504 (2014) Ω_{ccc}^*
arXiv:1502.01845 Ξ_{cc}^*

HADRON SCATTERING

PRD83 071504 (2011) $\pi\pi I = 2$
PRD86 034031 (2012) $\pi\pi I = 2$
PRD87 034505 (2013) $\pi\pi I = 1, \rho$
PRL113 182001 (2014) $\pi K, \eta K$
PRD91 054008 (2015) $\pi K, \eta K$

“TECHNOLOGY”

PRD79 034502 (2009) lattices
PRD80 054506 (2009) distillation
PRD85 014507 (2012) $\vec{p} > 0$

MATRIX ELEMENTS

arXiv:1501.07457 $M' \rightarrow \gamma M$
PRD90 014511 (2014) f_{π^*}

Summary

- LQCD spectroscopy program maturing. First phase:
 - With only “single-hadron” operators obtain sketch of hadron spectrum
 - Suggests rich spectrum of mesons & baryons - exotic & non-exotic hybrids

Summary

- LQCD spectroscopy program maturing. First phase:
 - With only “single-hadron” operators obtain sketch of hadron spectrum
 - Suggests rich spectrum of mesons & baryons - exotic & non-exotic hybrids
- *Goal is to compute resonance information - decays & branching fractions*
 - Including multi-hadron operators leads to richer spectrum
 - Demonstrated viability of finite-volume methods
 - Work underway at lower pion masses (230 MeV)
 - S-matrix formalism increasingly important

E.G., JOINT-PHYSICS ANALYSIS CENTER

Summary

- LQCD spectroscopy program maturing. First phase:
 - With only “single-hadron” operators obtain sketch of hadron spectrum
 - Suggests rich spectrum of mesons & baryons - exotic & non-exotic hybrids
- *Goal is to compute resonance information - decays & branching fractions*
 - Including multi-hadron operators leads to richer spectrum
 - Demonstrated viability of finite-volume methods
 - Work underway at lower pion masses (230 MeV)
 - S-matrix formalism increasingly important
- Ultimately, determine underlying structure
 - Resonance scattering with external currents
 - Both vacuum overlaps (“wave-function”) & transition form-factors

E.G., JOINT-PHYSICS ANALYSIS CENTER

SHULTZ 2015

BRICENO & HANSEN 2015