偏極$^6Li^{3+}$イオン源の開発

A. Tamii

Research Center for Nuclear Physics,

Osaka University, Japan

Collaboration of the Polarized $^{6}Li^{3+}$ Ion Source Project

RCNP, Osaka University, Japan
K. Hatanaka, A. Tamii, Y. Sakemi, Y. Shimizu, K. Fujita, and Y. Tameshige

CYRIC, Tohoku University, Japan
H. Okamura

Kyushu University, Japan
T. Wakasa

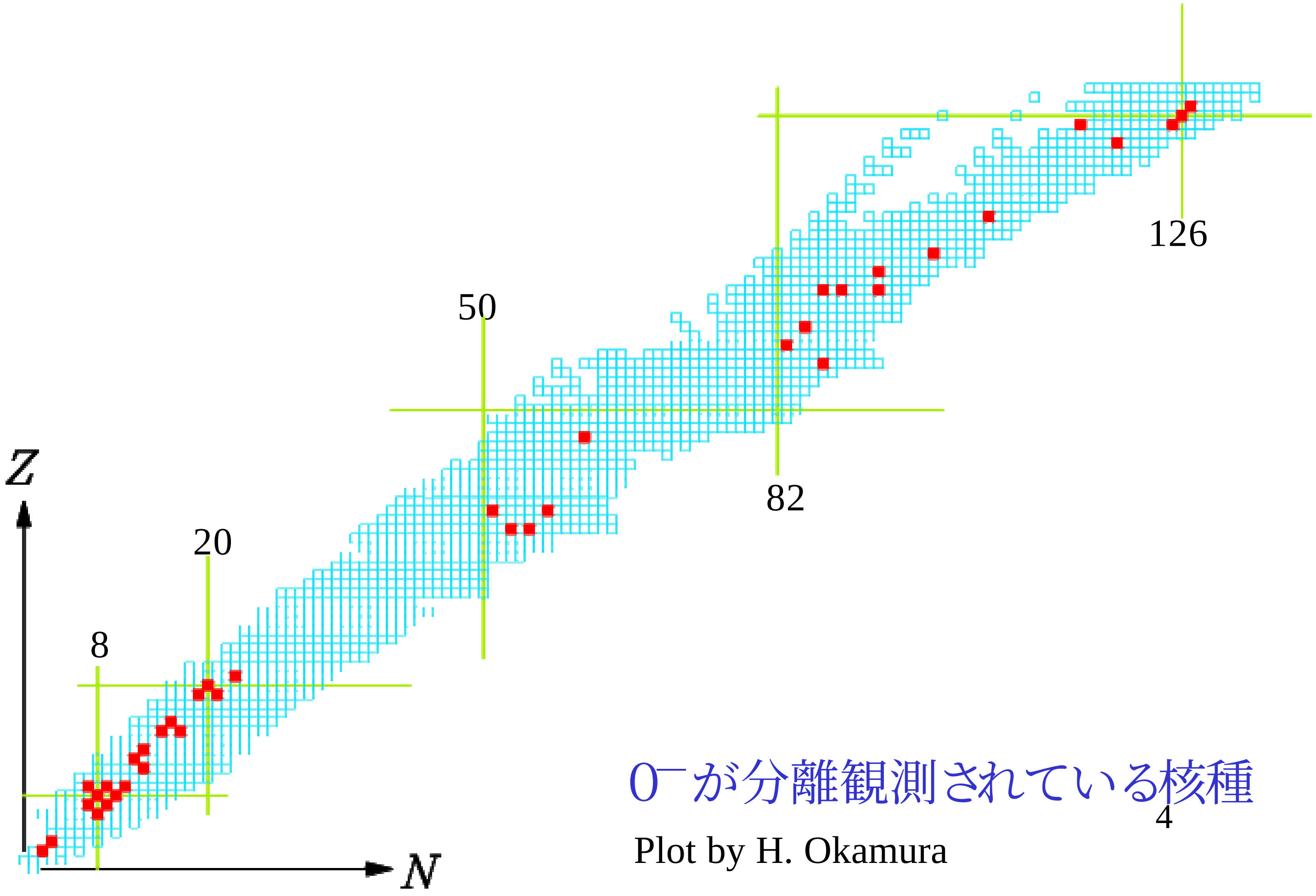
CNS, University of Tokyo, Japan
T. Uesaka and T. Wakui

RIKEN, Japan
T. Nakagawa

Physics and Experimental Method

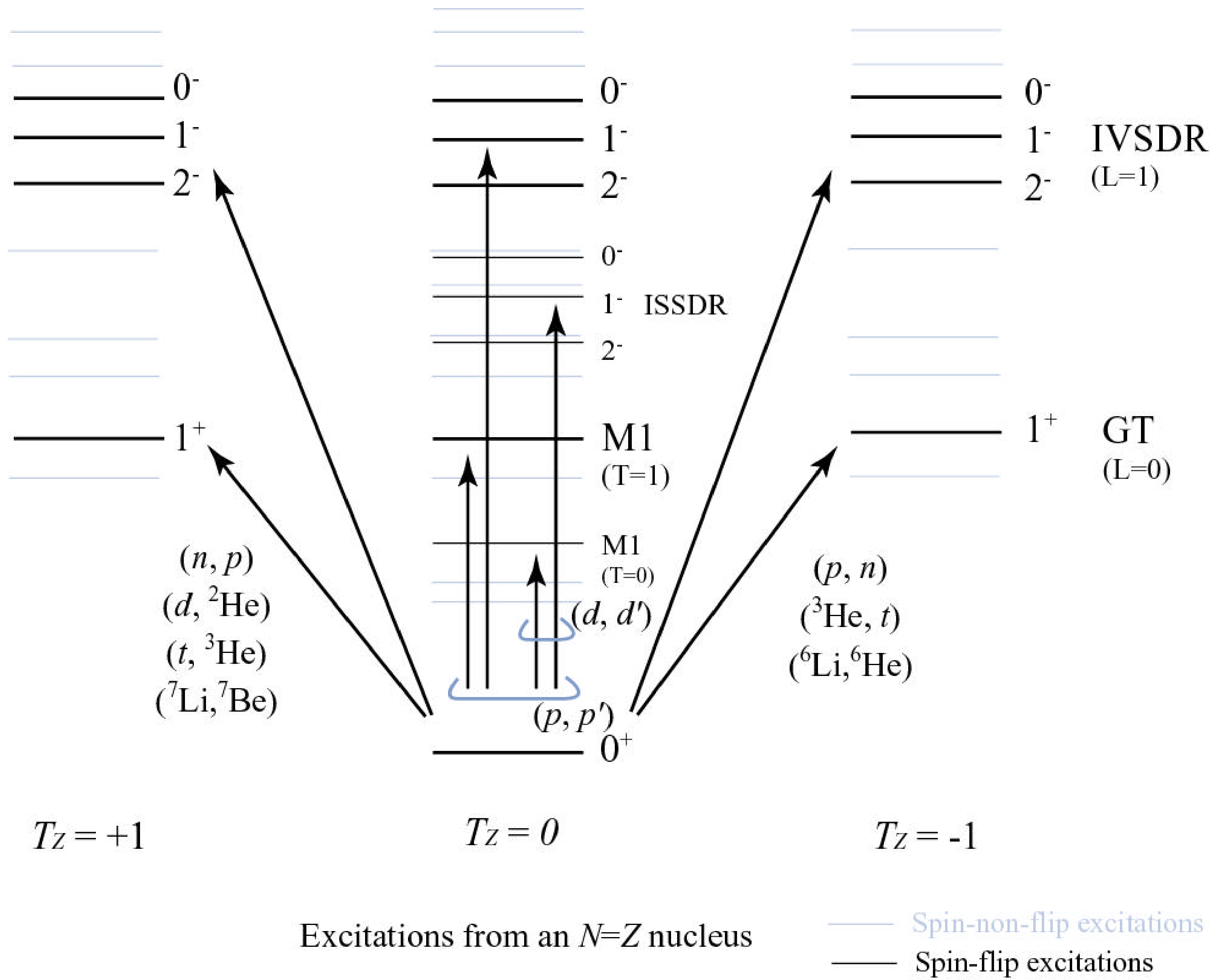
Study of (Isovector) 0^- states

- Chiral Partner of the 0+ ground state
- Role of the π-meson and the tensor force in the nuclear structure



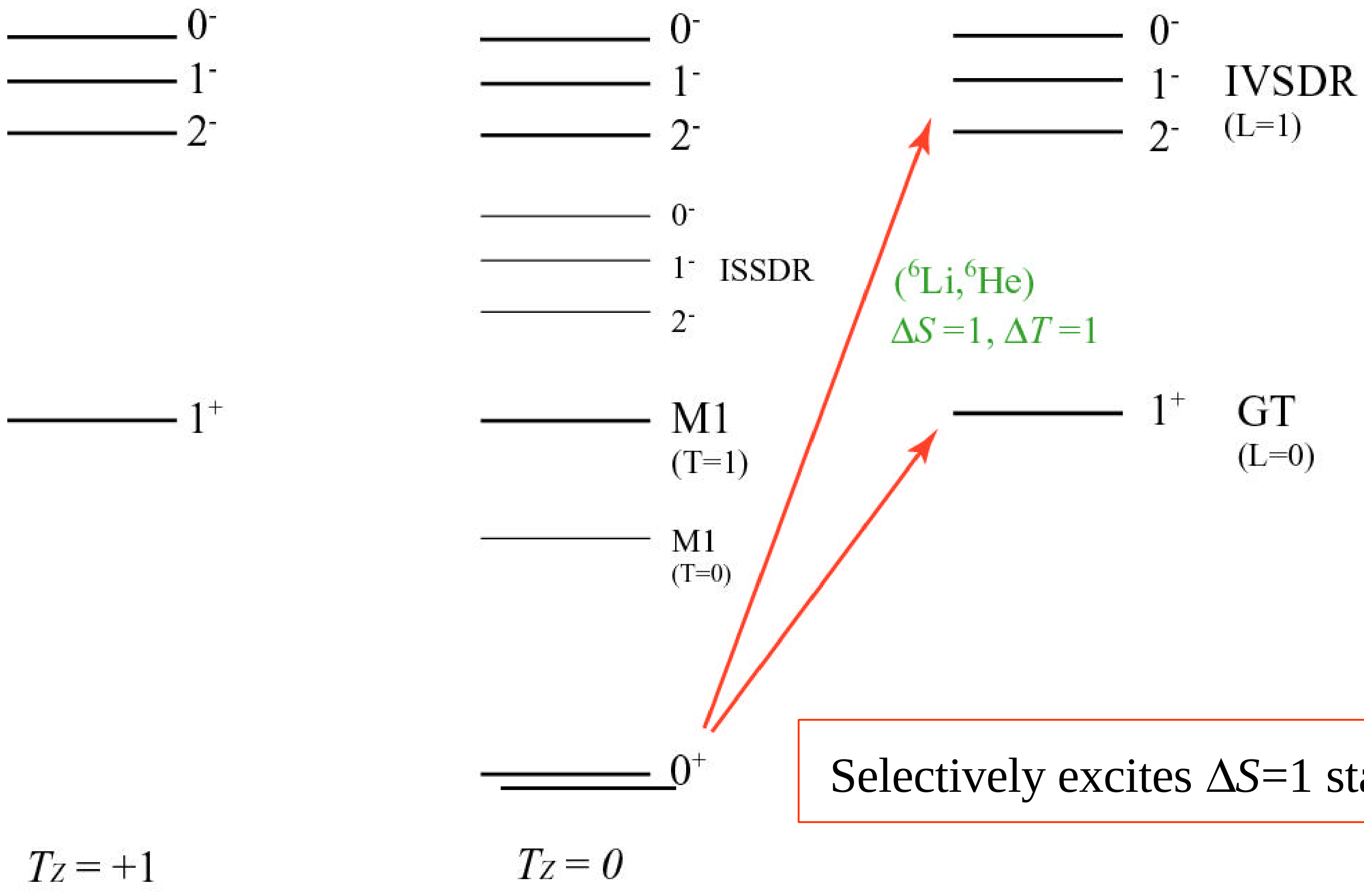
0^-が分離観測されている核種

Plot by H. Okamura



Excitations from an N=Z nucleus

Study of Nuclear Structures by the ($^{6}\vec{\mathrm{Li}}$,^{6}He) Reaction at 100 MeV/U



Selectively excites ΔS=1 states

Spin-Flip Excitations from an N=Z nucleus

Study of Nuclear Structures by the ($^{6}\vec{Li}$,^{6}He) Reaction at 100 MeV/U

One-step reaction is dominant at 100 MeV/U

Selective excitation of ΔT=1, ΔS=1

1^+ for ΔL=0

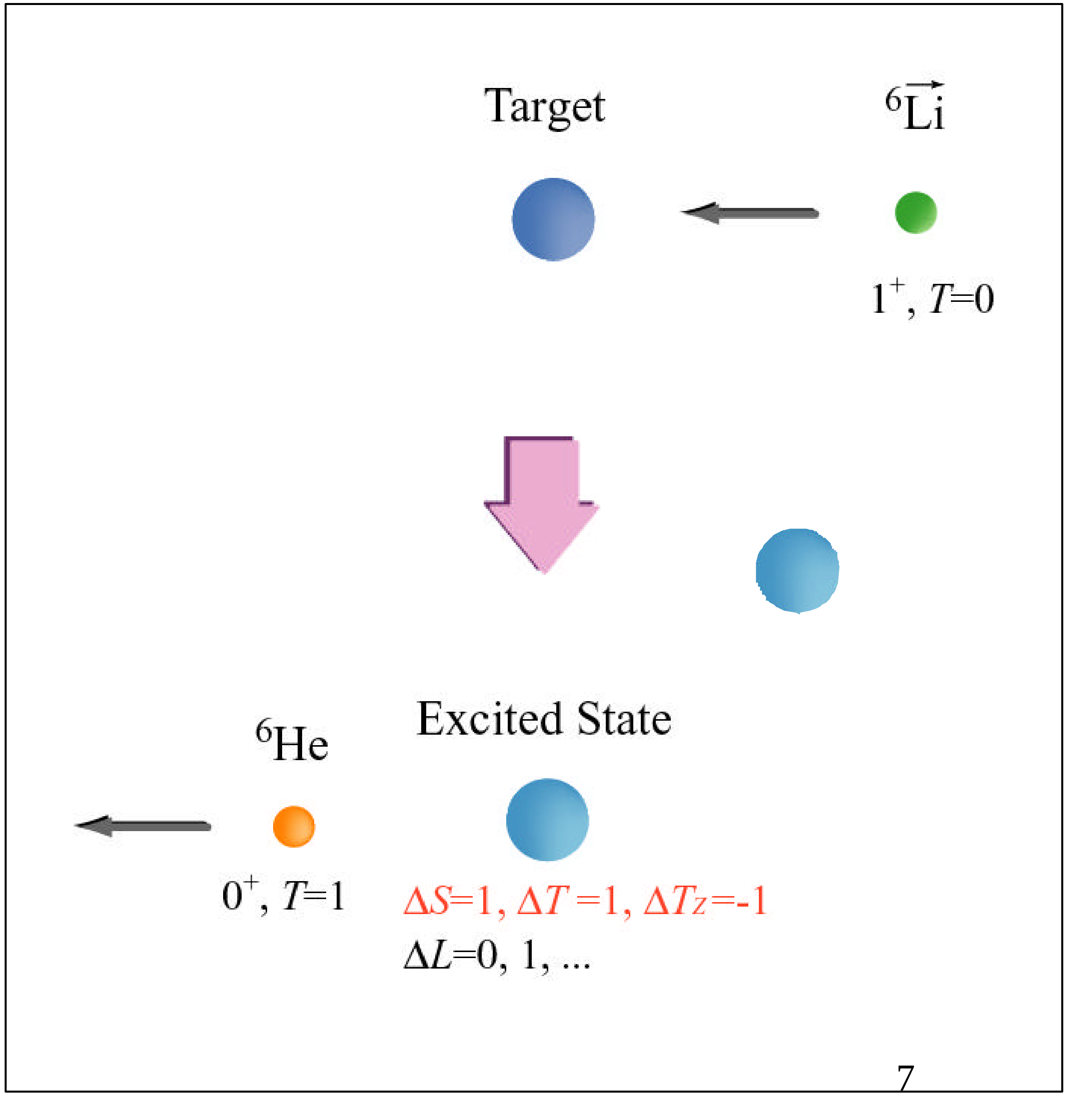
0^-, 1^- and 2^- for ΔL=1

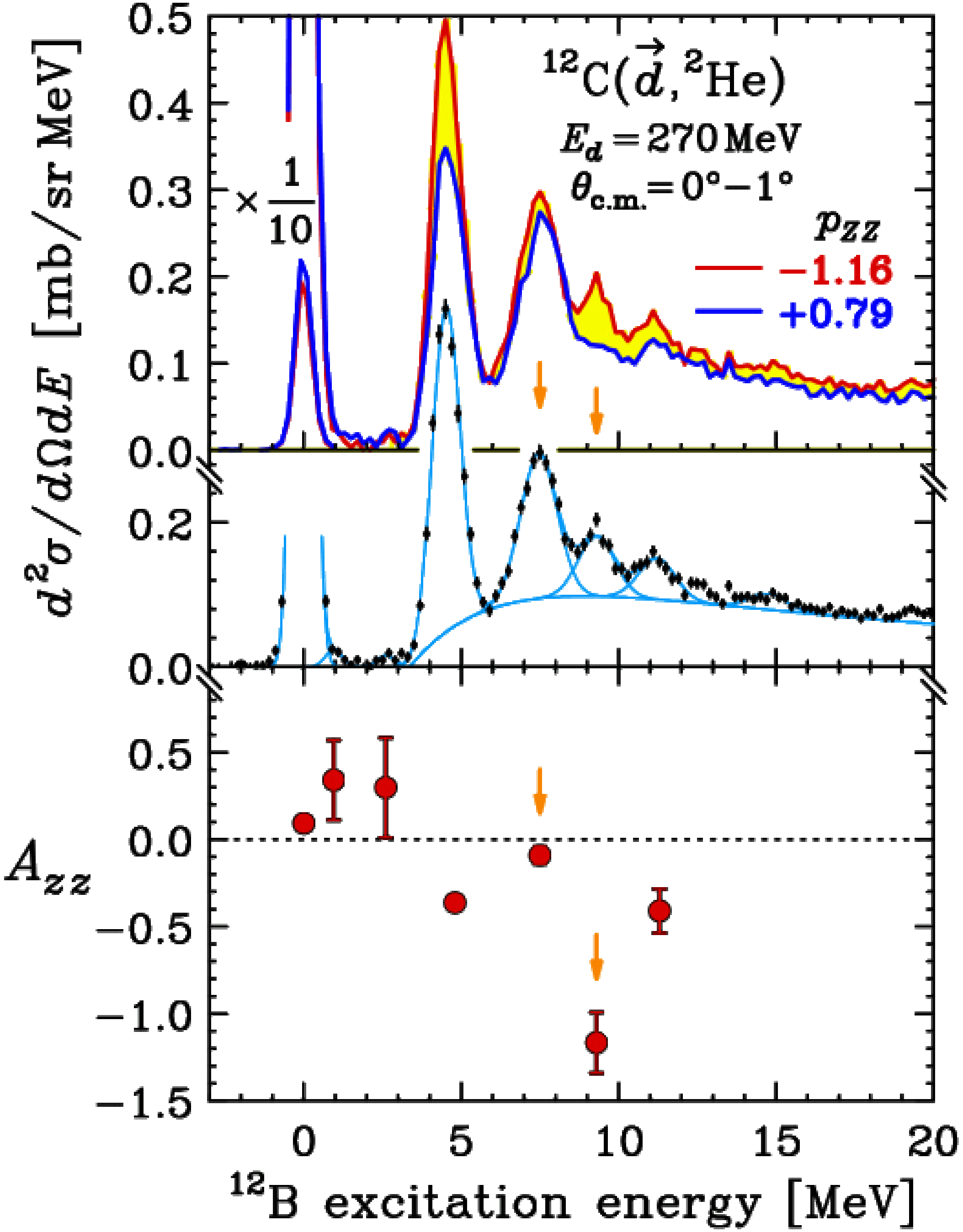
Tensor analyzing power at 0 °

Sensitive to J of 0^-,1^-, and 2^- states

High resolution measurements by the dispersion matching method

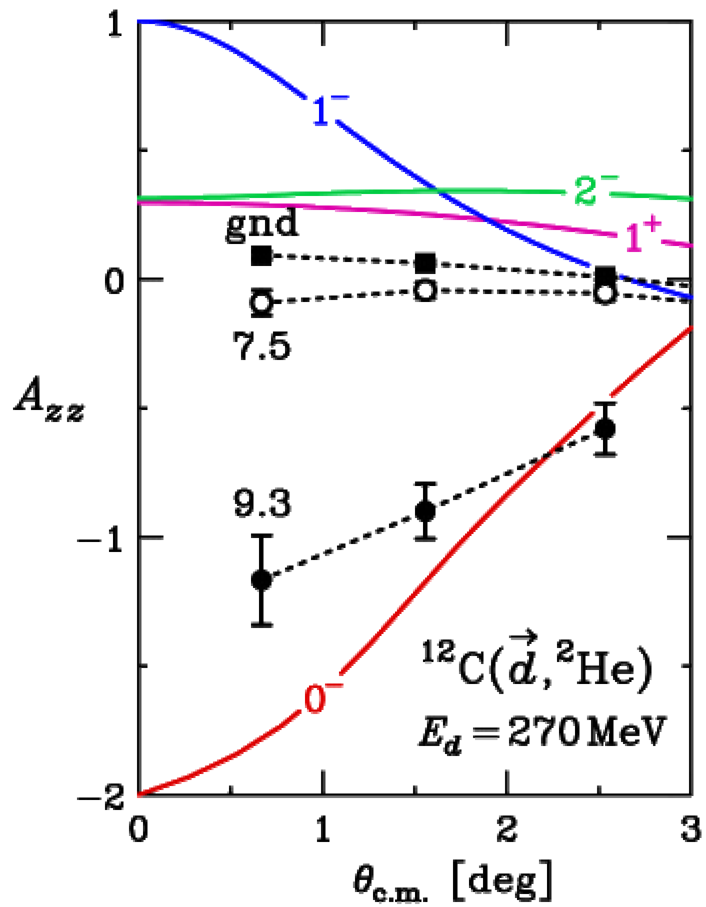
⇔(d,^{2}He), (p,n)





$$\sigma = \sigma_0\left(1 + \tfrac{1}{2} P_{ZZ} A_{ZZ}\right)$$

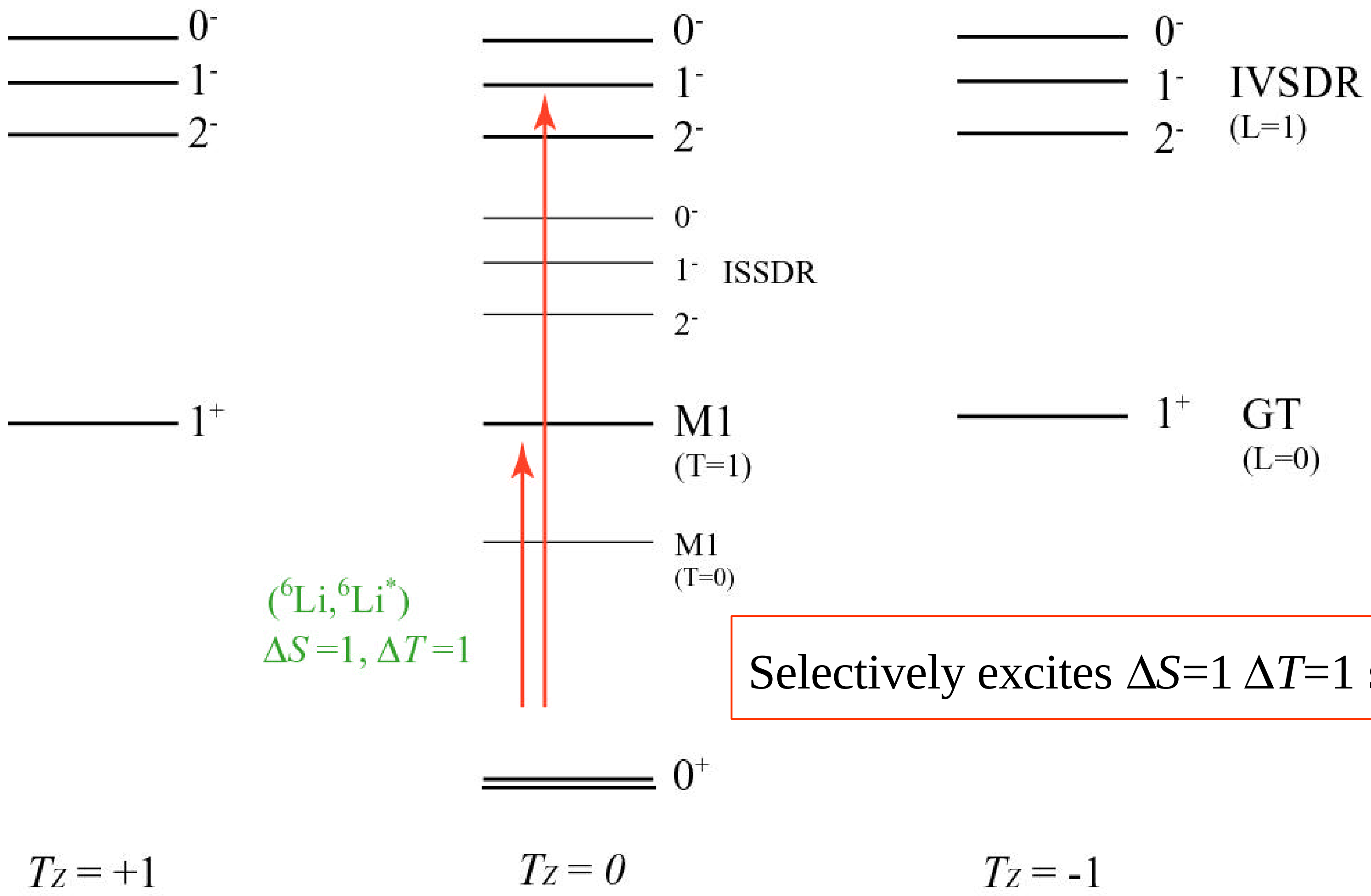
$$\xrightarrow{A_{ZZ}(0^-)=-2} \begin{cases} 3\sigma_0 & (P_{ZZ} = -2) \\ 0 & (P_{ZZ} = +1) \end{cases}$$



A = 12 系で初めて 0^- を発見

H. Okamura *et al*. Phys. Rev. C 66 (2002) 054602

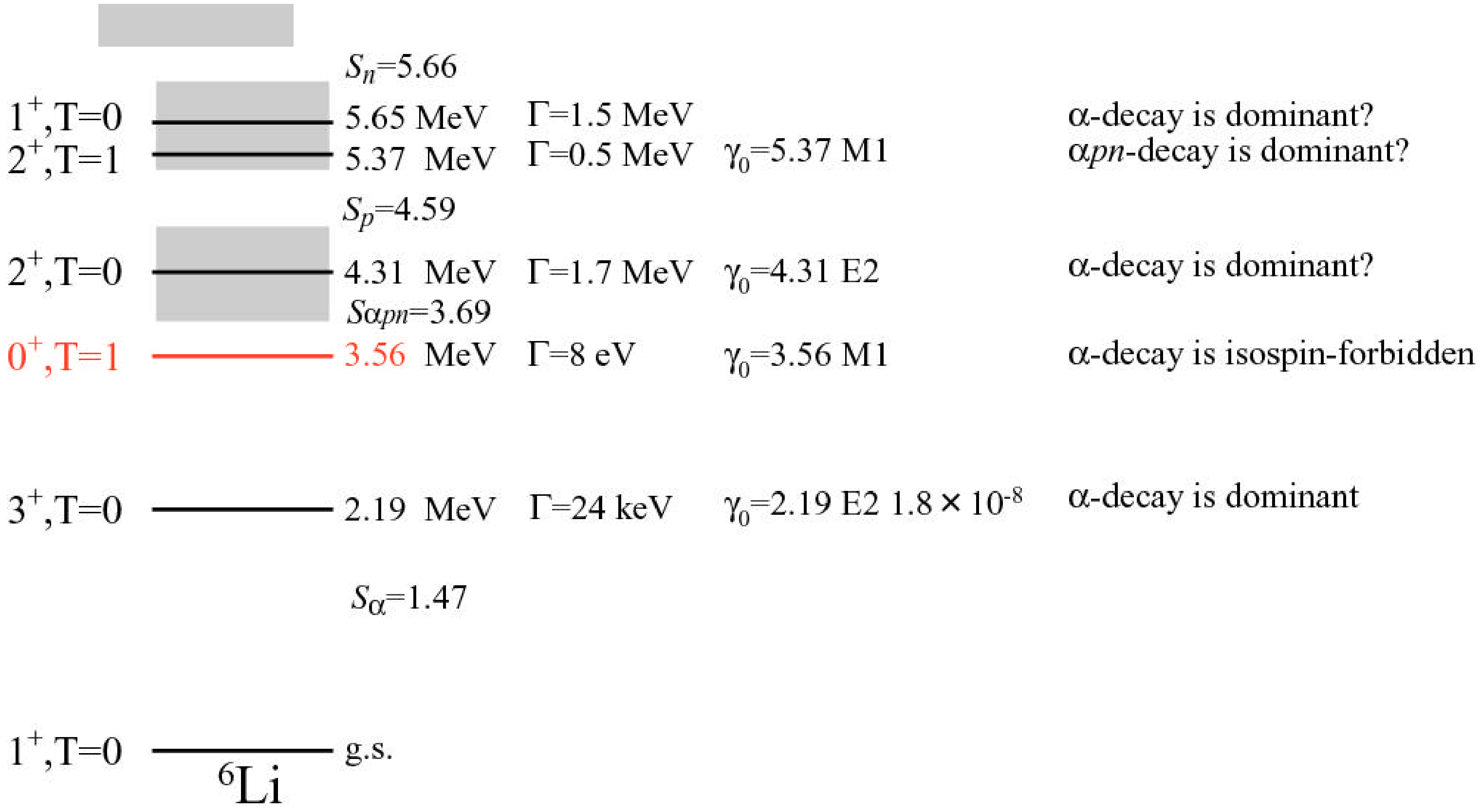
Study of Nuclear Structures by the $(^6\vec{\mathrm{Li}}, {}^6\mathrm{Li}^*(0^+, T=1; 3.56\mathrm{MeV}))$ Reaction at 100 MeV/U



Selectively excites $\Delta S=1$ $\Delta T=1$ states

Spin-Flip Excitations from an $N=Z$ nucleus

Energy Level of ^6Li



Study of Nuclear Structures by the $({}^{6}\vec{\mathrm{Li}}, {}^{6}\mathrm{Li}^{*}(0^{+}, T=1; 3.56\mathrm{MeV}))$ Reaction at 100 MeV/U

Selective excitation of ΔT=1, ΔS=1 and

ΔT_z=0

ΔT=1 M1 transition

IVSDR

Tensor analyzing power at 0 °

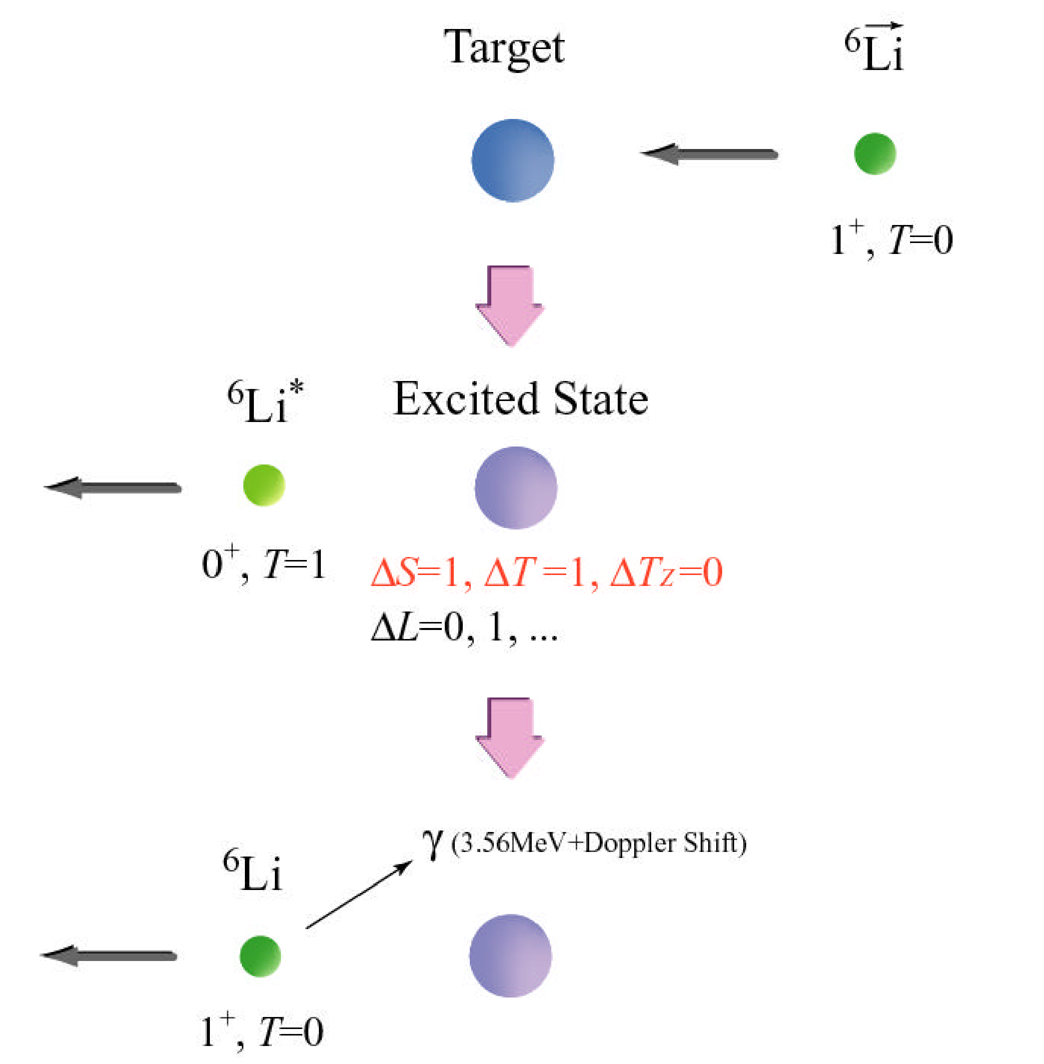
Selectivity for the 0^-, 1^-, and 2^- states

Relatively high resolution measurement

$\Leftrightarrow (d, d_{S=0})$

$d_{S=0} \rightarrow pn$

coincidence measurement



By K. Suda

$^{16}O(d, d^*_{singlet})^{16}O^*$(SDR) 反応測定

スピン双極子共鳴状態 (Spin Dipole Resonance)

$J^\pi = 0^-, 1^-, 2^-$ (ΔL=1, ΔS=1, ΔT=1)

cf. Gamow-Teller 共鳴, $J^\pi = 1^+$

(ΔL=0, ΔS=1, ΔT=1)

天体物理学への応用

超新星ニュートリノ測定 (特に ν_μ と ν_τ) に用いる提案

$^{16}O(\nu, \nu')^{16}O^*$(SDR) の粒子崩壊からの γ 線を
スーパーカミオカンデによって検出

Langanke et al., Phys. Rev. Lett. 76, 2629 (1996)

まず、$^{16}O^*$(SDR) についてのデータ (断面積) が必要
→ ハドロンプローブによる測定

SDR 励起に有効なプローブ

$(d, d^*_{singlet})$ 反応 · · · $d^*_{singlet}$: p-n [$^1S_0$], 非束縛系

選択性 ΔS=1, ΔT=1

ΔL=1 は角度分布から判定

開発した装置

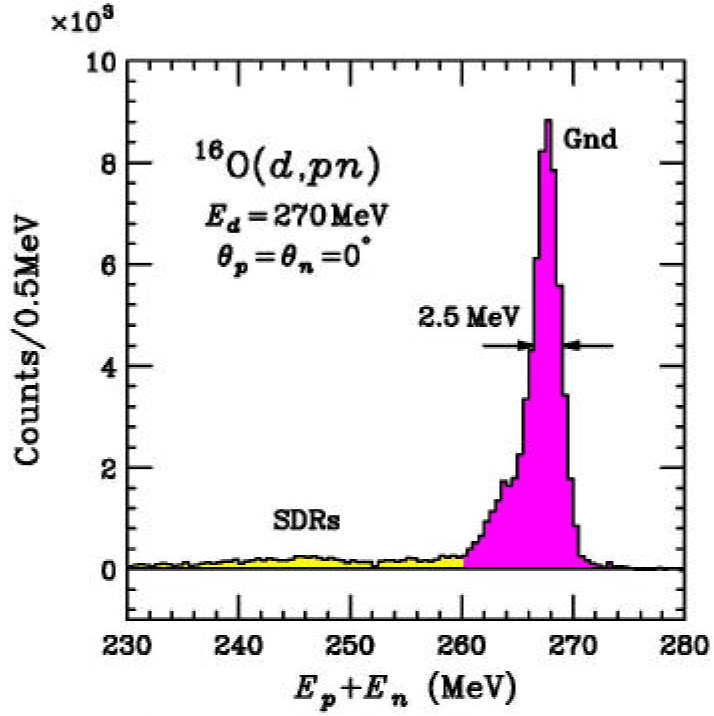
固体酸素標的 および 中性子検出器

実験

2002年6月 固体酸素標的のテスト実験

7月 本測定

典型的なスペクトル



エネルギー分解能の改善策

中性子 TOF の"START" の取り方を変える
加速器の RF 信号 → 陽子の飛行時間

目標分解能1MeV(解析進行中)

Dependence of the ^{6}Li energy on the angle of the decay gamma

$E^{lab}_{^6Li^*} = 600$ MeV

θ^{lab}_γ (deg)	E^{lab}_γ (MeV)	$E^{lab}_{^6Li}$ (MeV)	
0	5.630	597.930	
90	3.215	600.345	
			0.223
100	2.992	600.568	
			0.187
110	2.805	600.755	
			0.158
120	2.647	600.913	
			0.113
130	2.520	601.040	
			0.101
140	2.419	601.141	
			0.075
150	2.344	601.216	
			0.053
160	2.291	601.269	
			0.031
170	2.260	601.300	
			0.010
180	2.250	601.310	

"High-resolution" measurement is not feasible unless precisely detecting the angle or energy of decay gamma rays. → Ge Detector?

- Study of the reaction mechanism of composite particles
  - Elastic Scattering, inelastic scattering, (^{6}Li, ^{6}He) Reaction
  - diff. cross section and analyzing power

- Study of the Coulomb/nuclear break up mechanism of a polarized ^{6}Li nucleus

- Study of the spin structure of ^{6}Li

  ⇔ cluster 3-body calculation

Polarized $^{6}Li^{3+}$ Ion Source

Development of Polarized ^{6}Li ion Sources at Other Laboratories.

- Max Plank Institute, Heidelberg, Germany

  Optical Pumping + Surface Ionizer (+ Charge Exchange+Tandem)

  $^6Li^{1+}$: 20-30*m*A
- Florida State University, USA

  Optical Pumping + Surface Ionizer (+ Charge Exchange+Tandem+LINAC)
- Daresbury, UK

  Stern-Gerlach+RF + Surface Ionizer (+ Charge Exchange+Tandem)
- Wisconsin, USA

  Stern-Gerlach+RF (+ Cs-Beam-Bombard+Tandem)

  $^6Li^{0+}$: 3×10^{15}, $^6Li^{1-}$: 0.18*m*A
- Saturne, France

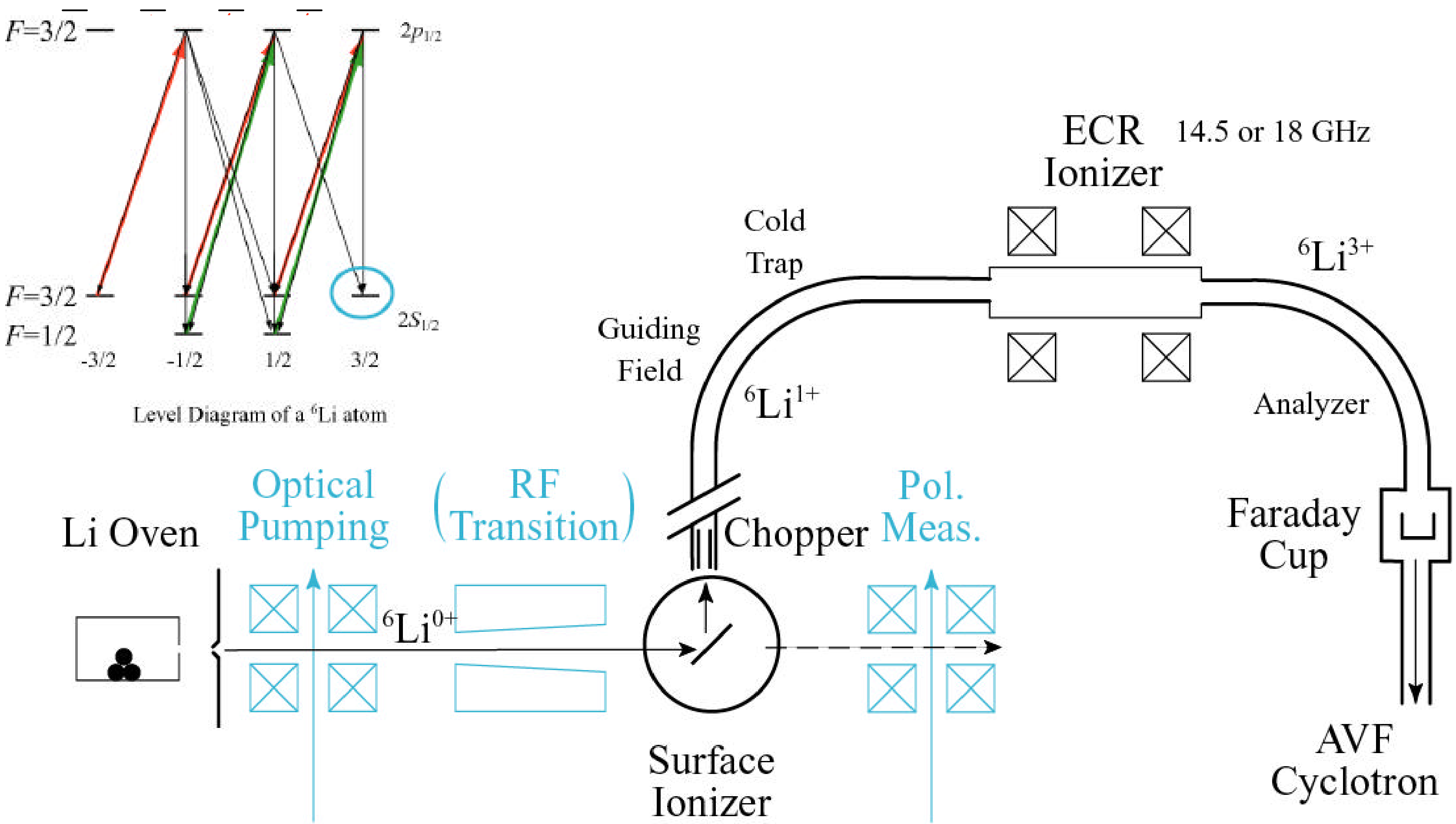
  Optical Pumping + Surface Ionizer (+ EBIS+Accum. Ring+Synchrotron)

  $^6Li^{1+}$: 20-35*m*A

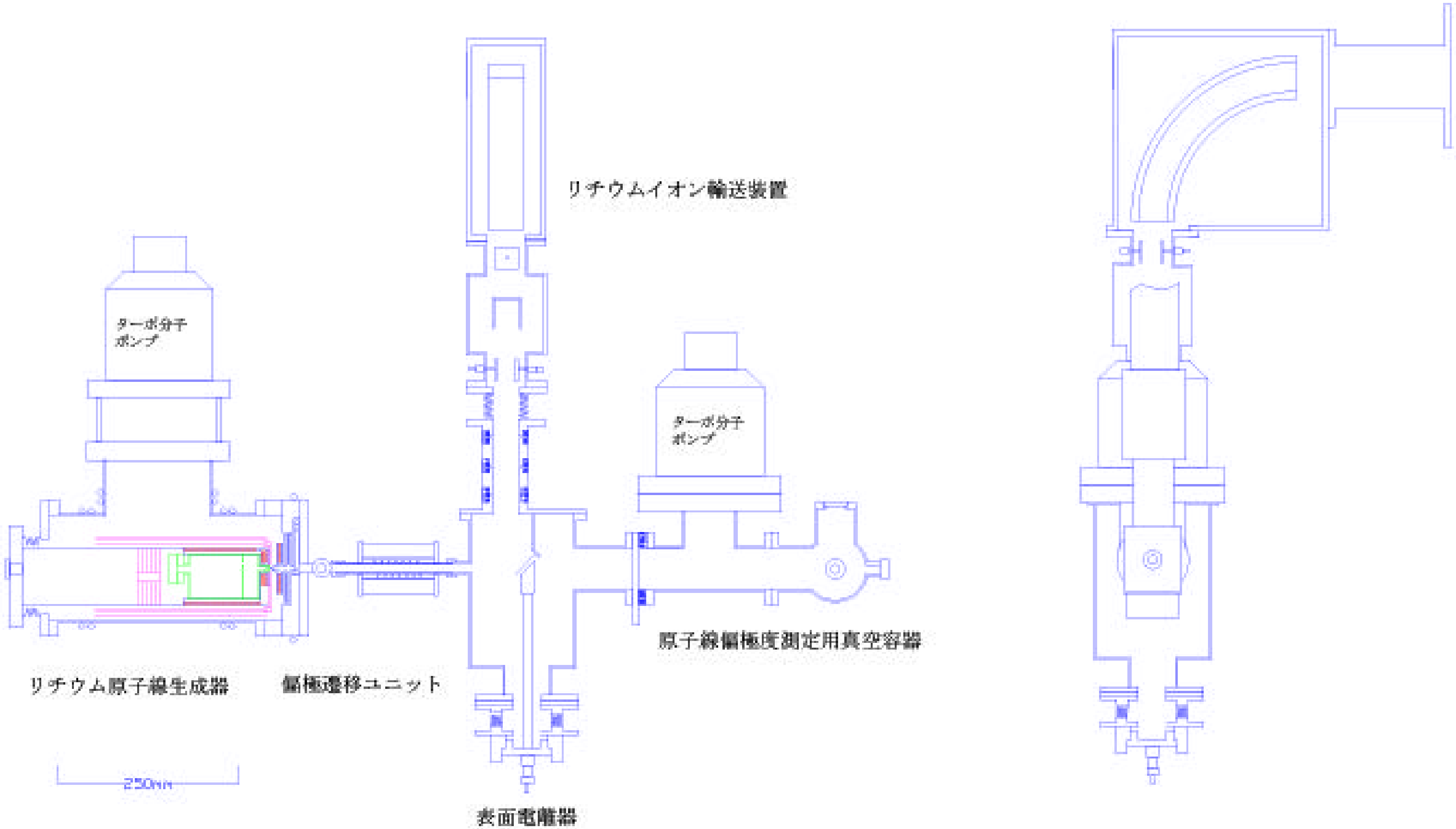
  $^6Li^{3+}$: 7×10^8 particles/spill

  P_{zz} = 70% at 187.5 keV/A

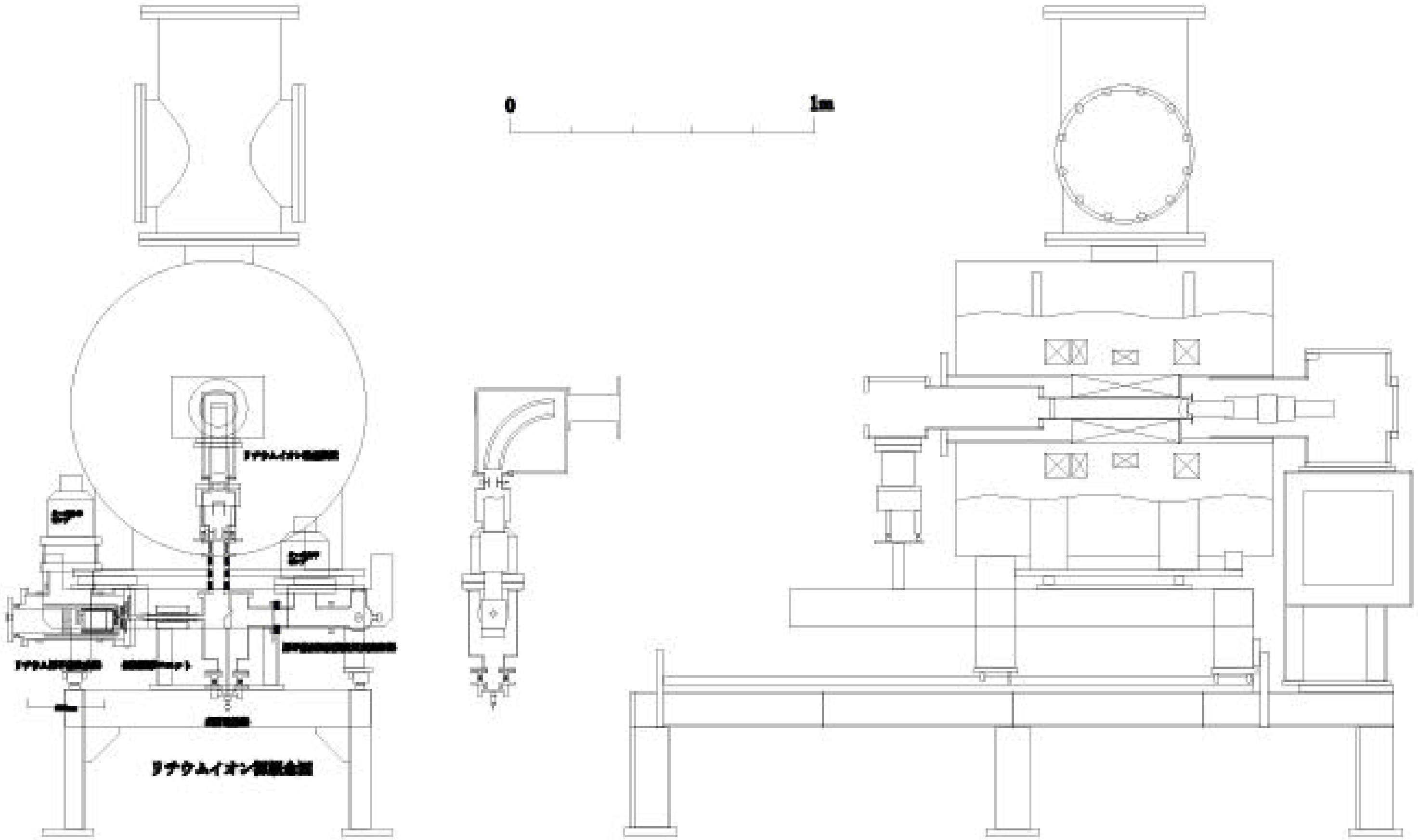
Plan of the polarized ^{6}Li ion source



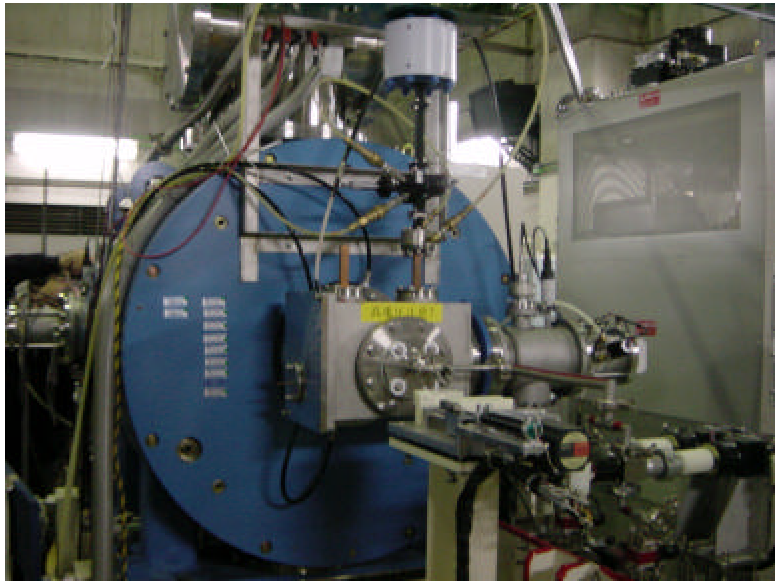
$^{6}Li^{1+}$: 20-30mA
Pol. 80-90%
at Florida State Univ.

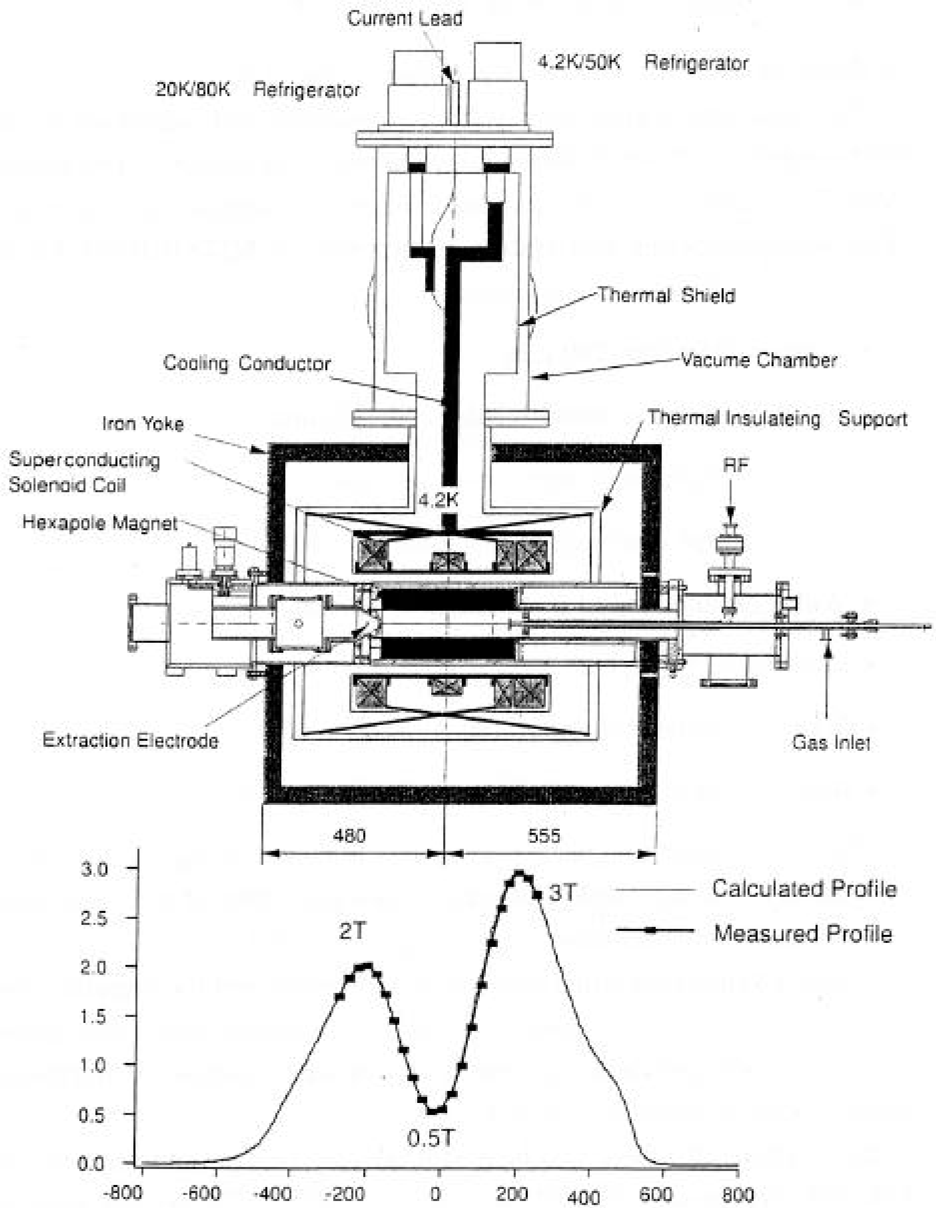


リチウムイオン源概念図


0
1m

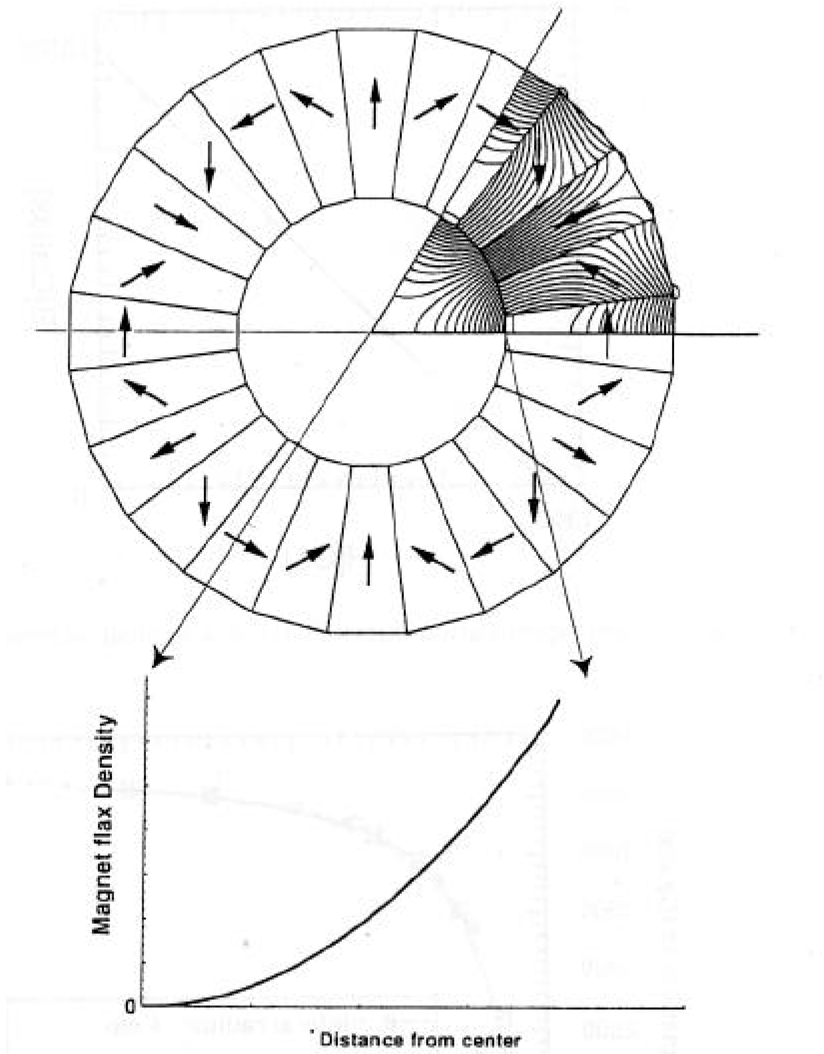
ECR Ionizer



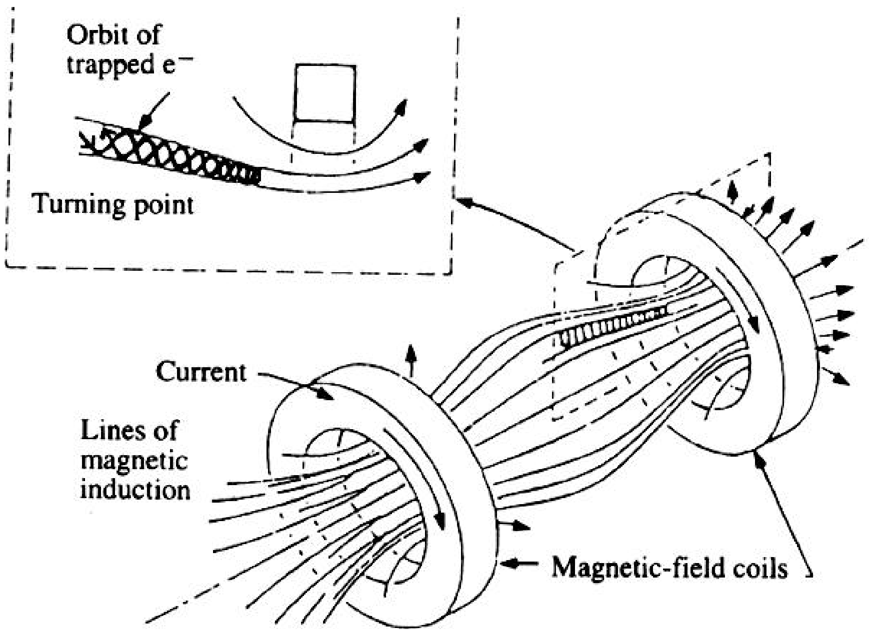

Current Lead
4.2K/50K Refrigerator
20K/80K Refrigerator
Thermal Shield
Cooling Conductor
Vacume Chamber
Iron Yoke
Thermal Insulateing Support
Superconducting Solenoid Coil
RF
4.2K
Hexapole Magnet
Extraction Electrode
Gas Inlet
480
555
3.0
2.5
2.0
1.5
1.0
0.5
0.0
3T
2T
0.5T
Calculated Profile
Measured Profile
-800
-600
-400
-200
0
200
400
600
800

ECR Ionizer

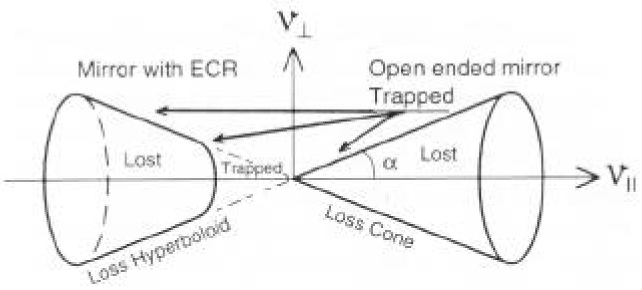
A hexapole magnet made of 24 pieces of permanent magnets



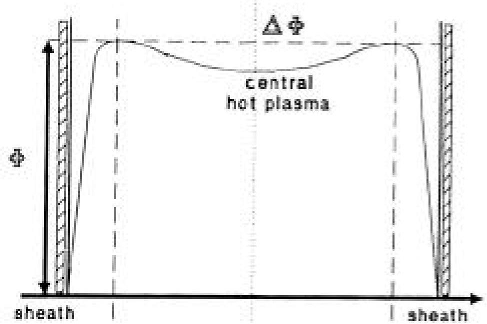
ECR Ionizer



The magnetic 'mirror', also called the simple mirror.



Losses in velocity space: loss cone in a simple mirror and loss hyperboloid



Critical Magnetic Field

Calc. by H. Okamura

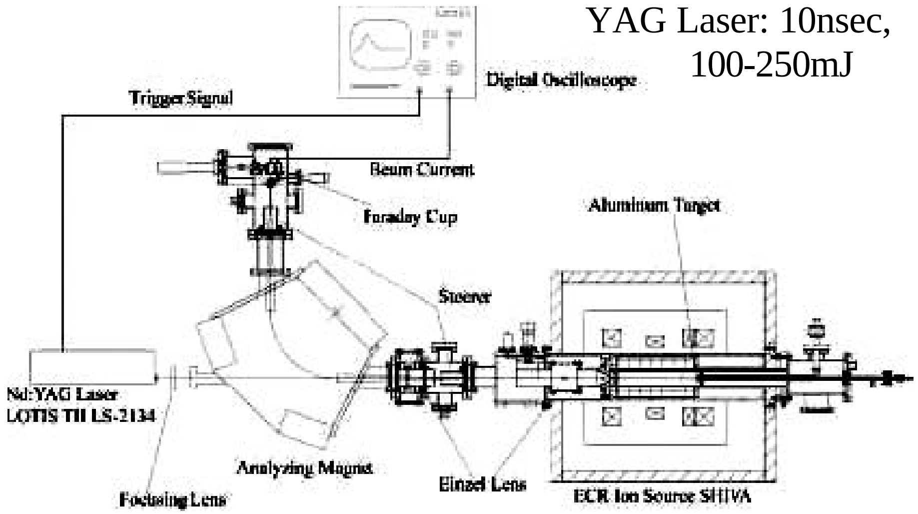
atom	state	calc. ν	calc. B_C	exp. ν	exp. B_C	μ_I/μ_N
^{1}H	1s	1422.586	508.204	1420.406	507.591	+2.7928
	2s	177.823	63.525	177.557	63.450	
^{2}H	1s	327.564	117.019	327.384	116.842	+0.8574
	2s	40.945	14.627	40.924	14.605	
^{3}H	1s	1517.387	542.071	1516.702	542.059	+2.9790
	2s	189.673	67.759	189.594	67.759	
$^3He^+$	1s	8669.430	3097.062			−2.1275
	2s	1083.679	387.133			
$^6Li^{2+}$	1s	8479.169	3029.093			+0.8220
	2s	1059.896	378.637			
		(MHz)	(Gauss)	(MHz)	(Gauss)	

Assumption of the Plasma Condition

The following plasma condition is assumed according to the empirical analysis of the laser abraded Al ion intensities from a 14.5 GHz ECR ionizer (SHIVA).

(M. Imanaka, PhD thesis, Univ. of Tsukuba)

YAG Laser: 10nsec, 100-250mJ



Buffer Gas: Oxygen

RF Power: 250 W

Neutral Oxygen Gas Density (n_{gas}): 1.44×10^{10} cm^{-3}

Electron Density (n_e): 2.23×10^{11} cm^{-3}

Electron Temperature (T_e) : 582 eV

Ion Temperature (T_i) : 5 eV

Ionization Rate: Voronov's empirical Fit

Charge Exchage Rate: Muller and Saltzborn

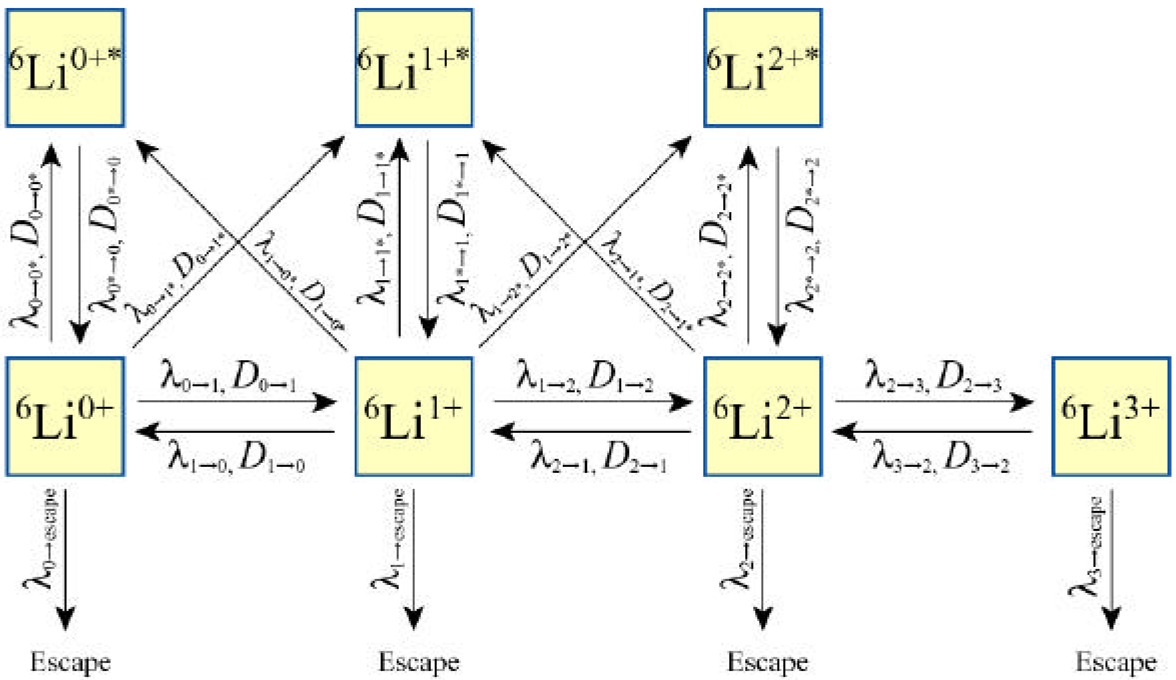
Confinement time of Al: $t_i \propto \frac{i}{i_{\max}} t_c$ for the i+ ions, τ_c=10msec

n_e, T_e, t_c, T_i are fitted to the data.

Simulation of the Depolarization in the ECR Ionizer

(extension of the simulation by Prof. M. Tanaka)

- Fractions and polarizations of escaped ions are calculated by assuming the initial conditions, transition rates, and magnetic-substate transition matrix.
- The rate equations are analytically solved.



$\lambda_{i\to j}$: Transition Rate from i to j [s^{-1}]

$D_{i\to j}$: Transition Matrix of Magnetic Substates from i to j ($0 \leqq D_{ji} \leqq 1$)

Results of the simulation

The result of the simulation is

$$P_{3+,escape} = \begin{pmatrix} 0.0165 & 0.0010 & 0.0000 \\ 0.0010 & 0.0148 & 0.0017 \\ 0.0000 & 0.0017 & 0.00157 \end{pmatrix} P_{1+,in}$$

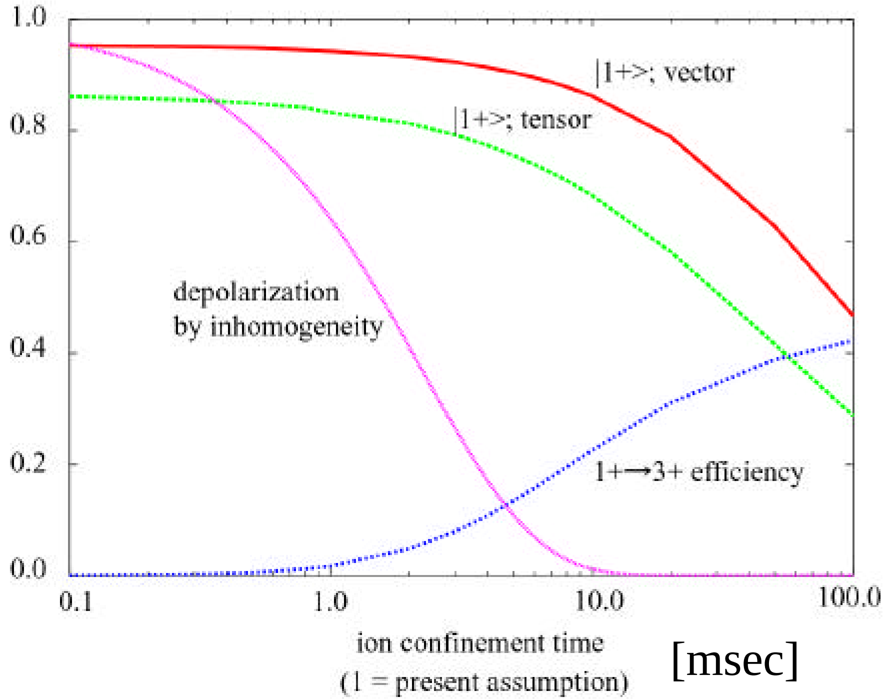
The polarization of escaped 3+ ions when we feed 1+ ions with pure magnetic substate population is summarized as follows

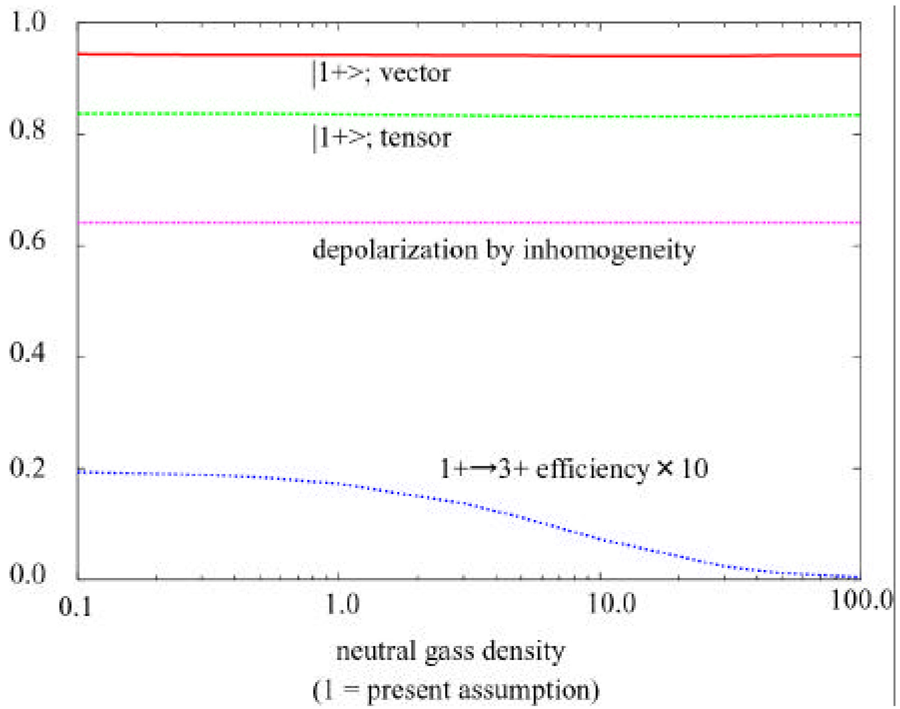
Table 2: Calculated depolarization and efficiency for the $^6Li^{1+} \rightarrow {}^6Li^{3+}$ ionization in the ECR ionizer.

Initial State ($^6Li^{1+}$)			Final Sate ($^6Li^{3+}$)		
state	vector pol.	tensor pol.	vector pol.	tensor pol.	efficiency
pure $\|+1>$	1.00	1.00	0.94	0.84	0.017
pure $\|\ 0>$	0.00	-2.00	-0.05	-1.54	0.017
pure $\|-1>$	-1.00	1.00	-0.90	0.70	0.017

Note that depolarization due to the inhomogeneous magnetic field is not included in the Present calculation.

Result of the simulation
(parameter dependence)

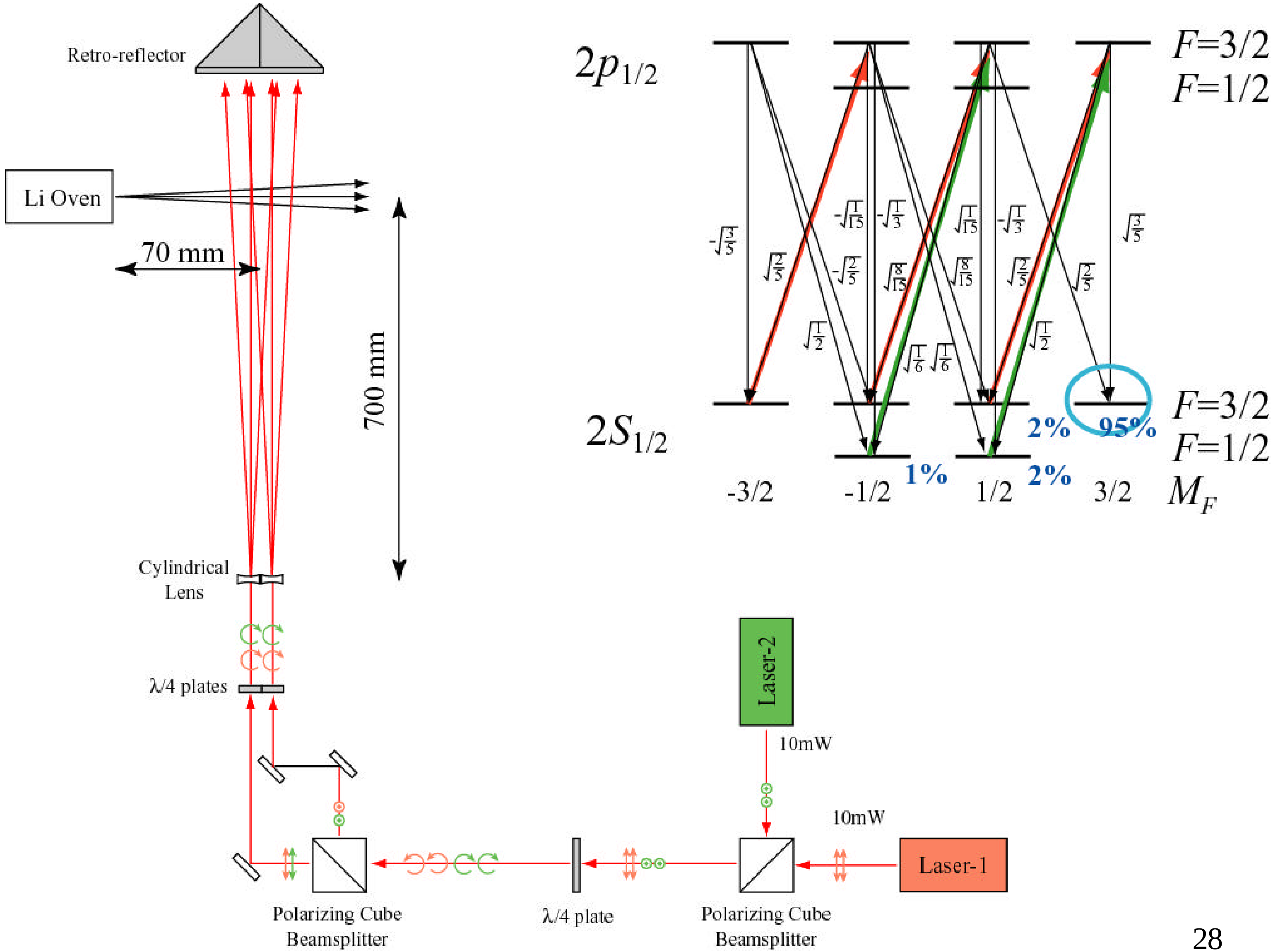




1=1msec for $^{6}Li^{3+}$

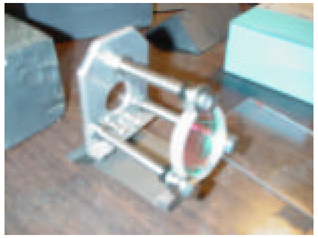
1 = Oxygen 1.44×10^{10} cm^{-3}

Simulation of the Optical Pumping

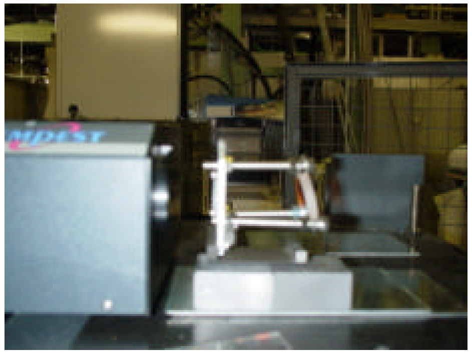


Study of the Confinement Time of Li ions by the Laser Ablation method

18GHz SC-ECRIS

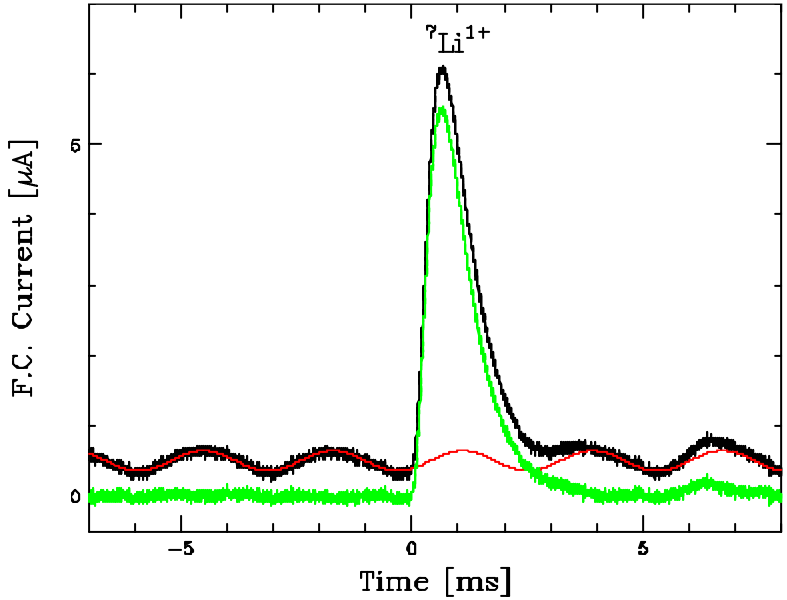


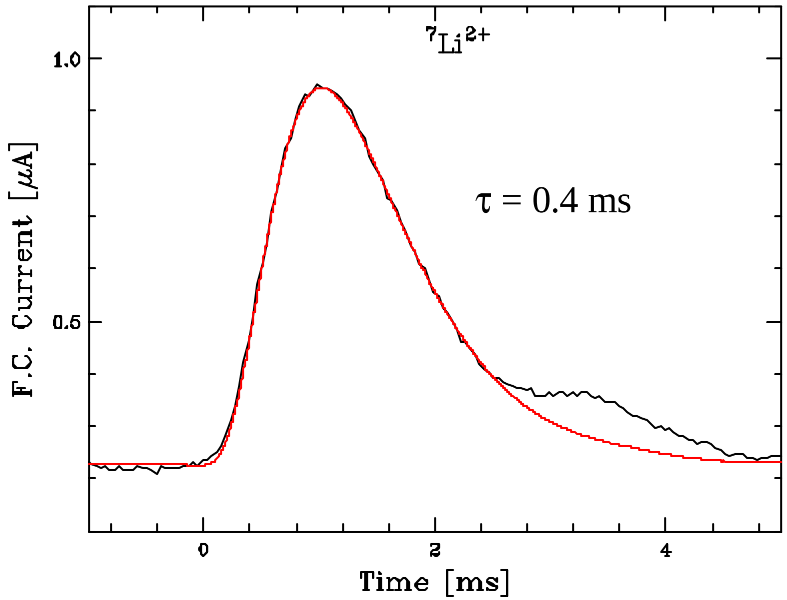
Lens, Mirror and
LiF rod



YAG 523nm 5ns
Max 100mJ/pulse

Laser ablation test
in atmosphere





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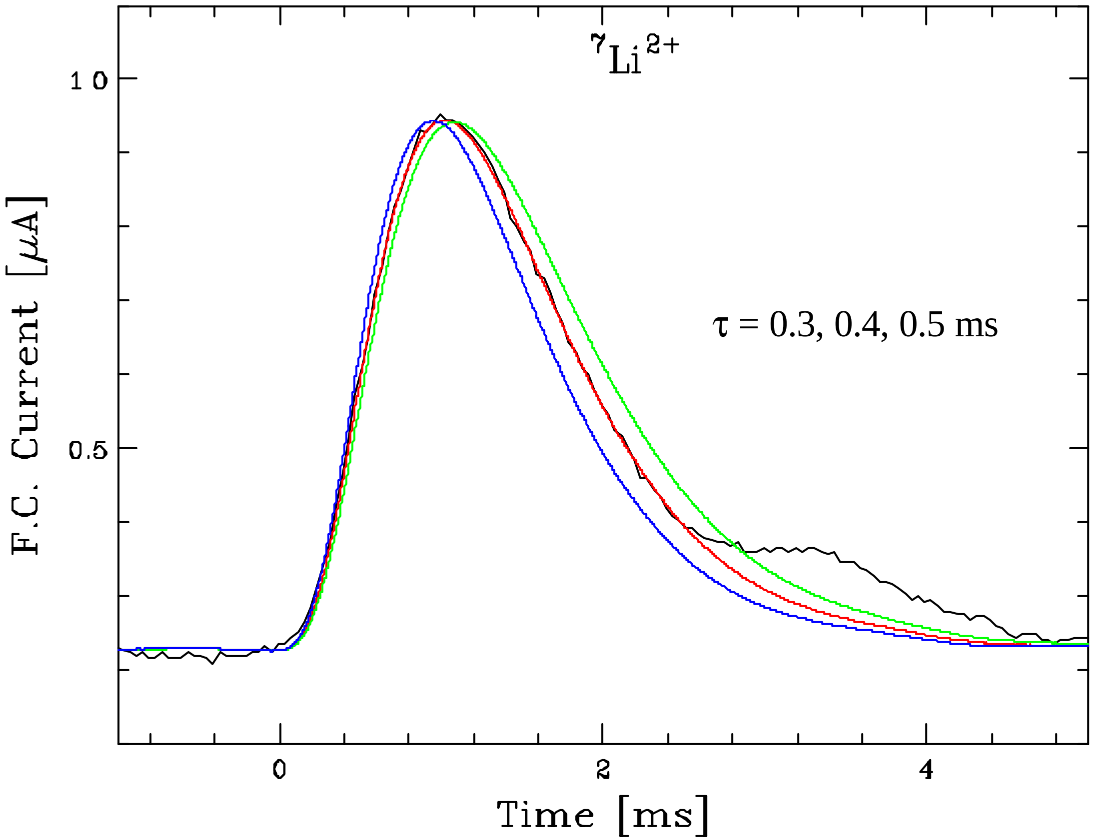
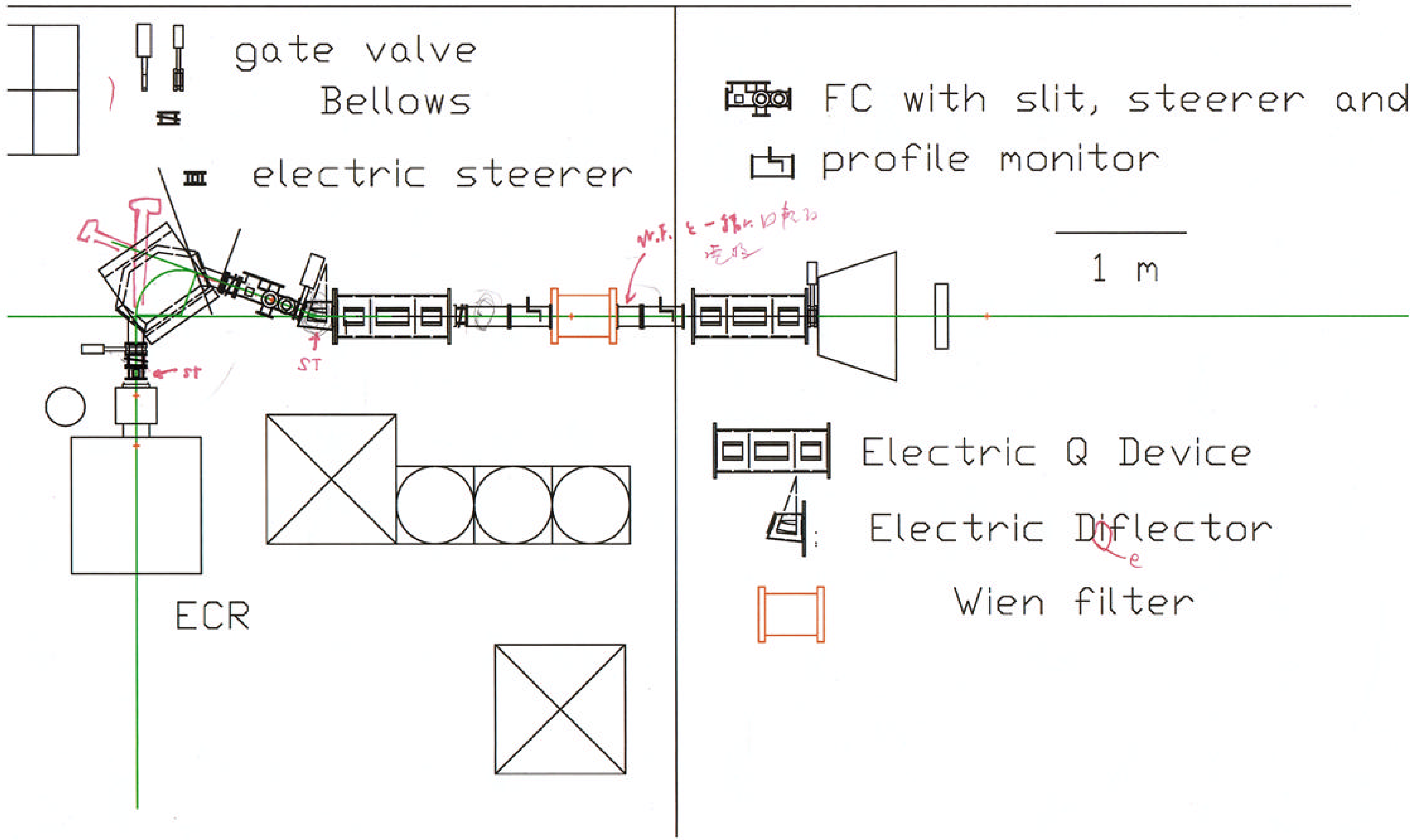


Note: the ECRIS operation has not tuned to $^{6}Li^{3+}$

Beam line of the new ECR ionizer



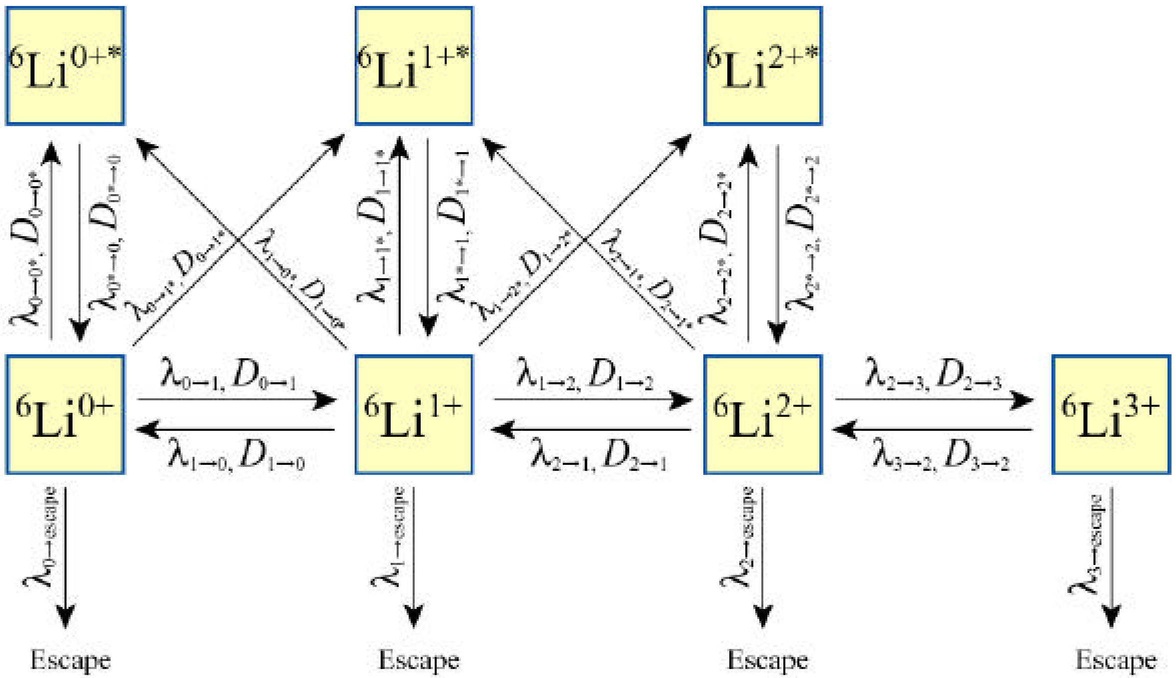
Physics and Experimental Method

Appendices

Simulation of the Depolarization in the ECR Ionizer

(extension of the simulation by Prof. M. Tanaka)

- Fractions and polarizations of escaped ions are calculated by assuming the initial conditions, transition rates, and magnetic-substate transition matrix.
- The rate equations are analytically solved.



$\lambda_{i\to j}$: Transition Rate from i to j [s^{-1}]

$D_{i\to j}$: Transition Matrix of Magnetic Substates from i to j ($0 \leqq D_{ji} \leqq 1$)

Magnetic-Substate Transition Matrix (1/2)

(according to the calc. of 3He by M. Tanaka and Y. Plis)

- The wave functions $\Psi_i(t)$ of the electron-nucleus system in a magnetic field system are written as a linear conbination of |IJ> states as

$$\Psi_{\mathrm{I}}(0) = \left|\uparrow +1\right\rangle$$
$$\Psi_{\mathrm{II}}(0) = \sin b_+ \left|\uparrow 0\right\rangle + \cos b_+ \left|\downarrow +1\right\rangle$$
$$\Psi_{\mathrm{III}}(0) = \sin b_- \left|\uparrow -1\right\rangle + \cos b_- \left|\downarrow 0\right\rangle$$
$$\Psi_{\mathrm{IV}}(0) = \left|\downarrow -1\right\rangle$$
$$\Psi_{\mathrm{V}}(0) = -\cos b_- \left|\uparrow -1\right\rangle + \sin b_- \left|\downarrow 0\right\rangle$$
$$\Psi_{\mathrm{VI}}(0) = -\cos b_+ \left|\uparrow 0\right\rangle + \sin b_+ \left|\downarrow +1\right\rangle$$

$$\sin b_\pm \equiv \sqrt{\frac{1}{2}(1+d_\pm)} \qquad \cos b_\pm \equiv \sqrt{\frac{1}{2}(1-d_\pm)}$$

$$d_\pm \equiv \frac{\pm\frac{1}{3}+x}{\sqrt{1\pm\frac{2}{3}x+x^2}}$$

$$x \equiv \frac{B}{B_c} \qquad B_c \text{ :critical magnetic field}$$

- The time revolution of the |↓ +1> state is

$$\begin{aligned}\left|\downarrow +1\right\rangle_t &= \cos b_+ \Psi_{\mathrm{II}}(t) + \sin b_+ \Psi_{\mathrm{IV}}(t) \\ &= \cos b_+ \Psi_{\mathrm{II}}(0)\exp(-iE_{\mathrm{II}}t) + \sin b_+ \Psi_{\mathrm{IV}}(0)\exp(-iE_{\mathrm{IV}}t) \\ &= \cos b_+ \left(\sin b_+ \left|\uparrow 0\right\rangle + \cos b_+ \left|\downarrow +1\right\rangle\right)\exp(-iE_{\mathrm{II}}t) \\ &\quad + \sin b_+ \left(-\cos b_+ \left|\uparrow 0\right\rangle + \sin b_+ \left|\downarrow +1\right\rangle\right)\exp(-iE_{\mathrm{IV}}t)\end{aligned}$$

- The probability to find |↓ +1> and its time average (after sufficient time) is

$$\begin{aligned}P(t) &= \left|\cos^2 b_+ \exp(-iE_{\mathrm{II}}t) + \sin^2 b_+ \exp(-iE_{\mathrm{IV}}t)\right|^2 \\ &= \cos^4 b_+ + \sin^4 b_+ + 2\cos^2 b_+ \sin^2 b_+ \cos\left((E_{\mathrm{II}} - E_{\mathrm{IV}})t\right)\end{aligned}$$
$$\overline{P} = \cos^4 b_+ + \sin^4 b_+ = \frac{1}{2}(1+d_+^2)$$

Magnetic-Substate Transition Matrix (2/2)

- By similar calculations we obtain

$$\begin{pmatrix} |\uparrow +1\rangle' \\ |\uparrow 0\rangle' \\ |\uparrow -1\rangle' \\ |\downarrow -1\rangle' \\ |\downarrow 0\rangle' \\ |\downarrow +1\rangle' \end{pmatrix} = \begin{pmatrix} 1 & & & & & \\ & \frac{1}{2}\left(1+d_+^2\right) & & & & \frac{1}{2}\left(1-d_+^2\right) \\ & & \frac{1}{2}\left(1+d_-^2\right) & & \frac{1}{2}\left(1-d_-^2\right) & \\ & & & 1 & & \\ & & \frac{1}{2}\left(1-d_-^2\right) & & \frac{1}{2}\left(1+d_-^2\right) & \\ & \frac{1}{2}\left(1-d_+^2\right) & & & & \frac{1}{2}\left(1+d_+^2\right) \end{pmatrix} \begin{pmatrix} |\uparrow +1\rangle \\ |\uparrow 0\rangle \\ |\uparrow -1\rangle \\ |\downarrow -1\rangle \\ |\downarrow 0\rangle \\ |\downarrow +1\rangle \end{pmatrix}$$

- We are not interested in the electron spin.

  In the case that the orientation of the electron spin is random at t=0, by taking the average for the initial state and sum for the final state concerning the electron spin, we obtain

$$\begin{pmatrix} |+1\rangle' \\ |0\rangle' \\ |-1\rangle' \end{pmatrix} = \begin{pmatrix} \frac{1}{4}\left(3+d_+^2\right) & \frac{1}{4}\left(1-d_+^2\right) & 0 \\ \frac{1}{4}\left(1-d_+^2\right) & \frac{1}{4}\left(2+d_+^2+d_-^2\right) & \frac{1}{4}\left(1-d_-^2\right) \\ 0 & \frac{1}{4}\left(1-d_-^2\right) & \frac{1}{4}\left(3+d_-^2\right) \end{pmatrix} \begin{pmatrix} |+1\rangle \\ |0\rangle \\ |-1\rangle \end{pmatrix}$$

- When x=5/3, the matrix is

$$D_{\mathrm{dep}} = \begin{pmatrix} 0.955 & 0.045 & 0 \\ 0.045 & 0.871 & 0.083 \\ 0 & 0.083 & 0.917 \end{pmatrix}$$

Depolarization due to the electron spin resonance (ESR) effect

We take SHIVA as a model case.

If micro-wave with a power of 250W is applied in a (non-resonating) cylinder with a diameter of 72mm.

$$u = \frac{W}{pr^2 c} = 2.0 \times 10^{-10}\ \mathrm{J/cm3}$$

$$B_1 = \sqrt{m_0 u} = 0.16\ \mathrm{Gauss}$$

The thickness of the ESR region is

$$\Delta R = 4.0\,mm \quad \text{at } R = 5.0\text{cm (in axial direction)}$$

$$\Delta R = 0.9\,mm \quad \text{at } R = 1.9\text{cm (in radial direction)}$$

The effective thickness averaged for isotropic ion velocity distribution and averaged half-length between the ECR points are

$$L \cong \frac{4.0 + 0.9 \times 2}{3} \times \frac{1}{2}\left(1 + \ln\frac{2R}{\Delta R}\right) = 12\,mm$$

$$\overline{\mathrm{R}} = \frac{1}{2}\frac{5.0 + 1.9 \times 2}{3} = 1.5\,cm$$

The spin rotation angle of the electron calculated with random-walk approximation is

$$w = \Delta w \times \sqrt{N} = g_e B_1 \frac{L}{v} \times \sqrt{\frac{v}{\overline{R}} t_i} = 6.2 \times 10^{-2}\ \mathrm{rad} = 3.6°$$

The nuclear depolarization is caused by the hyper-fine coupling between the electron and the nucleus.

Hence depolarization is negligible. Note that the calculation depends on the assumed plasma parameters.

Depolarization due to the inhomogeneous magnetic field

The *T1* relaxation is calculated by the following formula by Schearer et al., Phys. Rev. 139 (1965) A1398.

$$\frac{1}{T1} = \frac{2}{3}\frac{v^2}{g_I^2 t_c H_0^4}\left(\frac{\partial H_y}{\partial y}\right)^2$$

For ions by putting the following numbers we obtain

$$g_I = 3.94\times10^7 \text{ rad}/s/T$$
$$t_c = 1.2\times10^{-6} \text{ sec}$$
$$v = 1.3\times10^6 \text{ sec}$$
$$H_0 = 0.5T$$
$$\frac{\partial H_y}{\partial y} = 0.15T/\text{cm}$$

$$T1 = 4.5 \text{ msec for ions}$$

For neutral lithium atoms, by putting the numbers we obtain

$$g_I = 3.94\times10^7 \text{ rad}/s/T$$
$$t_c = 3.7\times10^{-5} \text{ sec}$$
$$v = 9.7\times10^4 \text{ sec}$$
$$H_0 = 0.5T$$
$$\frac{\partial H_y}{\partial y} = 0.3T/\text{cm}$$

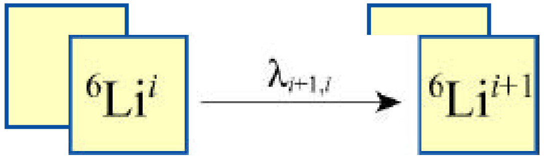
$$T1 = 6.3 \text{ for neutral atoms}$$

The *T1* relaxation time for ions has large depolarization effect when we consider the confinement time of $^6Li^{3+}$ (1 msec) and should be carefully taken care of.

Ionization Rate by Electron Impact

Voronov's empirical fit

G.S. Voronov, Atom. Data and Nucl. Data Tables 65 (1997)1.



$$c_{i\to i+1} = \langle s v_e \rangle = A \frac{1+PU^{1/2}}{X+U} U^{K} e^{-U} \quad [\mathrm{cm^3 s^{-1}}]$$

$$U = \frac{I_i}{T_e}$$

I_i: Ionization Energy

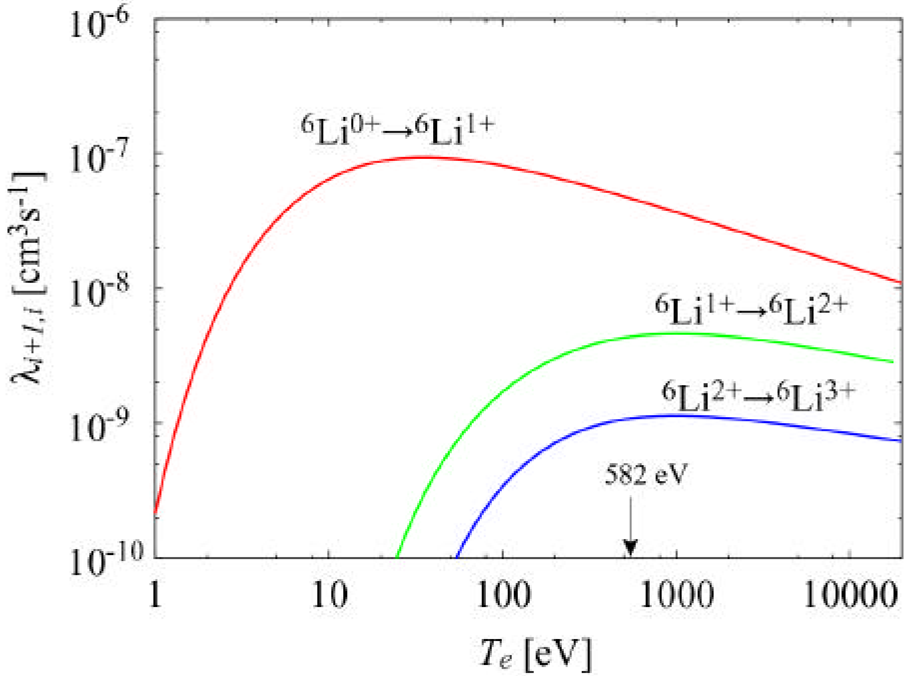
T_e: Electron Temperature

A, P, X, K: Fitting Parameters

$^{6}Li^{0+} \longrightarrow {}^{6}Li^{1+}$: 4.52×10^{-8} cm^{3}s^{-1}
$^{6}Li^{1+} \longrightarrow {}^{6}Li^{2+}$: 3.26×10^{-9} cm^{3}s^{-1}
$^{6}Li^{2+} \longrightarrow {}^{6}Li^{3+}$: 7.53×10^{-10} cm^{3}s^{-1}

$$l_{i\to i+1} = c_{i\to i+1} n_e$$

n_e: 2.23×10^{11} cm^{-3}



Charge Exchange Reaction Rate with the Neutral Gas

Muller and Saltzborn Empirical Fit

A. Muller and E. Saltzborn, Phys. Lett. A62 (1977) 391.

$$s = 1.43\times10^{-12} i^{1.17} I_{gas}^{-2.76} \quad [\mathrm{cm}^2]$$

$$z_{i\to i-1} = \langle s v_i \rangle = 3.15\times10^{-6} i^{1.17} I_{gas}^{-2.76} \sqrt{\frac{T_i}{A_i}} \quad [\mathrm{cm}^3\mathrm{s}^{-1}]$$

I_{gas}: Ionization Energy of the Neutral Gas (Oxygen: 13.6 eV)

T_i: Ion Temperature (5 eV)

A_i: Ion Mass in AMU

$^{6}Li^{1+} \to {}^{6}Li^{0+}$: 2.14×10^{-9} cm^{3}s^{-1}
$^{6}Li^{2+} \to {}^{6}Li^{1+}$: 4.81×10^{-9} cm^{3}s^{-1}
$^{6}Li^{3+} \to {}^{6}Li^{2+}$: 7.72×10^{-9} cm^{3}s^{-1}

$$l_{i\to i-1} = V_{i\to i-1} n_{gas}$$

n_{gas}: 1.44×10^{10} cm^{-3}

Atomic Excitation Rate by Electron Impact (1/2)

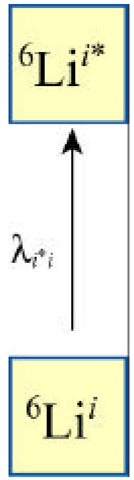
- $^{6}Li^{0+} \longrightarrow {}^{6}Li^{0+*}$ 2s→2p (including cascade)

  D. Leep and A. Gallagher, Phys. Rev. A 10 (1974) 1082.

  $$s \sim 3.5 p a_0^2 = 3.1\times10^{-16} \quad [\mathrm{cm}^2]\,\mathrm{at}\,T_e \sim 600\,\mathrm{eV}$$

  $$s v_e = 4.5\times10^{-7} \quad [\mathrm{cm}^3\mathrm{s}^{-1}] \qquad l_{0\to0*} = s v_e n_e$$

  a factor of ~10 larger than the ionization rate coefficient



- $^{6}Li^{1+} \longrightarrow {}^{6}Li^{1+*}$ 1s→2p

  assume that a factor of ~5 larger than the ionization rate coefficient

$$s v_e = 1.6\times10^{-8} \quad [\mathrm{cm}^3\mathrm{s}^{-1}] \qquad l_{1\to1*} = s v_e n_e$$

Atomic Excitation Rate by Electron Impact (2/2)

- $^{6}Li^{2+} \rightarrow {}^{6}Li^{2+*}$ 1s→2p

  Fisher *et al.*, Phys. Rev. A 55 (1997) 329.

  Empirical fit of 1s→2p excitation cross sections of hydrogen-like atoms

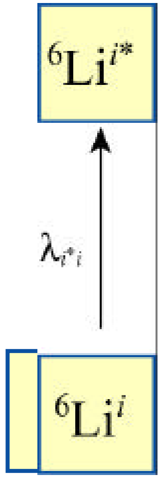
  $$\sigma \sim 1.0\pi a_0^2 Z_i^{-4} = 1.1\times10^{-18} \quad [\mathrm{cm}^2]\ \mathrm{at}\ T_e \sim 550\,\mathrm{eV}$$

  $$\sigma v_e = 1.6\times10^{-9} \quad [\mathrm{cm}^3\mathrm{s}^{-1}] \qquad \lambda_{2\rightarrow2^*} = \sigma v_e n_e$$

  Summing up transitions 1s→2,…,6 and taking the Boltzmann distribution

  $$\langle \sigma v_e \rangle = 1.82\times10^{-9} \quad [\mathrm{cm}^3\mathrm{s}^{-1}]$$

  a factor of ~2 larger than the ionization rate coefficient



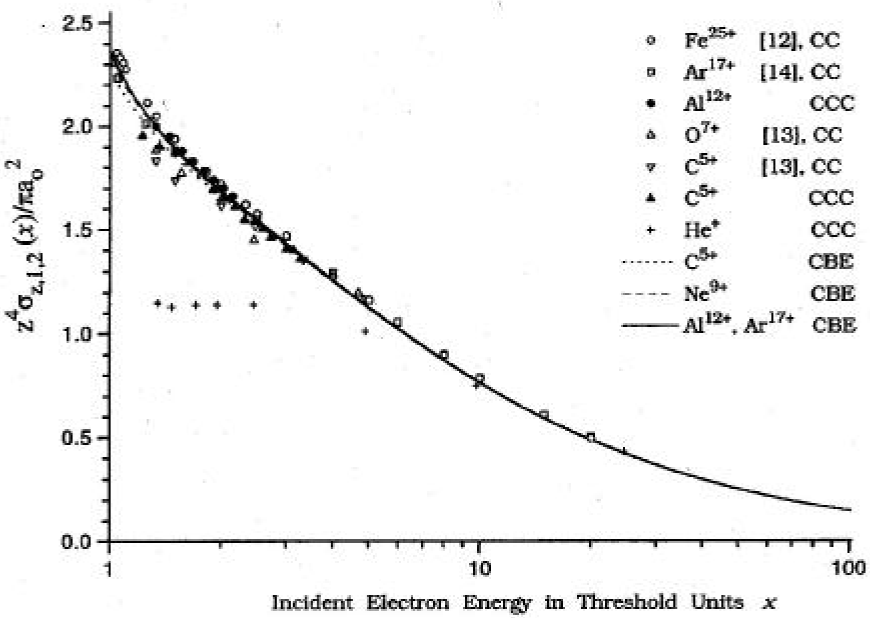
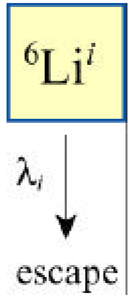


FIG. 3. Cross sections for transition 1→2 in hydrogenlike ions with $z=2-26$.

Confinement Time of The Ions



- It is very difficult to estimate the confinement time of ions in an ECR plasma.

  If we assume (M.Imanaka, PhD Thesis; Shirkov, CERN/PS 94-13)

  $$t_i \propto i\sqrt{A_i}$$

  and scale the value of τ_{3+}=2.3msec, which was fitted to the Al data,

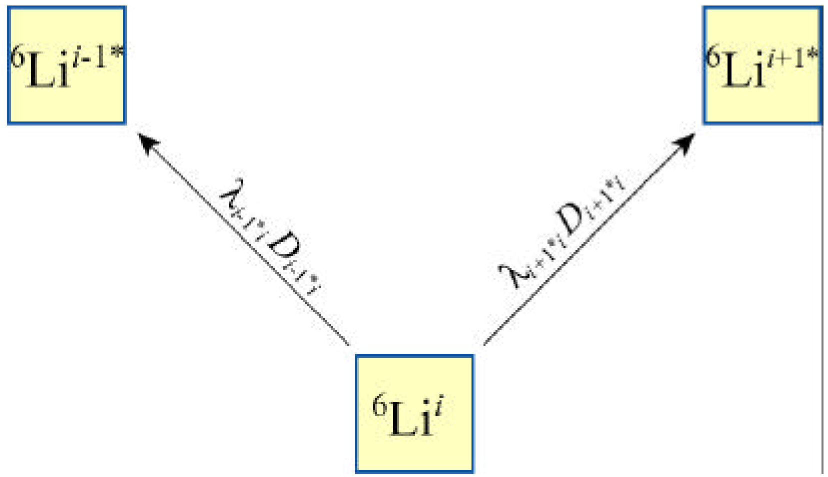
  $$t_{1+} = 0.33[\mathrm{ms}]$$
  $$t_{2+} = 0.66[\mathrm{ms}]$$
  $$t_{3+} = 0.99[\mathrm{ms}]$$

  $$l_i = t_i^{-1}$$

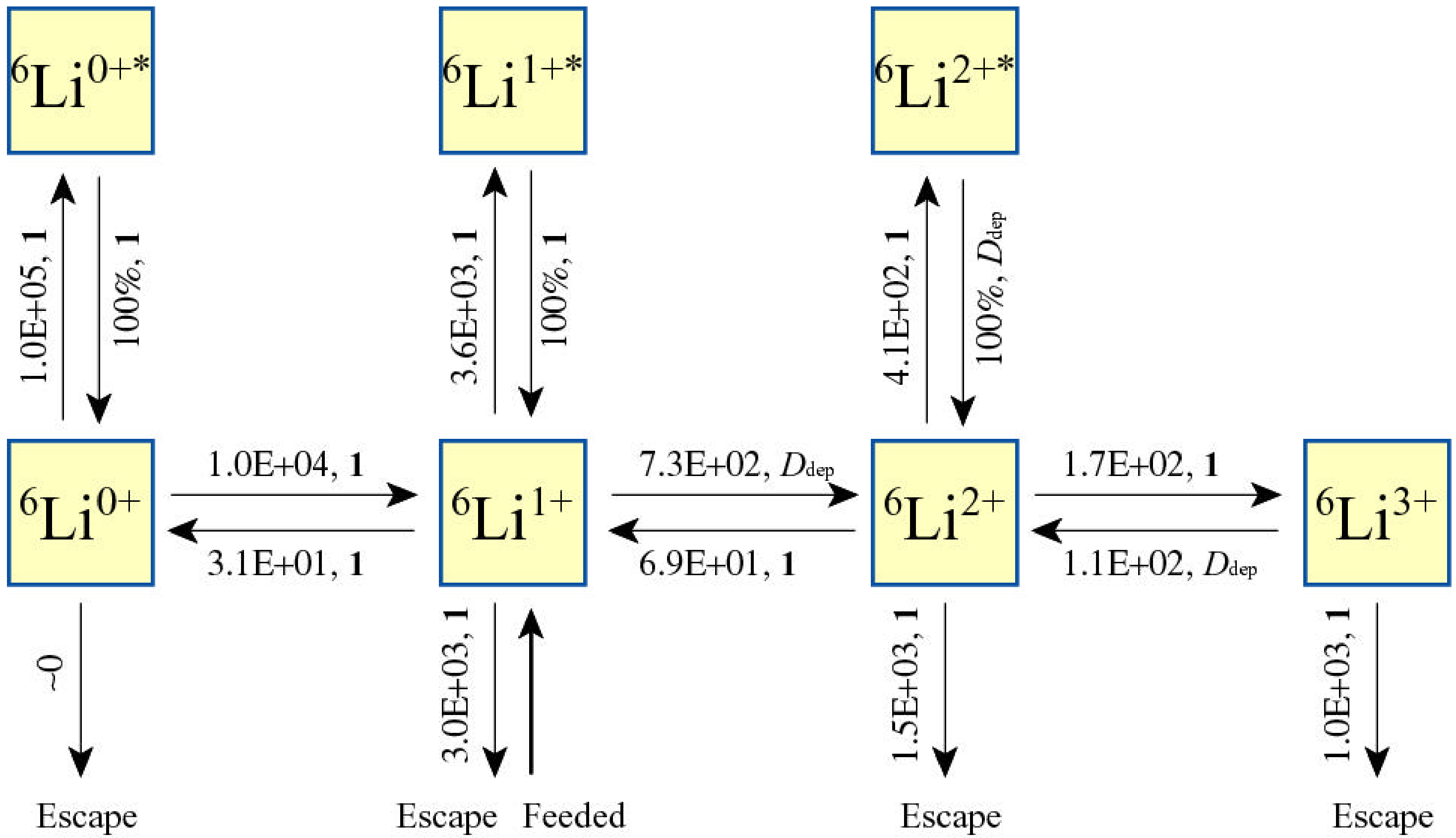
Other processes

Inelastic Ionization and Radiative Capture Processes

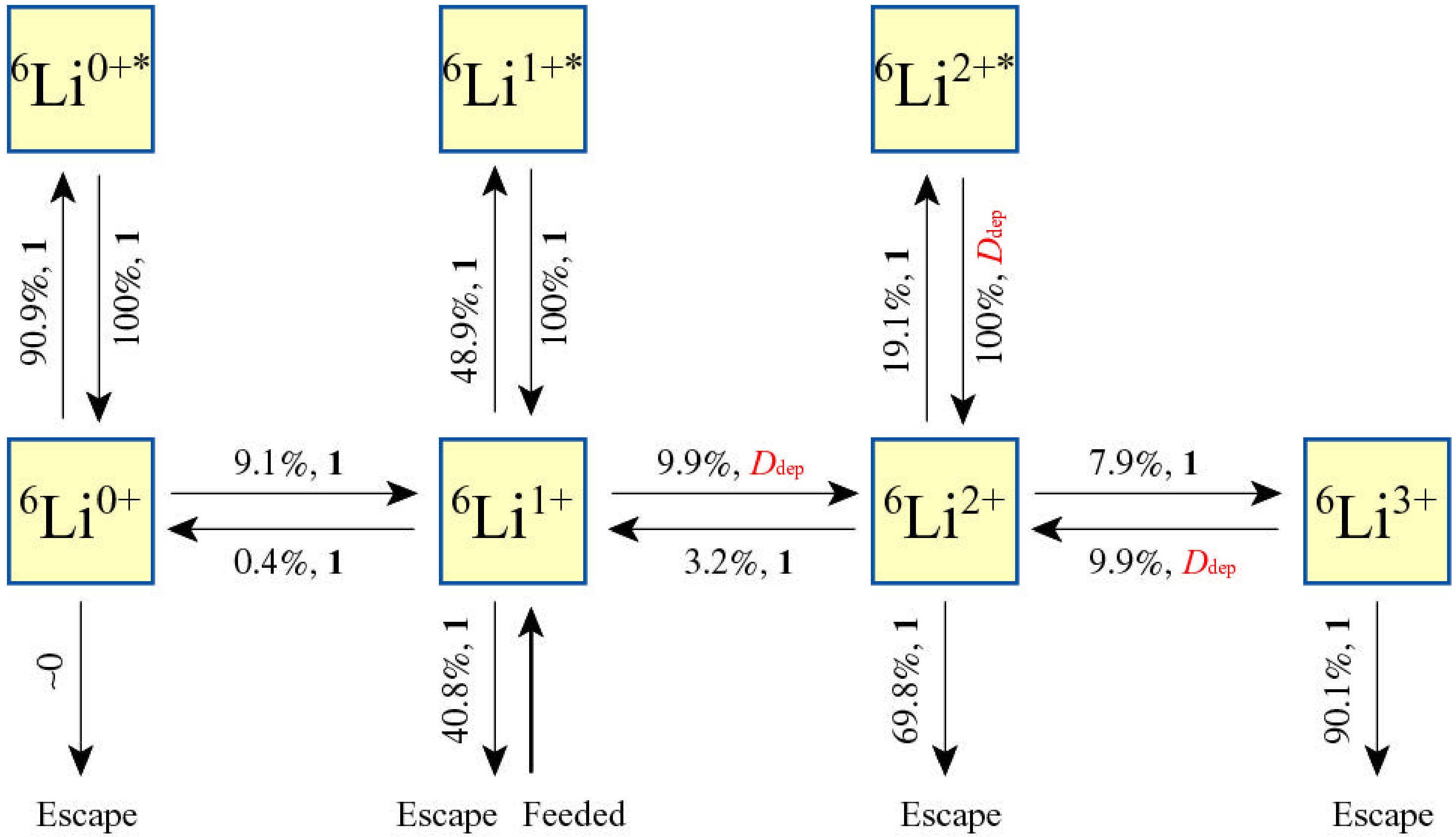
In the present calculation, these processes has no (or negligible) effect.

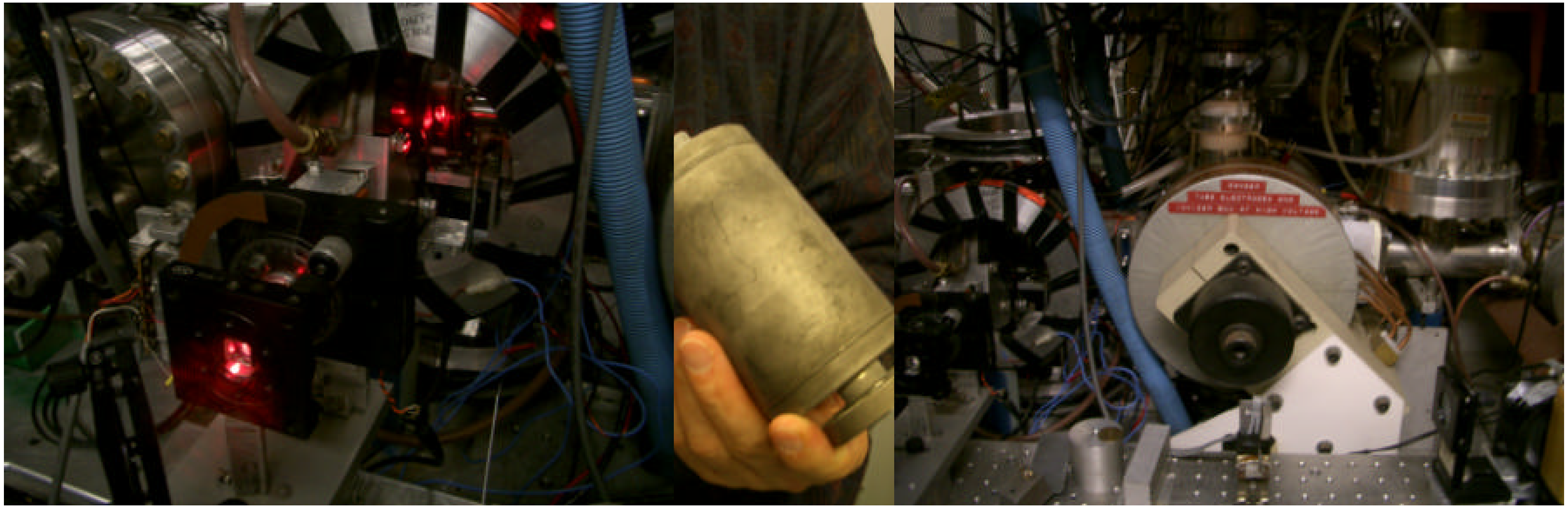


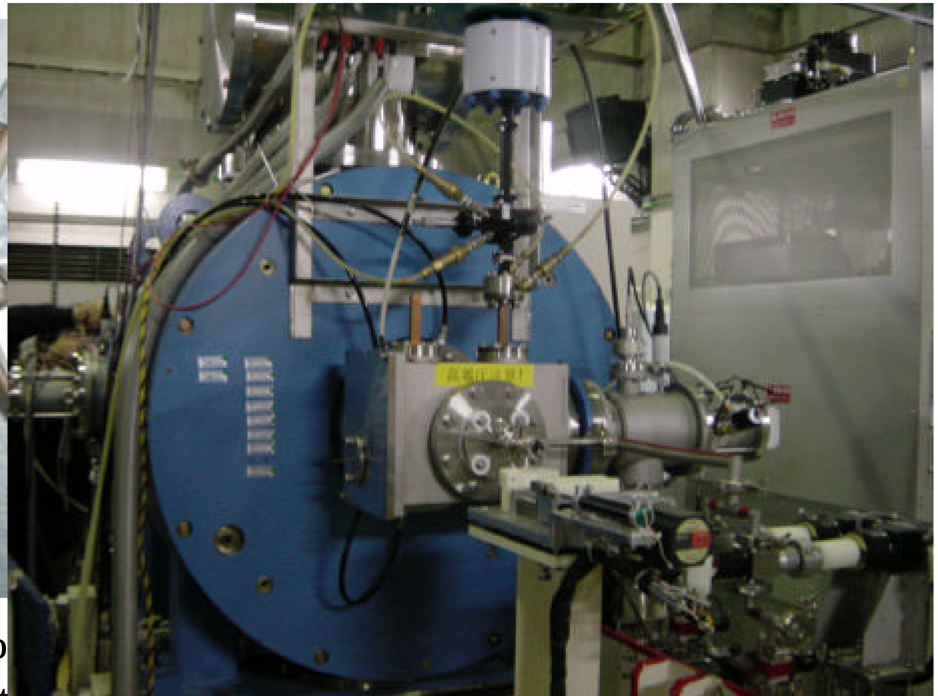
Summary of the Processes in the ECR Ionizer



Summary of the Processes in the ECR Ionizer







Pol. 80-90%
at Florida State Univ.

Feasibility Test Plan

- Study of confinement time and ionization efficiency of Li is planed by using the 18GHz superconducting ECR ion source at RIKEN and the laser abration method.

- Optimization of the plasma condition:

  Mirror ratio, neutral gass density, RF power

- Development of the Li-oven, surface ionizer for testing the beam current.

- Laser pumping system for testing the polarization of the $^{6}Li^{3+}$ beam

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- Further simulation with more realistic parameters is required.