

Studies of Gamow-Teller transitions using Weak and Strong Interactions



High-resolution
Spectroscopy
&
Tensor Interaction

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Neptune driving Waves

Neptune =
weak interaction
→ β decay



Powerful Waves = strong interaction)
→ Charge-Exchange Reaction

Neptune and the waves, or "steeds," he rides.

— Walter Crane, 1892

Gamow-Teller transitions

Mediated by $\sigma\tau$ operator: both **S & W int.** has this *Op.*

$$\Delta S = -1, 0, +1 \quad \text{and} \quad \Delta T = -1, 0, +1$$

($\Delta L = 0$, no change in radial w.f.)

→ no change in spatial w.f.

Accordingly, transitions among $j_>$ and $j_<$ configurations

$$j_> \rightarrow j_>, \quad j_< \rightarrow j_<., \quad j_> \leftrightarrow j_<$$

example $f_{7/2} \rightarrow f_{7/2}$, $f_{5/2} \rightarrow f_{5/2}$, $f_{7/2} \leftrightarrow f_{5/2}$

Note that **Spin** and **Isospin** are
unique quantum numbers in atomic nuclei !

→ GT transitions are sensitive to Nuclear Structure !

→ GT transitions in each nucleus are **UNIQUE** !

**Basic common understanding of β -decay and Charge-Exchange reaction

β decays :

Absolute $B(\text{GT})$ values,

but usually the study is limited to low-lying states

(p,n) , $(^3\text{He},t)$ reaction at 0° :

Relative $B(\text{GT})$ values, but **Highly Excited States**

** Both are important for the study of GT transitions!

β -decay & Nuclear Reaction

* β -decay GT tra. rate = $\frac{1}{t_{1/2}} = f \frac{\lambda^2}{K} B(\text{GT})$

$B(\text{GT})$: reduced GT transition strength
 $\propto (\text{matrix element})^2 = |\langle f | \sigma \tau | i \rangle|^2$

*Nuclear (CE) reaction rate (cross-section)

= reaction mechanism

(\otimes) operator

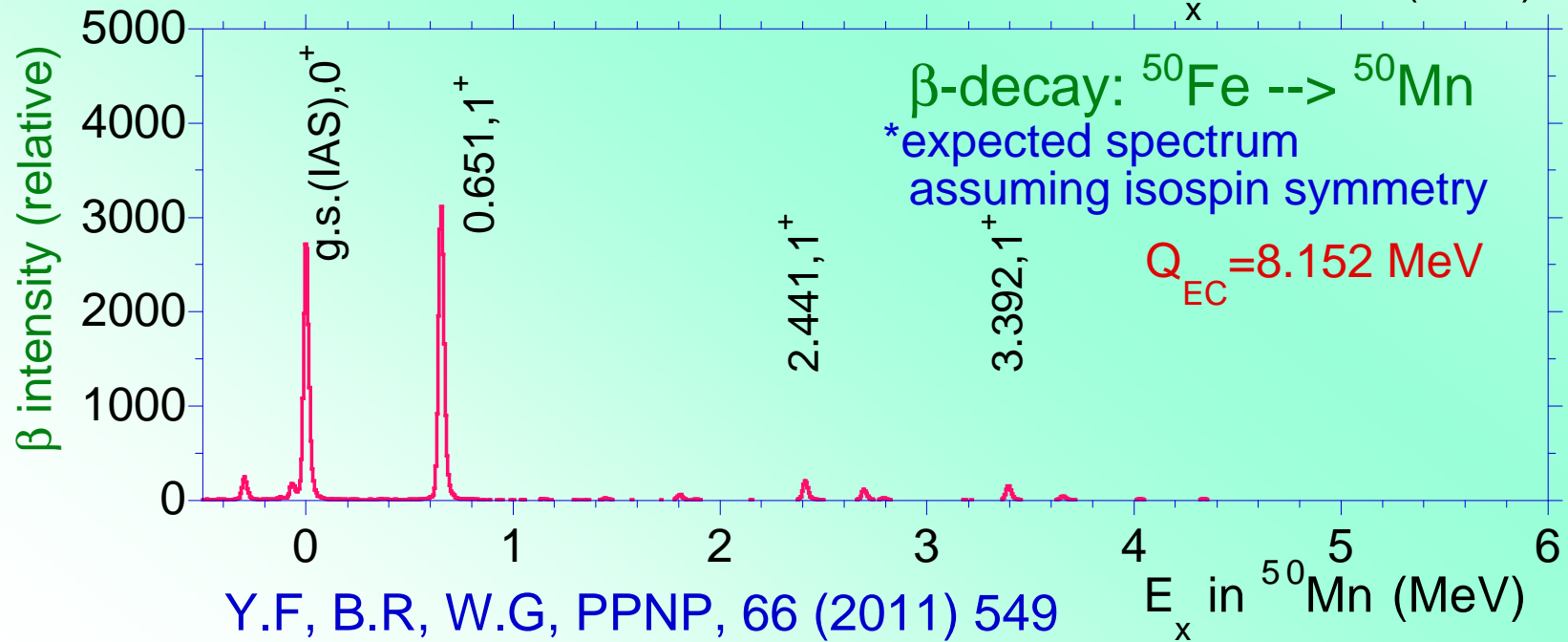
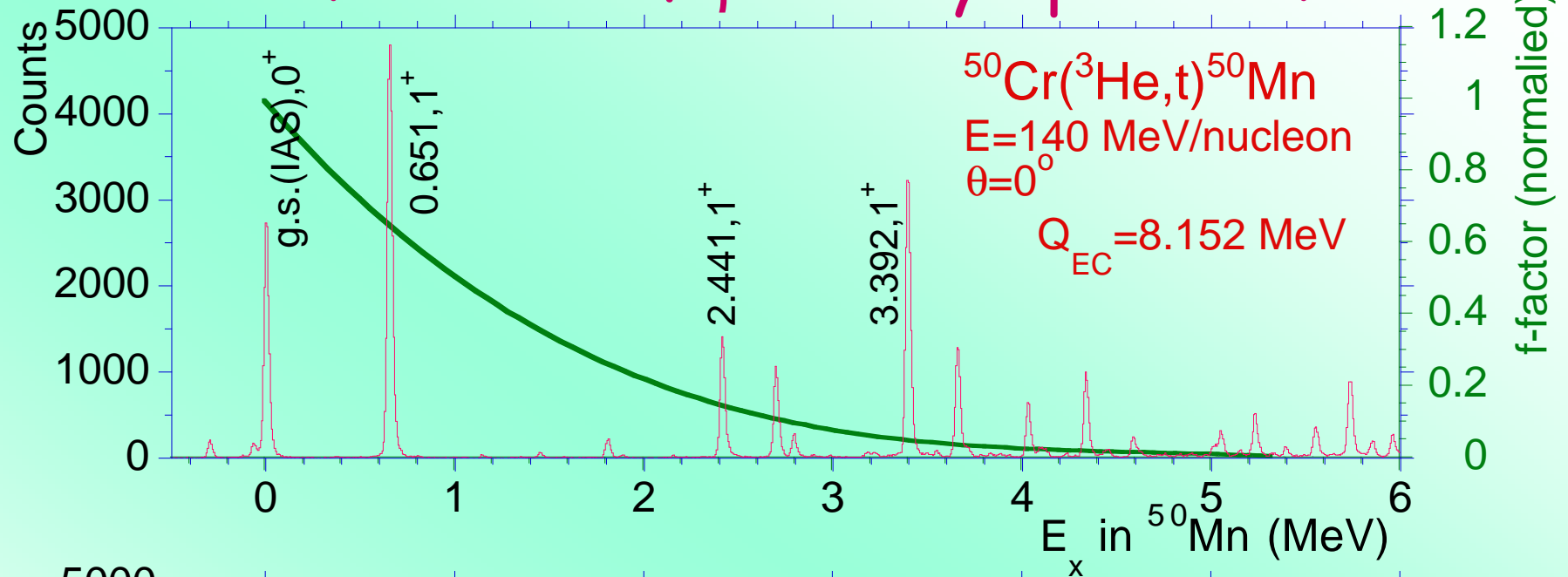
(\otimes) structure

=(matrix element)²

*At intermediate energies ($100 < E_{\text{in}} < 500$ MeV)

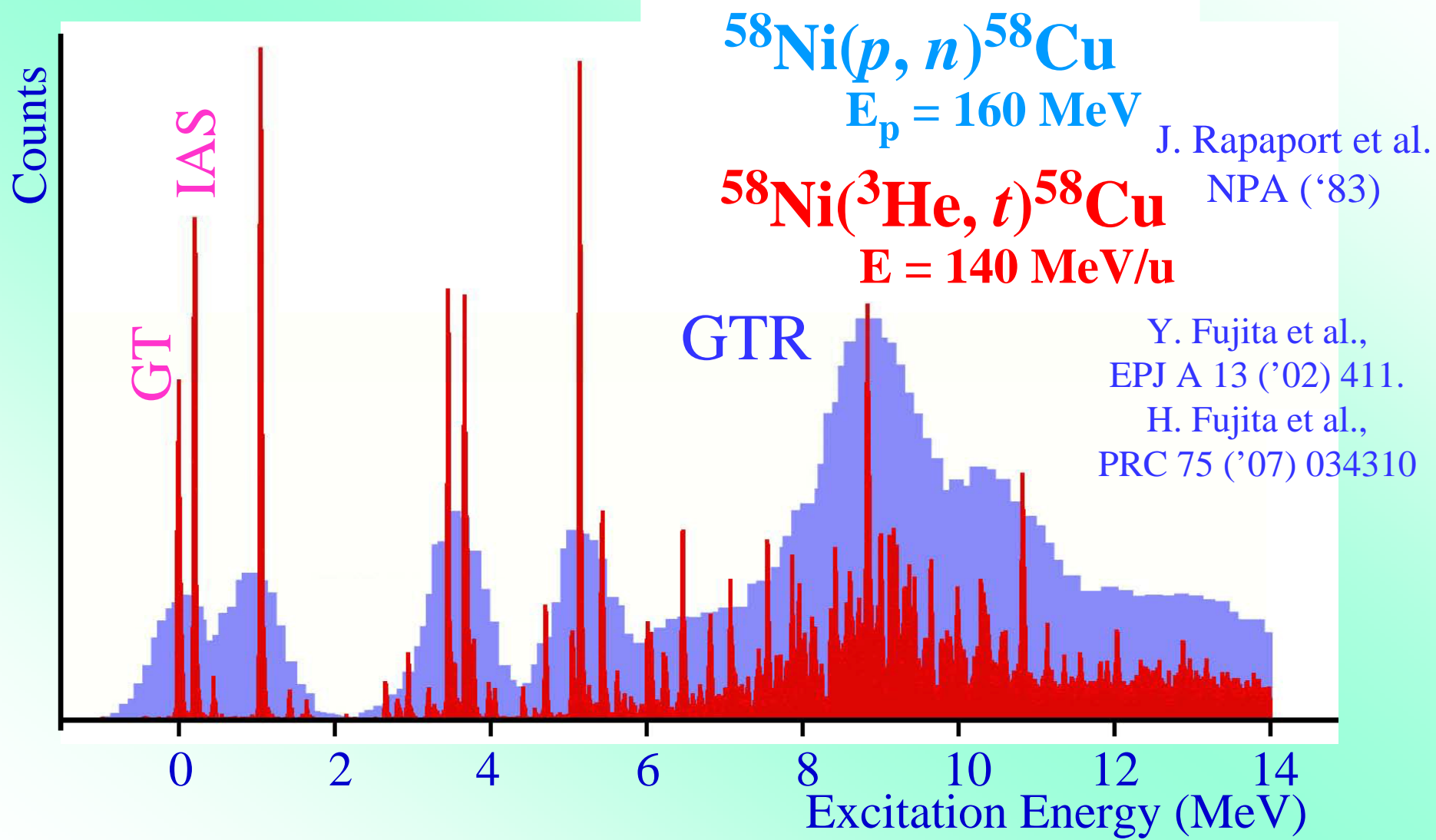
→ $d\sigma/d\omega(q=0)$: proportional to $B(\text{GT})$

Simulation of β -decay spectrum

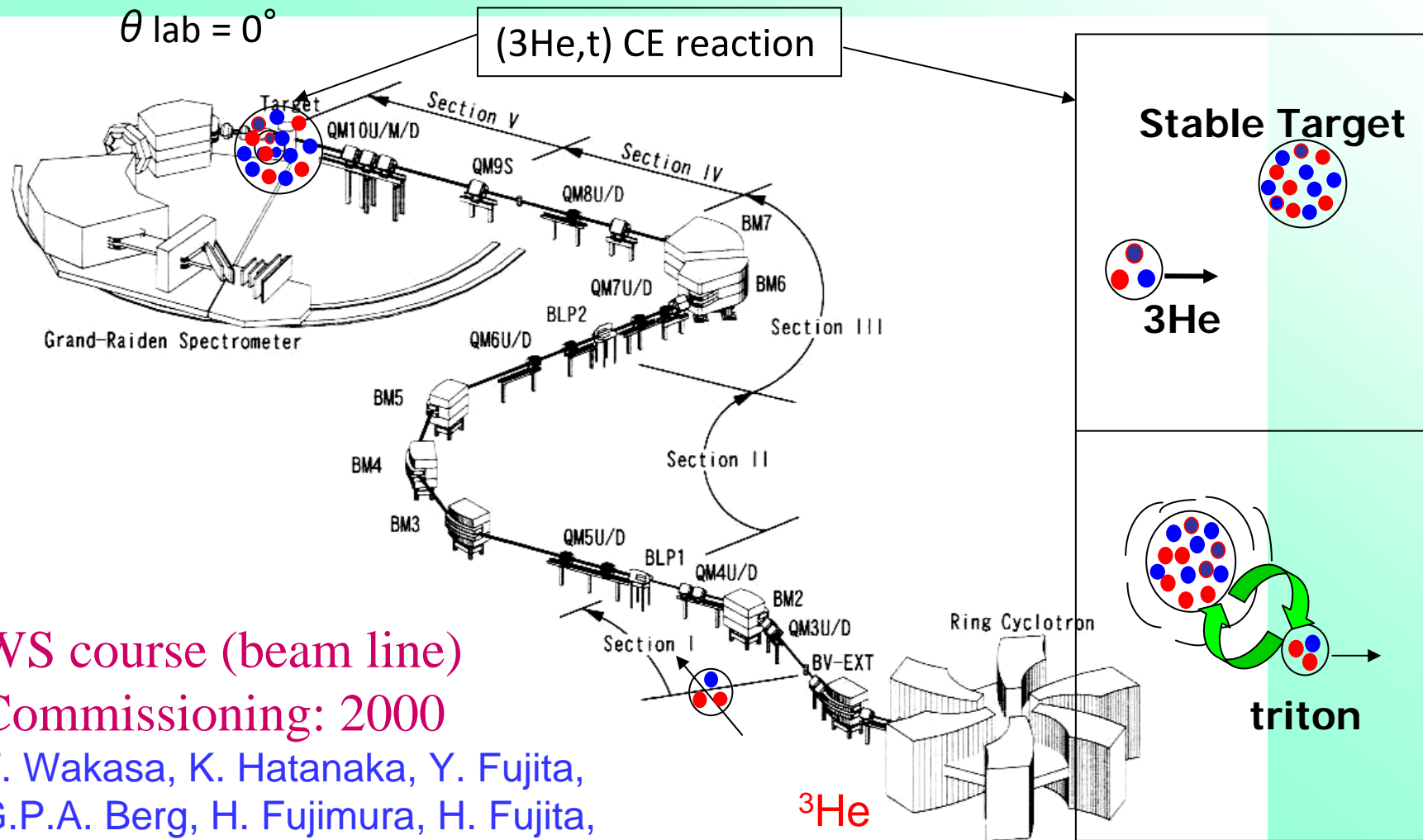


Y.F, B.R, W.G, PPNP, 66 (2011) 549

Comparison of (p, n) and (³He, t) 0° spectra



$(^3\text{He}, t)$ CE Reactions @ RCNP (Osaka)



WS course (beam line)

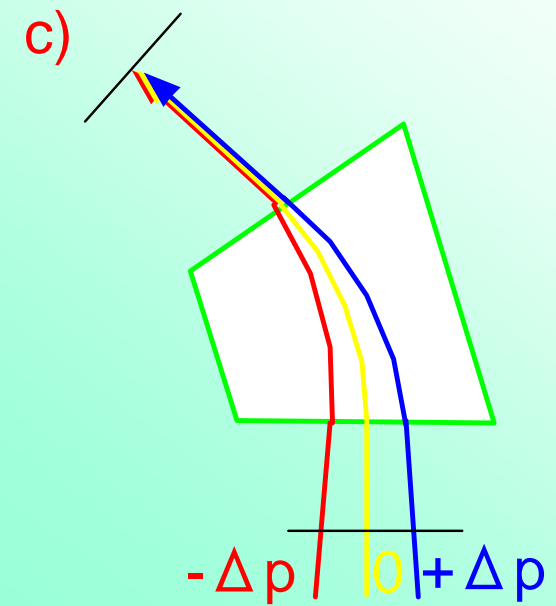
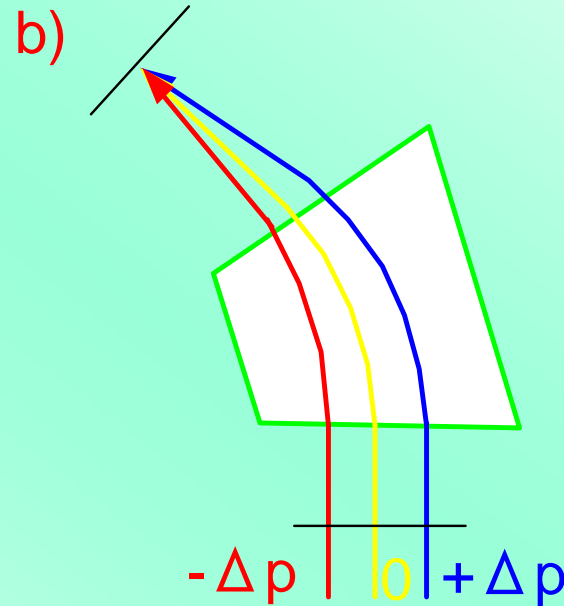
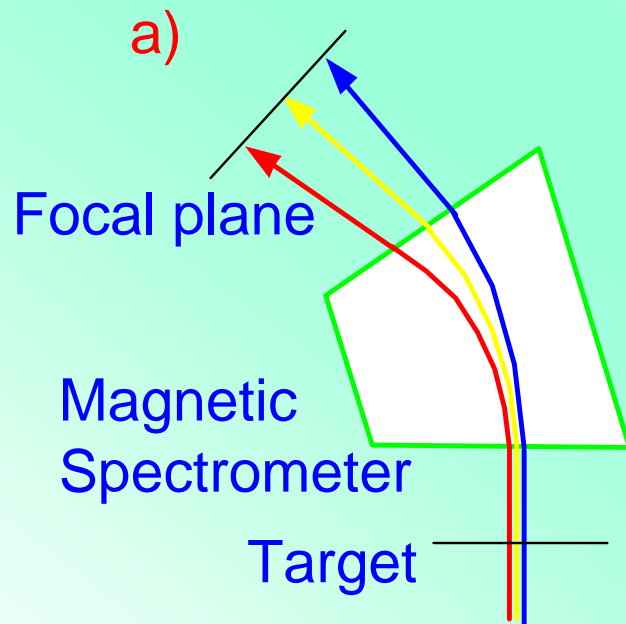
Commissioning: 2000

T. Wakasa, K. Hatanaka, Y. Fujita,
G.P.A. Berg, H. Fujimura, H. Fujita,
M. Itoh, J. Kamiya, T. Kawabata et al.,
N.I.M. A 482 (2002) 79.

Matching Techniques

Y. Fujita et al., N.I.M. B 126 (1997) 274.

H. Fujita et al., N.I.M. A 484 (2002) 17.



*Achromatic beam
transportation*

$\Delta E \sim 200$ keV
for 140MeV/u ^3He beam

*Lateral dispersion
matching*

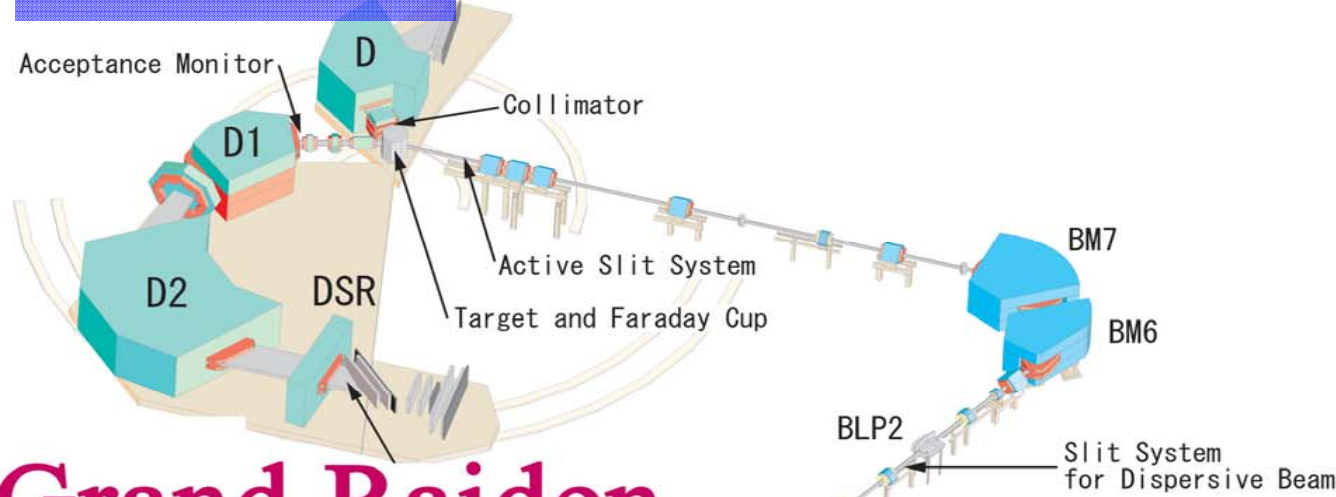
$\Delta E \sim 35$ keV
Horiz. angle resolution
 $\Delta\theta_{\text{sc}} > 15$ mrad

*Angular dispersion
matching*

$\Delta\theta_{\text{sc}} \sim 5$ mrad

$\Delta E = 30 \text{ keV}$

RCNP, Osaka Univ.



Grand Raiden

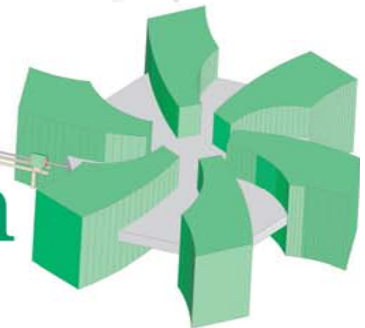
WS Beam Line

Dispersion Matching Techniques were applied!

Y. Fujita et al, NIM B 126 (1997) 274.
H. Fujita et al, NiM A 484 (2002) 17.

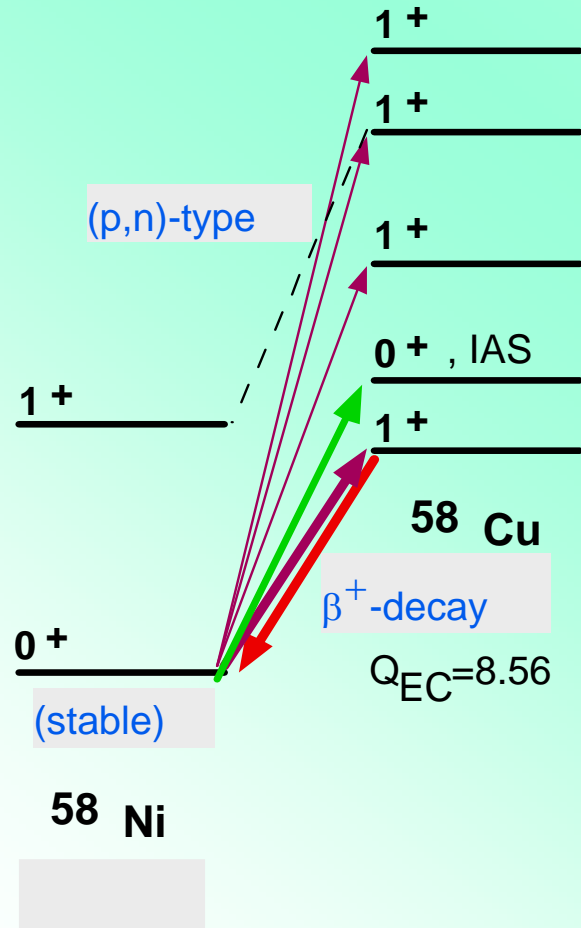
$\Delta E = 150 \text{ keV}$

Ring Cyclotron

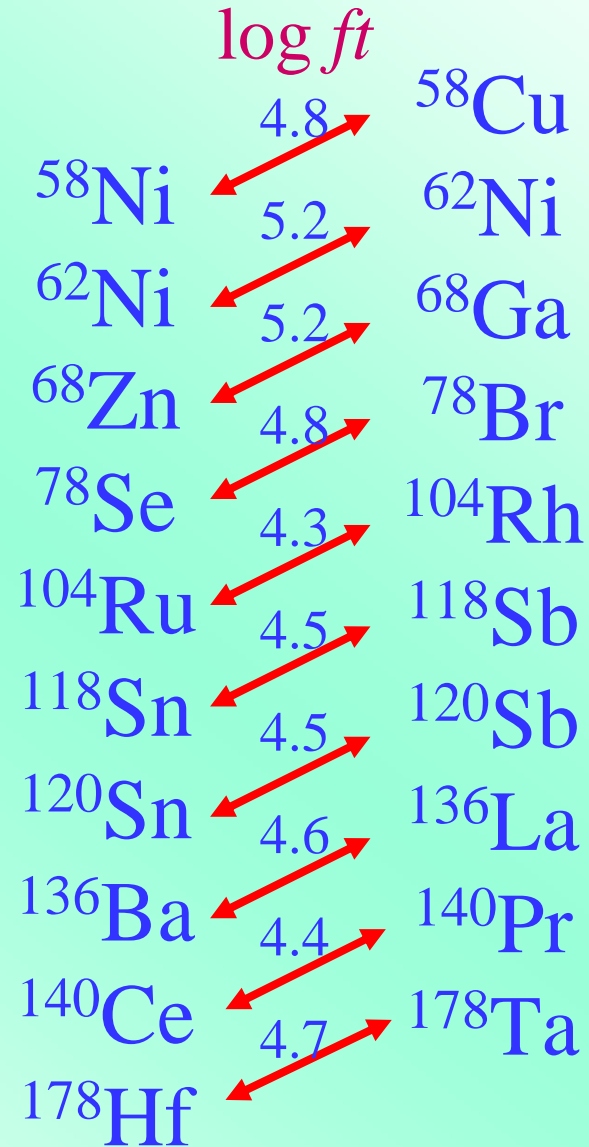


Connection: Charge Exchange & β decay

$$0^+ \leftrightarrow 1^+$$



** 0^+ & 1^+ relationship of g.s.



***Isospin Symmetry

an important idea to see the connection of
decays and excitations caused
by Strong, EM and Weak interactions !

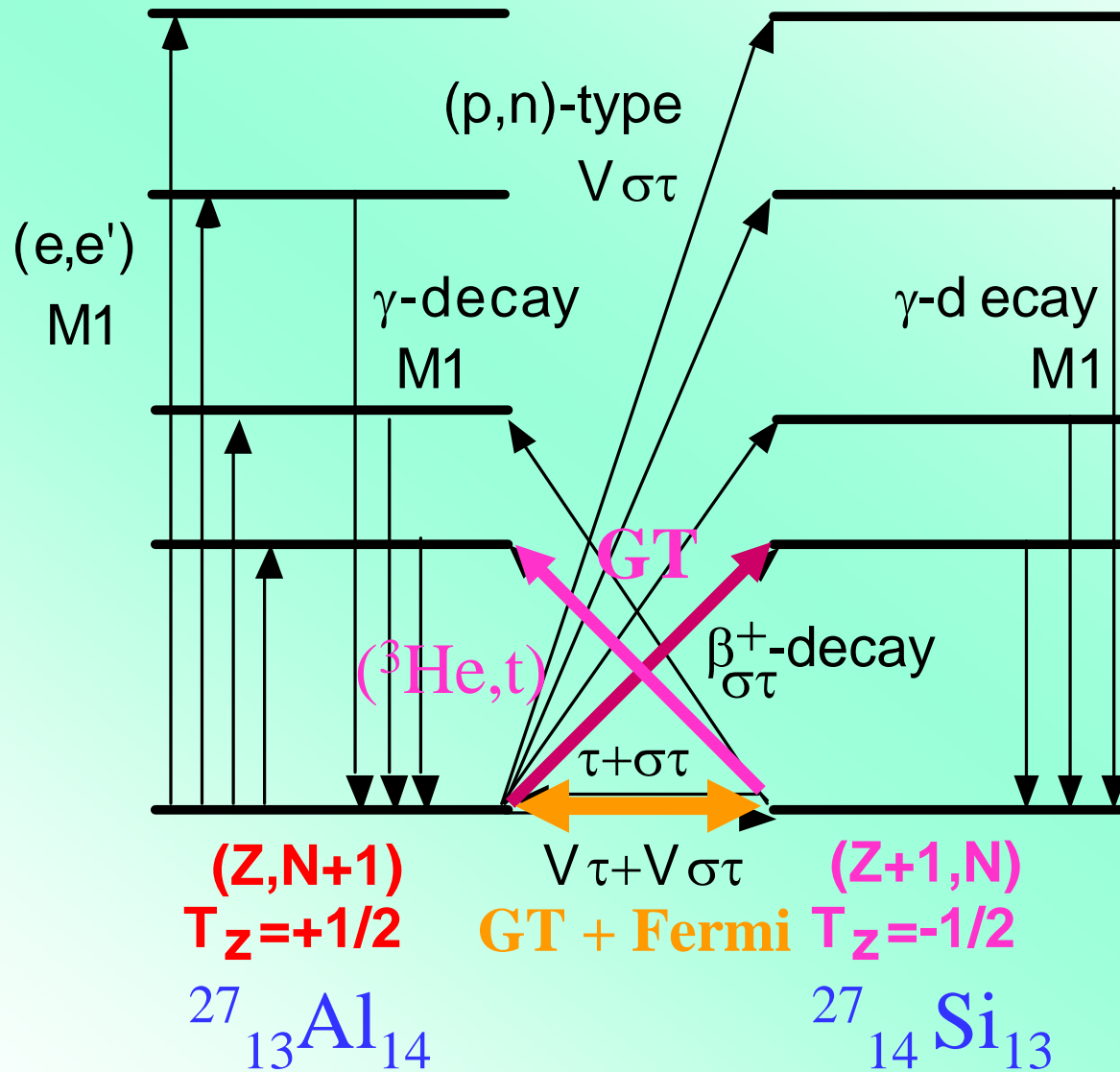
There are many cases that the “operators” are the same
in transitions caused by “strong,” “EM” and “weak” int.

$T=1/2$
Isospin
Symmetry

Koelner Dom
Koeln, Germany
(157m high)

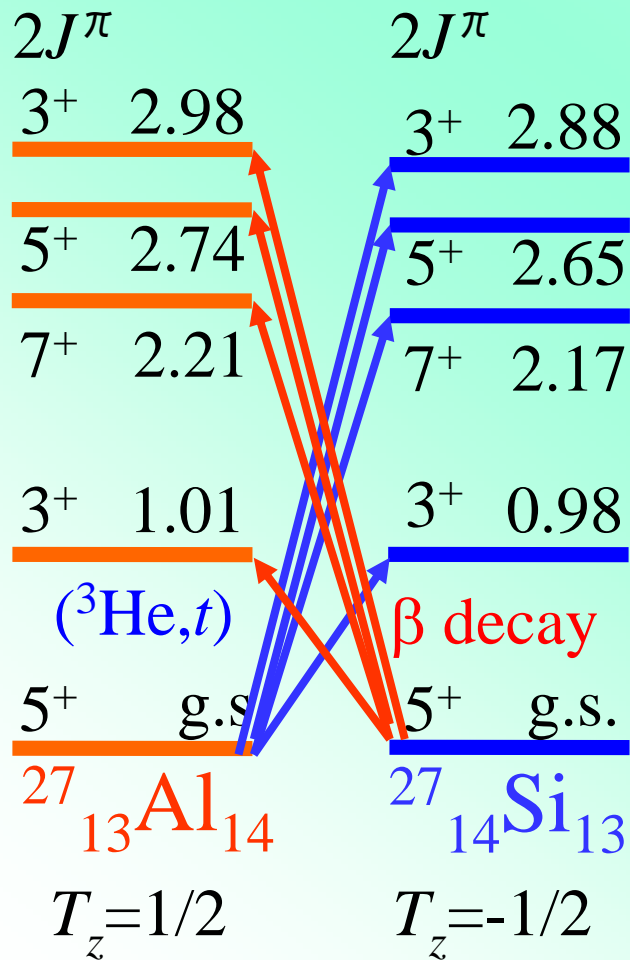


$T=1/2$ Mirror Nuclei : Structures & Transitions

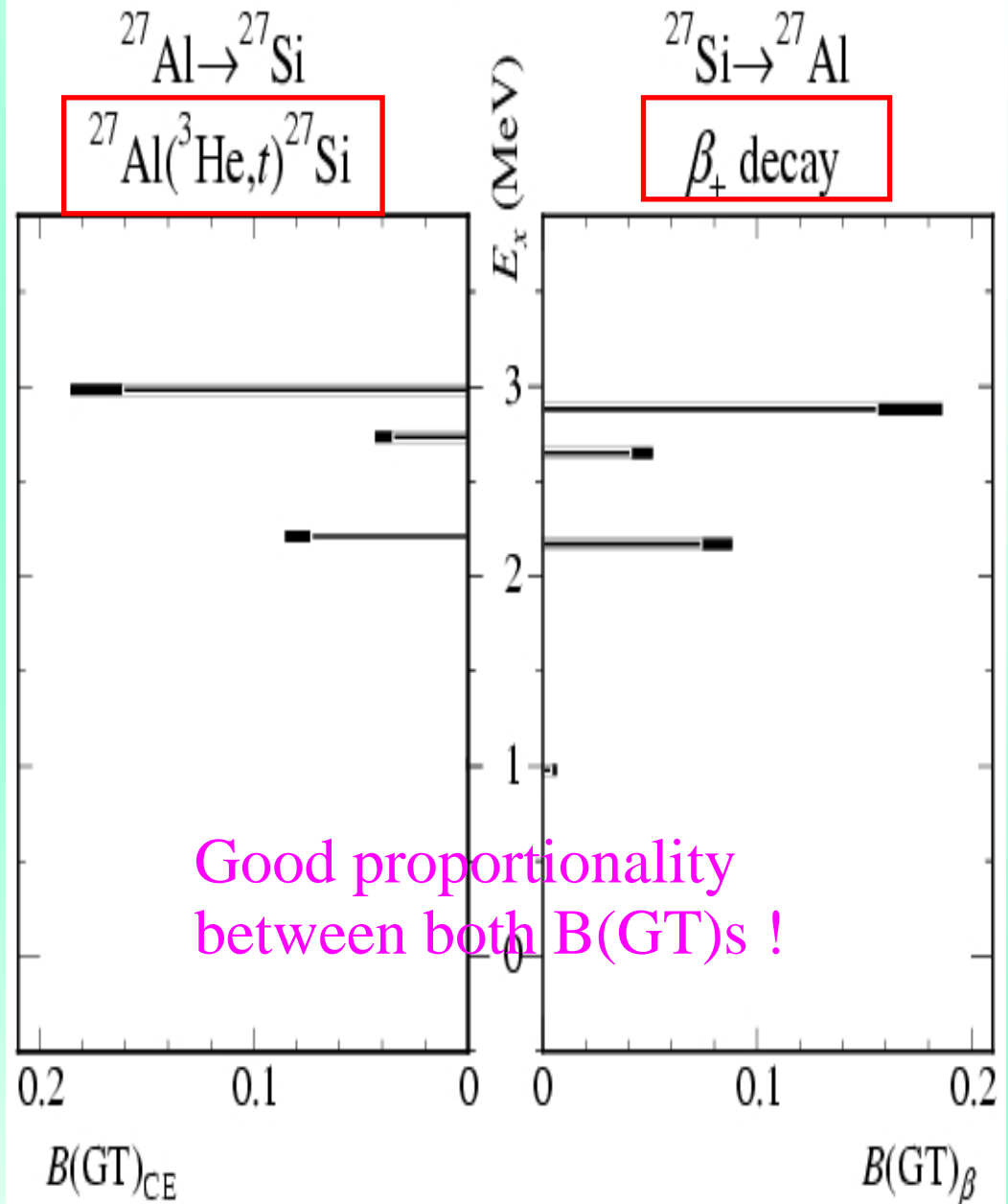


Symmetry in A=27 System

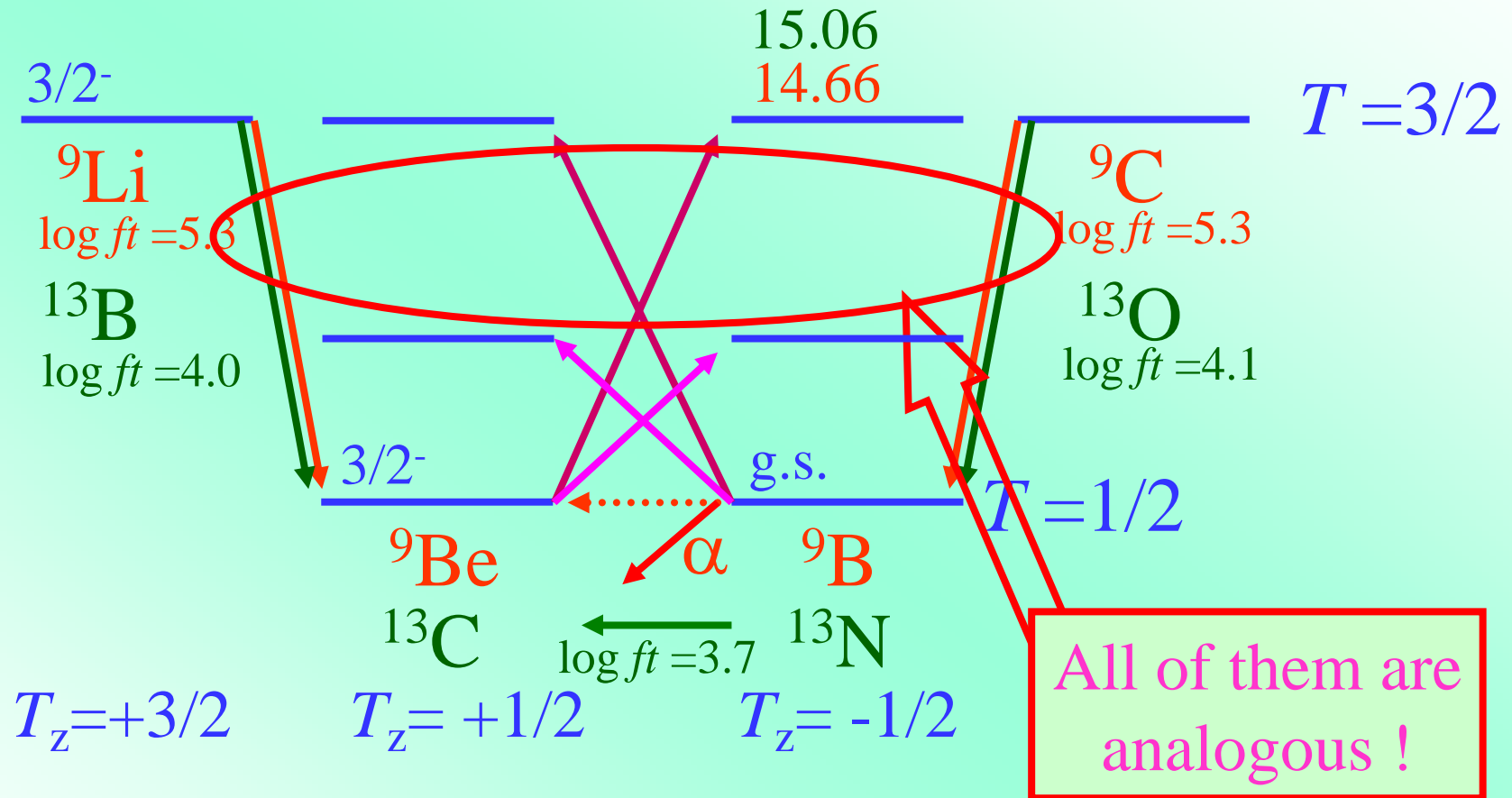
$$\frac{d\sigma}{d\Omega}(q \rightarrow 0) = KM |J|^2 B(\text{GT})$$



Experiments



Analogous relationship: A=9, 13 system



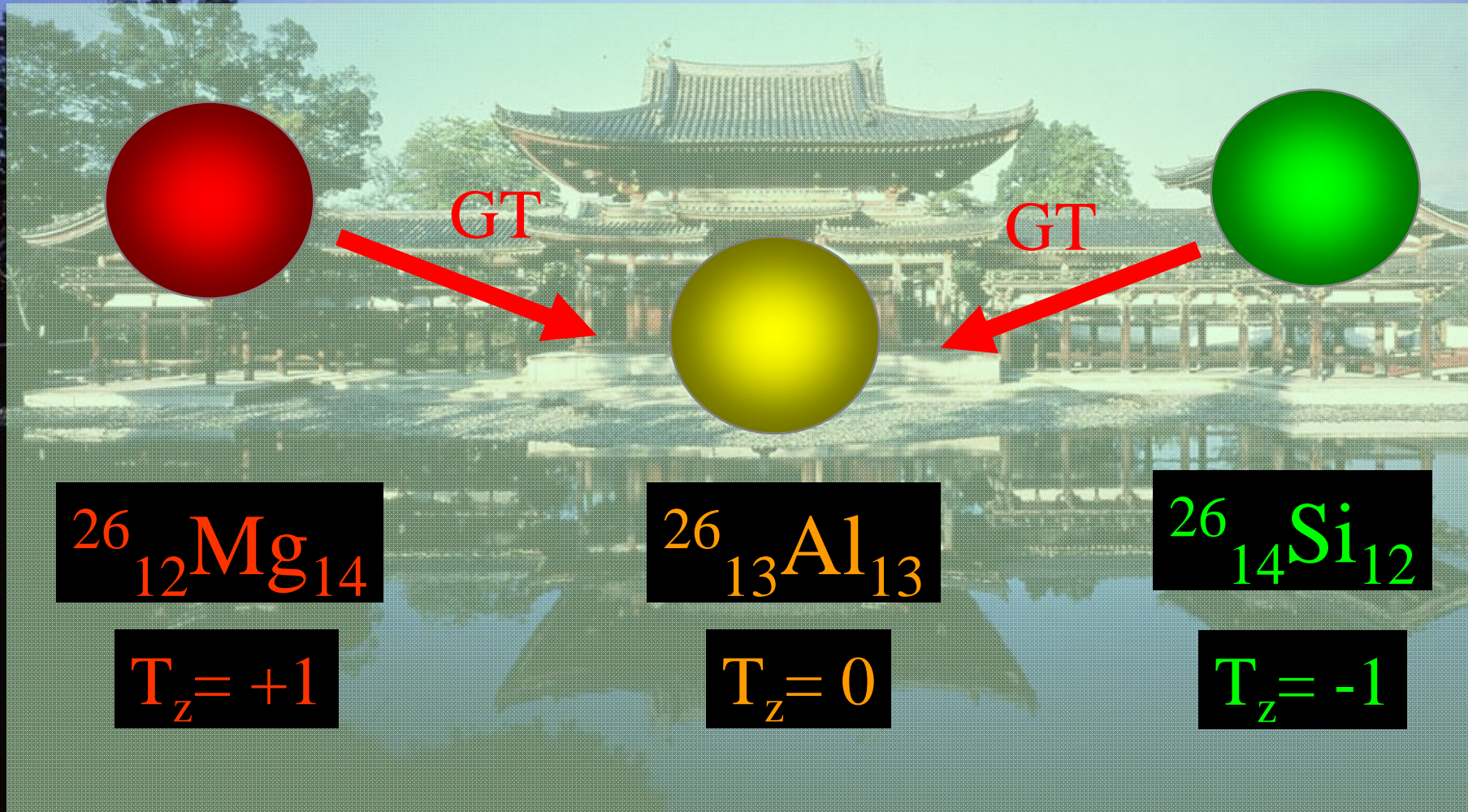
*Small isospin asymmetry can be seen
for $T_z=+3/2 \rightarrow +1/2$ and $T_z=-1/2 \leftarrow -3/2$
GT transitions.

T=1 Isospin Symmetry

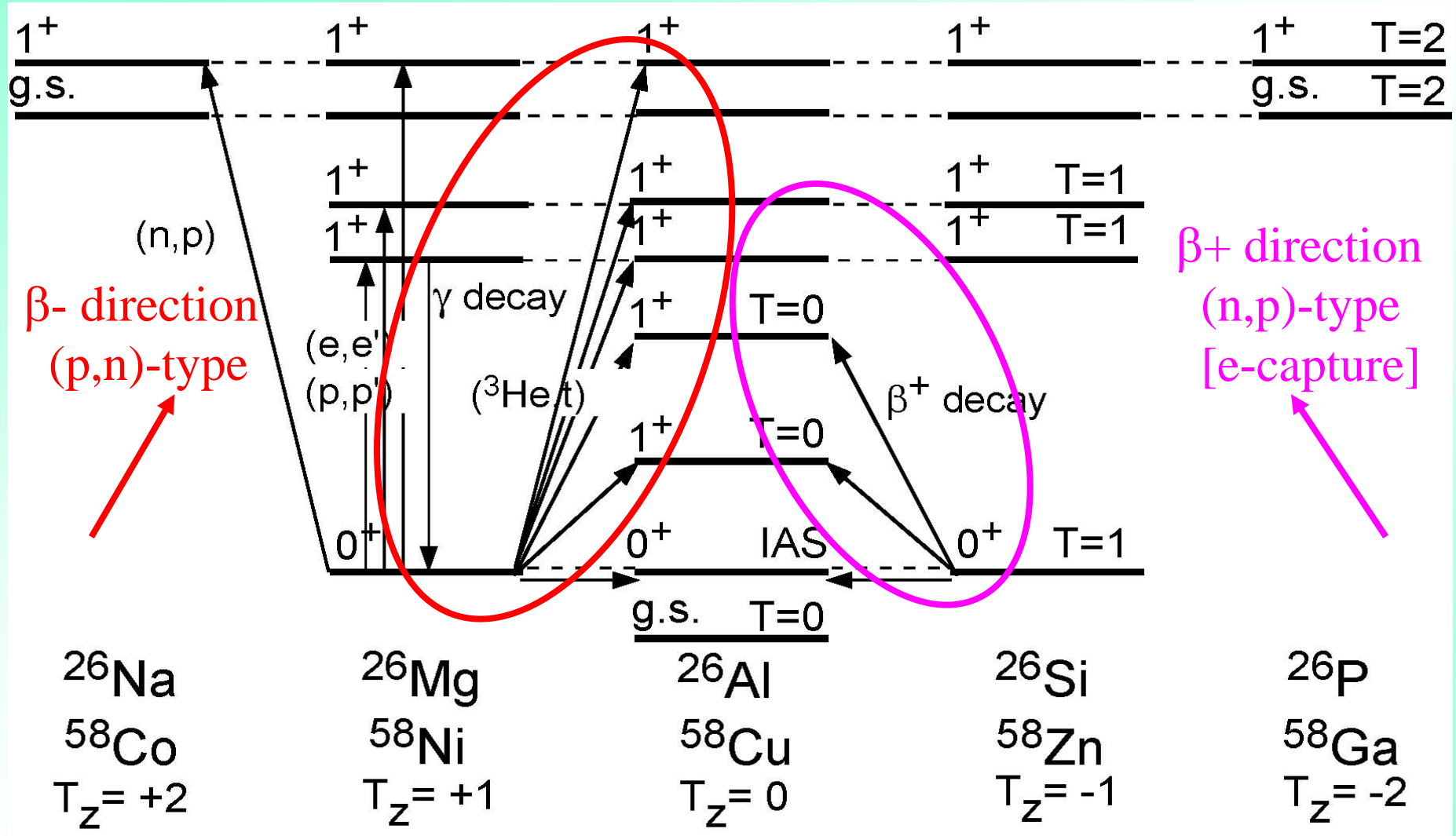


Byodoin-temple,
Uji, Kyoto

T=1 Isospin Symmetry

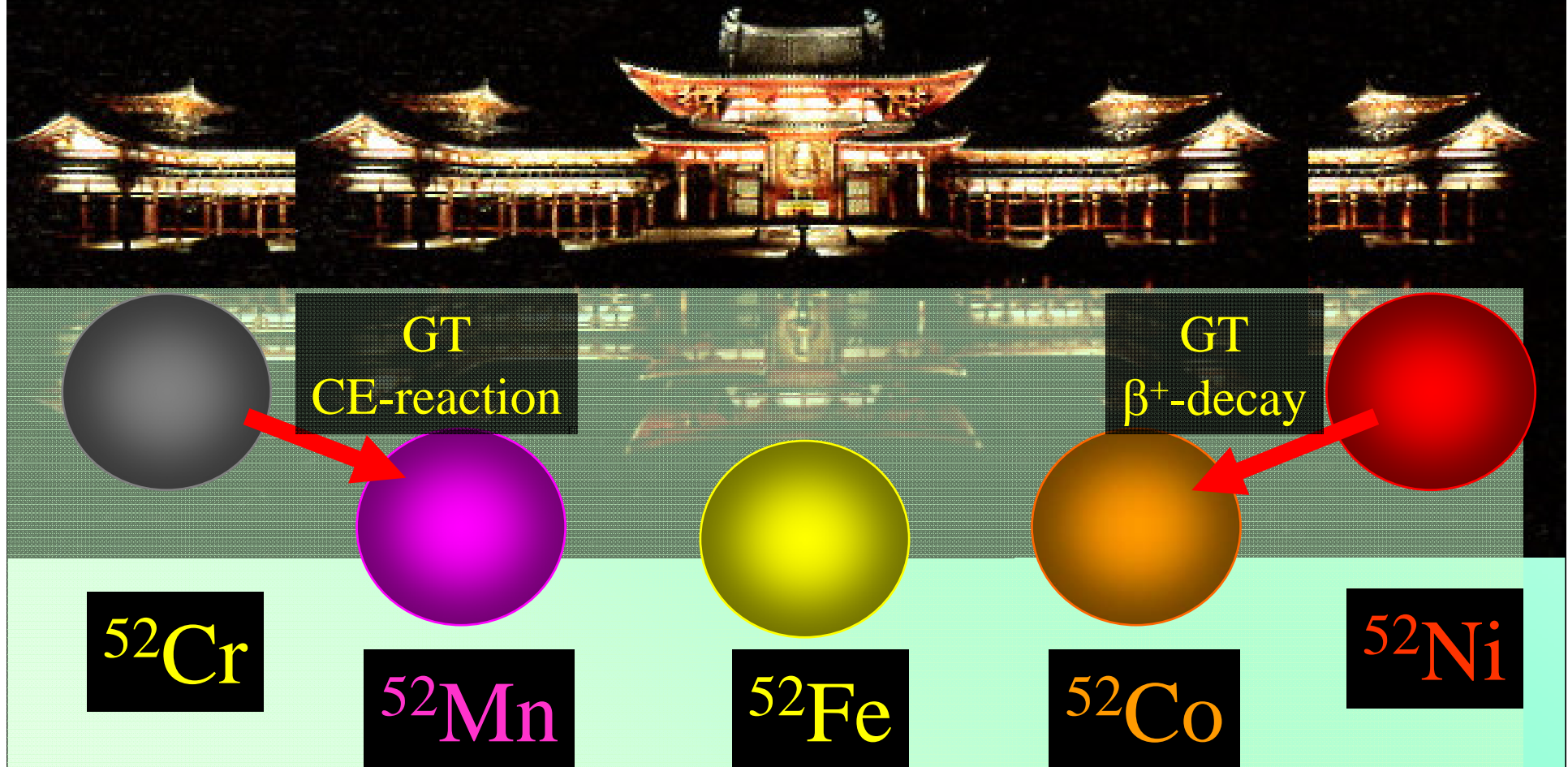


$T_z = +1 \rightarrow 0 \leftarrow -1$ Symmetry

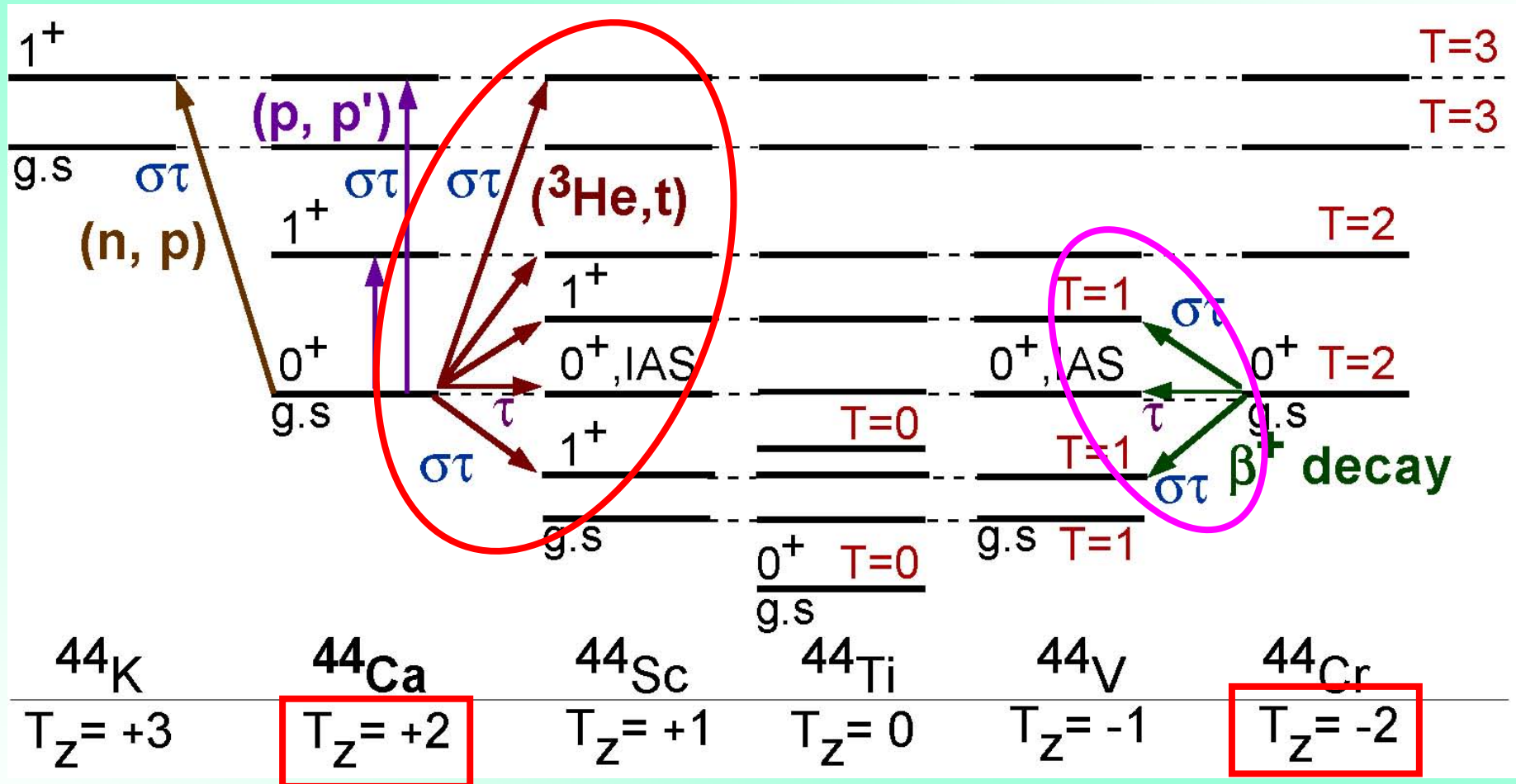


Super-Byodoin 平等院

T=2 Isospin Symmetry



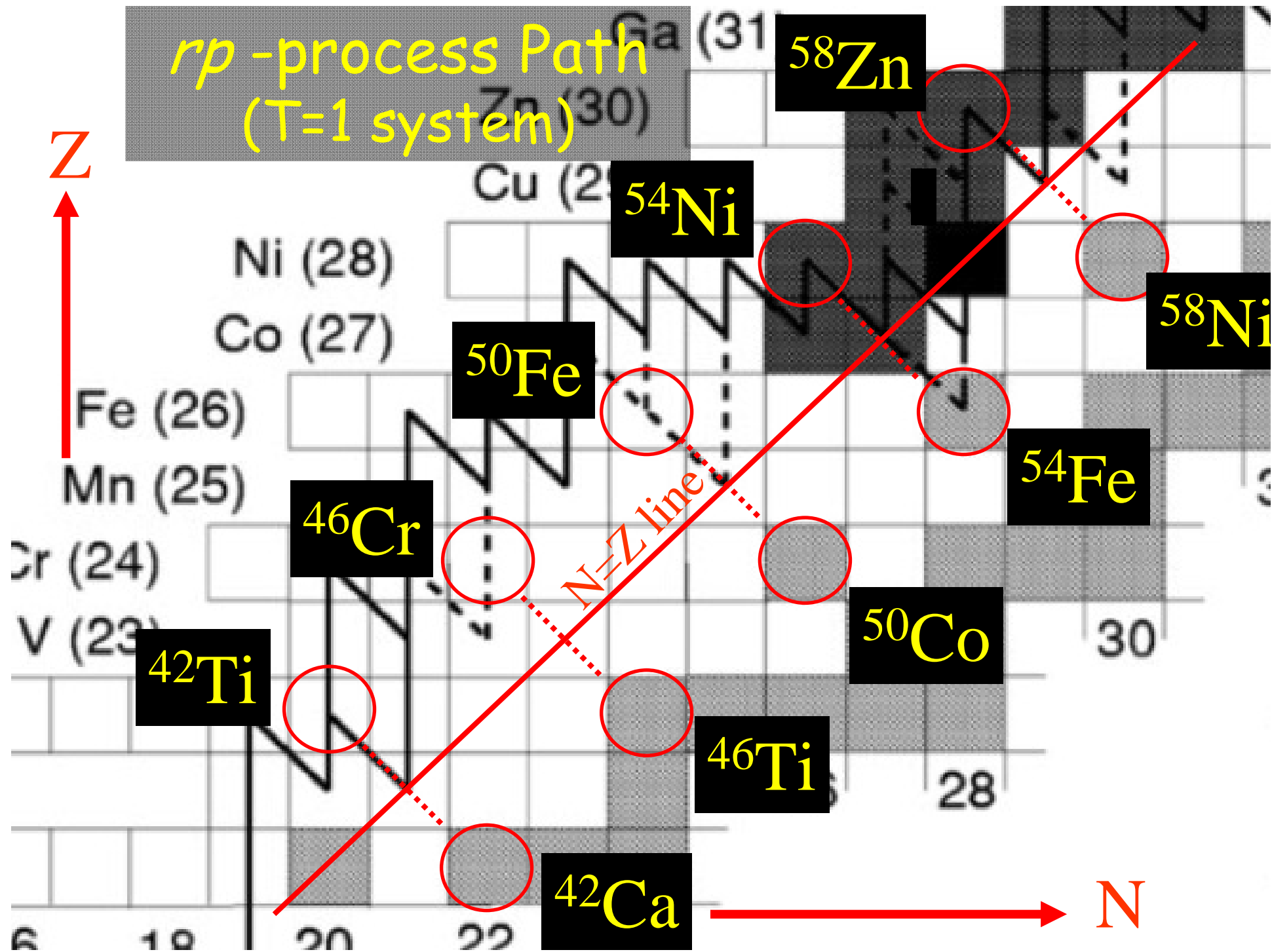
Isospin Structure of T=2 system



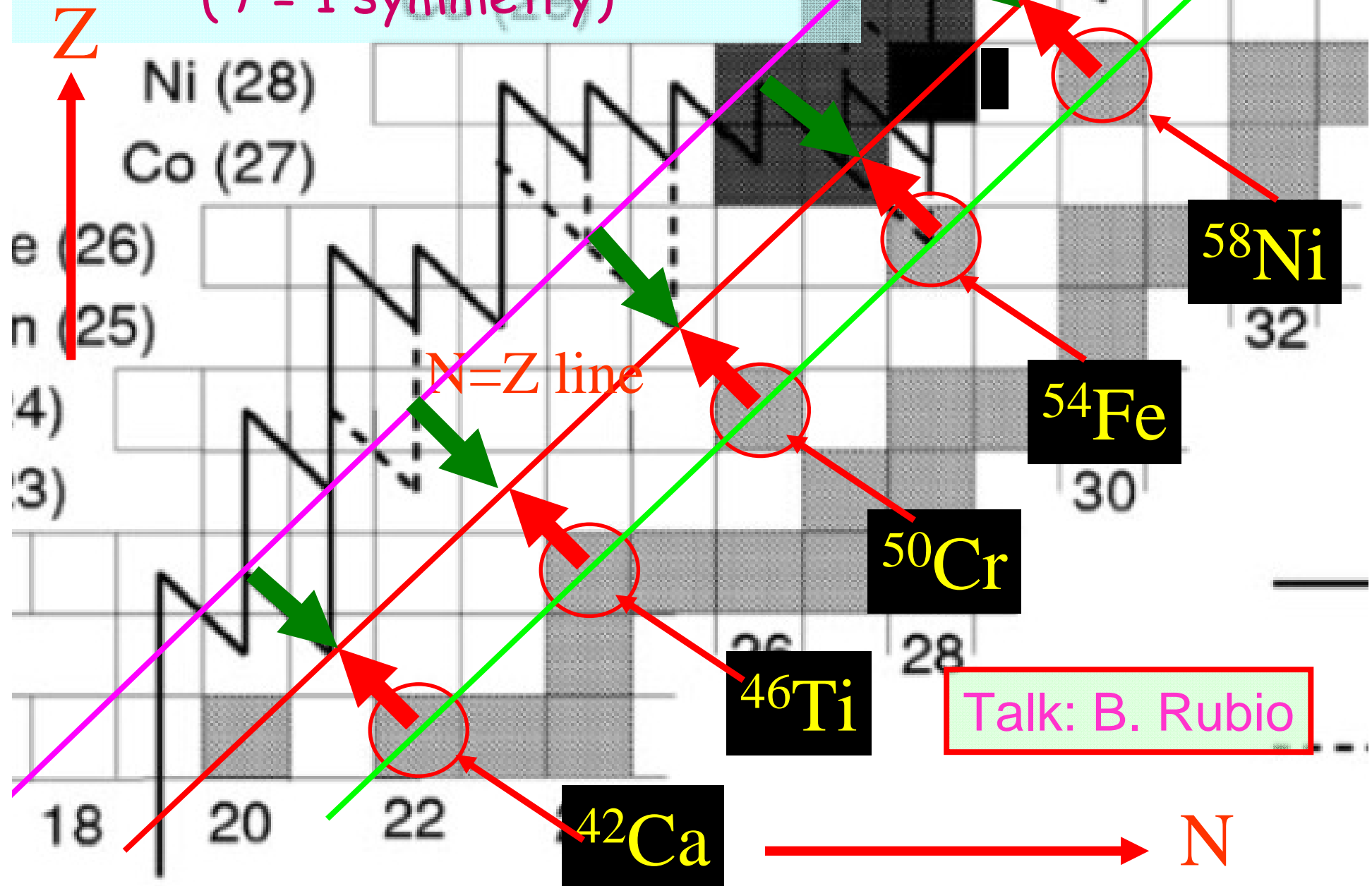
Talk by S. Orrigo: ⁴⁸Fe, ⁵²Ni, ⁵⁶Zn β decay

****GT transitions in each nucleus are
UNIQUE!**

- *pf*-shell nuclei -

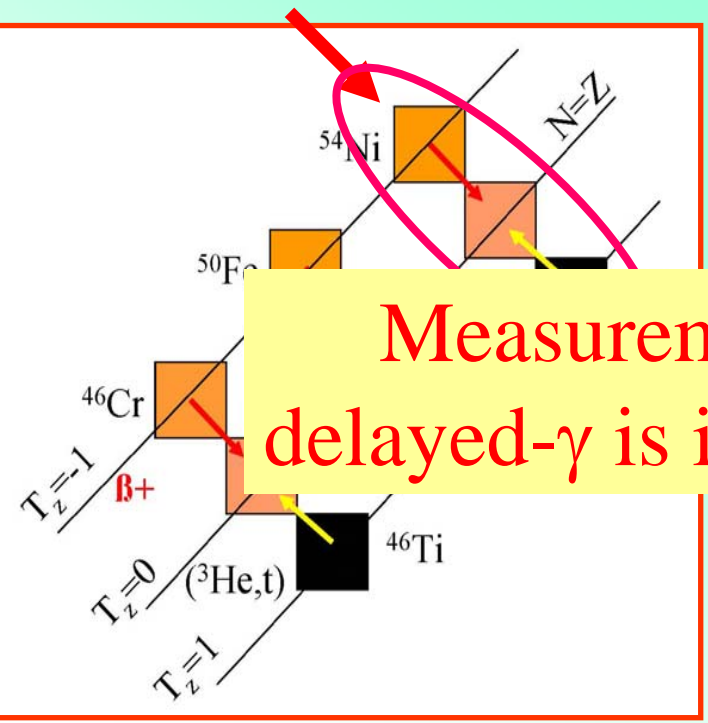


rp-process path nuclei
($T = 1$ symmetry)

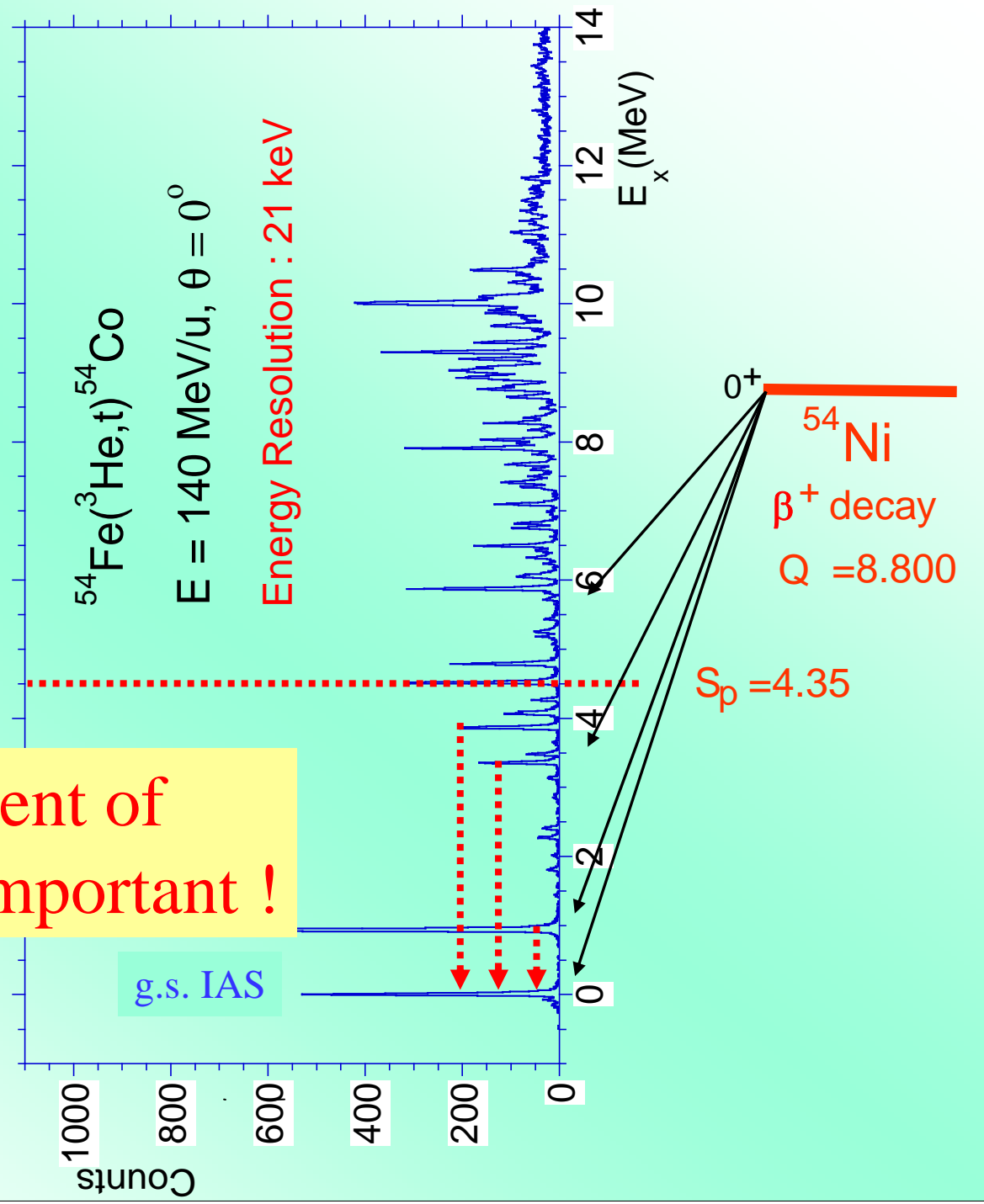


^{54}Ni β -decay measurement

- at GSI (FRS facility)
- RISING (stopped beam campaign)



Measurement of delayed- γ is important !



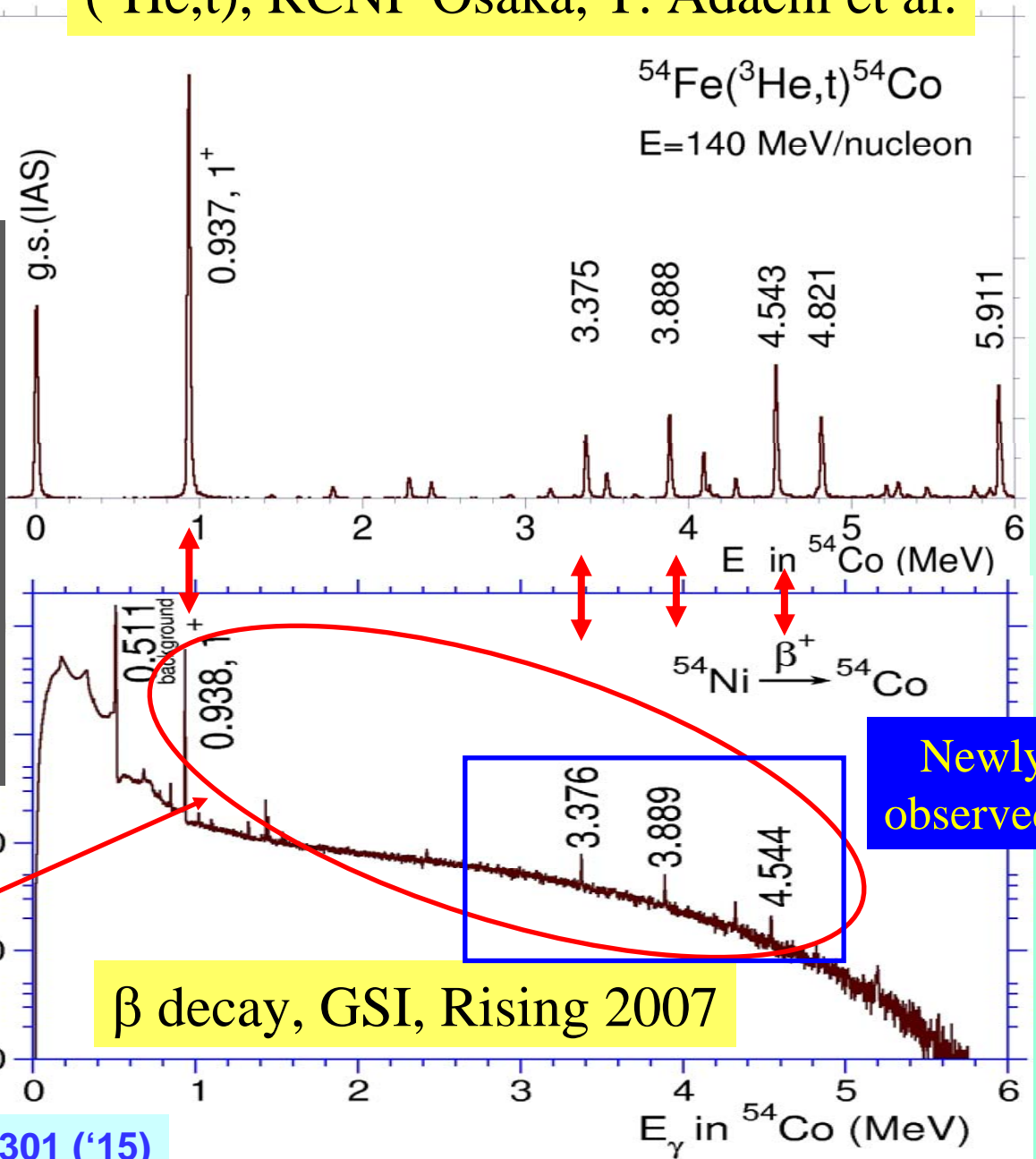
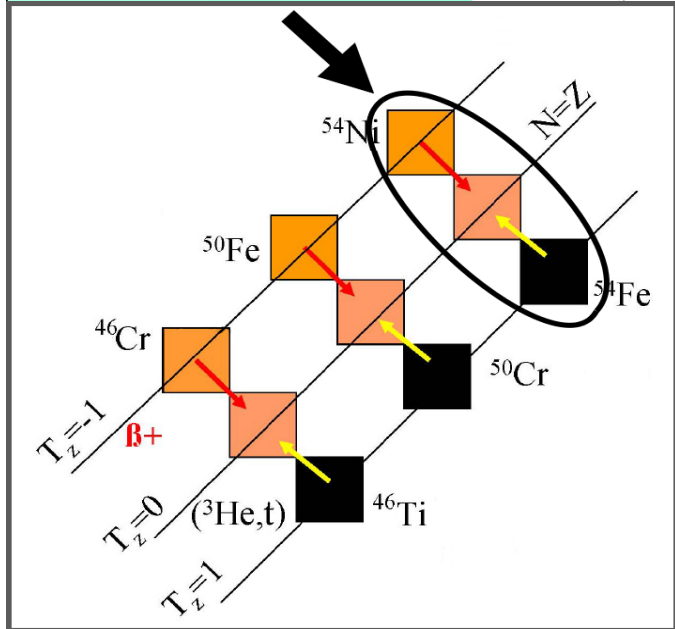
GSI RISING set up



Active Beam Stopper Campaign
July-August, 2007

$(^3\text{He},t)$, RCNP Osaka, T. Adachi et al.

Talk: B. Rubio

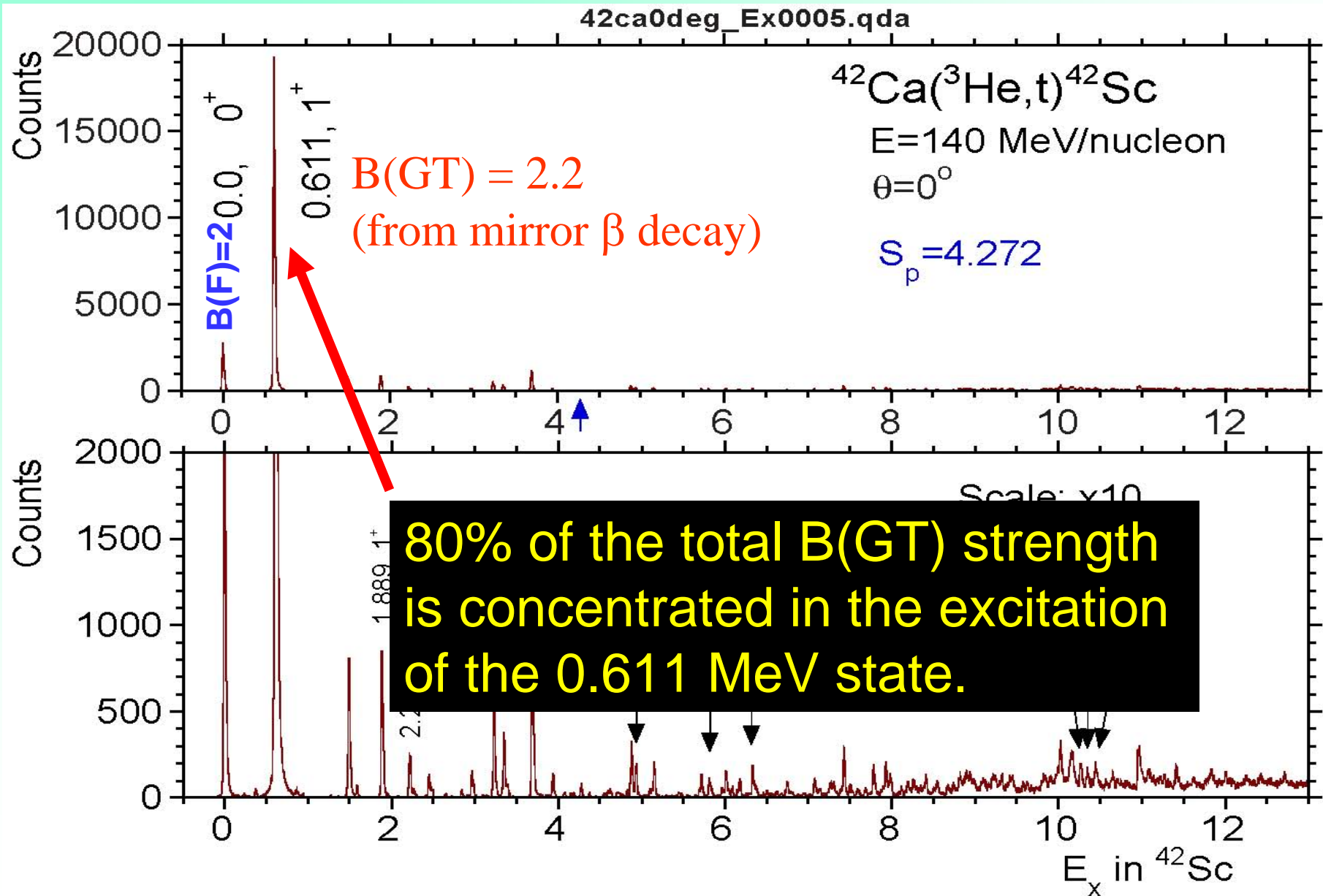


Corresponding Transitions were observed in a wide E_x range !

β decay, GSI, Rising 2007

F. Molina et al., PRC 91, 014301 ('15)

$^{42}\text{Ca}(^3\text{He},t)^{42}\text{Sc}$ in 2 scales



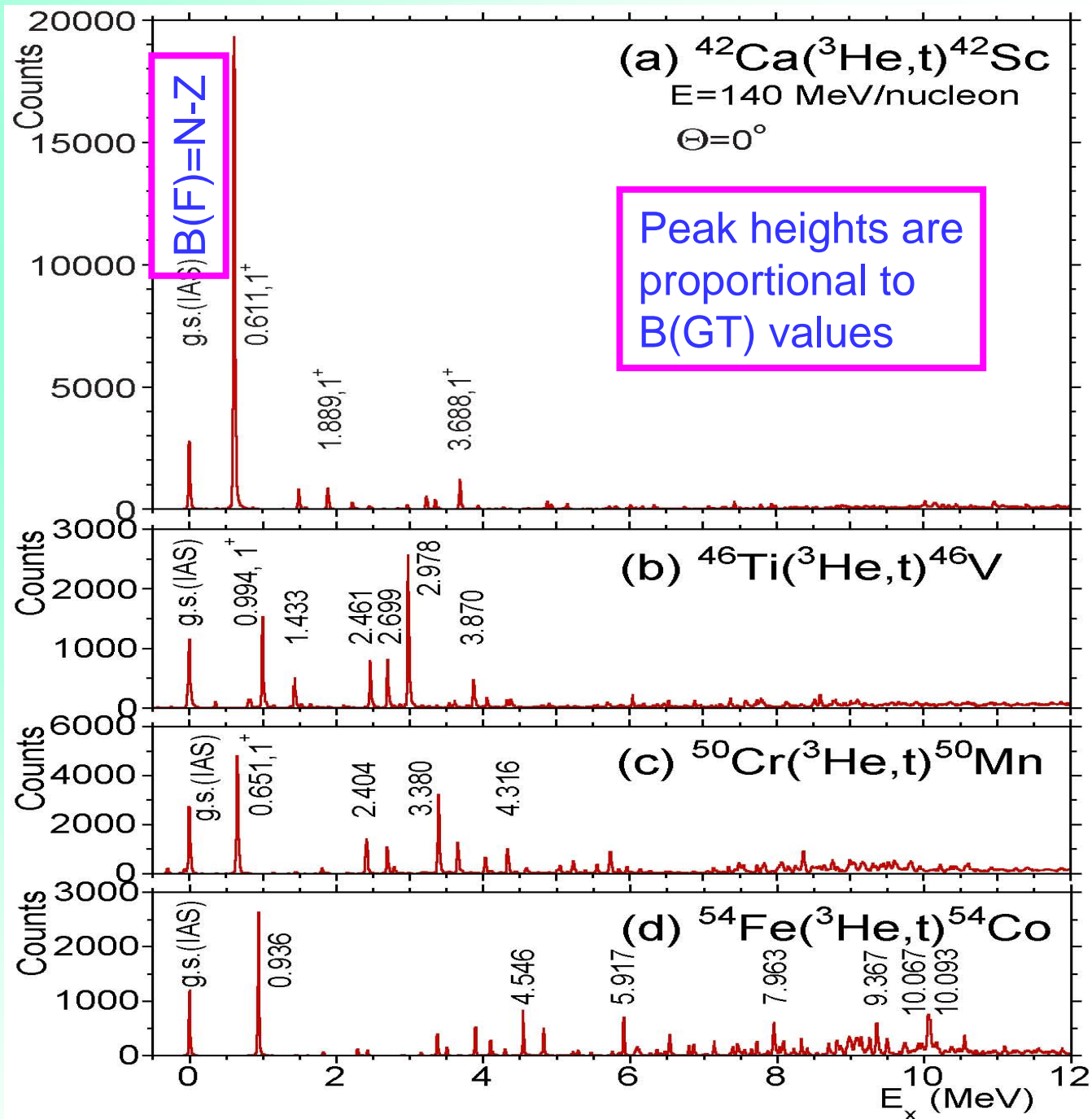
GT states in $A=42-54$ $T_z=0$ nuclei

Y. Fujita et al.
PRL 2014
PRC 2015

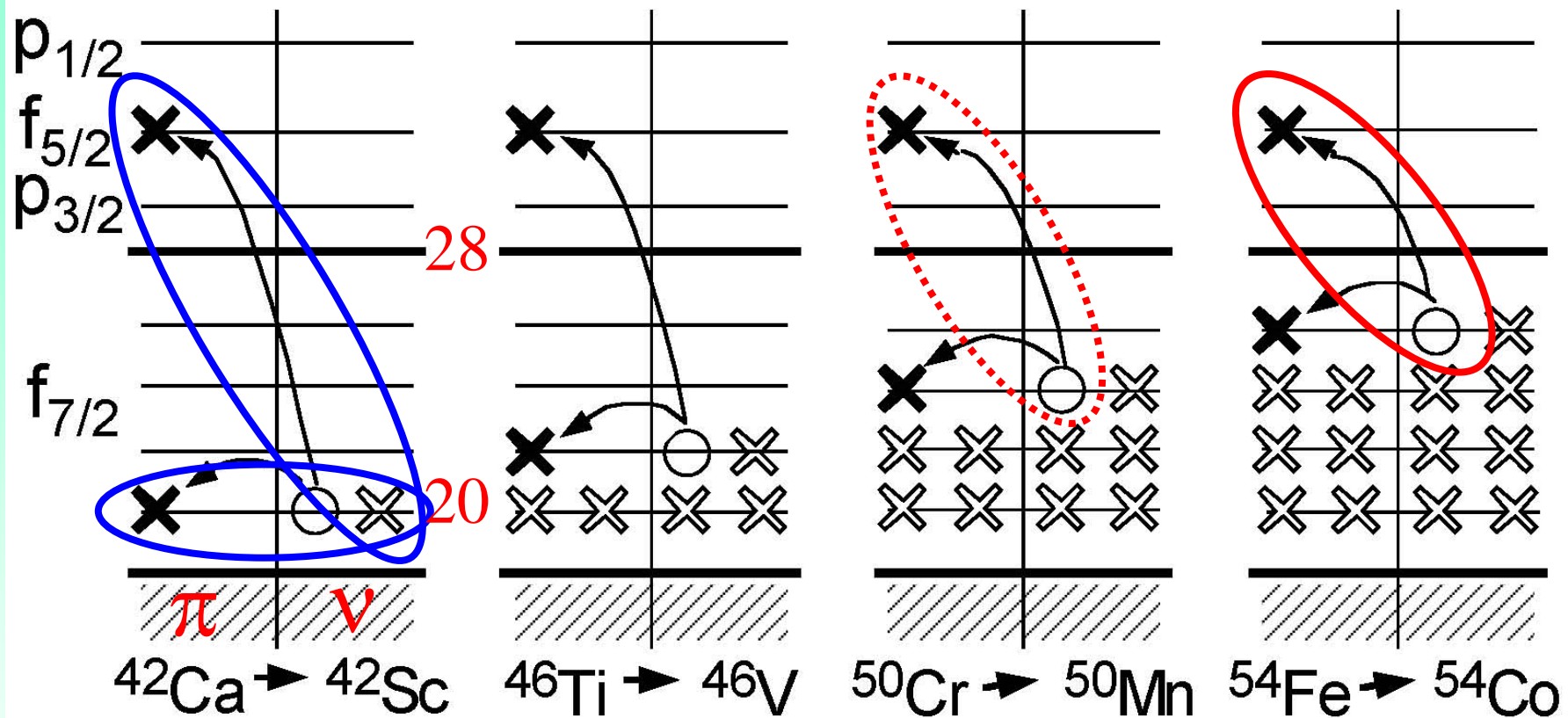
T. Adachi et al.
PRC 2006

Y. Fujita et al.
PRL 2005

T. Adachi et al.
PRC 2012



SM Configurations of GT transitions



particle-hole configuration
 + IV-type int.
 = REPULSIVE

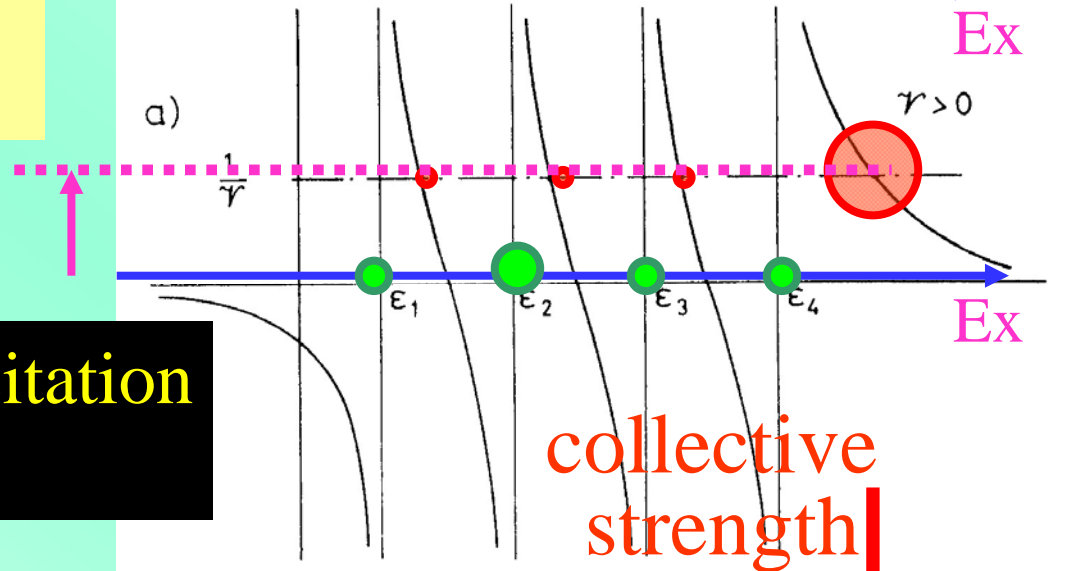
Role of Residual Int. (repulsive)

Single particle-hole strength distribution



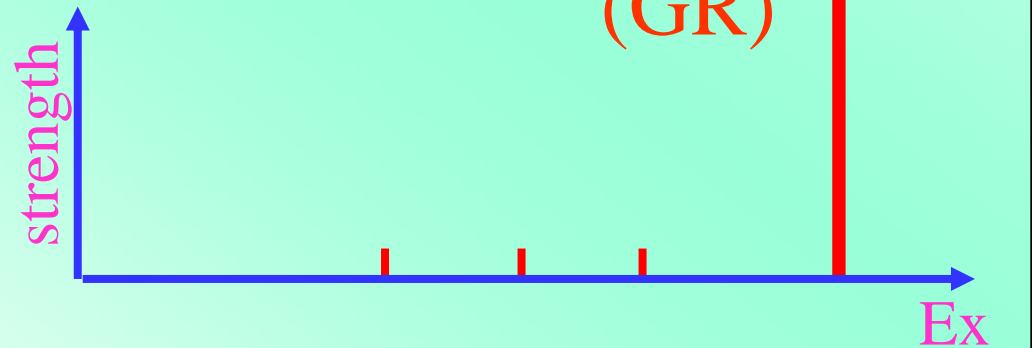
Graphical solution of the RPA dispersive eigen-equation

positive = repulsive



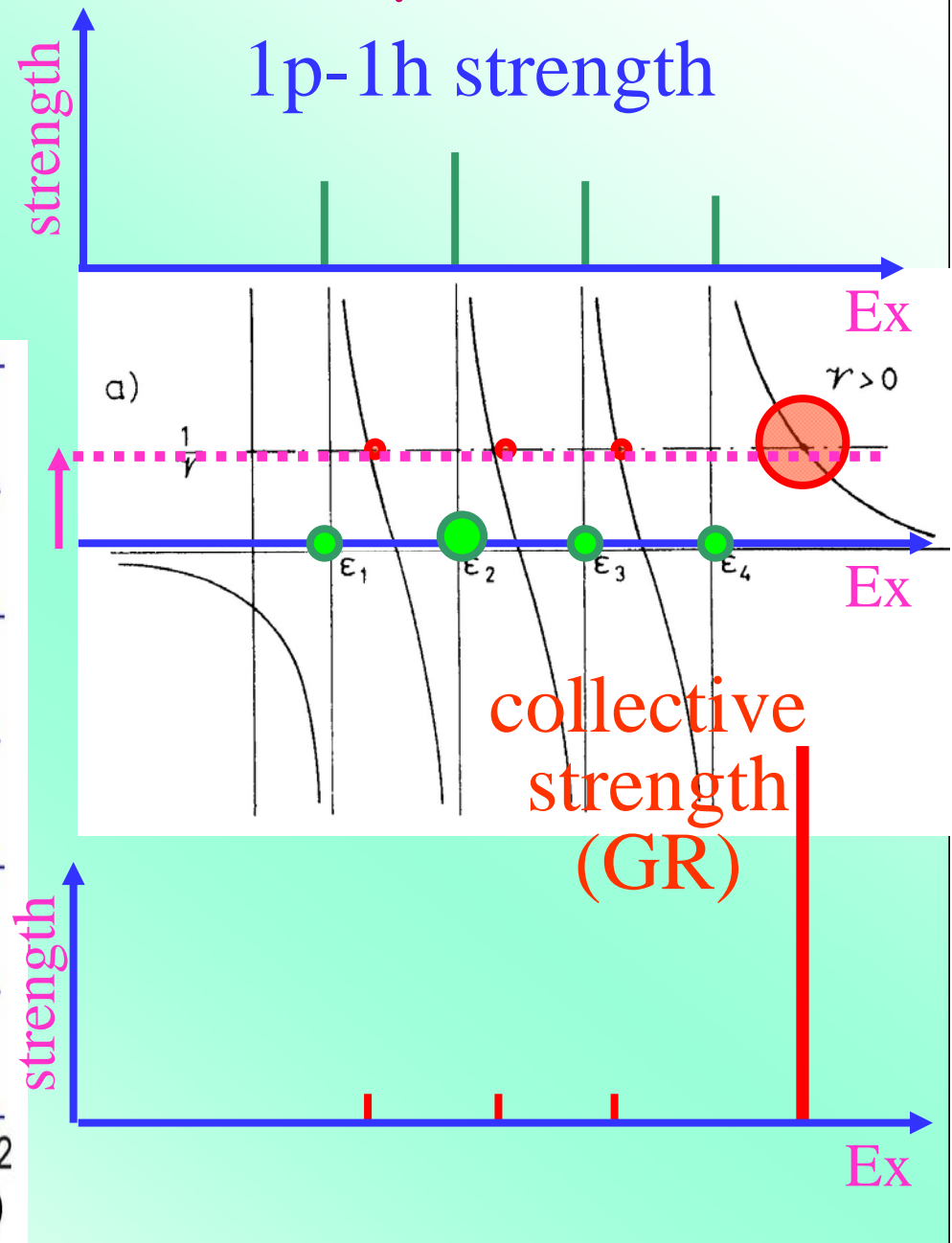
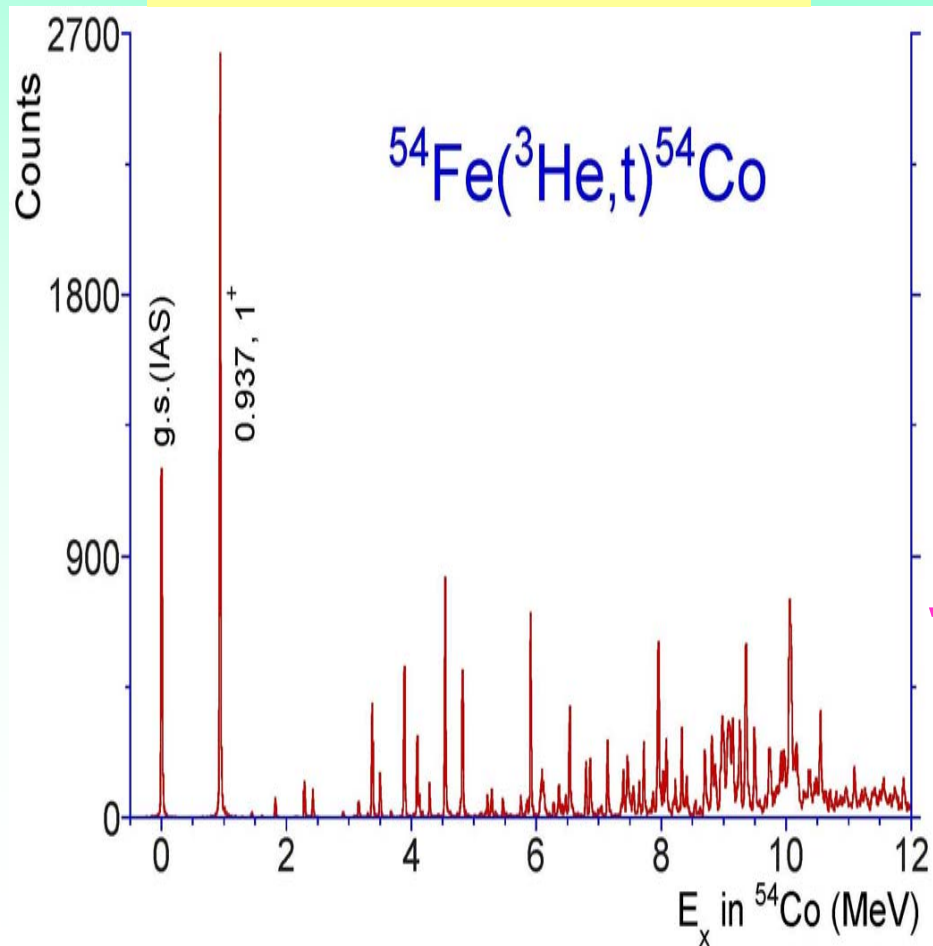
p-h configuration + IV excitation = repulsive

Collective excitation formed by the repulsive residual interaction

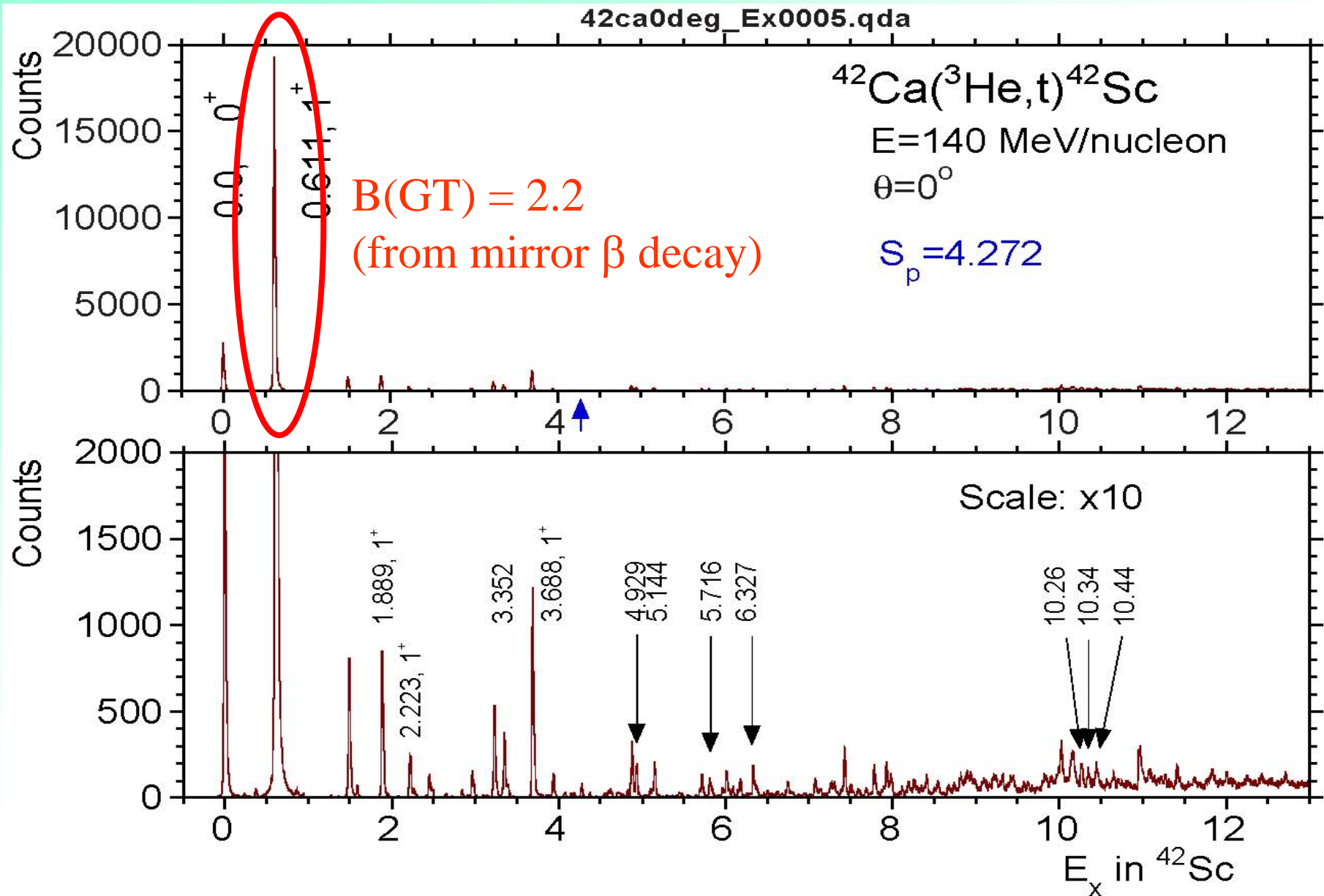


Role of Residual Int. (repulsive)

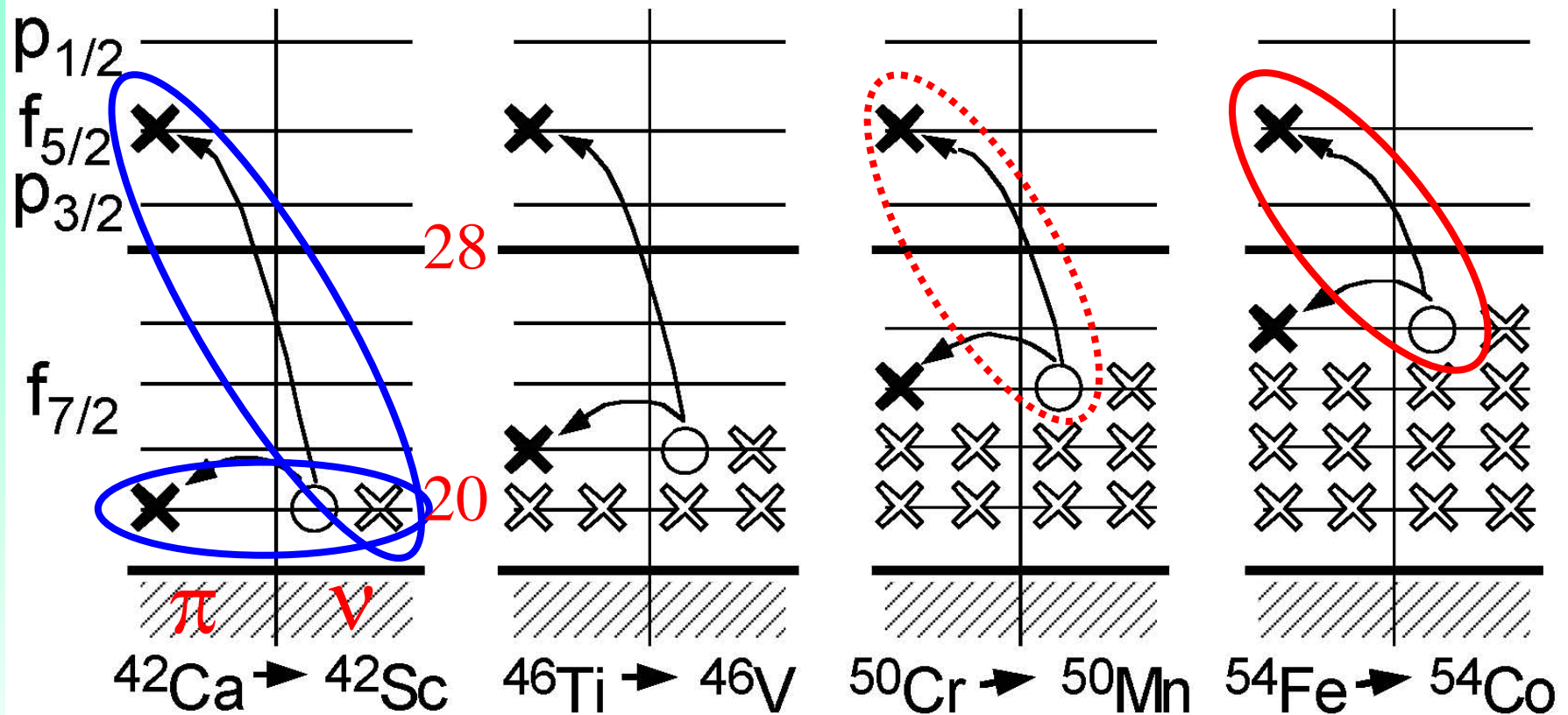
Collective excitation formed by the **repulsive** residual interaction



$^{42}\text{Ca}(^3\text{He},t)^{42}\text{Sc}$ in 2 scales



SM Configurations of GT transitions



π -p - ν -p configurations
sensitive to IS pairing int.

→ attractive

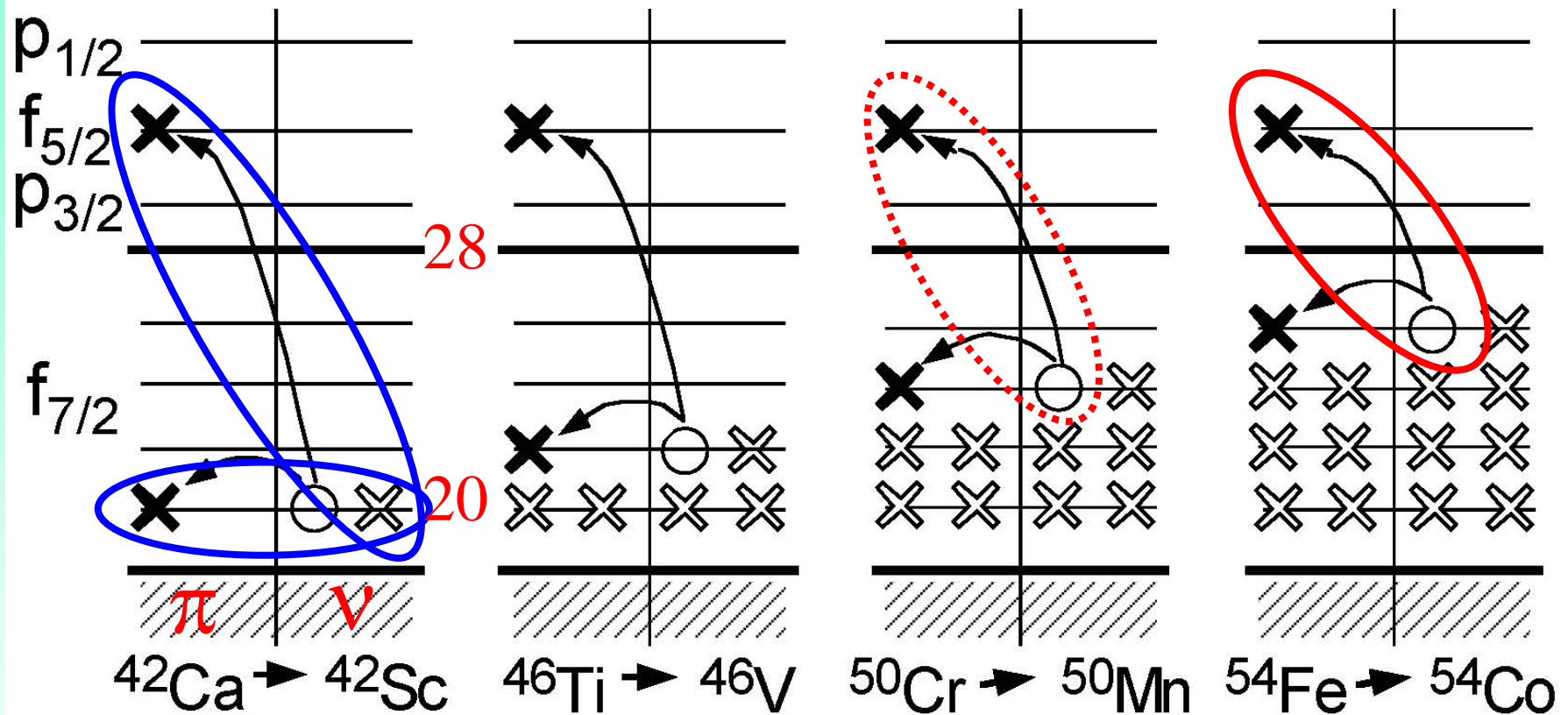
(spin-triplet, IS int. is stronger
than spin-singlet, IV int.)

particle-hole configurations
+ IV-type excitation ($\sigma\tau$)

→ repulsive

by Engel, Bertsch, Macchiavelli

SM Configurations of GT transitions



particle-particle int. (attractive)

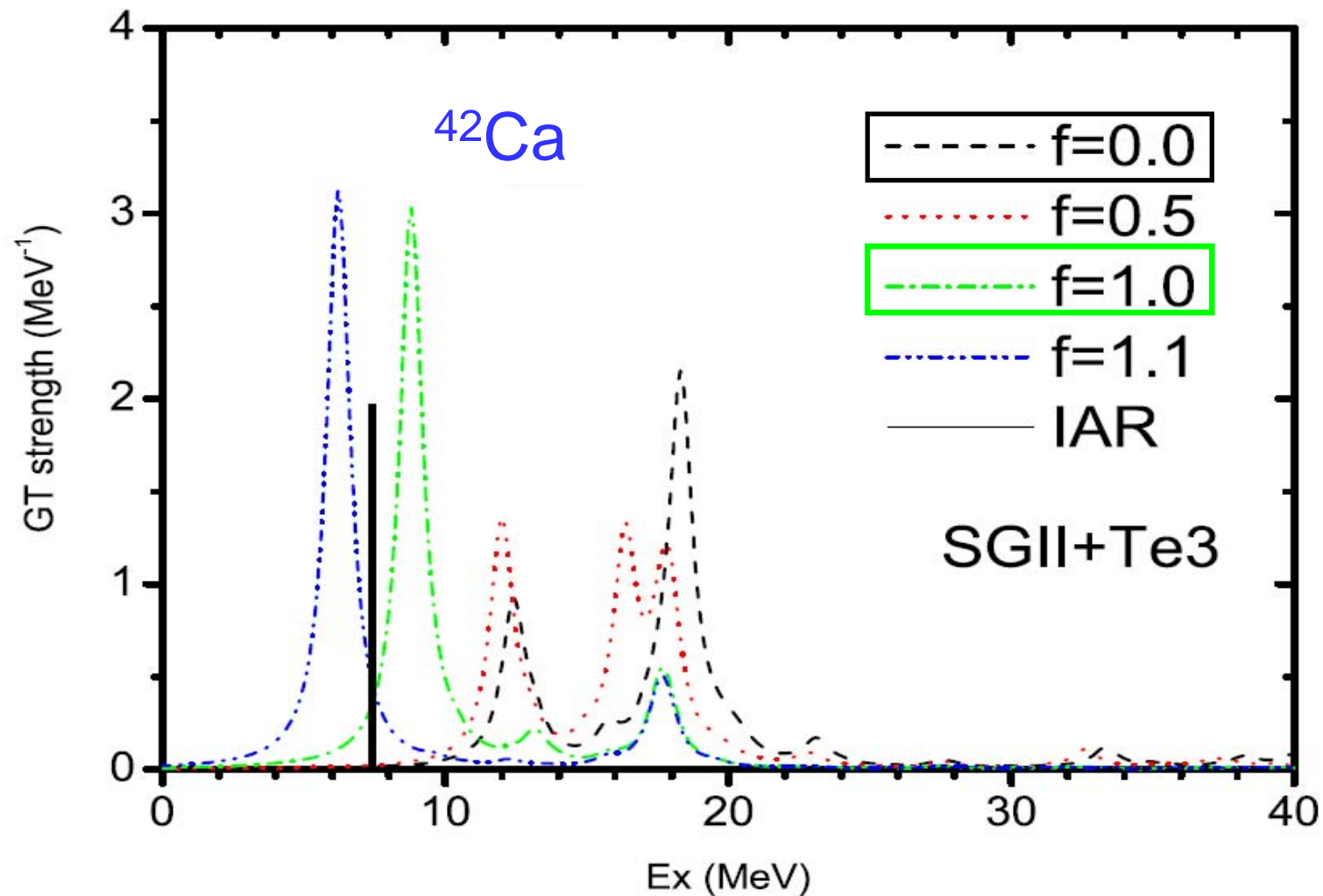
(IS p-n int. is attractive)

particle-hole int. (repulsive)

**Isoscalar interaction
can play important roles !**

QRPA-cal. GT-strength (with IS-int.)

Bai, Sagawa, Colo et al., PRC 90 (2014) 054335



⁴²Ca → ⁴²Sc (Q-value)

QRPA cal. including IS int.

Bai, Sagawa, Colo et al., PRC 90 (2014) 054335

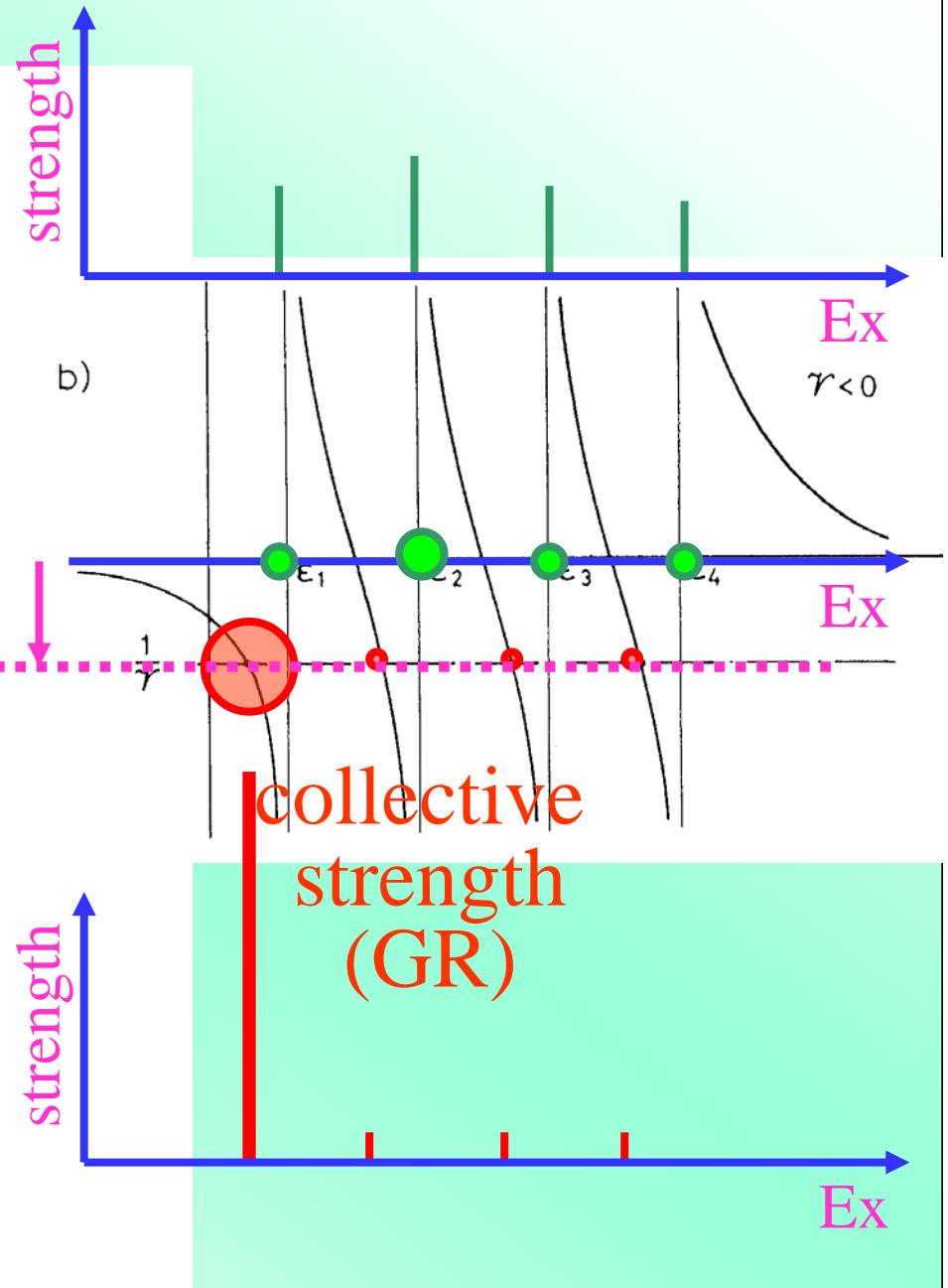
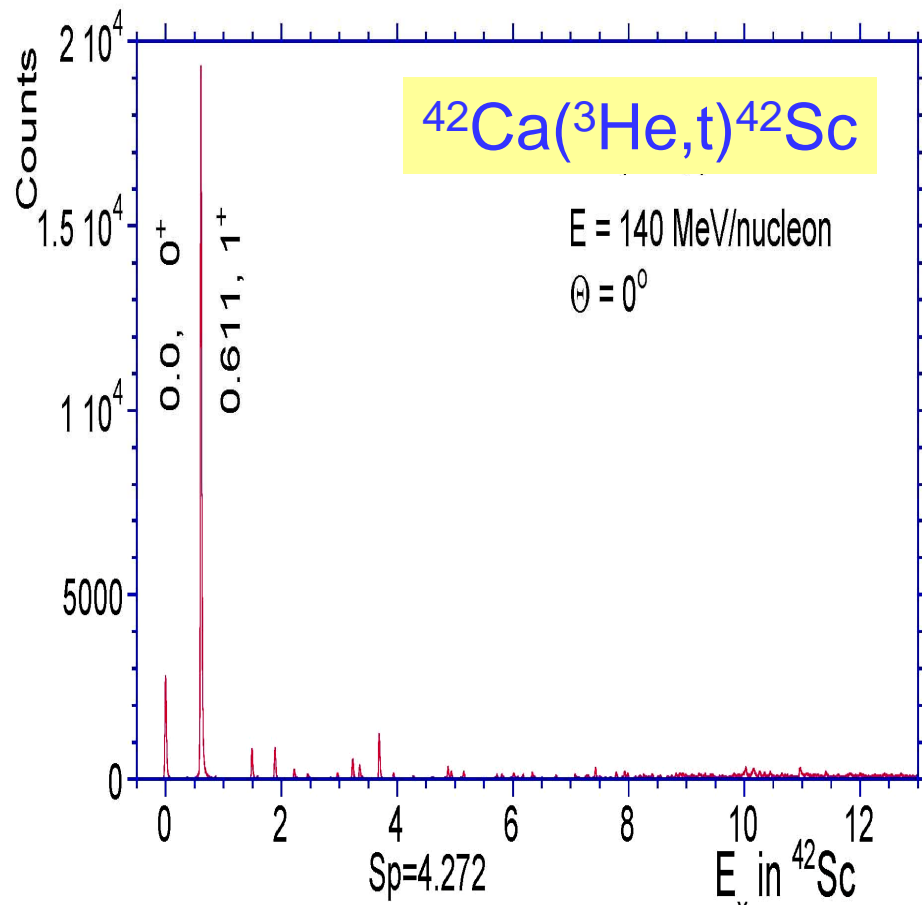
f	B_{np}	neutron	proton	$(X_{upvn}+Y_{unvp})$	$(X_{upvn}+Y_{unvp}) * \langle p GT n \rangle$
0	1.34	1f7/2	1f7/2	0.427	1.3689
0.5	2.051	1f7/2	1f7/2	0.432	1.384
1	4.75	1f5/2	1f7/2	0.053	0.2158
		1f7/2	1f5/2	0.129	0.474
		1f7/2	1f7/2	0.33	1.059

Configurations are in phase!

Low-energy collective GT excitation !
(collectivity is from IS p-n int. !)

Role of Residual Int. (attractive)

Collective excitation formed by the attractive IS residual interaction



42Ca → 42Sc: Shell Model Cal.: Transition Matrix Elements

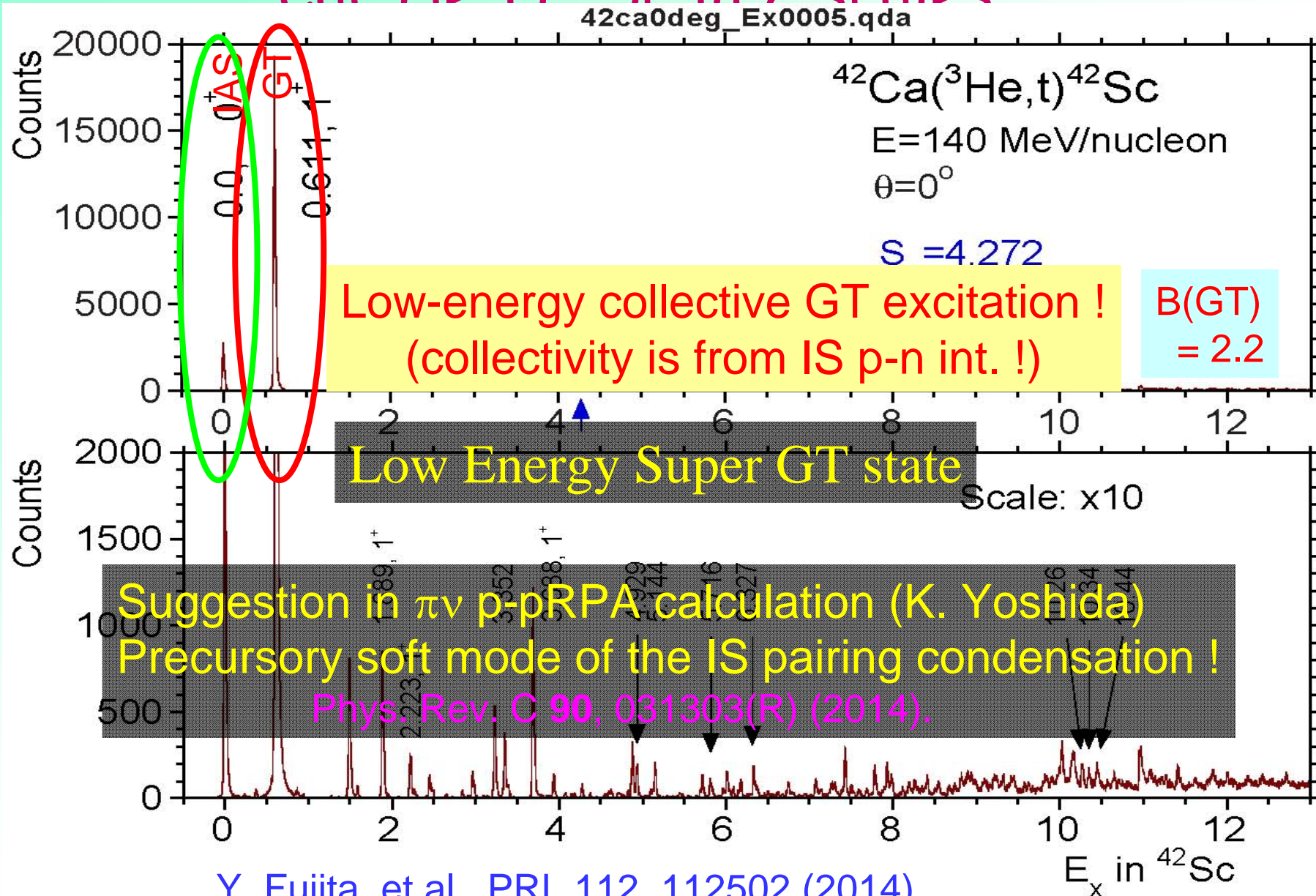
TABLE VI. Results of the pf-shell SM calculation using the GXPF1J interaction. The matrix elements $M(\text{GT})$ of GT transitions exciting individual $J^\pi = 1^+$ GT states in ^{42}Sc from the g.s. of ^{42}Ca are shown for each configuration. The results are shown for all excited GT states predicted in the region up to 9.82 MeV. The notation $f7 \rightarrow f7$, for example, stands for the transition with the $\nu f_{7/2} \rightarrow \pi f_{7/2}$ type and $p3 \rightarrow p3$ the $\nu p_{3/2} \rightarrow \pi p_{3/2}$. The summed value of the matrix elements is denoted by $\Sigma M(\text{GT})$ and its squared value is the $B(\text{GT})$, where the $B(\text{GT})$ values do not include the quenching factor of the SM calculation.

SM cal: M. Honma

States in ^{42}Sc		Configurations						Transition strengths		
E_x (MeV)	T	$f7 \rightarrow f7$	$f7 \rightarrow f5$	$f5 \rightarrow f7$	$p3 \rightarrow p3$	$p3 \rightarrow p1$	$p1 \rightarrow p3$	$\Sigma M(\text{GT})$	$B(\text{GT})$	
0.33	1⁺₁	0	1.383	0.548	0.063	0.031	0.024	0.016	2.07	4.28
4.41	0	0	0.719	-0.742	-0.085	-0.079	-0.073	-0.048	-0.31	0.09
7.41	0	0	0.193	-0.788	-0.090	0.142	0.060	0.040	-0.44	0.19
8.62	0	0	-0.151	0.385	0.044	0.109	-0.071	-0.047	0.30	0.09
9.82	1	0	0.0	1.196	-0.137	0.0	-0.053	0.035	1.04	1.08

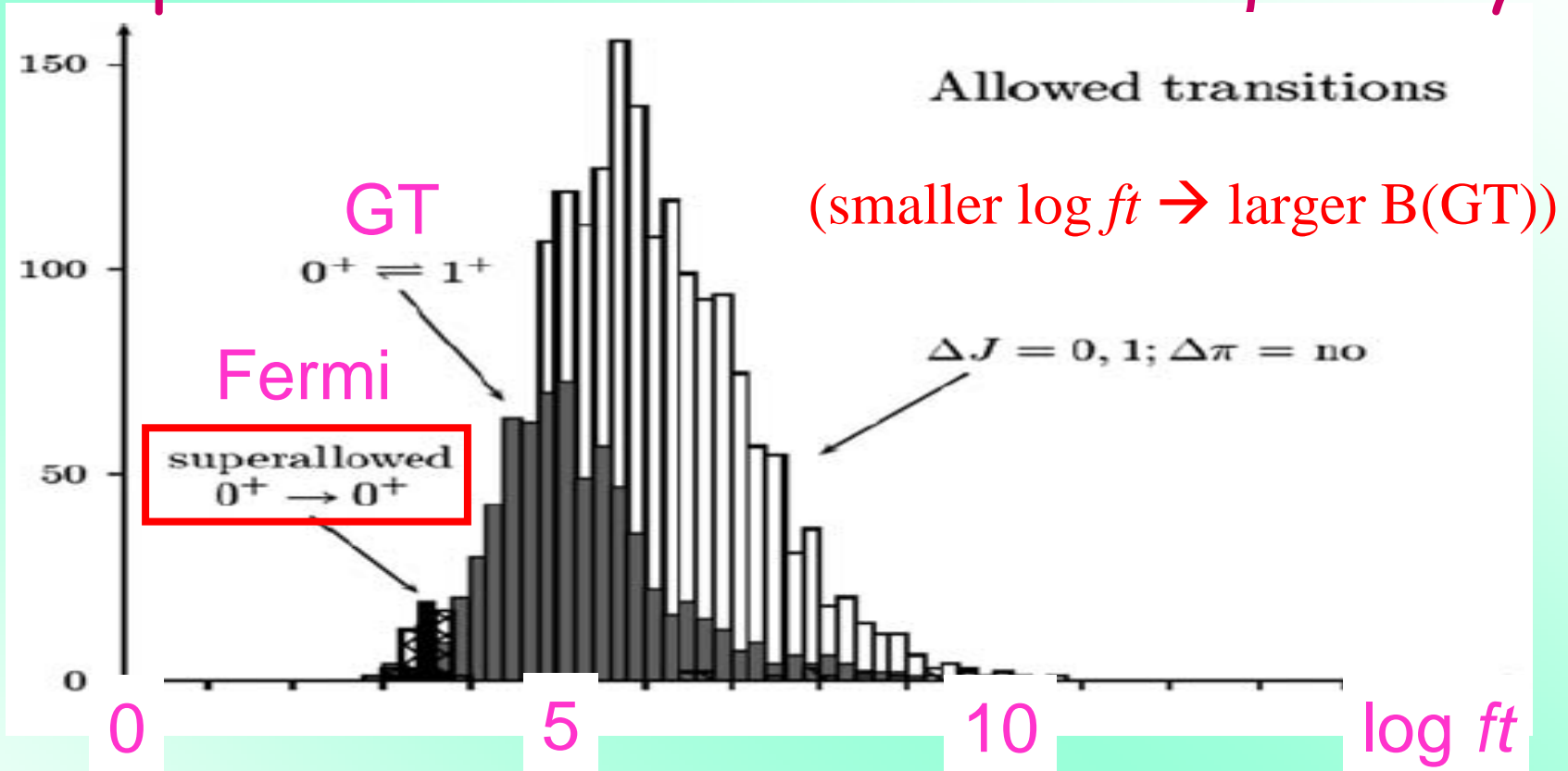
Matrix Elements are in-phase !

$^{42}\text{Ca}(^3\text{He}, t)^{42}\text{Sc}$ in 2 scales



Y. Fujita, et al., PRL 112, 112502 (2014).
 PRC 91, 064316 (2015).

Super-allowed GT transitions in β decay



${}^6\text{He}, 0^+ \rightarrow {}^6\text{Li}, 1^+$	$\log ft = 2.9$
${}^{18}\text{Ne}, 0^+ \rightarrow {}^{18}\text{F}, 1^+$	$\log ft = 3.1$
${}^{42}\text{Ti}, 0^+ \rightarrow {}^{42}\text{Sc}, 1^+$	$\log ft = 3.2$

Super-allowed
GT transitions

Super-Multiplet State

*proposed by Wigner (1937)

In the limit of null $L \cdot S$ force, SU(4) symmetry exists.

We expect:

- a) GT excitation strength is concentrated in a low-energy GT state.
- b) excitation energies of both the IAS and the GT state are identical.

→ *Super-Multiplet State*

In ^{54}Co , we see a broken SU(4) symmetry.

In ^{42}Sc , we see a good SU(4) symmetry.

→ attractive IS residual int. restores the symmetry !

→ 0.611 MeV state in ^{42}Sc has a character close to

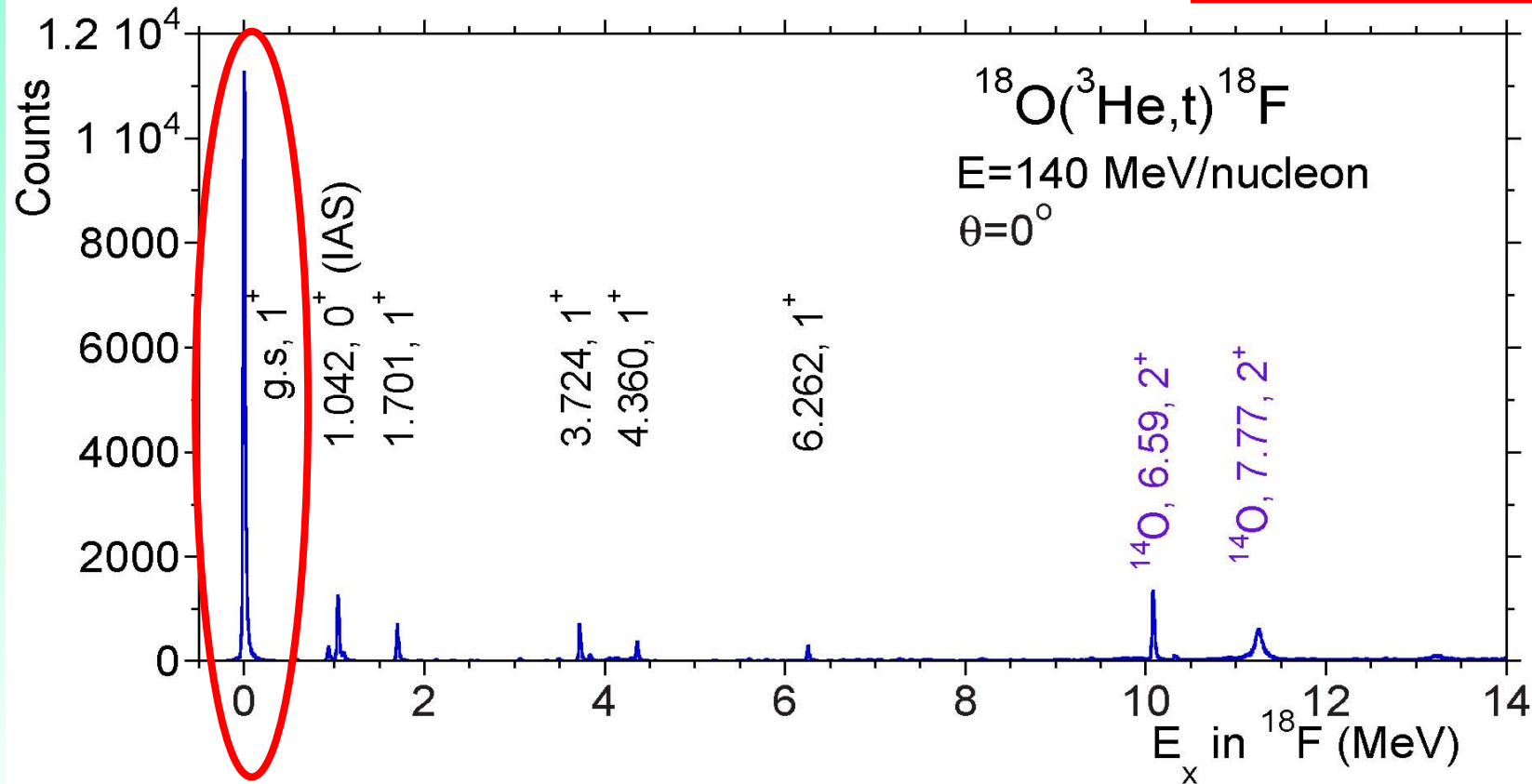
Super-Multiplet State !

We call this state the

Low-energy Super GT state !

$^{18}\text{O}(^3\text{He},t)^{18}\text{F}$ at 0°

Talk: H. Fujita



Low-energy collective GT excitation: $B(\text{GT})=3.1$

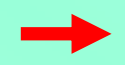
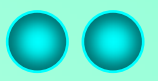
Low Energy Super GT state

GT transitions forming Low-Energy Super GT state

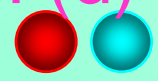
$J^\pi = 0^+ \rightarrow 1^+$

Σ (Sum rule) = $3 \times |N-Z| = 6$

$2n$



${}^2\text{H} (d)$

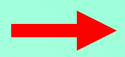
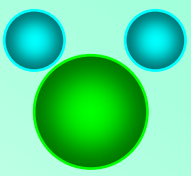


g.s.

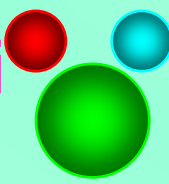
$B(\text{GT}) = 6.0 ?$

Large !

${}^6\text{He}$



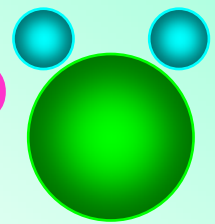
${}^6\text{Li}$



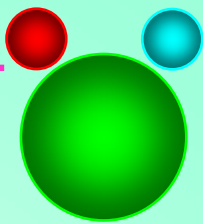
g.s.

$B(\text{GT}) = 4.73$

${}^{18}\text{O}$



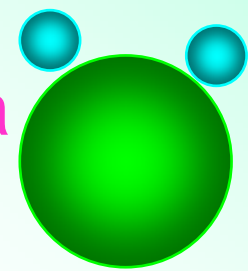
${}^{18}\text{F}$



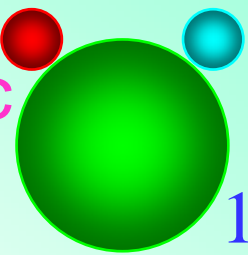
g.s.

$B(\text{GT}) = 3.09$

${}^{42}\text{Ca}$



${}^{42}\text{Sc}$



$1^{\text{st}} E_x$ state (IAS is the g.s.)

$B(\text{GT}) = 2.17$

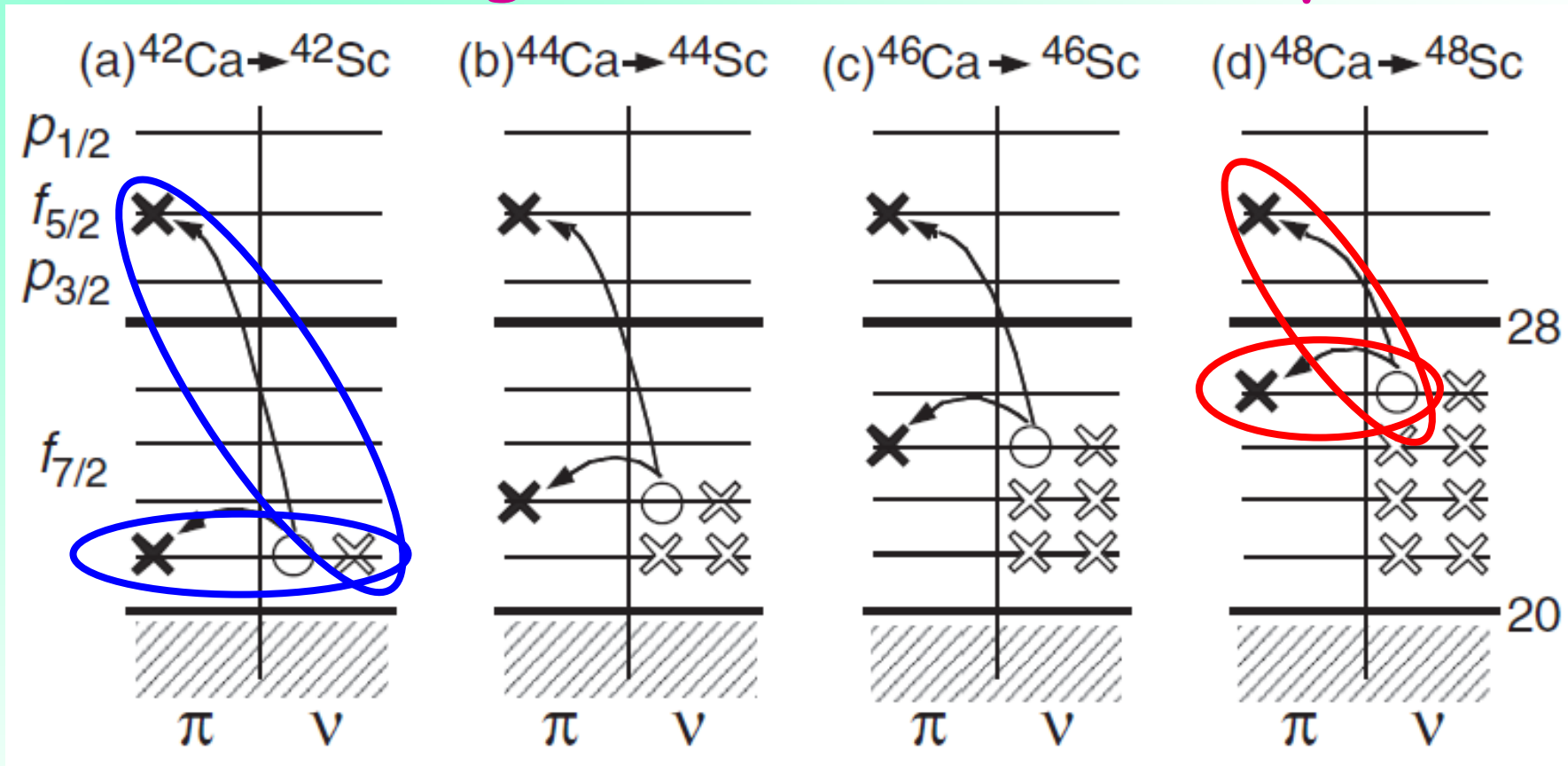
Smaller !



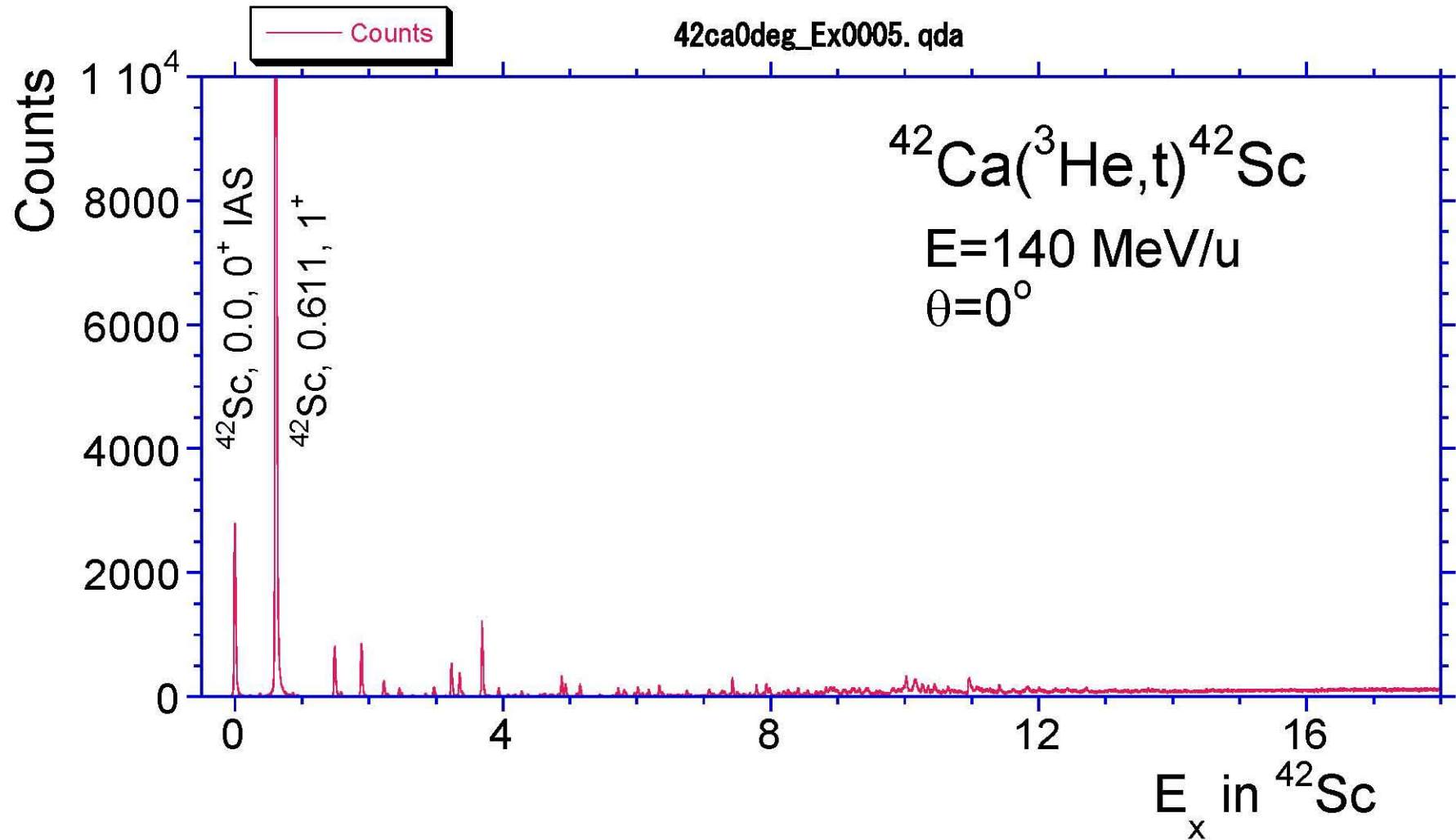
***from $p-p$ to $p-h$ configuration

LESGT stae \rightarrow GTR structure
in $A= 42$ to 48 Ca isotopes

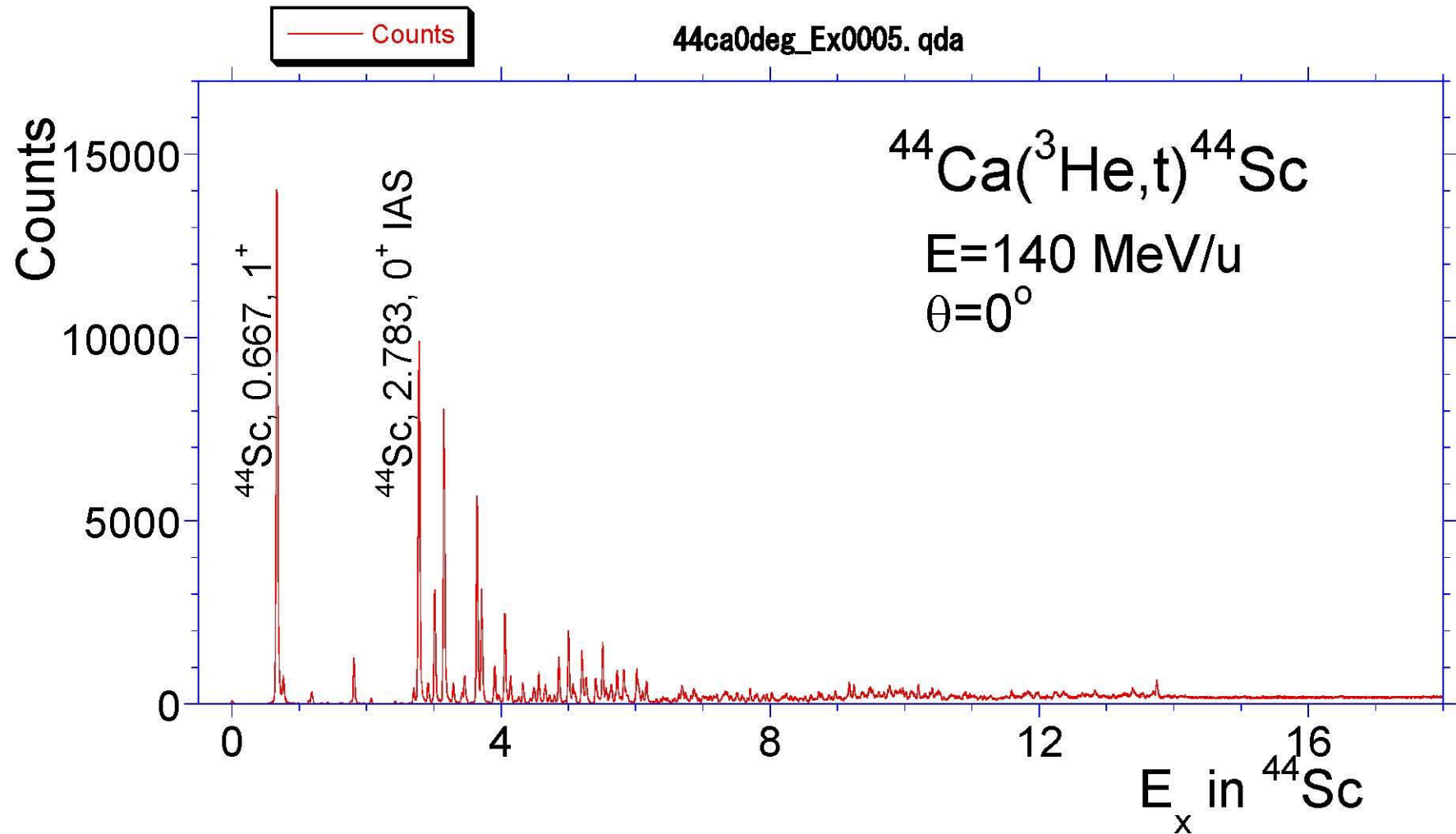
GT Configurations in Sc isotopes



particle-particle int. (attractive) \longrightarrow particle-hole int. (repulsive)

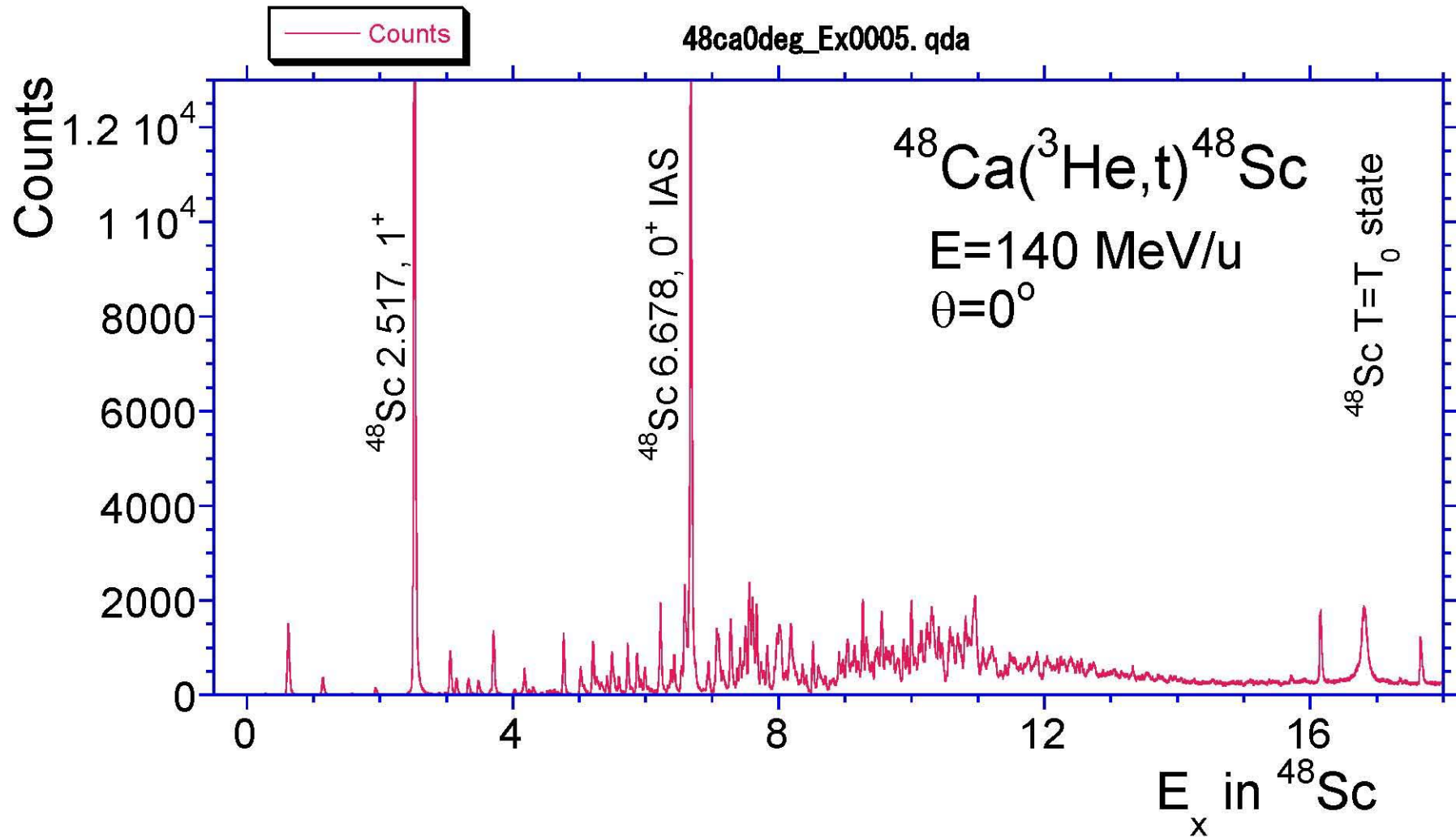


$^{44}\text{Ca}(^3\text{He},t)^{44}\text{Sc}$



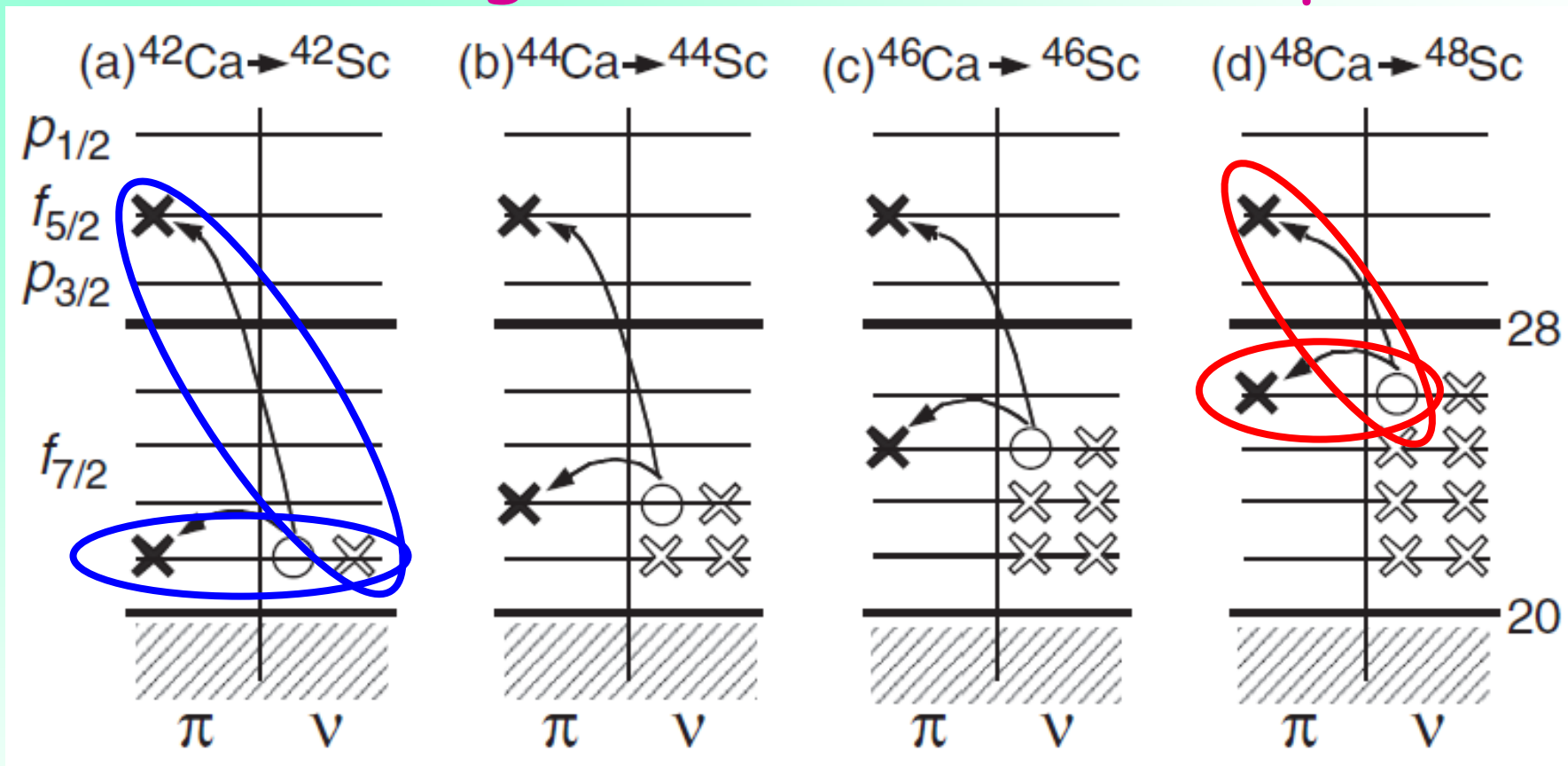
Y. Fujita et al., PRC 88, 014308 (2013)

$^{48}\text{Ca}(^3\text{He},t)^{48}\text{Sc}$



H.F analysis

GT Configurations in Sc isotopes



particle-particle int. (attractive)



particle-hole int. (repulsive)

Low-Energy
Super GT state
Is formed !

Gamow-Teller
Resonance
Is formed !

Summary

GT ($\sigma\tau$) operator : a simple operator !

- * GT transitions: sensitive to the structure of $|i\rangle$ and $|f\rangle$

High resolution of the ($^3\text{He},t$) reaction

- * Fine structures of GT transitions

Mirror β decays and Isospin Symmetry

- * Giving the Absolute GT strength

- GT transitions in each nucleus are UNIQUE !
- Low-energy Super GT state (LESGT state)
- Assuming T-symmetry → GT in unstable nuclei !

We can learn a lot by the comparison of analogous GT transitions !

GT-study Collaborations

Bordeaux (France) : β decay
GANIL (France) : β decay
Gent (Belgium) : (^3He , t), (d, ^2He), (γ , γ'), theory
GSI, Darmstadt (Germany) : β decay, theory
ISOLDE, CERN (Switzerland) : β decay
iThemba LABS. (South Africa) : (p, p'), (^3He , t)
Istanbul (Turkey): (^3He , t), β decay
Jyvaskyla (Finland) : β decay
Koeln (Germany) : γ decay, (^3He , t), theory
KVI, Groningen (The Netherlands) : (d, ^2He)
Leuven (Belgium) : β decay
LTH, Lund (Sweden) : theory
Milano : theory
Osaka University (Japan) : (p, p'), (^3He , t), theory
RIKEN : β decay, theory
Surrey (GB) : β decay
TU Darmstadt (Germany) : (e, e'), (^3He , t)
Valencia (Spain) : β decay
Michigan State University (USA) : theory, (t, ^3He)
Muenster (Germany) : (d, ^2He), (^3He ,t)
Univ. Tokyo and CNS (Japan) : theory, β decay

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Review

Spin–isospin excitations probed by strong, weak and electro-magnetic interactions

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