

Clustering aspects in neutron-rich nuclei studied by AMD

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Introduction: Island of Inversion

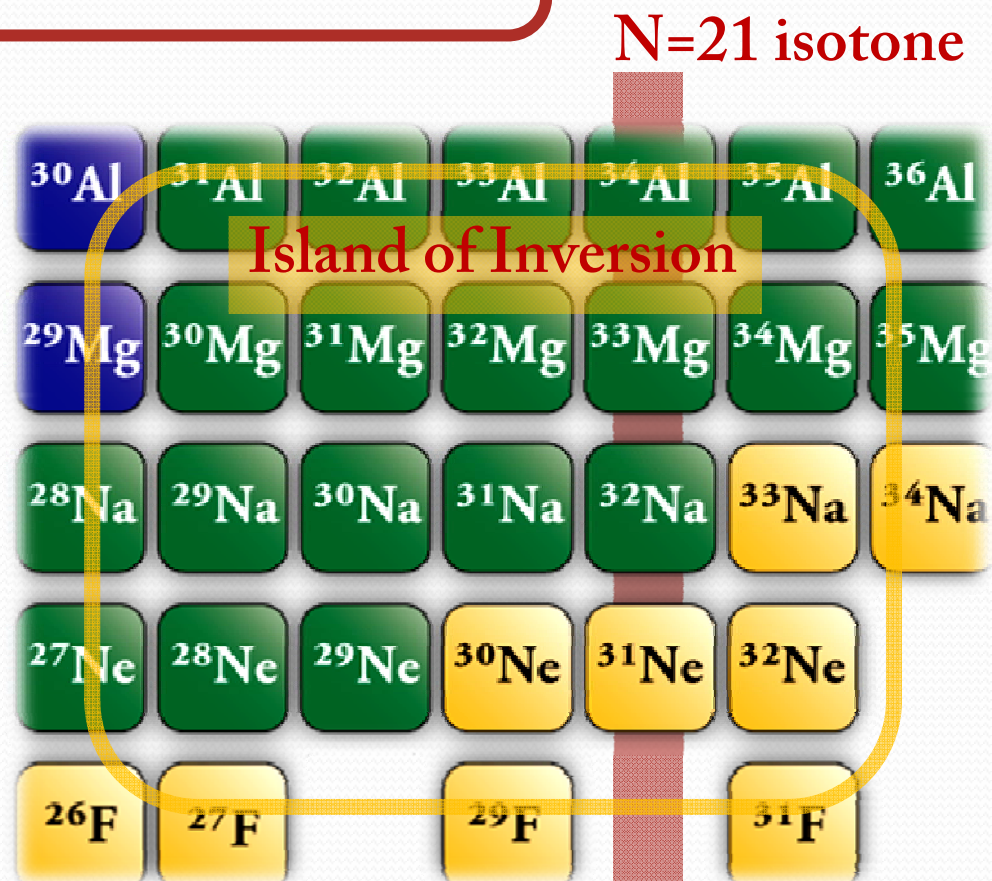
Island of Inversion: $N \sim 20$ Ne, Na, Mg isotopes

Topics: large deformation, mpmh config., shell gap quench

N=21 isotone: ^{31}Ne (binding limit), ^{33}Mg

Topics: spin-parity, p-wave halo, neutron orbital

- Odd-mass(neutron) system reveals neutron-orbital, shape coexistence in the Island
- The last neutron particle (hole) orbital has strong influence on nuclear properties
- In N=21 isotone, $0f_{7/2}$ and $0d_{3/2}$ orbitals take part in the game



Theoretical Framework: AMD

A-body Hamiltonian

$$\hat{H} = \sum_{i=1}^A \hat{t}_i - \hat{t}_{c.m.} + \sum_{i < j} \hat{v}_{ij}^{NN} + \sum_{i < j \in P} \hat{v}_{ij}^{Coulomb}, \quad \hat{v}^{NN} : \text{Gogny D1S}$$

Parity projected Slater determinant

$$\Psi_{\text{int}}^{\pm} = \frac{1 \pm \hat{P}_x}{2} \mathcal{A}\{\varphi_1, \varphi_2, \dots, \varphi_A\},$$

Single particle wave packets

$$\begin{aligned} \varphi_i &= \phi(\mathbf{r}) \chi_i \eta_i, & \phi(\mathbf{r}) &= \exp\{-(\mathbf{r} - \mathbf{Z}_i)M(\mathbf{r} - \mathbf{Z}_i)\}, \\ & & \chi_i &= \alpha_i \chi_{\uparrow} + \beta_i \chi_{\downarrow}, \end{aligned}$$

Variational parameters

$\mathbf{Z}_i$: centroid of wave packet

M : deformation of wave packet

α_i, β_i : spin direction

Theoretical Framework: AMD

Angular momentum projection

$$\Psi_{MK}^{J\pm} = \int d\Omega D_{MK}^{J*}(\Omega) \hat{R}(\Omega) \Psi_{\text{int}}^{\pm},$$

GCM

$$\Psi_{MK}^{J\pm}(\beta_1), \Psi_{MK}^{J\pm}(\beta_2), \dots, \Psi_{MK}^{J\pm}(\beta_N)$$

Generator Coordinate: quadrupole deformation β

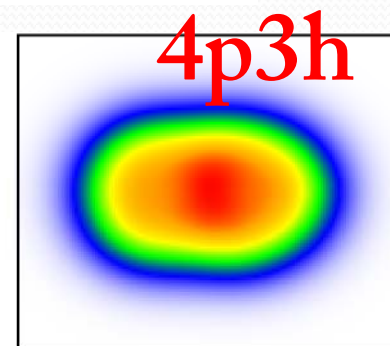
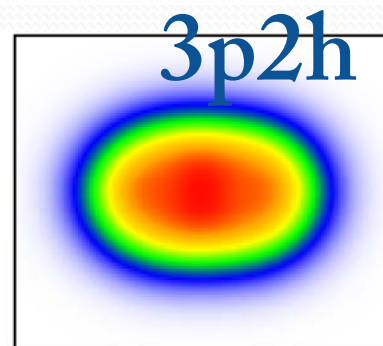
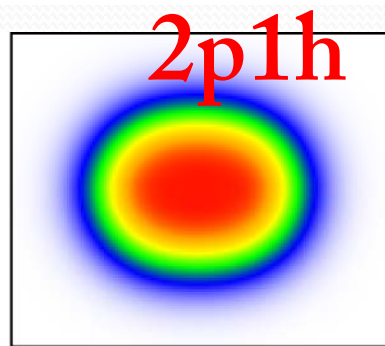
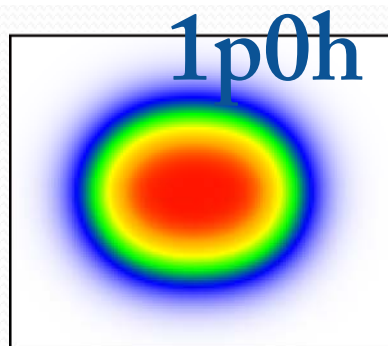
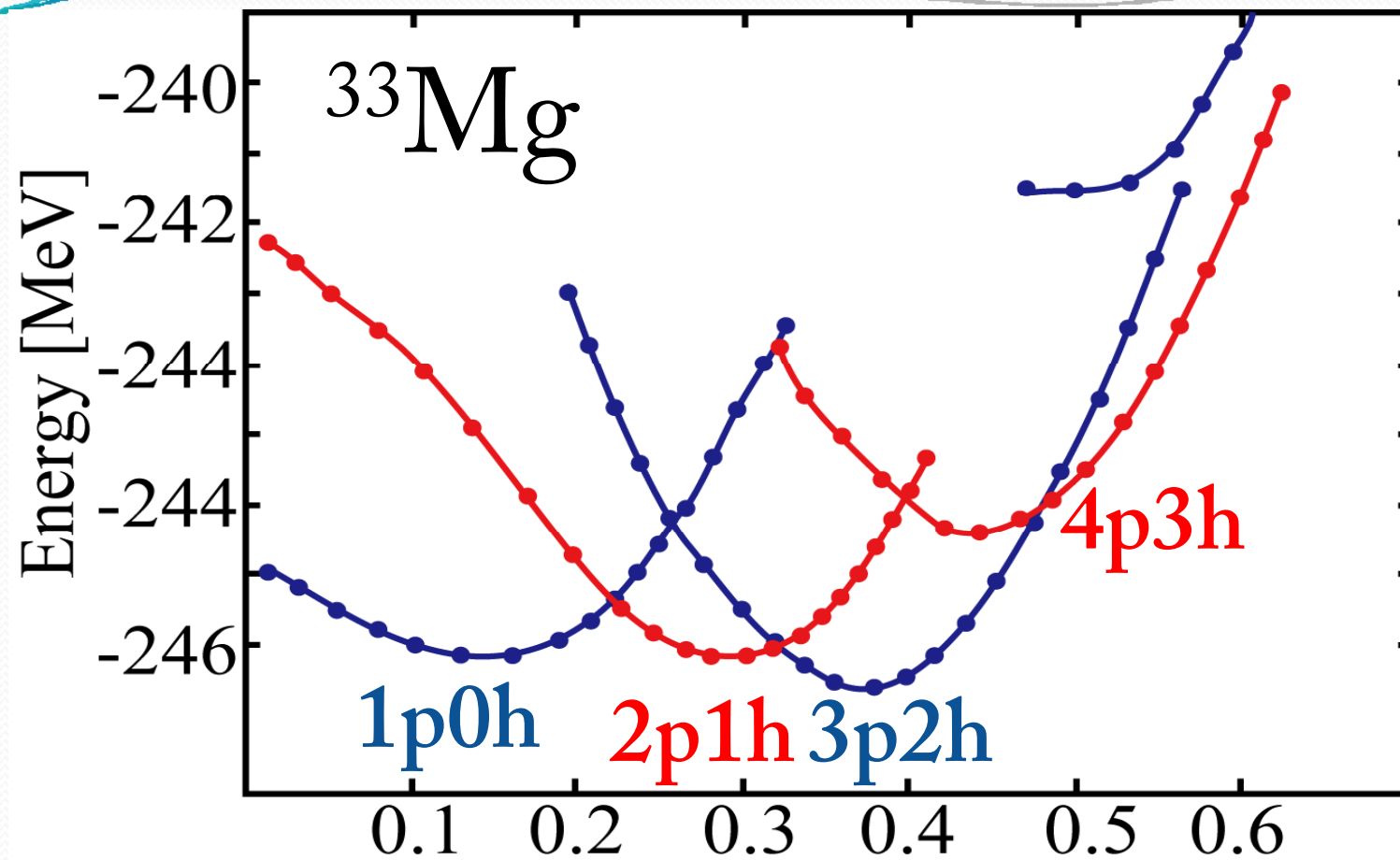
$$\Psi_{\alpha}^{J\pm} = \sum_{iK} c_{ik} \Psi_{MK}^{J\pm}(\beta_i),$$

Hill-Wheeler eq.

$$\sum_{jK'} H_{iKjK'} c_{jK',\alpha} = E_{\alpha} \sum_{jK'} N_{iKjK'} c_{jK',\alpha},$$

$$H_{iKjK'} = \langle \Psi_{MK}^{J\pm}(\beta_i) | \hat{H} | \Psi_{MK'}^{J\pm}(\beta_j) \rangle, \quad N_{iKjK'} = \langle \Psi_{MK}^{J\pm}(\beta_i) | \Psi_{MK'}^{J\pm}(\beta_j) \rangle$$

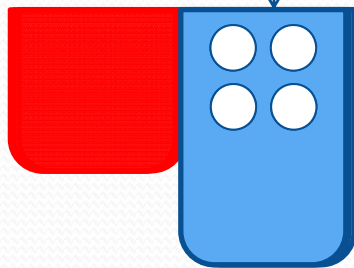
Results: Energy Curve and PH Config. of ^{33}Mg



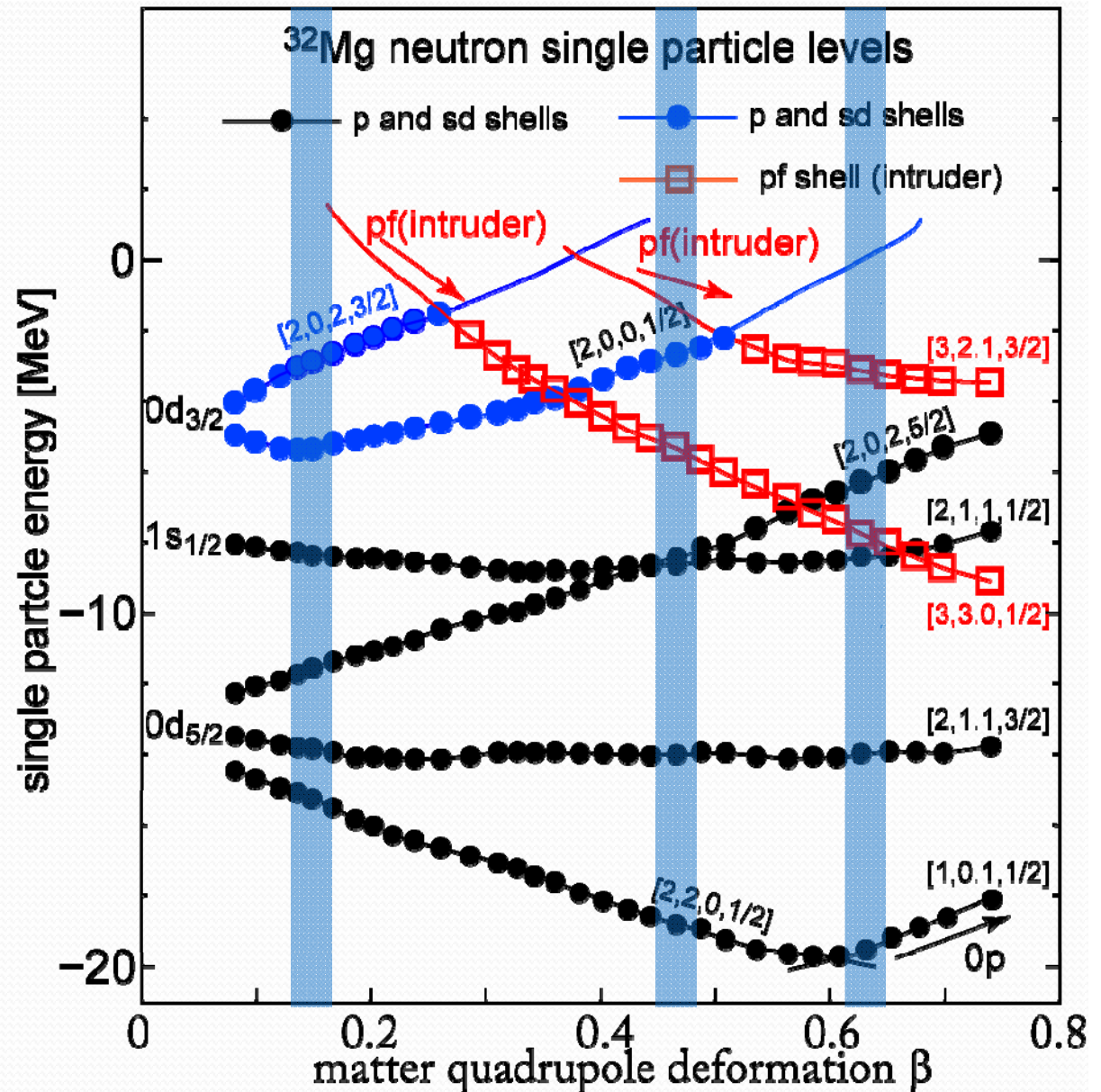
Consideration on the curves

0p0h (normal config.)

2p2h (intruder)
N=20 shell gap



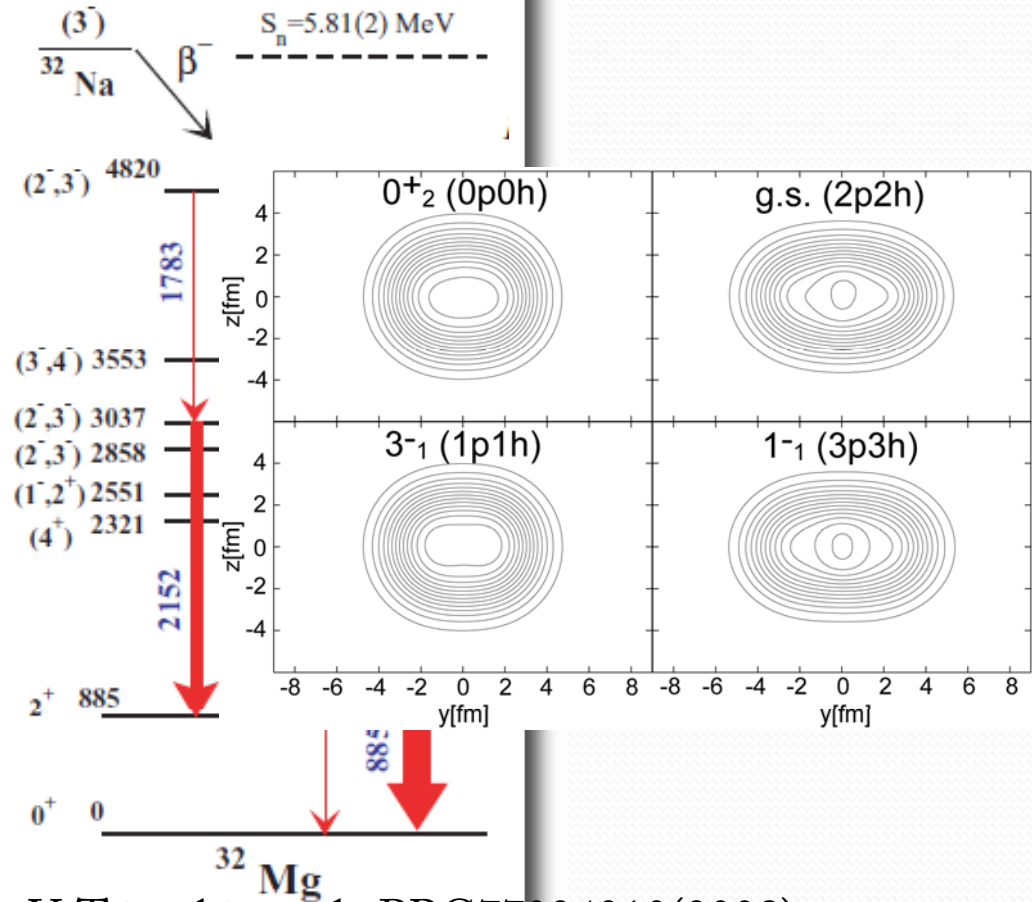
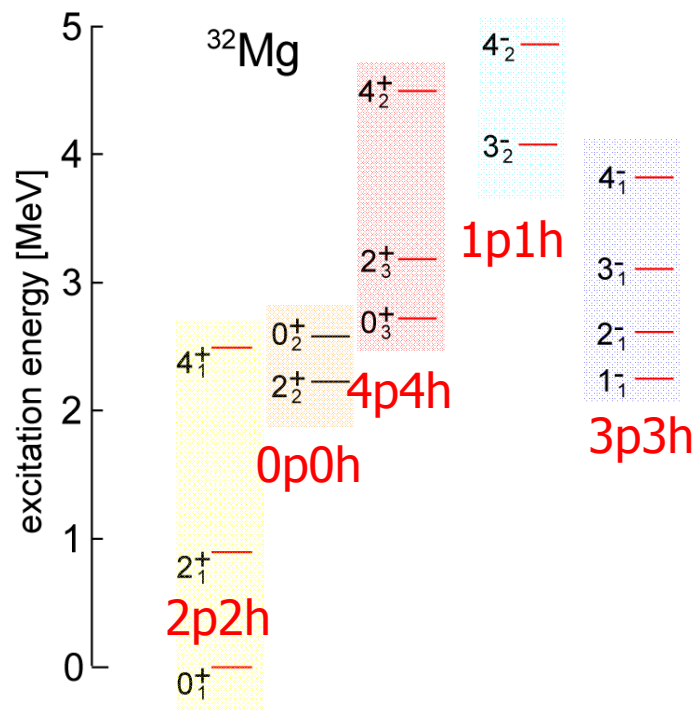
- 2 orbits of sd and 2 of pf shell participate in this game
- From this one can deduce a naive rule for ph configurations



Results: Behavior of negative parity states

level scheme of

DAMD+GCM

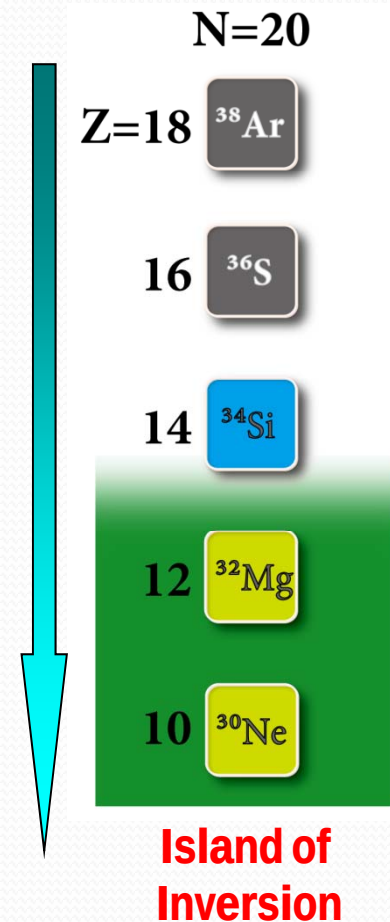
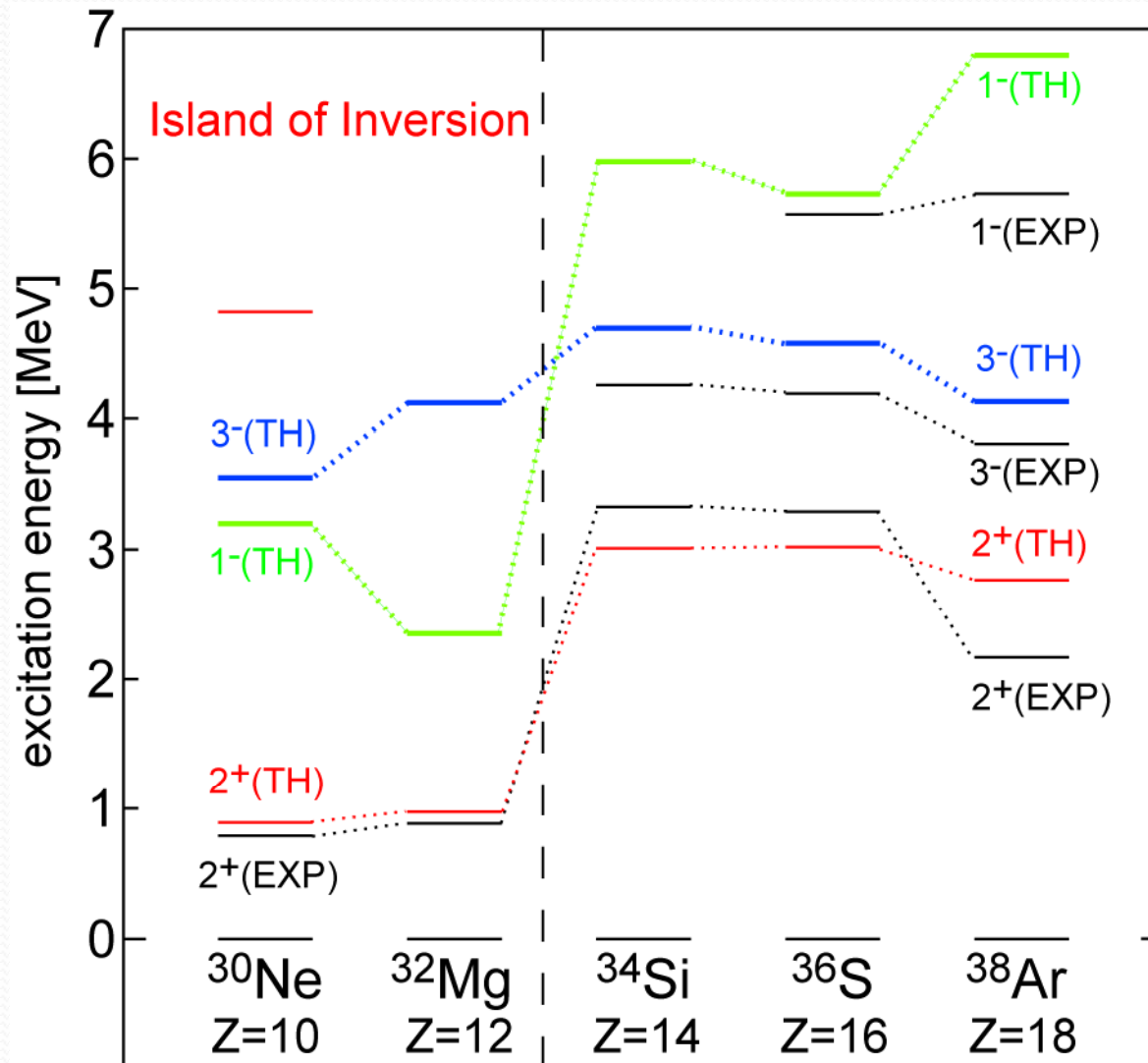


V. Tripathi, et. al., PRC77034310(2008).

- $1^-(3p3h)$ state is built on the deformed g.s. ($1p1h$ from $2p2h$)
 - $3^-(1p1h)$ state is built on the spherical 0^+ state ($1p1h$ from $0p0h$)
- Low lying 1^- state is a good signal of the inversion.

Results: Behavior of negative parity states

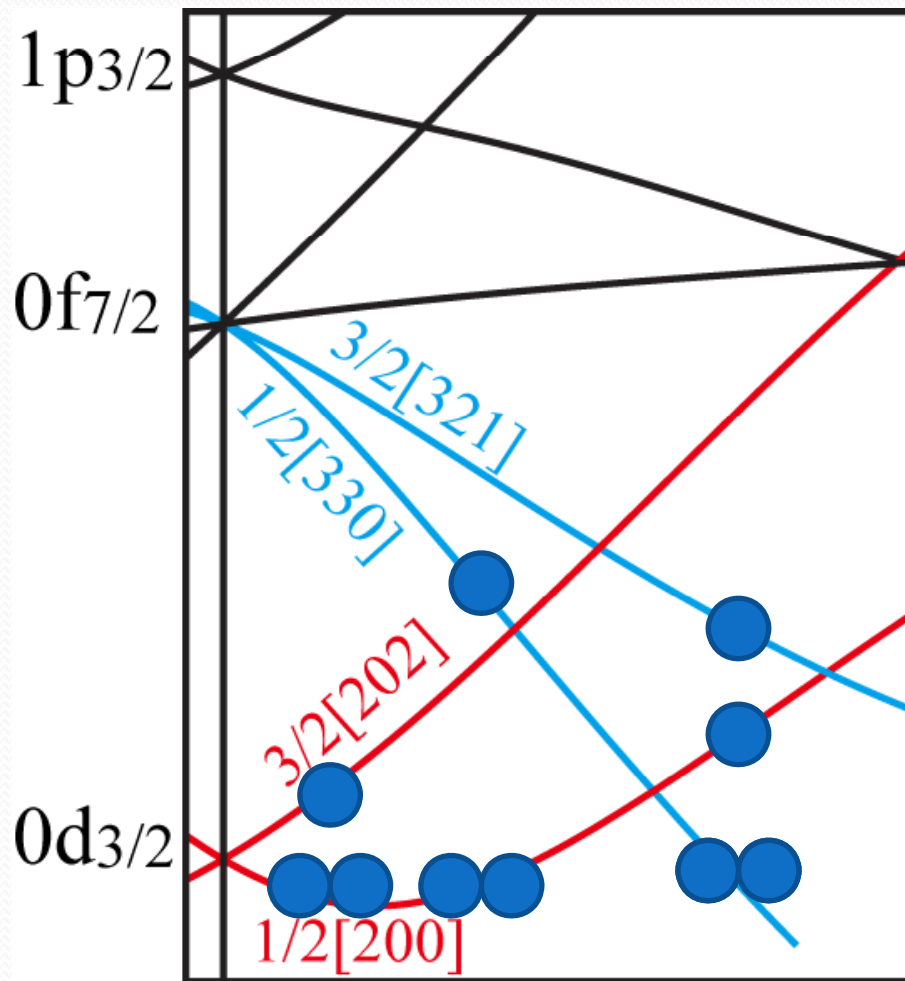
Case of the N=20 isotone



Introduction: Particle Hole Config. (N=19 System)

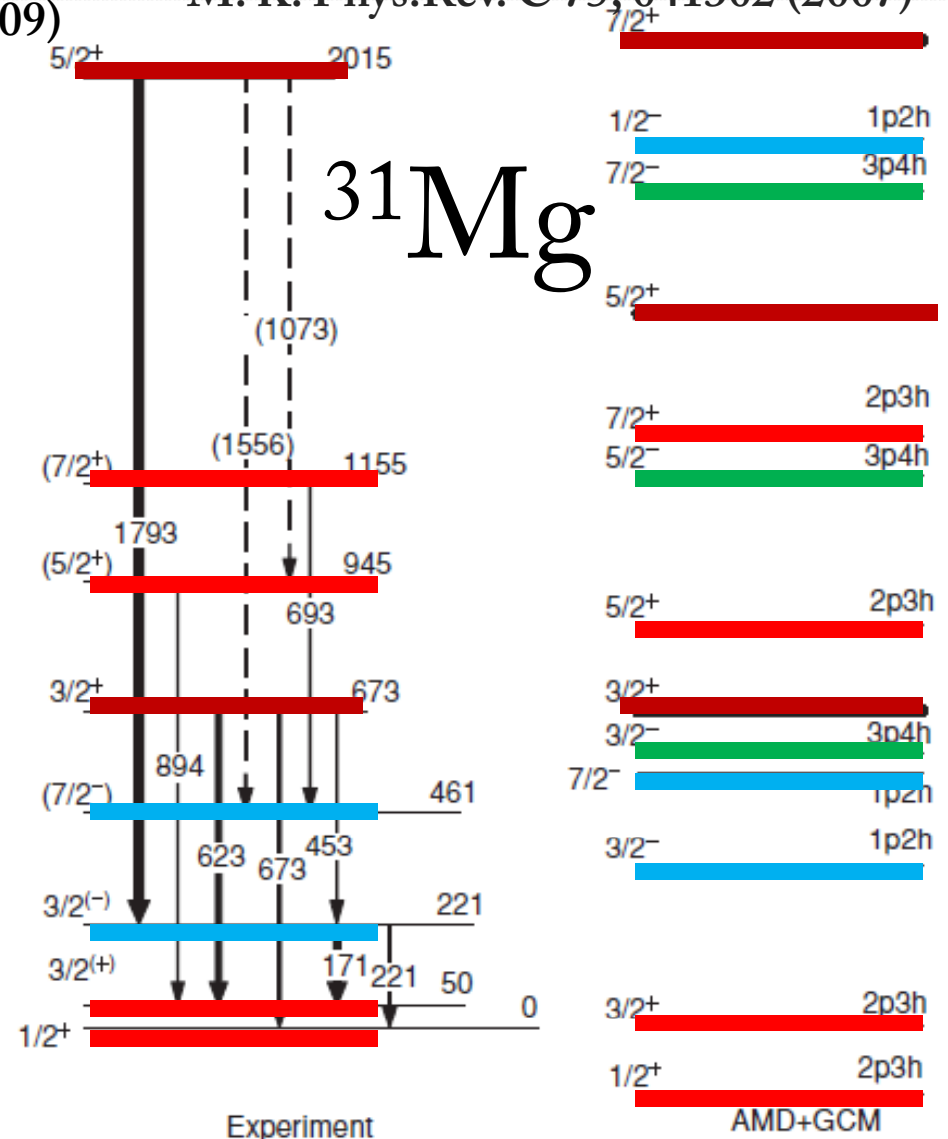
1p-removal from ^{32}Al

D. Miller, et.al., Phys. Rev. C 79, 054306 (2009)

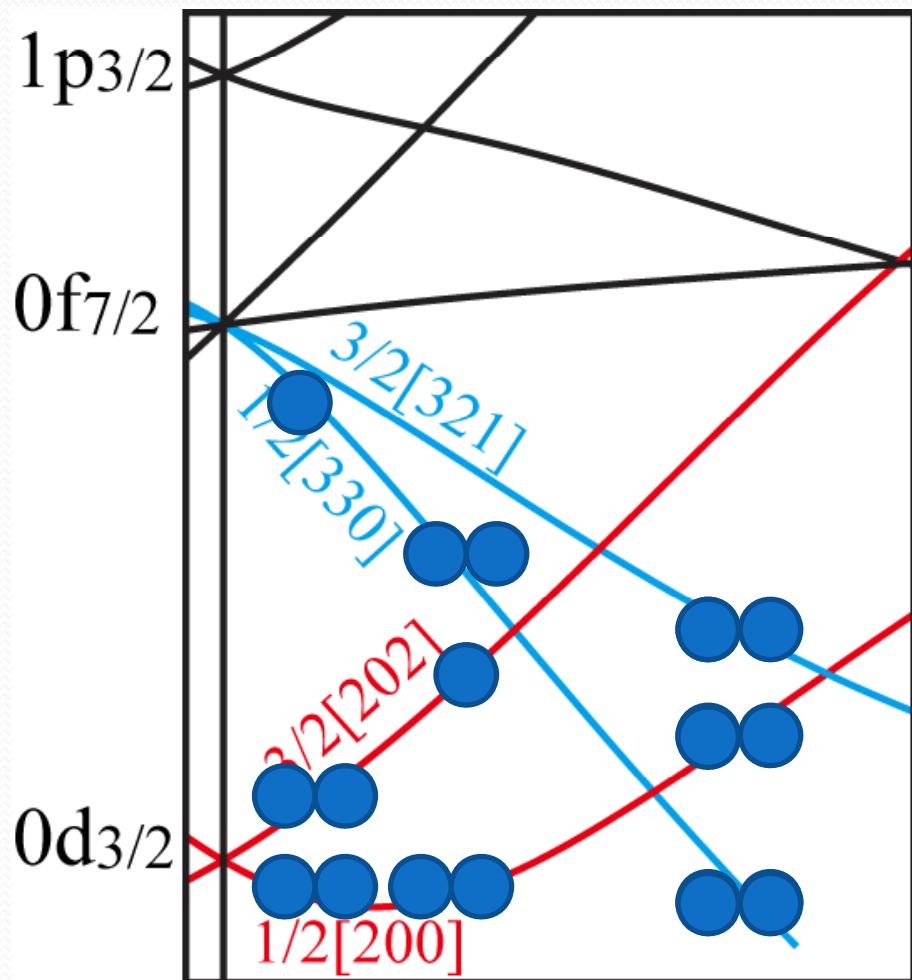


0p 1h 2h 3h 4h

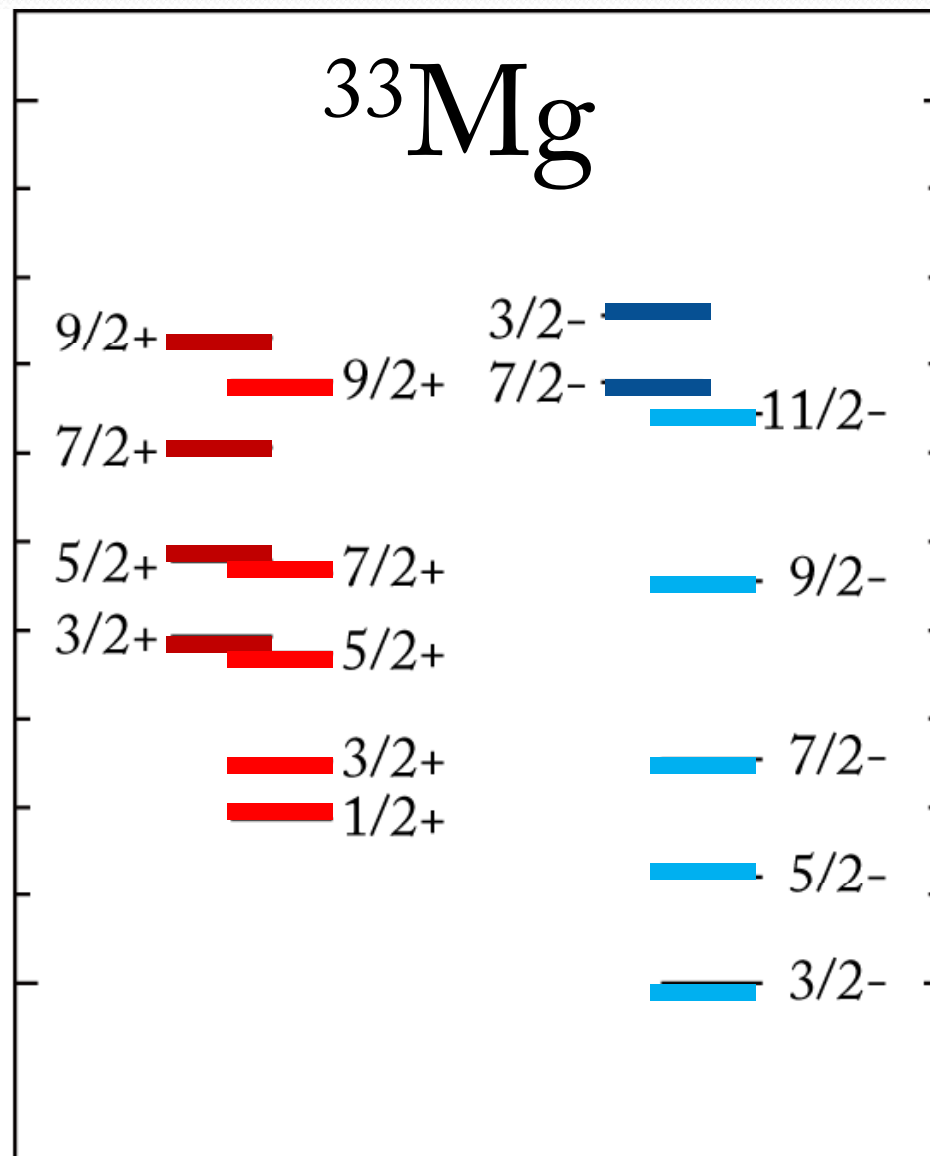
M. K. Phys. Rev. C 75, 041302 (2007)



Introduction: Particle Hole Config. (N=21 System)



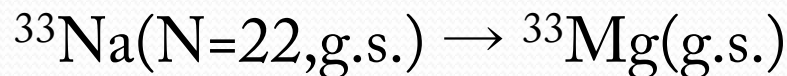
1p0 2p1h 3p2h 4p3h



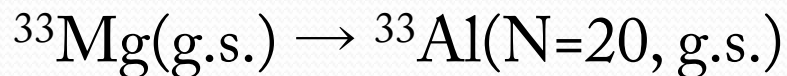
Introduction: Spin-parity of $^{33}\text{Mg}(N=21)$

- $J^\pi = (3/2, 5/2)^+$

β -decays



S. Nummela, et. al., PRC64, 054313 (2001)



V. Tripathi, et. al., PRL101,142504 (2008)

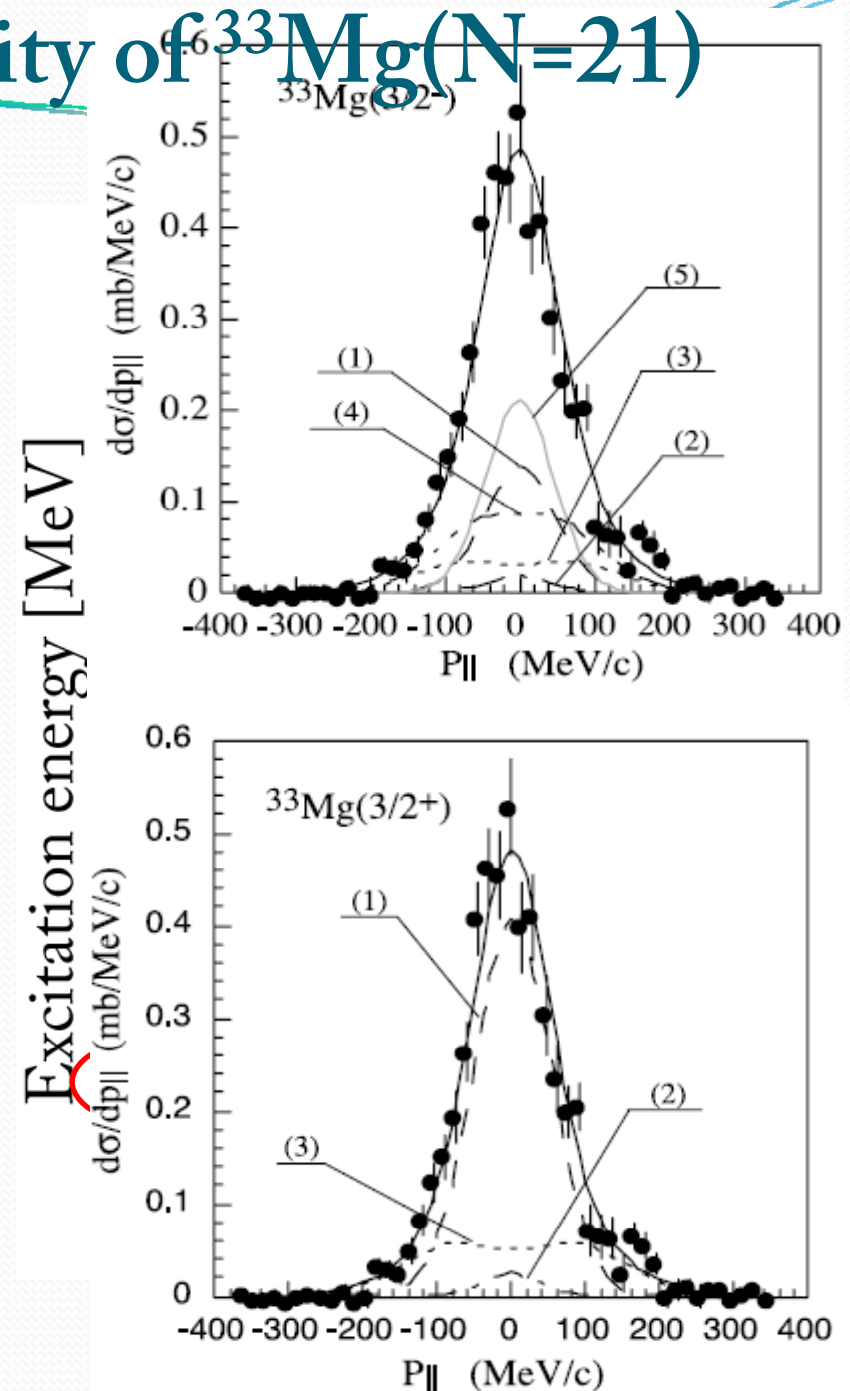
- $J^\pi = 3/2^-$

magnetic-moment: $\mu = -0.7456\mu_N$

D. T. Yordarov, et. al., PRL99, 212501 (2007)

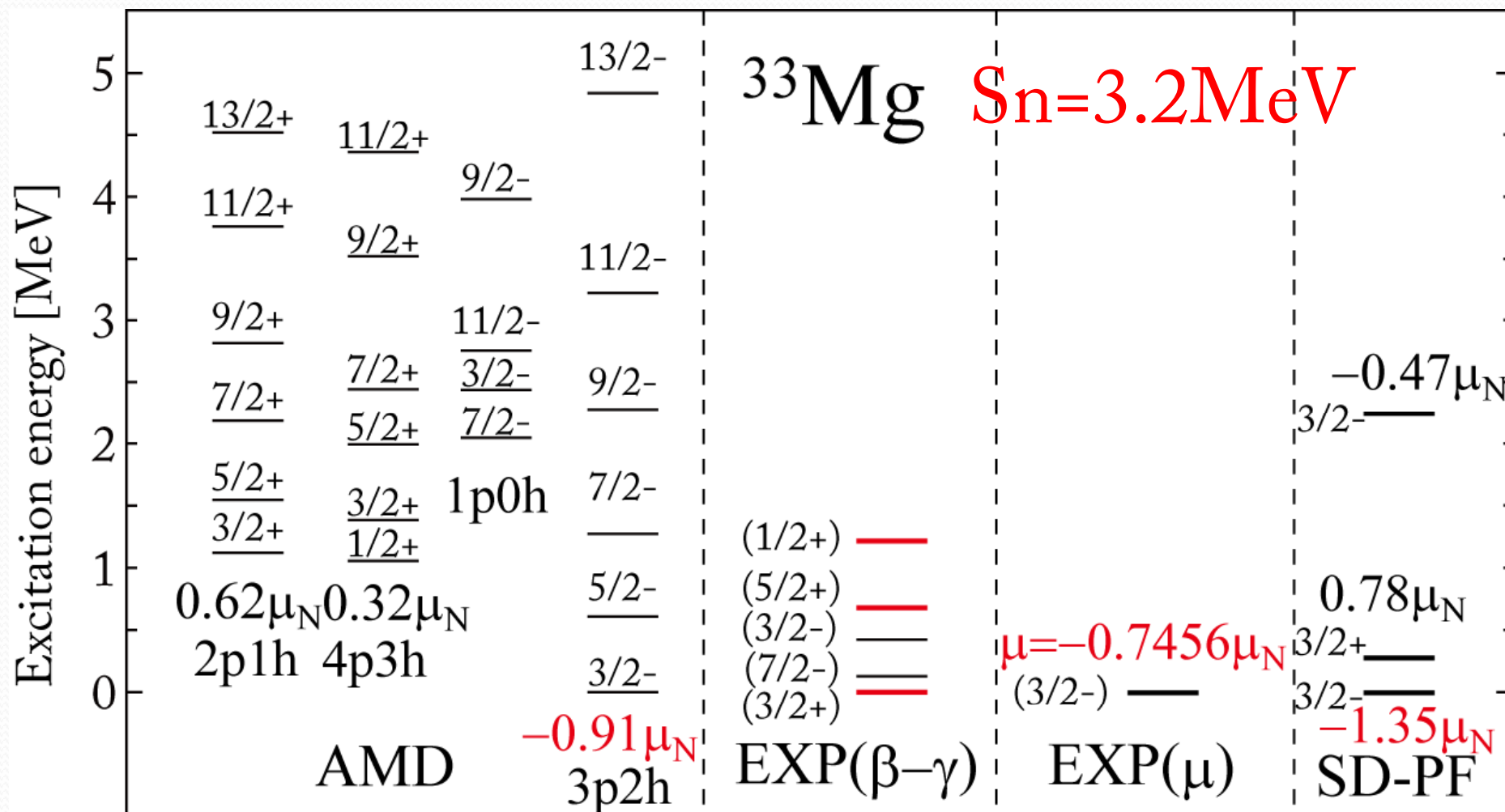
1n-removal: $p_{3/2}$

R. Kanungo, et. Al., PLB685, 253 (2010)



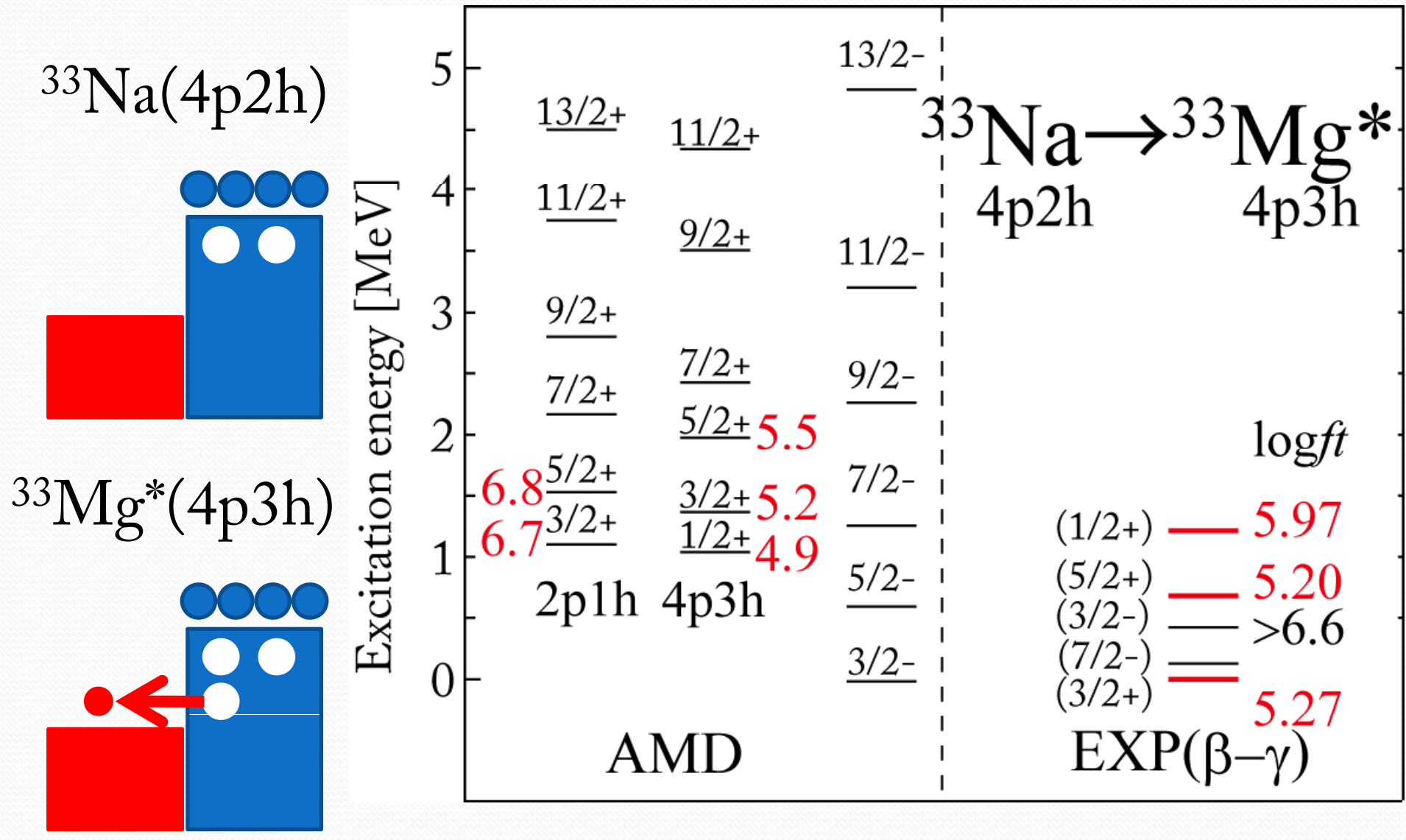
Results: Spectrum of ^{33}Mg

- Consistent with magnetic moment exp.
- AMD underestimates number of very low-lying states



Results: β -decay and $J\pi$ assignment

- $^{33}\text{Na}(4p2h)$ feeds $4p3h$ config.

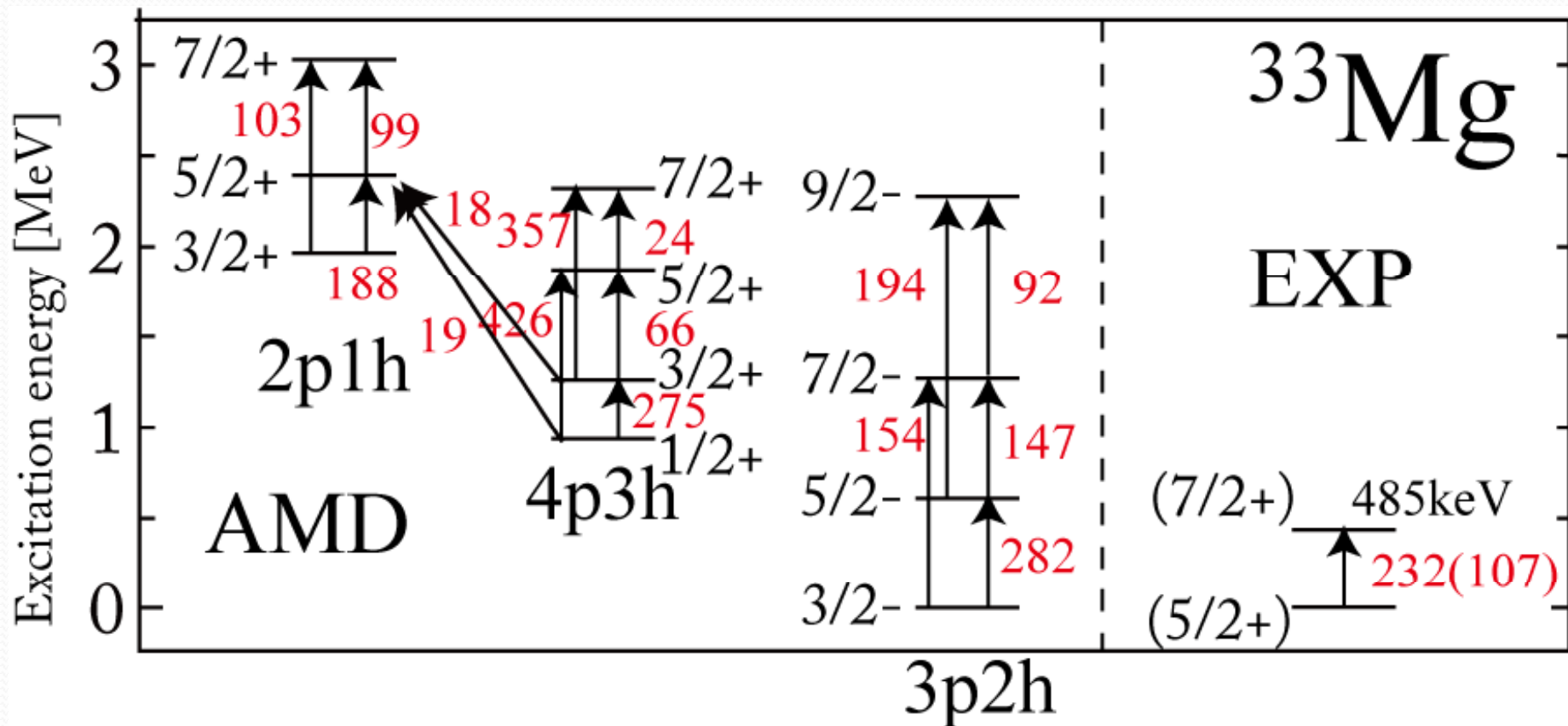


Results: B(E2) of ^{33}Mg

- **Coulomb Excitation Exp. :** B.V. Pritychenko, et. al., PRC65, 061304(R) (2002).

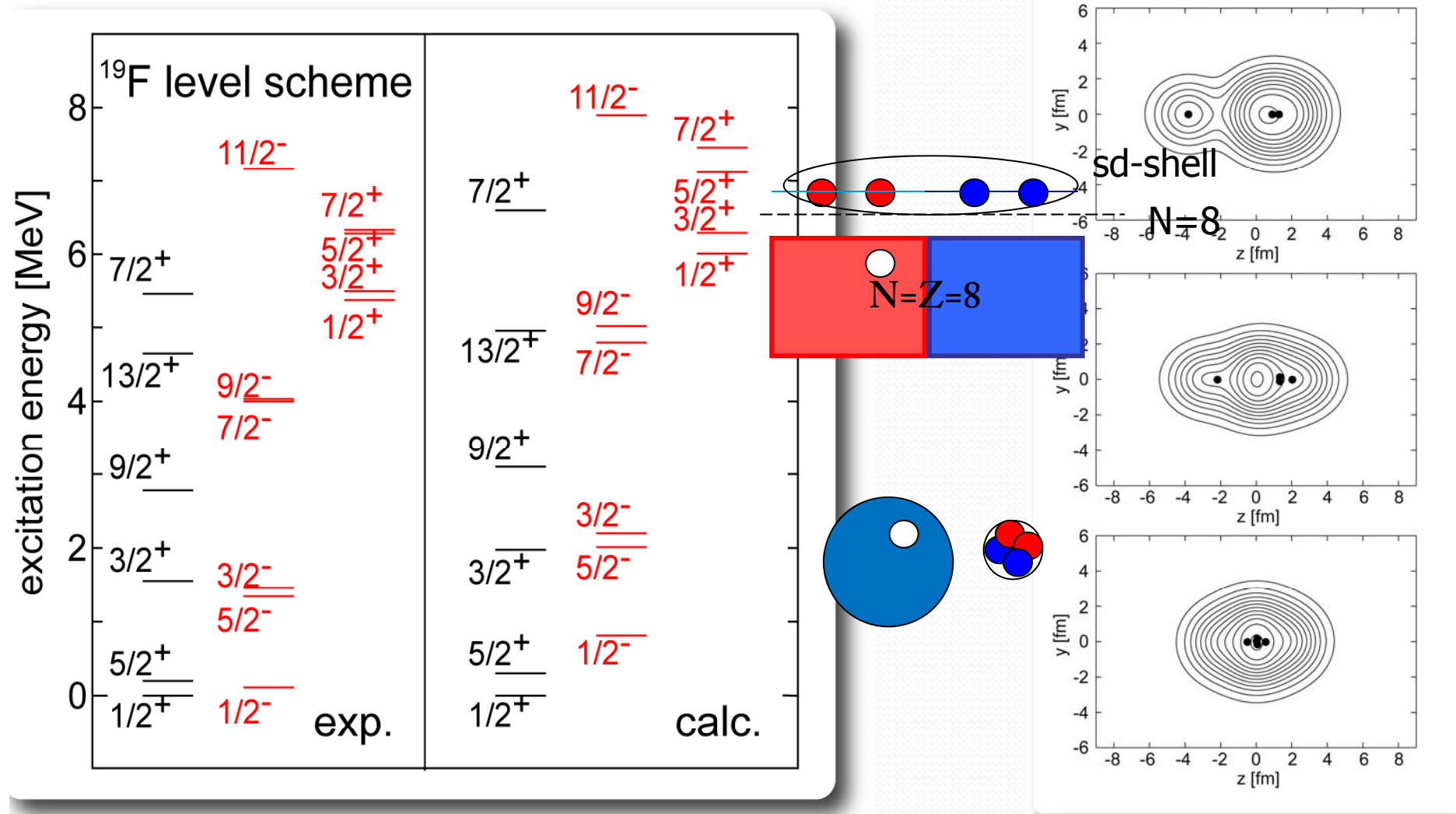
Large cross section for 485keV state

- Authors assumed $J\pi=5/2+$ for the ground state (no information was available at that time except for β -decay data that suggested positive parity)
- Large cross section does not contradict to AMD (Further exp. is required.)



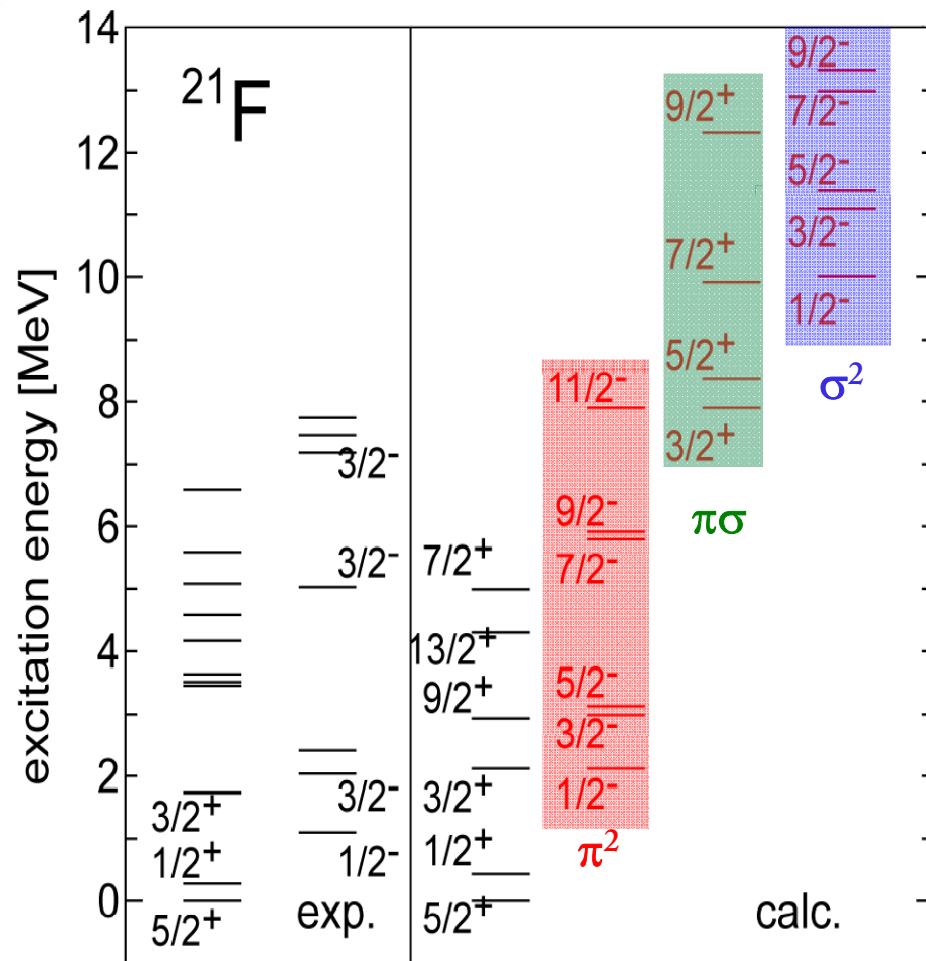
Results: alpha cluster structure in ^{19}F

Key point: Alpha cluster states of ^{19}F
has a proton hole in 0p orbital.

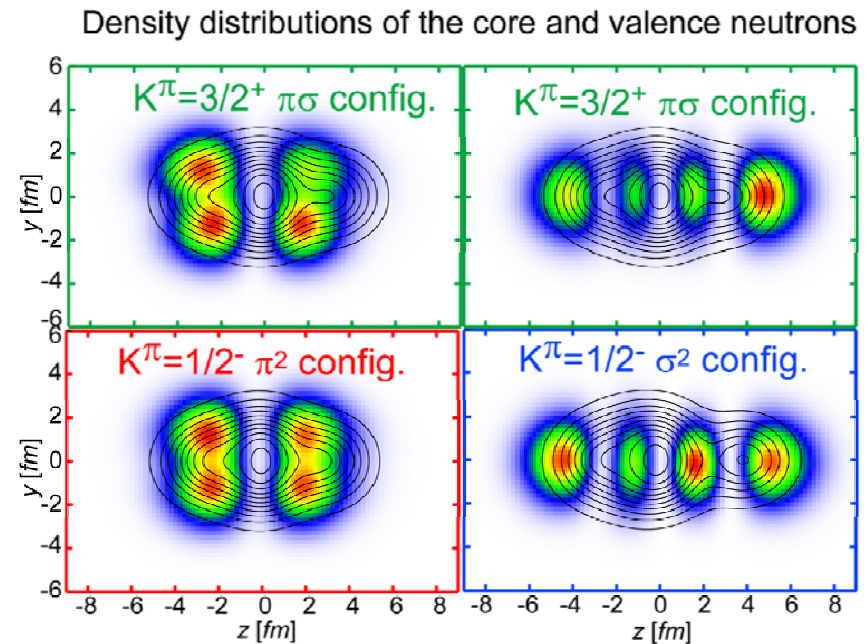


Results: Molecular-orbitals in ^{21}F

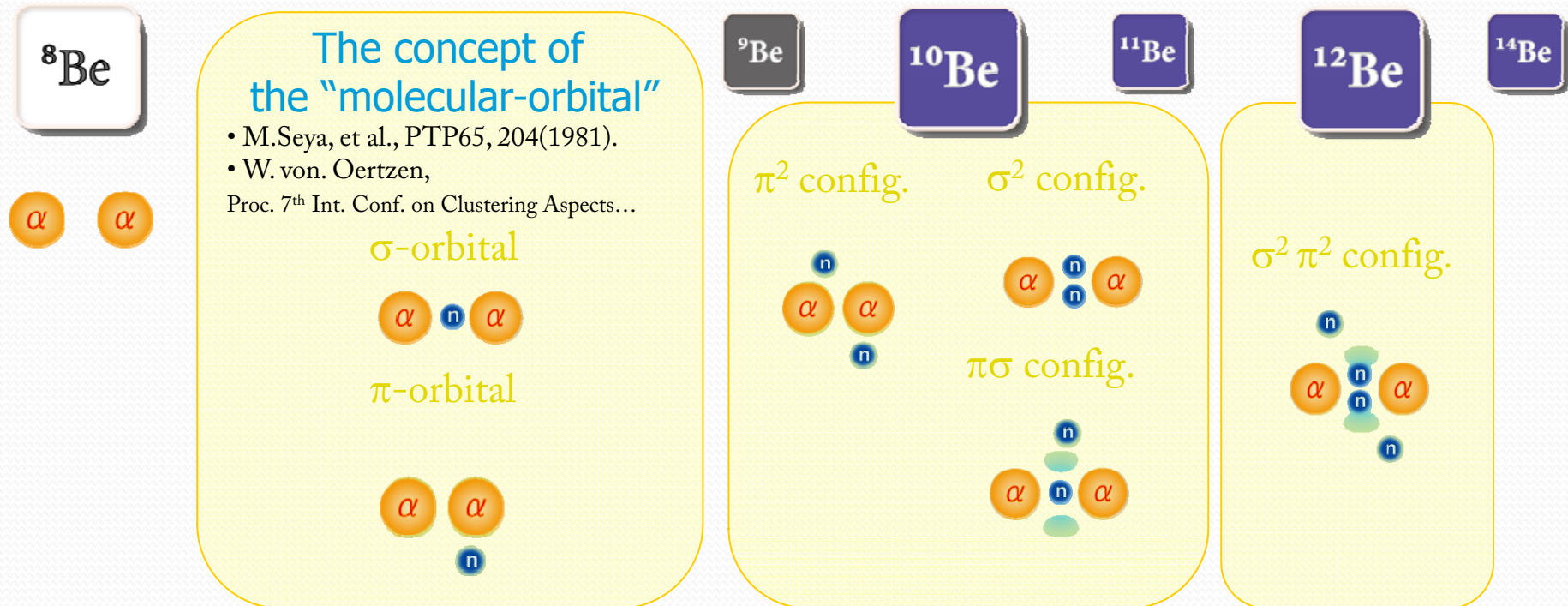
M. Kimura, submitted to PRC.



- proton hole states has the α cluster structure
- π^2 , $\pi\sigma$ and σ^2 configurations appear as $K^\pi=1/2^-$, $3/2^+$ and $1/2^-$ bands



Results: Analogue to molecular-orbitals of Be



- M. Seya, et al., PTP65, 204 (1981).
- Y. Kanada-En'yo, et al., PRC60, 064304(1999).
- N. Itagaki, et al., PRC62 034301, (2000). ...

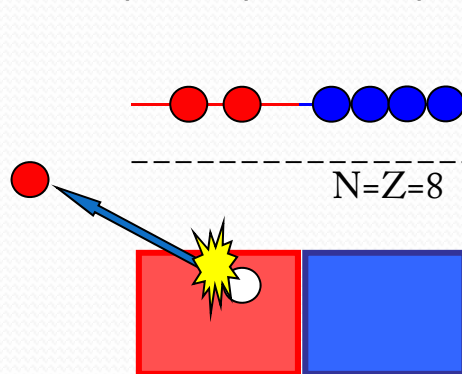
Molecular Structure in Be isotopes suggests us new aspects specific to $N \neq Z$ nuclei.

- A new binding mechanism.
- A new type of cluster structure.

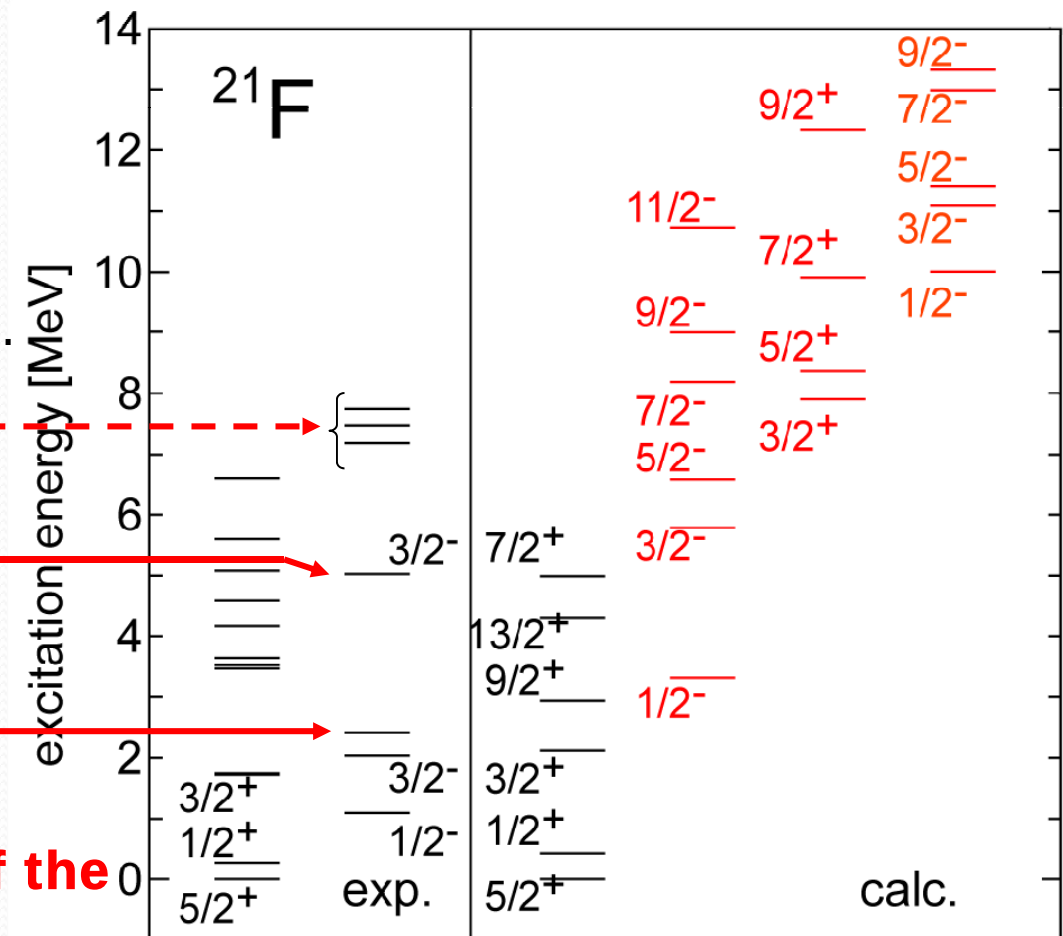
Results: Molecular-orbitals in ^{21}F

**Populated directly by only
1p knockout from ^{22}Ne**

- G. Th. Kaschl, NPA155, 417 (1970).
- G. Mairle, et al., NPA363, 413 (1981).

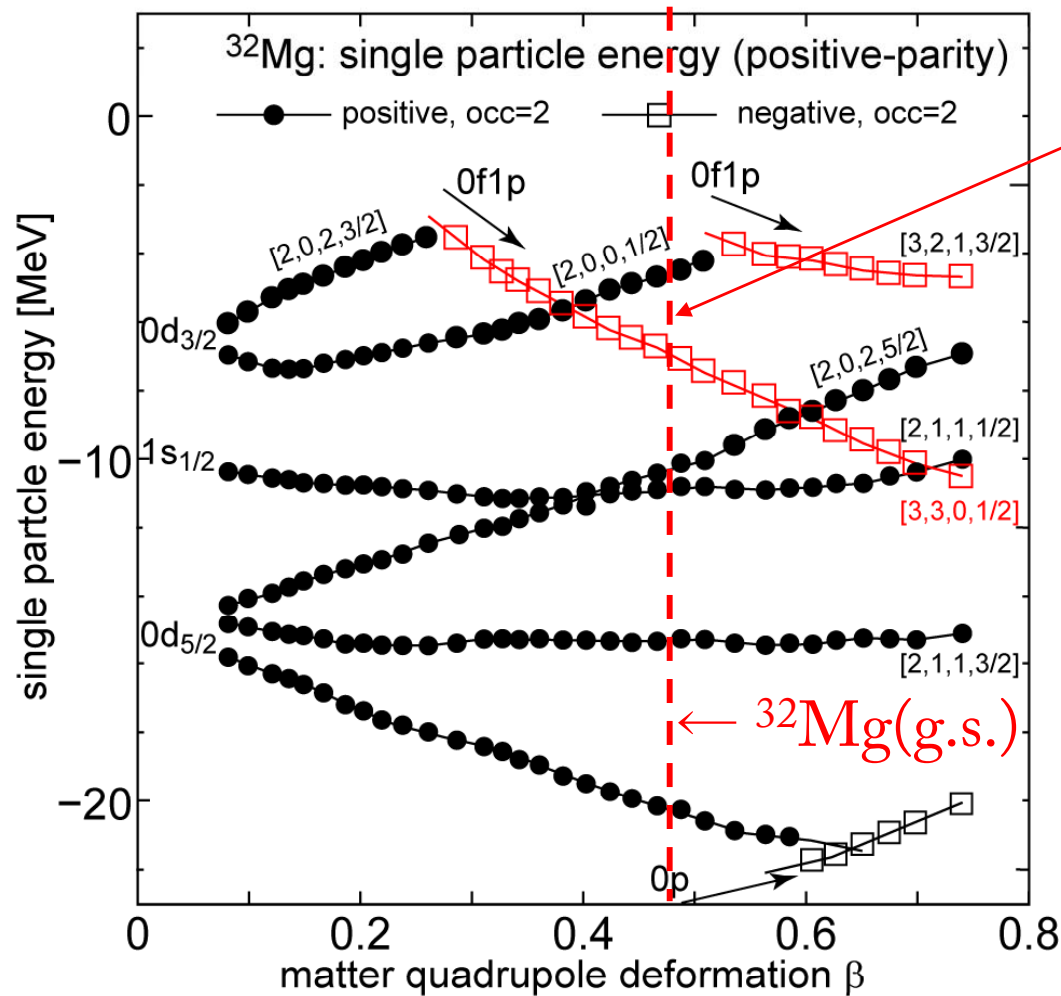


They are the candidates of the molecular states



Results: Evolution toward drip-line

Behavior of the α cluster states in E isotones



Intruder orbital from pf-shell has large overlap with σ -orbital

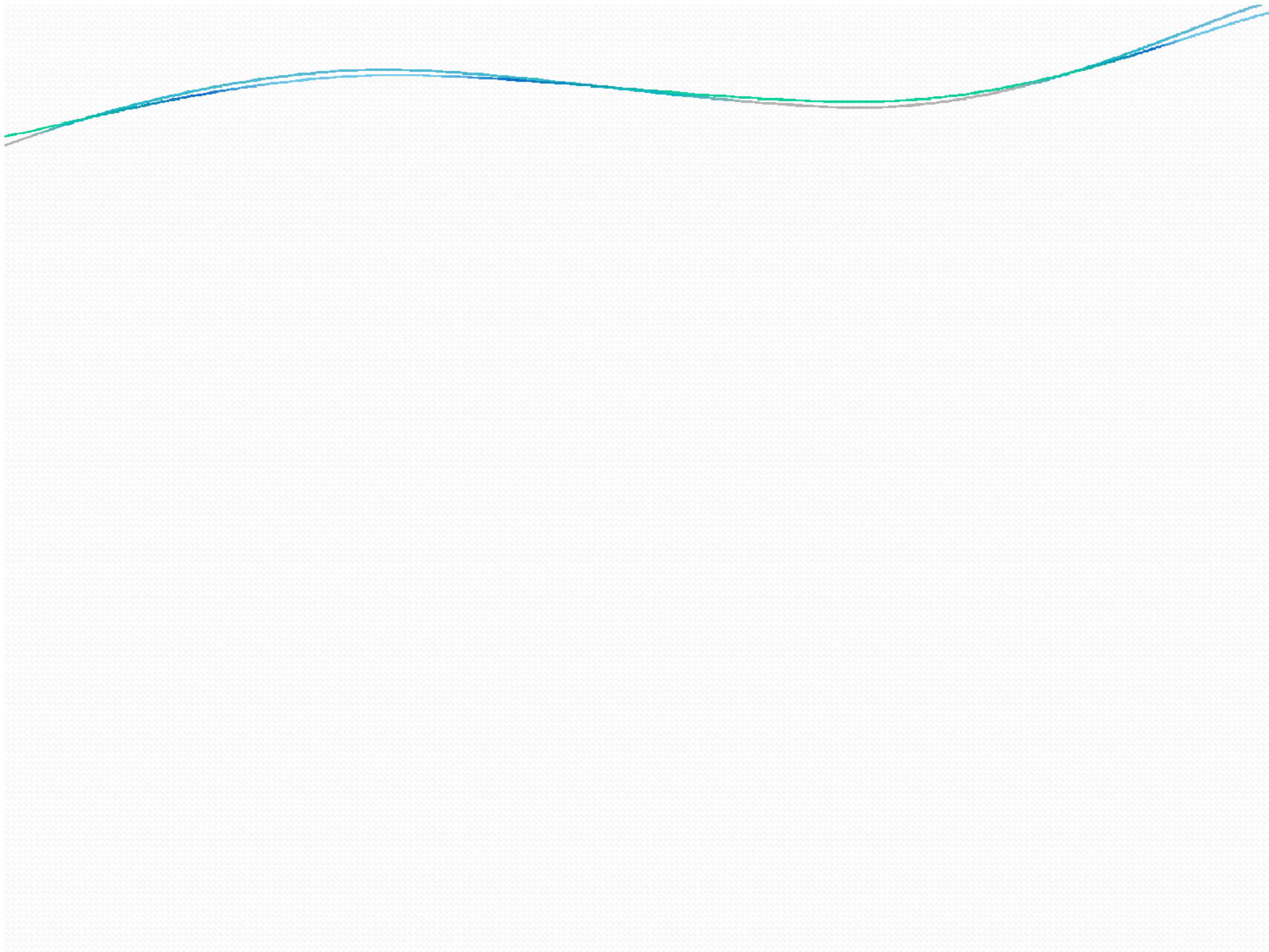
The origin of the reduction of ex. energy

1. proton excitation from p to sd shell does not cost much energy due to strong α correlation.
2. It induces the strong deformation of system.
3. Strong deformation lowers σ -orbital



Molecular states with valence neutron in σ -orbital comes down.

- α cluster states with neutron $0h_{9/2}$ configuration monotonously go up.



I. Introduction : ^{24}Mg の多様なクラスター構造

池田図:

^{24}Mg の高励起状態に多様なクラスター構造が存在することを示唆
 $(\alpha + ^{20}\text{Ne}, ^{12}\text{C} + ^{12}\text{C}, ^{16}\text{O} + 2\alpha, ^{12}\text{C} + 3\alpha, 6\alpha)$

高励起状態のクラスター構造:

- ^{24}Mg のクラスター状態は十分に明らかにはされていない。

(^{24}Mg の低励起状態では平均場構造が支配的)

- ^{24}Mg よりも軽い核のクラスター構造はよく理解されている。近年、 3α , 4α クラスター状態(α 凝縮状態)の理解が進展。

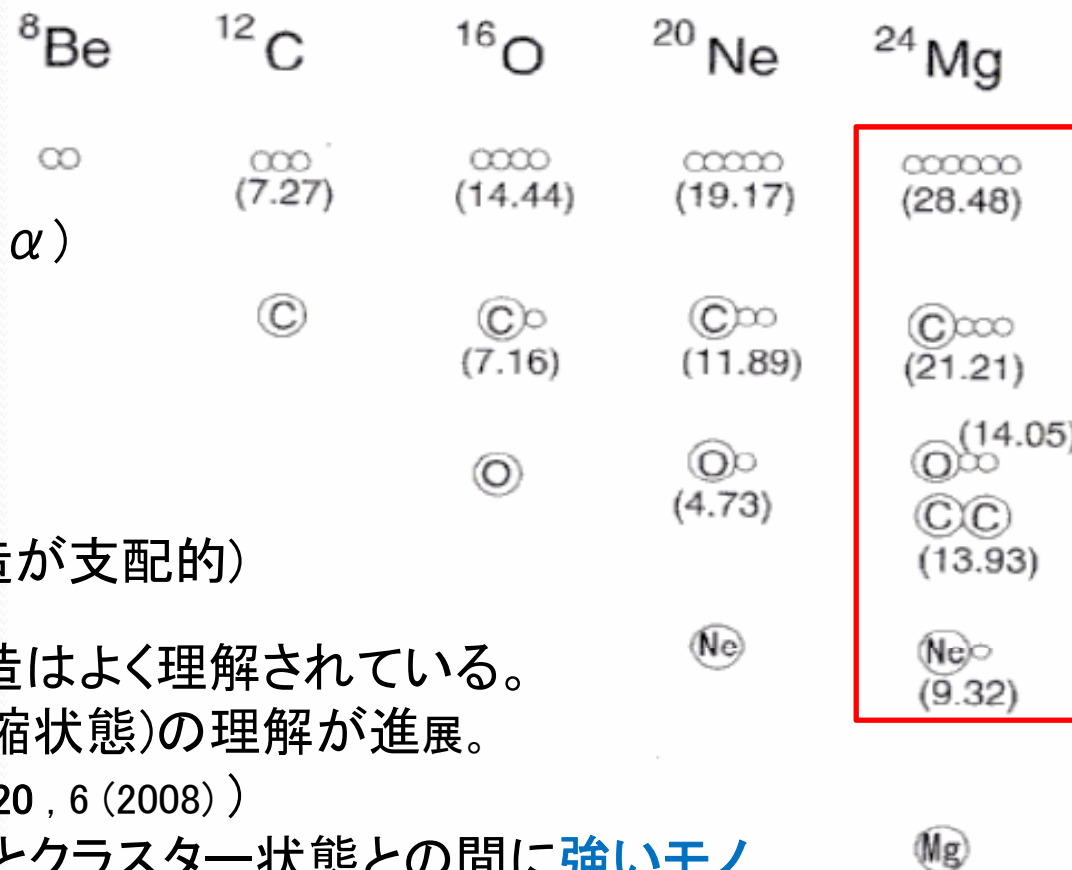
(Taiichi Yamada, Yasuro Funaki, et al PTP 120, 6 (2008))

- α 凝縮状態の研究では基底状態とクラスター状態との間に強いモノポール遷移があることがしめされた。

(Taiichi Yamada, Yasuro Funaki, et al PTP 120, 6 (2008))



^{24}Mg のクラスター状態を理解するための鍵となりうる。



• $^{24}\text{Mg}(\alpha, \alpha')$ 散乱のデータ

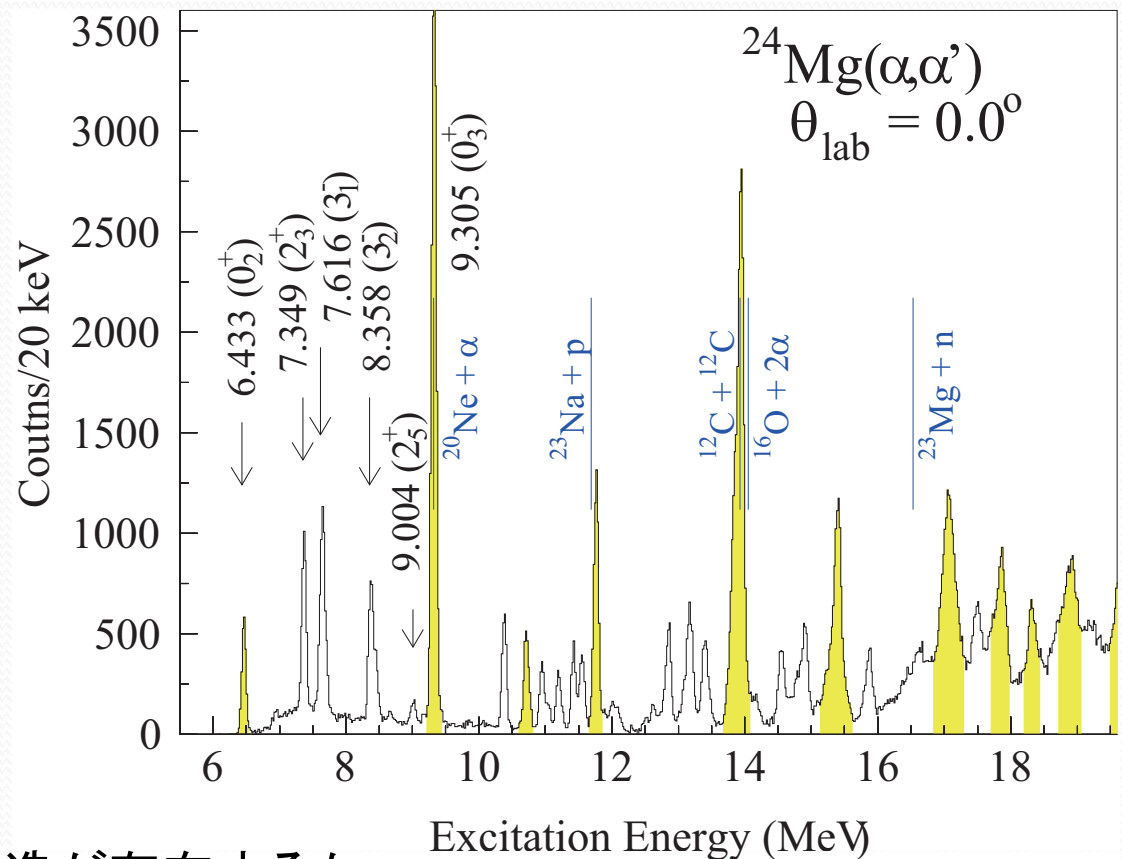
- 超前方の α 非弾性散乱(0^+ を選択的に励起)
- 10~20MeVの領域に強く励起される状態が多数見つかった。
- これらのピークは敷居値エネルギー近傍にあり、大きなモノポール遷移確率をもつ可能性が指摘されている。
- 観測されたこれらの状態は池田図に現れる多様なクラスター構造に対応しているのではないか。

[目的]

^{24}Mg の高励起状態にクラスター構造が存在するかどうか。どのようなクラスター構造が存在するか。

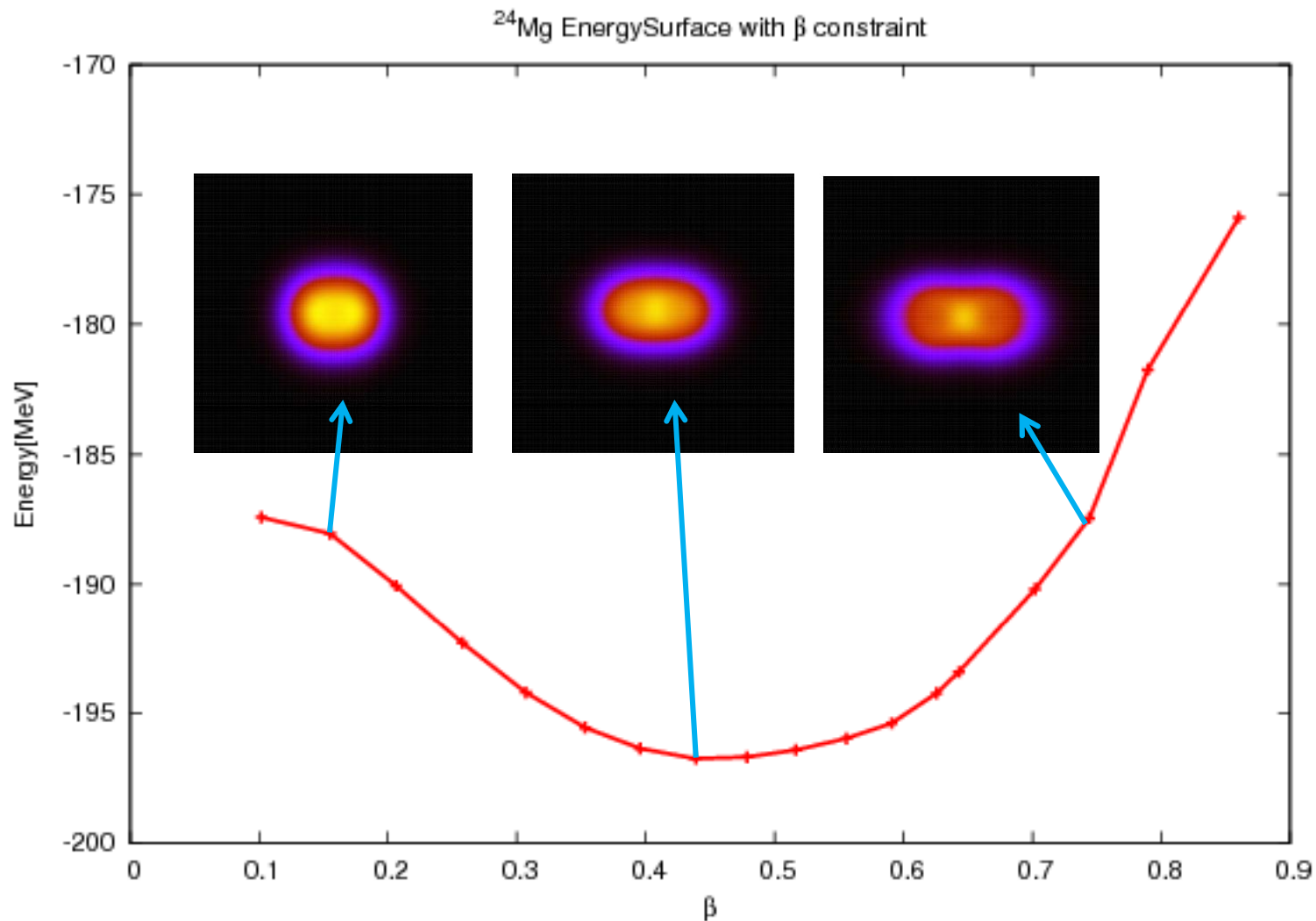
[方法]

平均場構造とクラスター構造を両方記述するためAMD模型を用い、 ^{24}Mg の高励起状態を調べる。



Takahiro Kawabata : ppt
Alpha Inelastic Scattering and
Cluster Structures in ^{13}C and ^{24}Mg より

- β 拘束でのエネルギー曲線



β 拘束での計算では平均場構造がえられ、クラスター構造が現れない。
 ^{24}Mg の低励起状態では平均場構造が支配的となっている。

• 主量子数Nの拘束

^{24}Mg の高励起状態は多粒子多空孔配位をもつと考えられる。

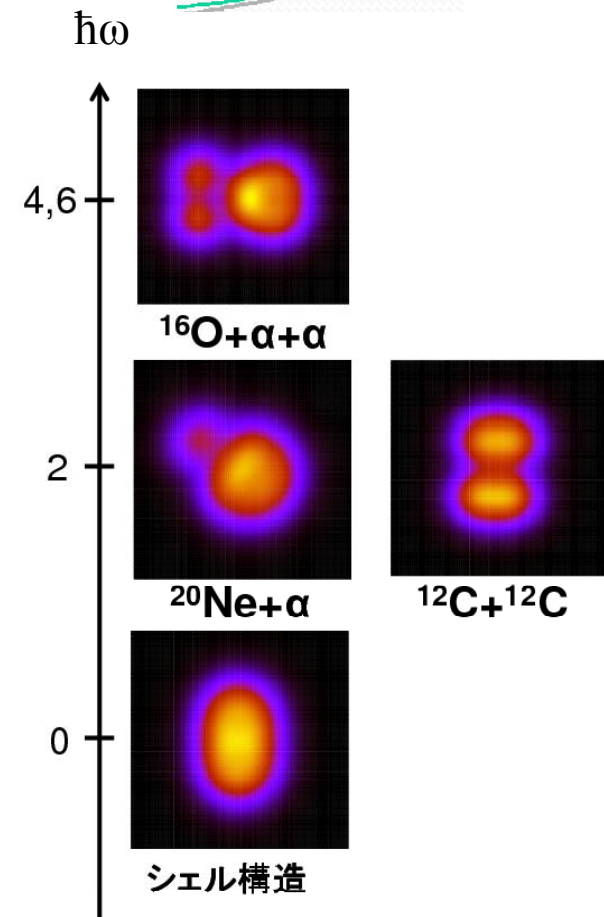
そこで、主量子数を拘束する。

$$\hat{N} = \hat{a}^\dagger \hat{a} \quad \hat{a}^\dagger = \sqrt{\nu} \hat{r} - \frac{i}{2\hbar\sqrt{\nu}} \hat{p}$$

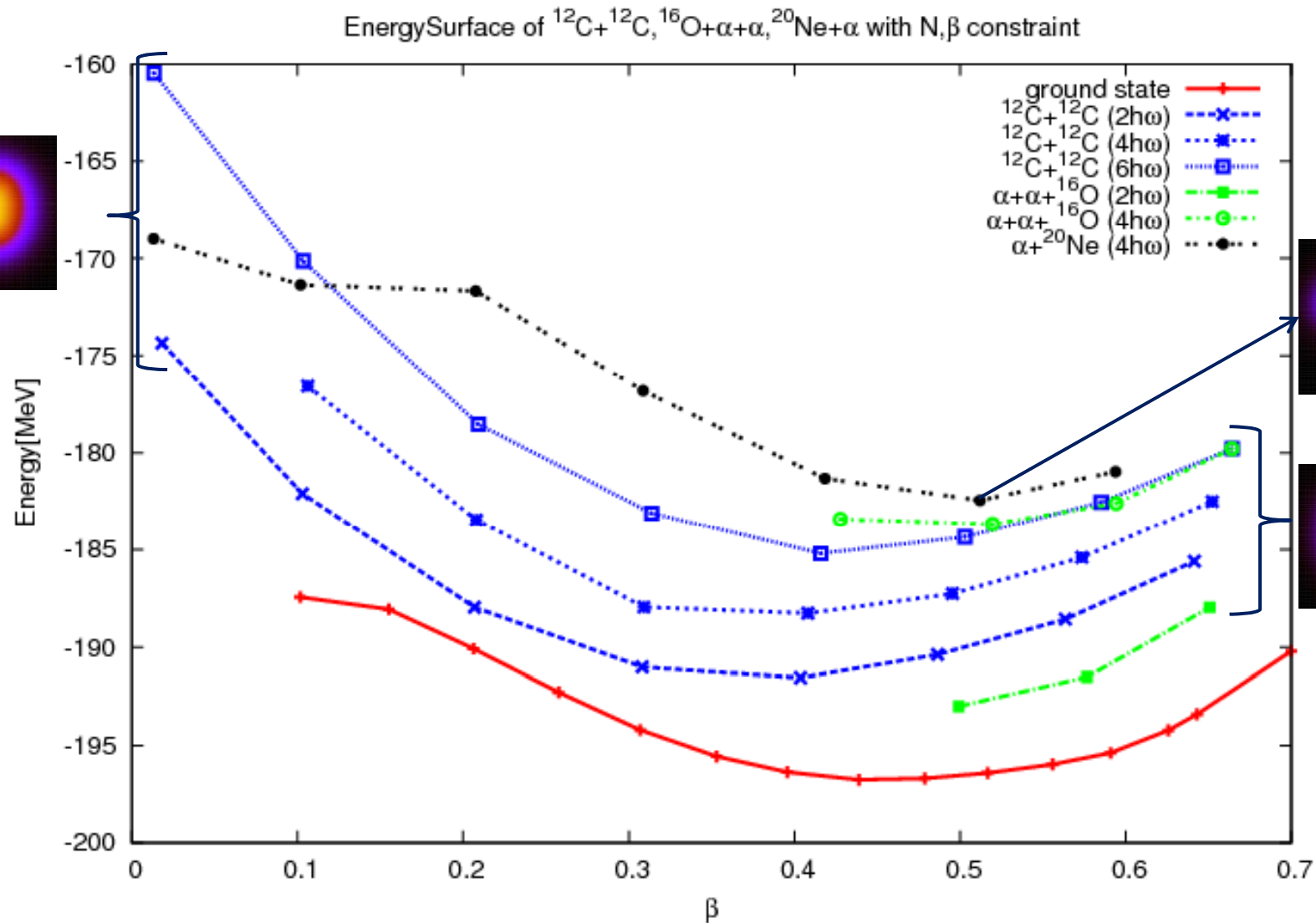
主量子数を拘束すると、高励起状態に多様なクラスター構造をもつ解が現れた。(β拘束とは対照的)

高励起状態では低励起状態とは逆にクラスター構造が支配的になっている可能性

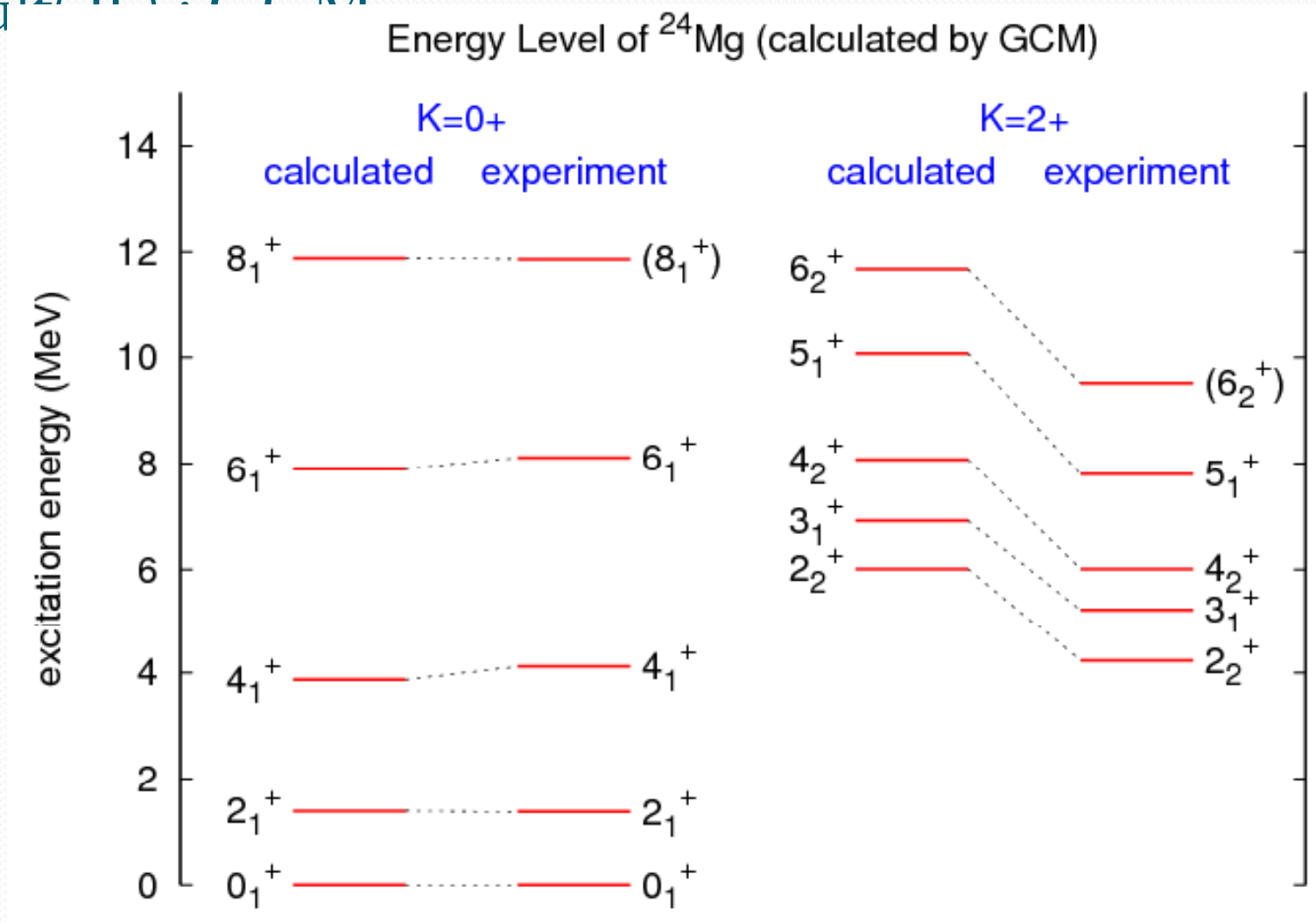
そこで主量子数と四重極変形度 β を同時に拘束した系統的な計算を行い、 ^{24}Mg の低-高励起状態を調べる。



III. Result(I) : Energy Surface ($^{12}\text{C}+^{12}\text{C}$, $^{16}\text{O}+\alpha+\alpha$, $^{20}\text{Ne}+\alpha$)

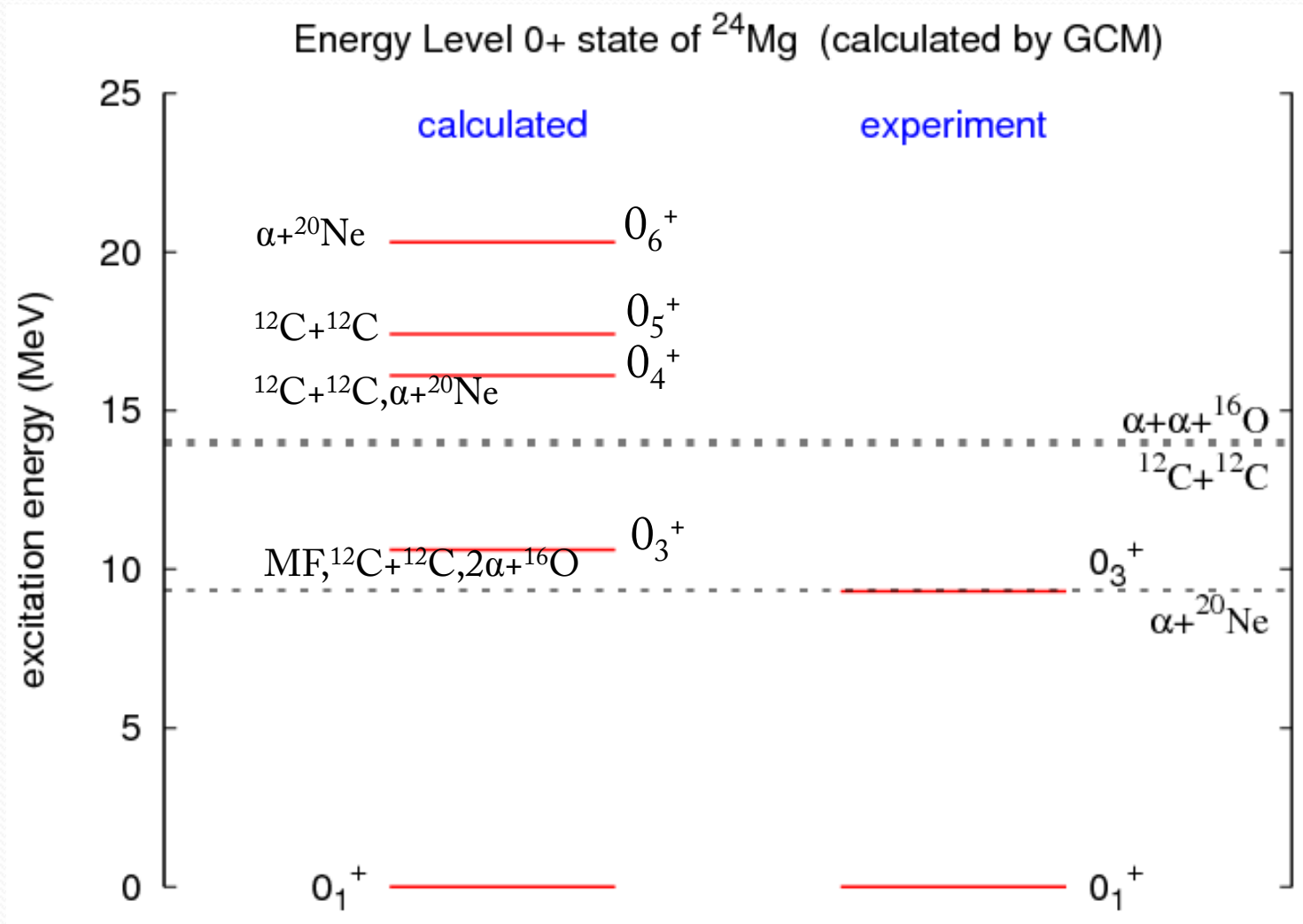


• Result (II) : GCM



0^+ : 198.6MeV (実験値 : 198.26MeV) K=0+状態は実験とほぼ一致している。

- Result(III) : GCM



実験データでみられたような10MeV~20MeV付近に多数の0+状態が現れた。これらの状態がクラスター状態をつくっていると予想される。