

Spectroscopic Study of Charmed Baryons at J-PARC

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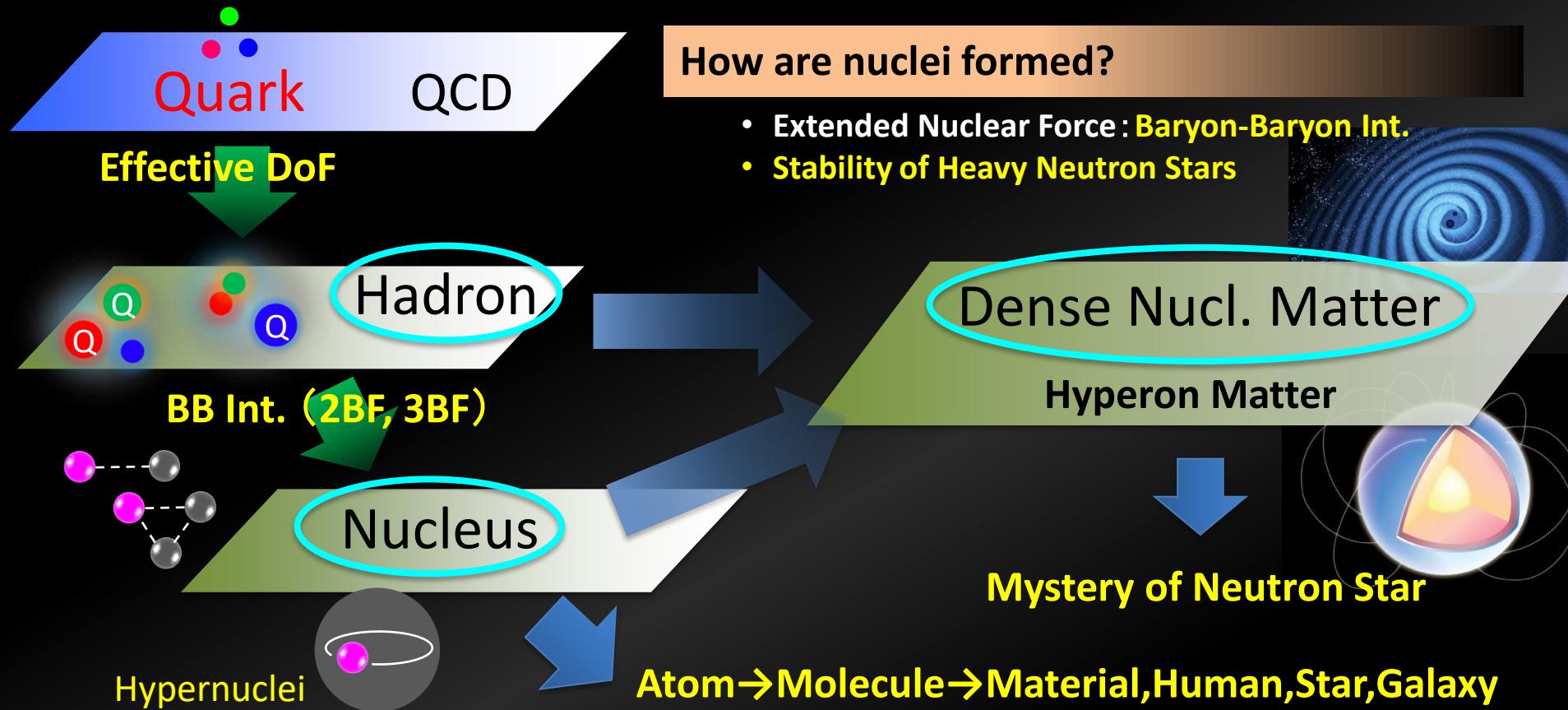
- I. Baryon Structure with Heavy Quarks
- II. Charmed Baryon Spectroscopy
- III. New Platform of Hadron Physics at J-PARC

Hierarchy of Matter in the Universe

Matter Evolution from Quark to Hadron, Nucleus, and Neutron Star

How QCD works in Hadron?

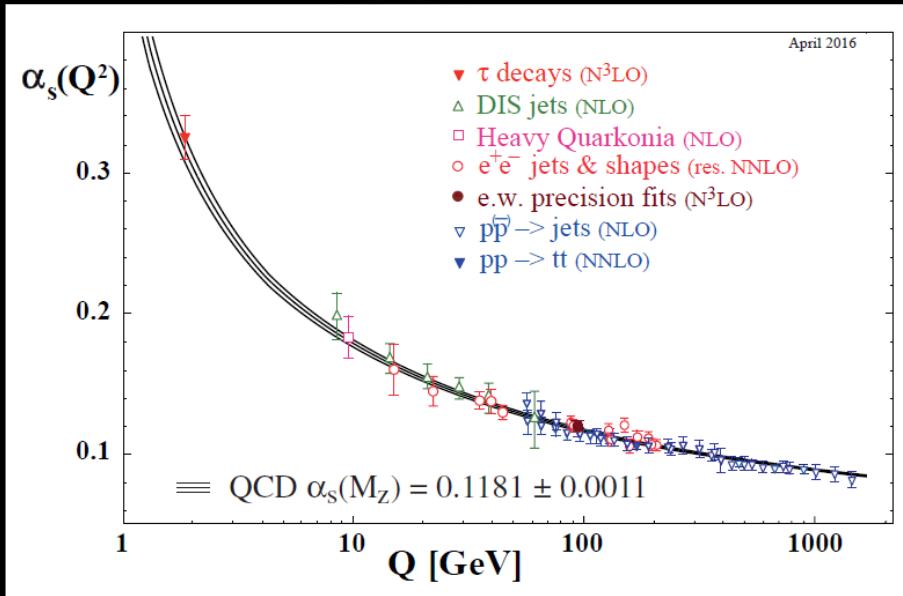
- Effective DoF (**building blocks**) to describe hadrons
- Change of Hadron Properties in Matter



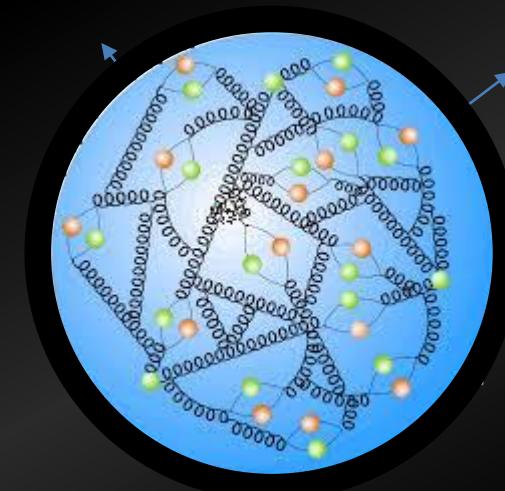
Form of Hadrons

Low E

High E



$\alpha_s = \infty$
at Λ_{QCD}

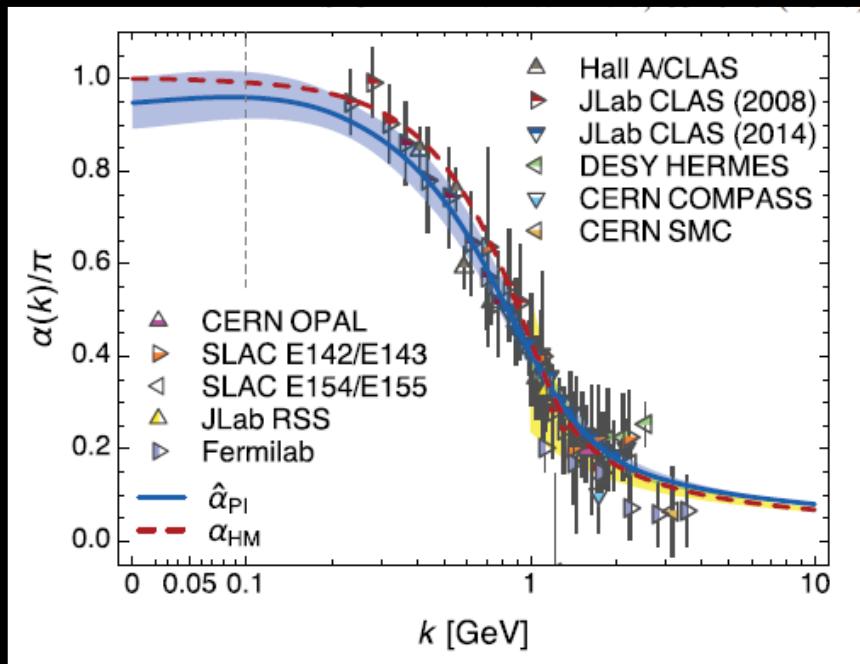


Quark-Gluon

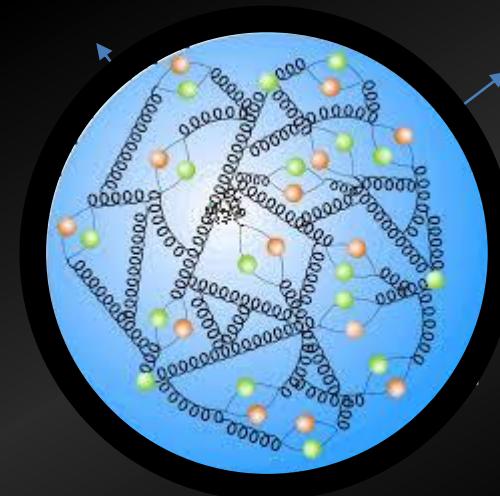
Form of Hadrons

Low E

High E

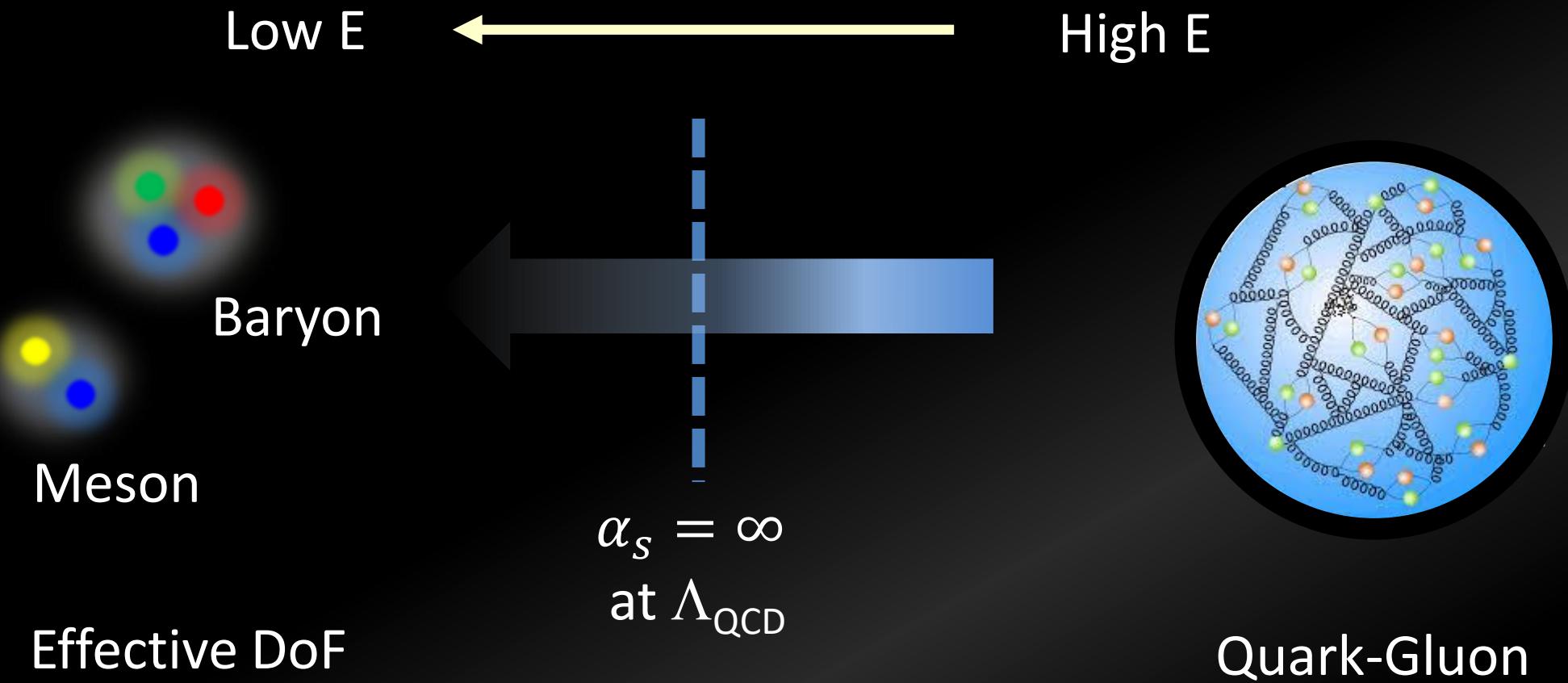


$\alpha_s = \infty$
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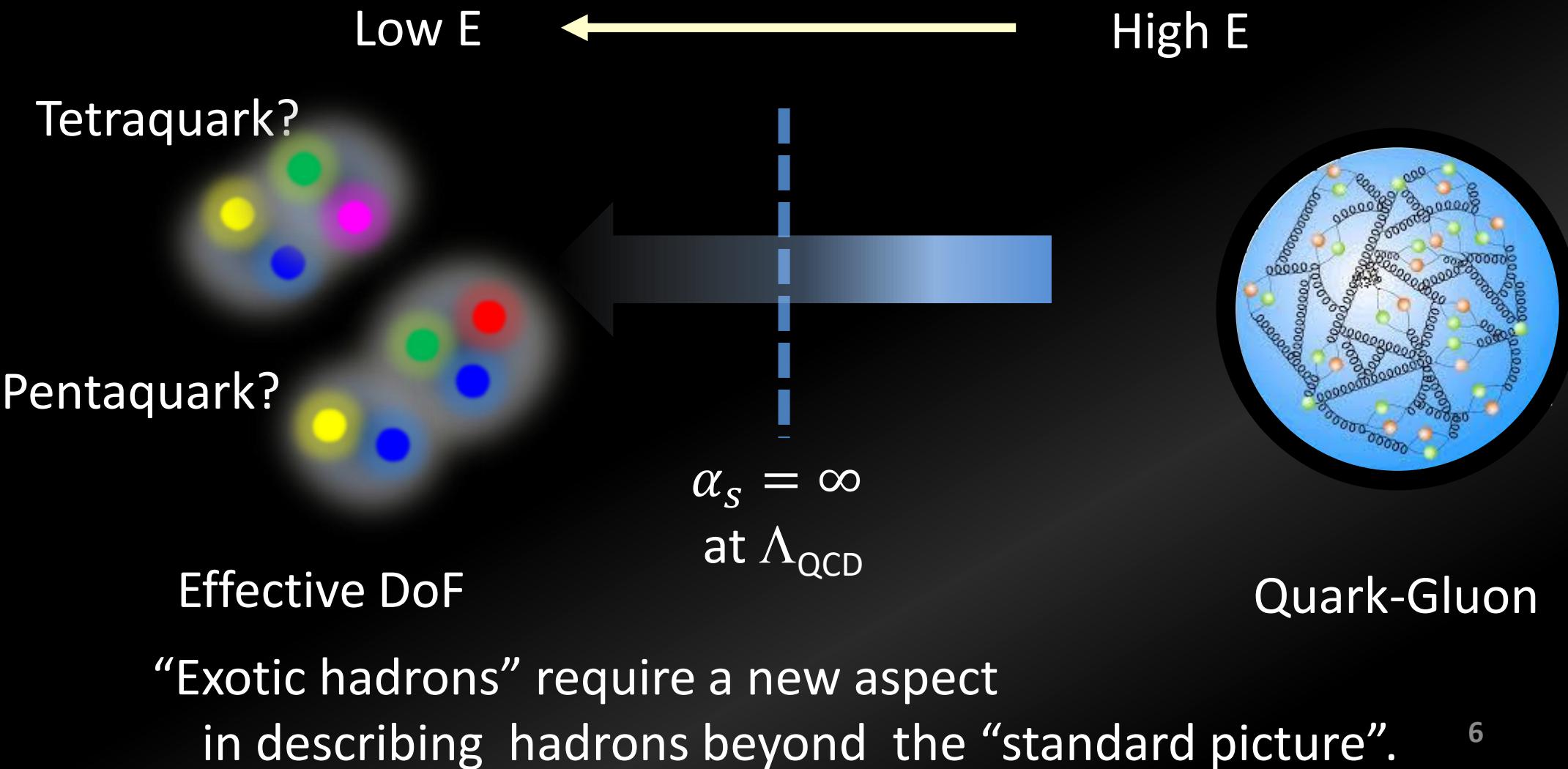


Quark-Gluon

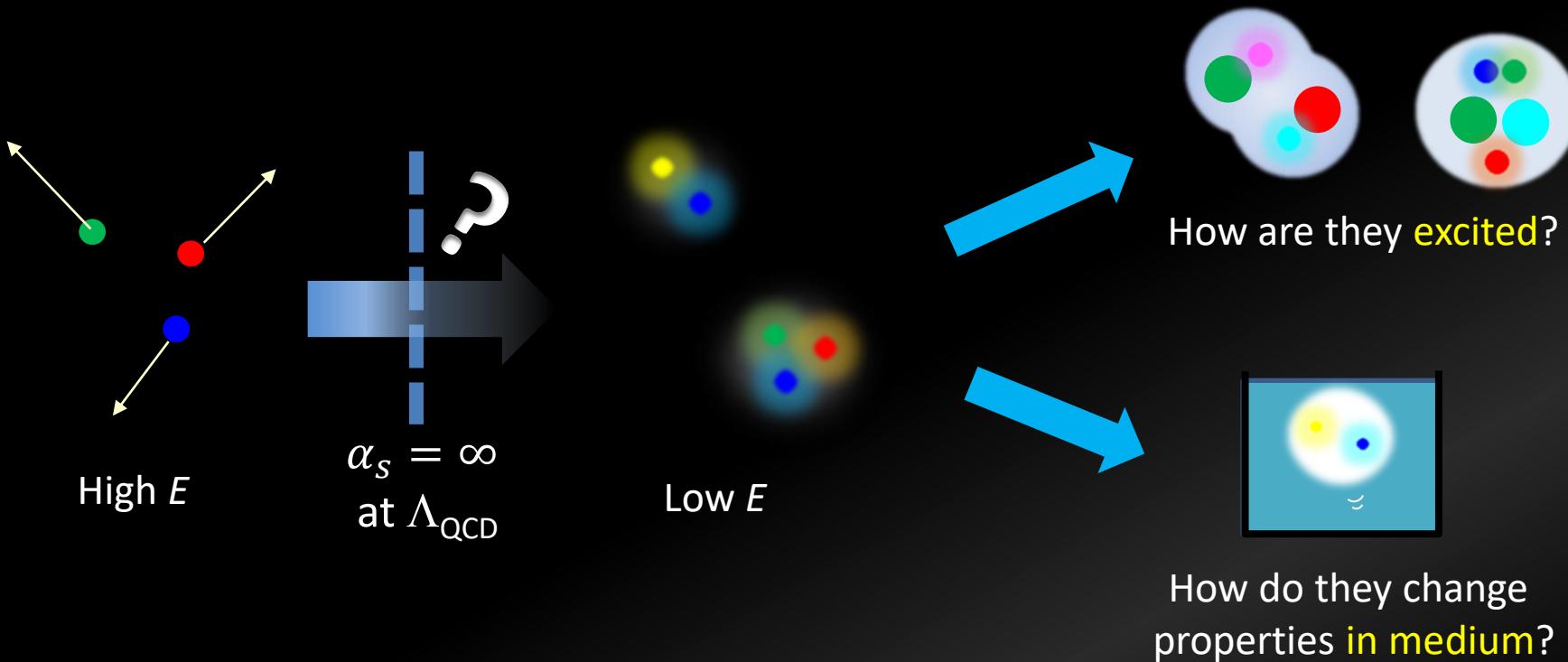
Form of Hadrons



Form of Hadrons

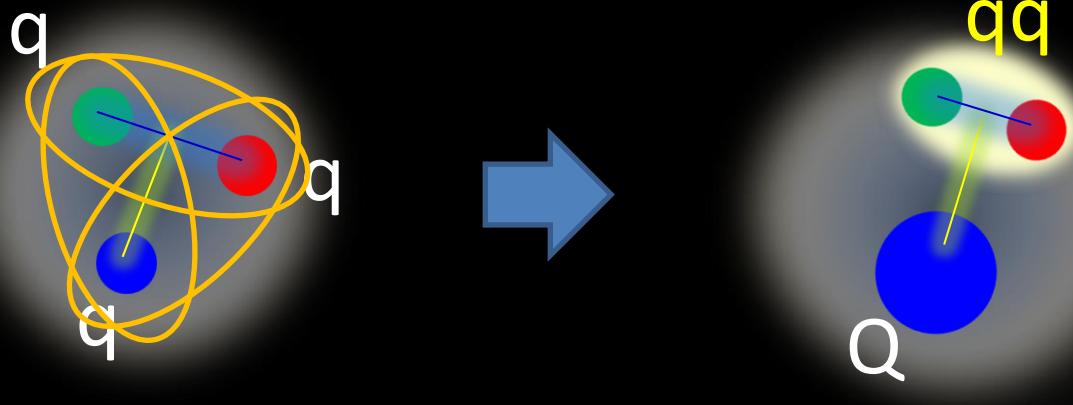


Hadron Physics at J-PARC



- Quasi-Particles (= Effective DoF) emerging at Low E describe hadron properties effectively.
- QP could play a role to form hadrons in early Universe.

Roles of Heavy Flavors



$$V_{CMI} \sim [\alpha_s/(m_i m_j)]^* (\lambda_i, \lambda_j) (\sigma_i, \sigma_j)$$

$\rightarrow 0$ if $m_{i,j} \rightarrow \infty$

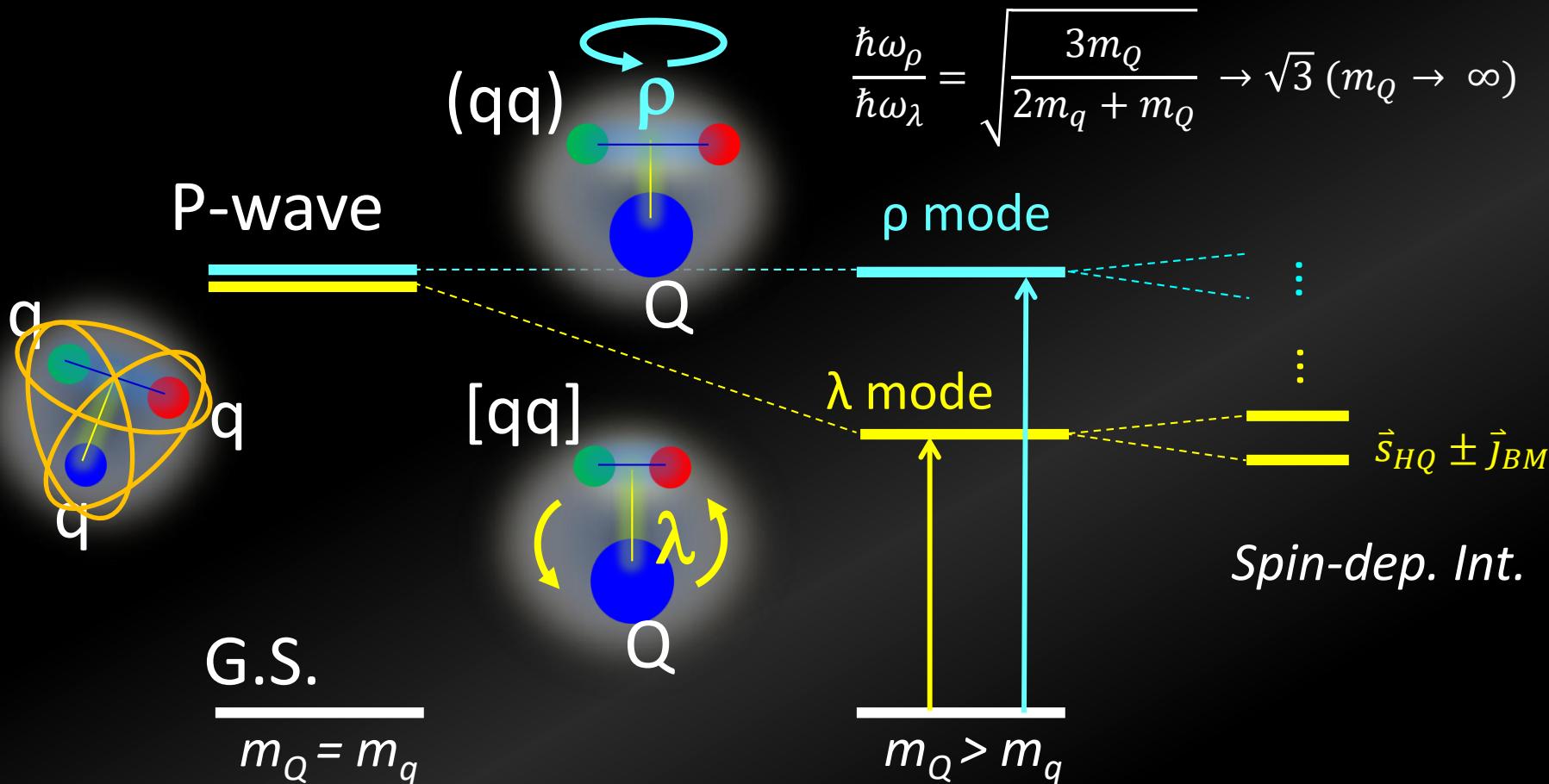
$$V_{CMI}(^1S_0, \bar{3}_c) = 1/2 * V_{CMI}(^1S_0, 1_c)$$

$[qq]$ $[\bar{q}\bar{q}]$

- Motion of “qq” is singled out by a heavy Q
 - **Diquark correlation**
- Level structure, Production rate, Decay properties
 - sensitive to the internal quark(diquark) WFs.
- Properties are expected to depend on a Q mass.

Schematic Level Structure of Heavy Baryons

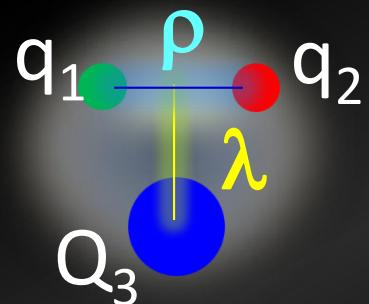
- λ and ρ motions split (Isotope Shift)
- HQ-spin multiplet ($\vec{s}_{HQ} \pm \vec{j}_{Brown\;Muck}$)



Essence of Non-relativistic Formulation in CQM

$$H = H_0 + V_c + V_{ss} + V_{SO} + V_T \dots$$

- H_0 : kinematic term
- V_c : confinement potential
- V_{ss} : spin-spin interaction
- V_{SO} : spin-orbit interaction (LS force)
- V_T : Tensor interaction



Coordinate: $\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3$
 $\rho = (\mathbf{r}_1 - \mathbf{r}_2)/\sqrt{2}$
 $\lambda = (\mathbf{r}_1 + \mathbf{r}_2 - 2\mathbf{r}_3)/\sqrt{6}$
 $\mu_\rho = m_q$
 $\mu_\lambda = 3m_q m_Q/(2m_q + m_Q)$

Confinement

- $H = H_0 + V_c + V_{ss} + \dots$
 $\Psi \sim \psi_\ell \chi_S \phi_I(\text{color}) \rightarrow \text{symmetrize (anti-symm.)}$
- $V_c = k/2 \sum r_{ij}^2 \rightarrow \text{analytic } (<-> \text{Cornell potential } \sum \left(\frac{a}{r_{ij}} + b r_{ij} \right))$

$$\omega_{\lambda,\rho} = \sqrt{3k/m_{\lambda,\rho}}, \quad \left(m_\lambda = \frac{3m_q m_Q}{2m_q + m_Q}, m_\rho = m_q \right)$$

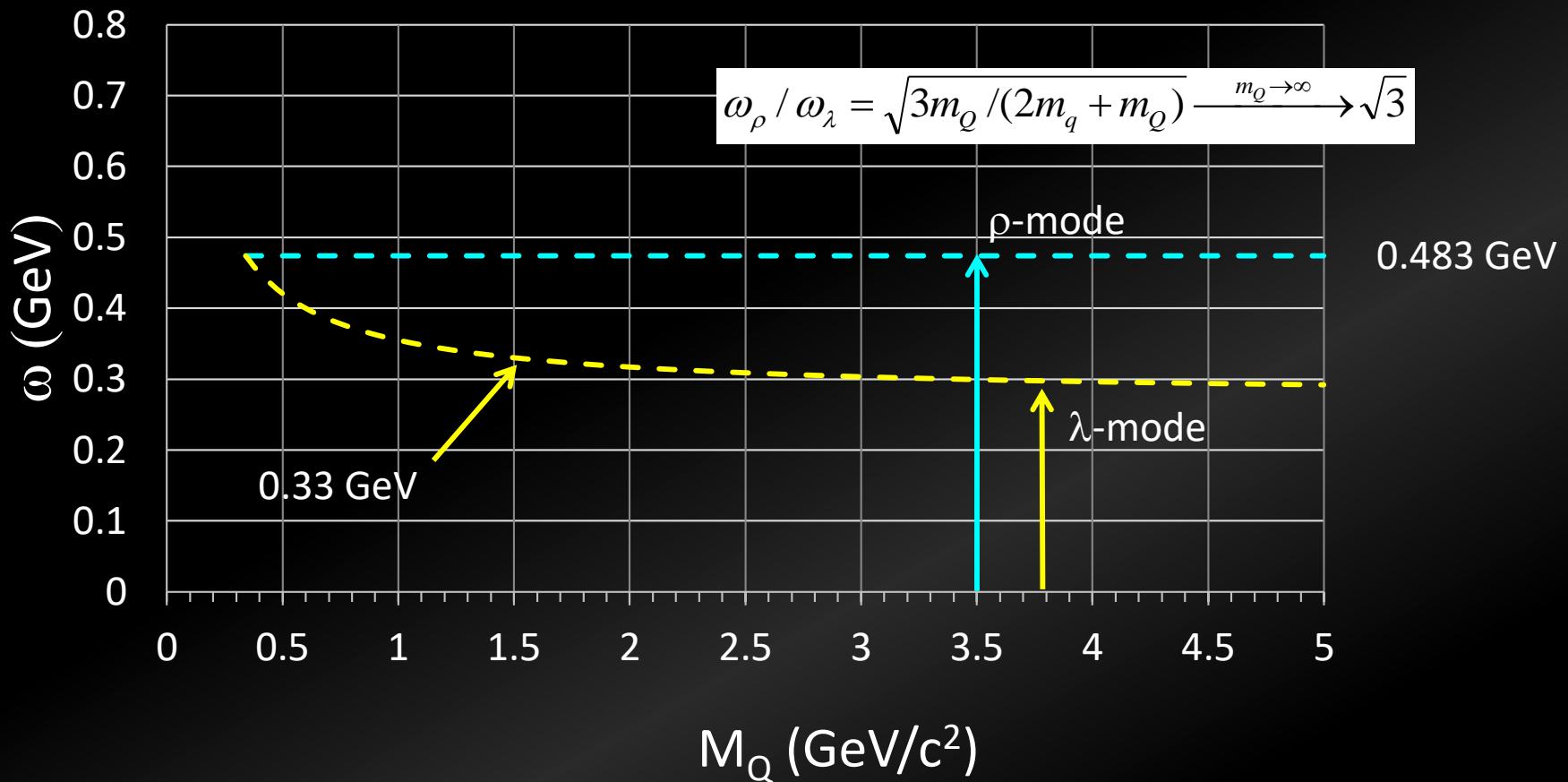
$$k = 0.33^2 m_\lambda / 3, \text{ at } m_Q = 1.5 \text{ GeV}/c^2$$

$$c.f. 1\hbar\omega_\lambda \sim \frac{\Lambda_c\left(\frac{1}{2}^-\right) + 2\Lambda_c\left(\frac{3}{2}^-\right)}{3} - \Lambda_c\left(\frac{1}{2}^+\right) \sim 0.33 \text{ GeV}/c^2 \text{ @ charm sector}$$

$$1\hbar\omega_\rho \sim \hbar\sqrt{3k/m_q} \sim 0.483 \text{ GeV} \text{ (independent of } m_Q)$$

P-wave (ρ , λ -mode excitations)

isotope shift



Spin-spin Interaction

- $H = H_0 + V_c + V_{ss} + \dots$

$\Psi \sim \psi_\ell \chi_S \phi_I(\text{color}) \rightarrow \text{symmetrize (anti-symm.)}$

- $V_c = k/2 \sum r_{ij}^2$

$$\omega_{\lambda,\rho} = \sqrt{3k/m_{\lambda,\rho}}, \quad \left(m_\lambda = \frac{3m_q m_Q}{2m_q + m_Q}, m_\rho = m_q \right)$$

- $V_{ss} = c_s \sum \frac{\sigma_i \cdot \sigma_j}{m_i m_j} \delta(r_{ij}) \quad \langle \chi_S | V_{ss} | \chi_S \rangle:$

$$\Lambda\left(\frac{1}{2}^+\right) = \omega_0 - 3c_s/m_q^2$$

(S, χ^ρ) : “qq”-spin anti-symmm.

$$\Sigma\left(\frac{1}{2}^+\right) = \omega_0 + c_s \left(\frac{1}{m_q^2} - \frac{4}{m_q m_Q} \right)$$

(S, χ^λ) : “qq”-spin symmm., $[q q Q]^{1/2}$

$$\Sigma^*\left(\frac{3}{2}^+\right) = \omega_0 + c_s \left(\frac{1}{m_q^2} + \frac{2}{m_q m_Q} \right)$$

(S, χ^s) : “qqQ” spin symmm.

Spin-spin Interaction

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$$\Sigma \left(\frac{1}{2}^+ \right) = \omega_0 + c_s \left(\frac{1}{m_q^2} - \frac{4}{m_q m_Q} \right)$$

$$\Sigma^* \left(\frac{3}{2}^+ \right) = \omega_0 + c_s \left(\frac{1}{m_q^2} + \frac{2}{m_q m_Q} \right)$$

}

$$\boxed{\rightarrow \frac{\Sigma + 2\Sigma^*}{3} - \Lambda = \mathcal{C}_S \frac{4}{m_q^2} \sim 0.2 \text{ GeV}/c^2}$$

Spin-spin Interaction

- $H = H_0 + V_c + V_{ss} + \dots$
 $\Psi \sim \psi_\ell \chi_S \phi_I(\text{color}) \rightarrow \text{symmetrize (anti-symm.)}$

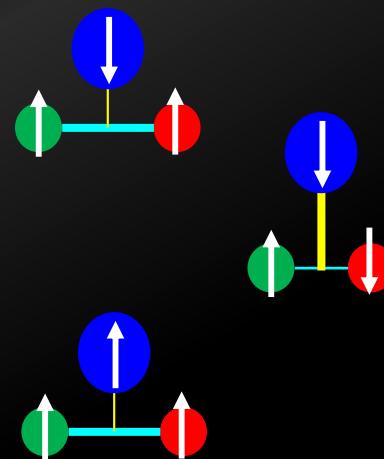
- $V_c = k/2 \sum r_{ij}^2$

$$\omega_{\lambda,\rho} = \sqrt{3k/m_{\lambda,\rho}}, \quad \left(m_\lambda = \frac{3m_q m_Q}{2m_q + m_Q}, m_\rho = m_q \right)$$

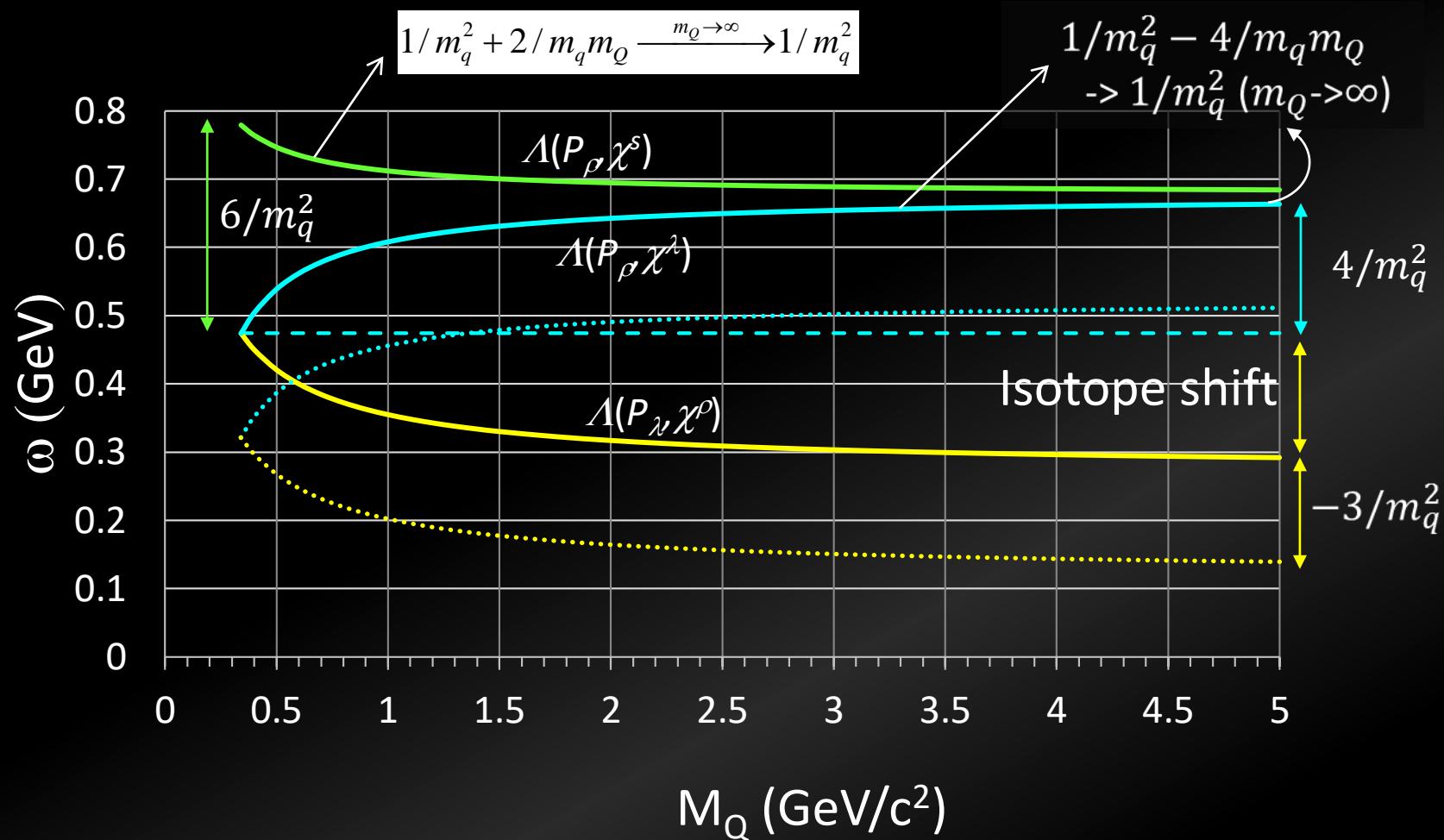
- $V_{ss} = c_s \sum \frac{\sigma_i \cdot \sigma_j}{m_i m_j} \delta(r_{ij})$ $\langle \chi_S | V_{ss} | \chi_S \rangle:$

$$\Lambda\left(\frac{1}{2}^-, \frac{3}{2}^-\right) = \begin{cases} \omega_\rho + c_s \left(\frac{1}{m_q^2} - \frac{4}{m_q m_Q} \right) (\ell_\rho = 1, \chi^\rho) \\ \omega_\lambda - 3c_s/m_q^2 \quad (\ell_\lambda = 1, \chi^\lambda) \end{cases}$$

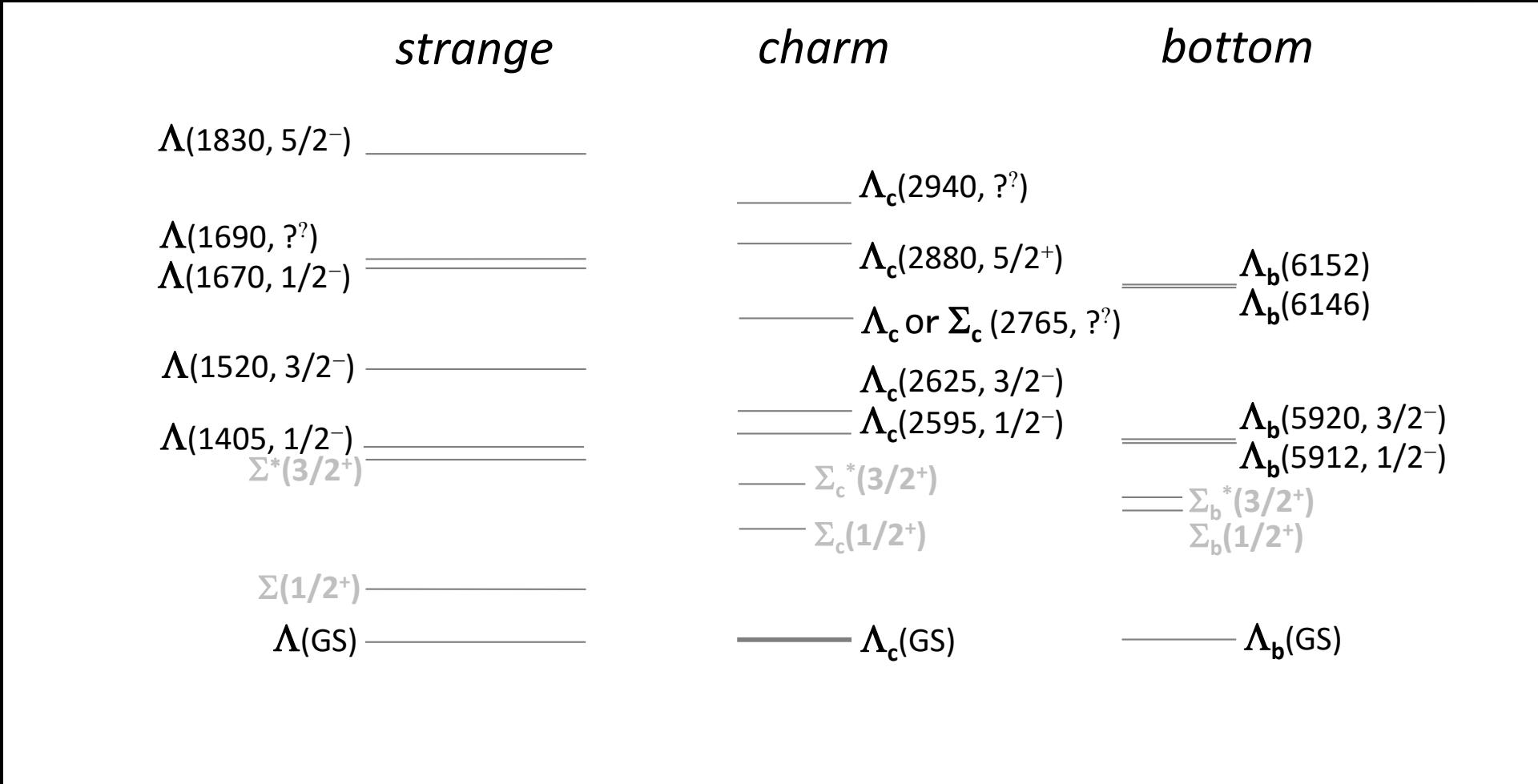
$$\Lambda\left(\frac{1}{2}^-, \frac{3}{2}^-, \frac{5}{2}^-\right) = \omega_\rho + c_s \left(\frac{1}{m_q^2} + \frac{2}{m_q m_Q} \right) (\ell_\rho = 1, \chi^s)$$



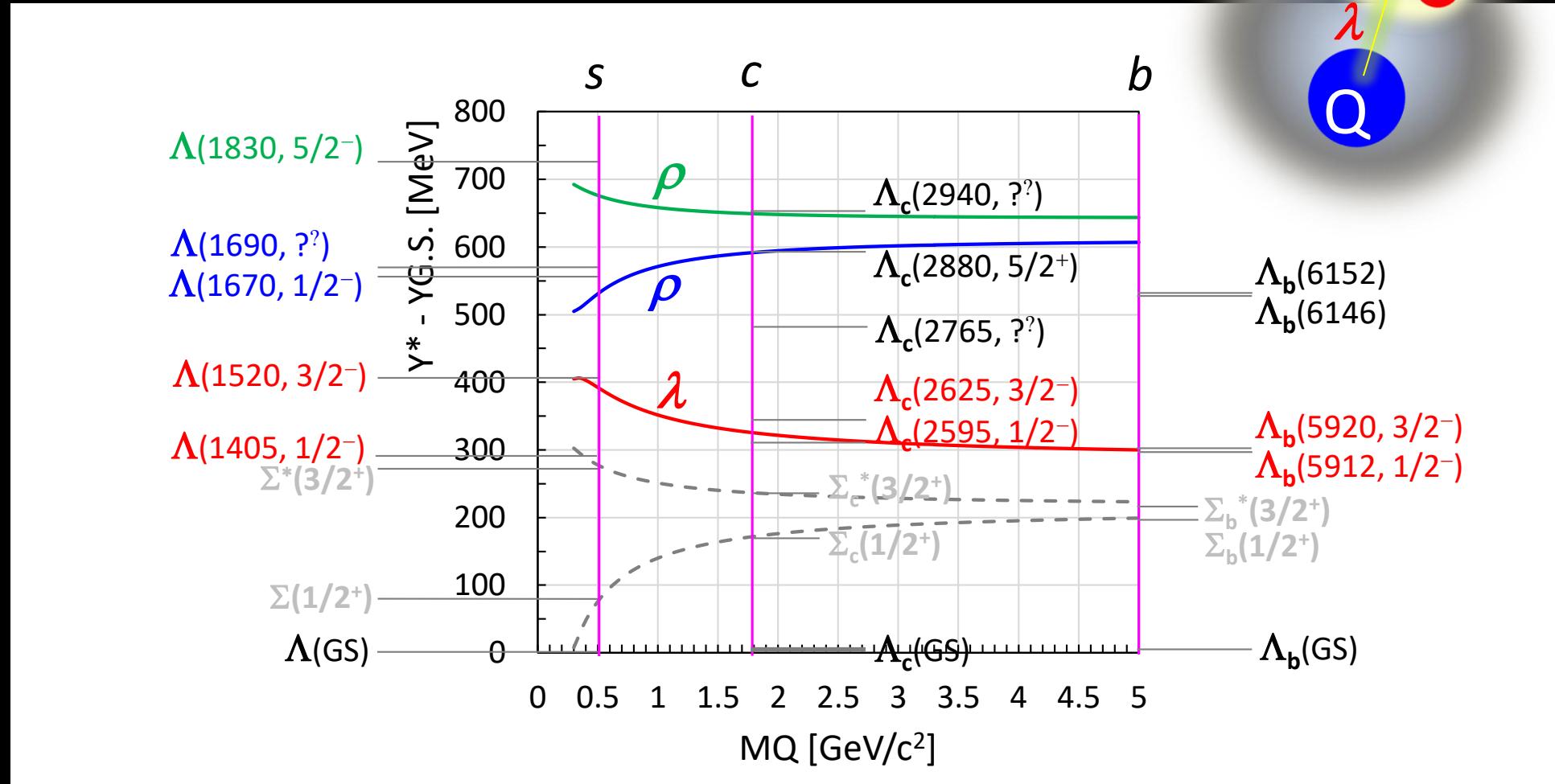
$\Lambda(P,\chi)$ (ρ, λ -mode excitations w/ V_{ss})



Lambda Baryons

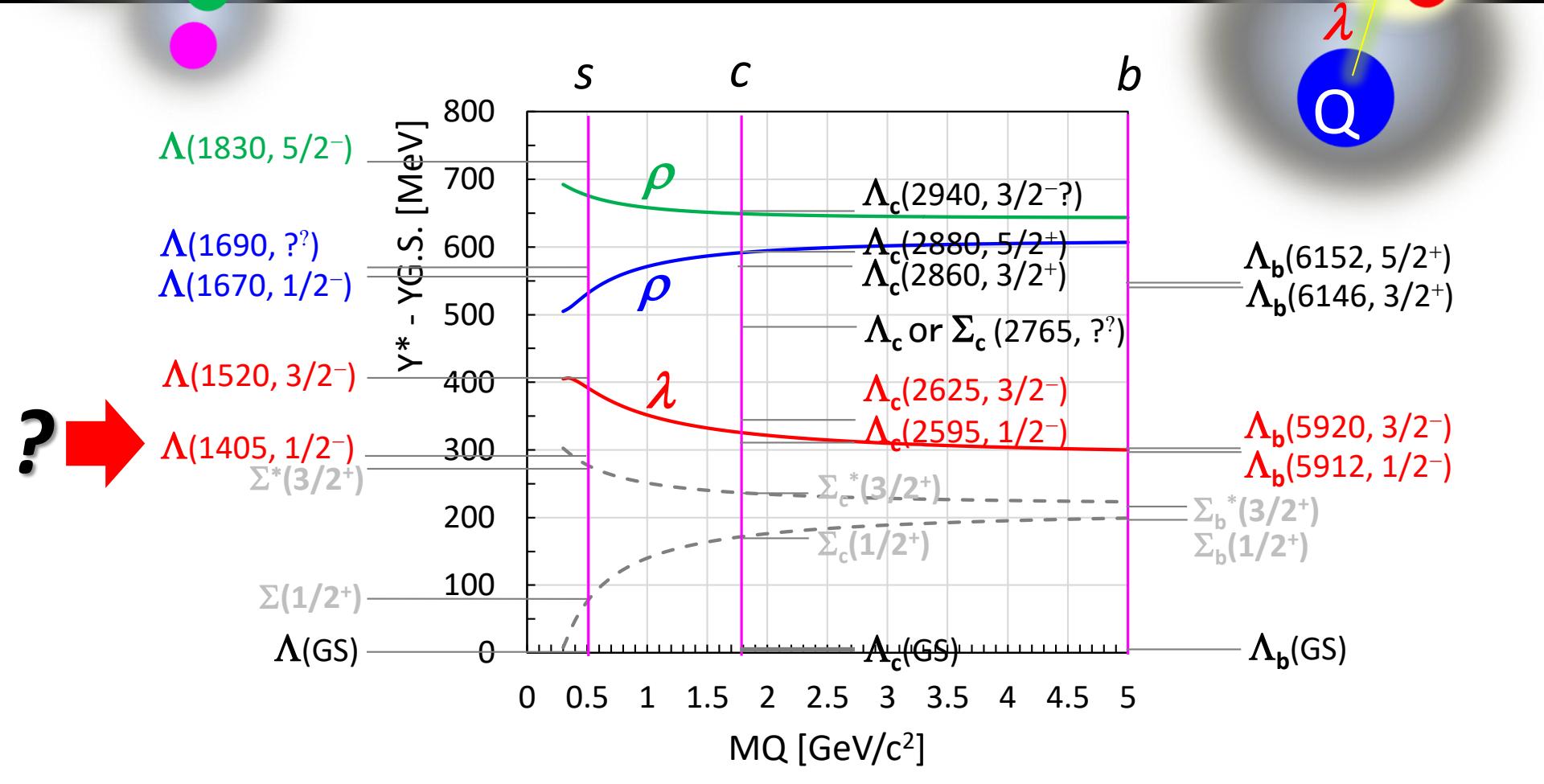


Lambda Baryons



T. Yoshida et al.,
 Phys. Rev. D92, 114029(2015)

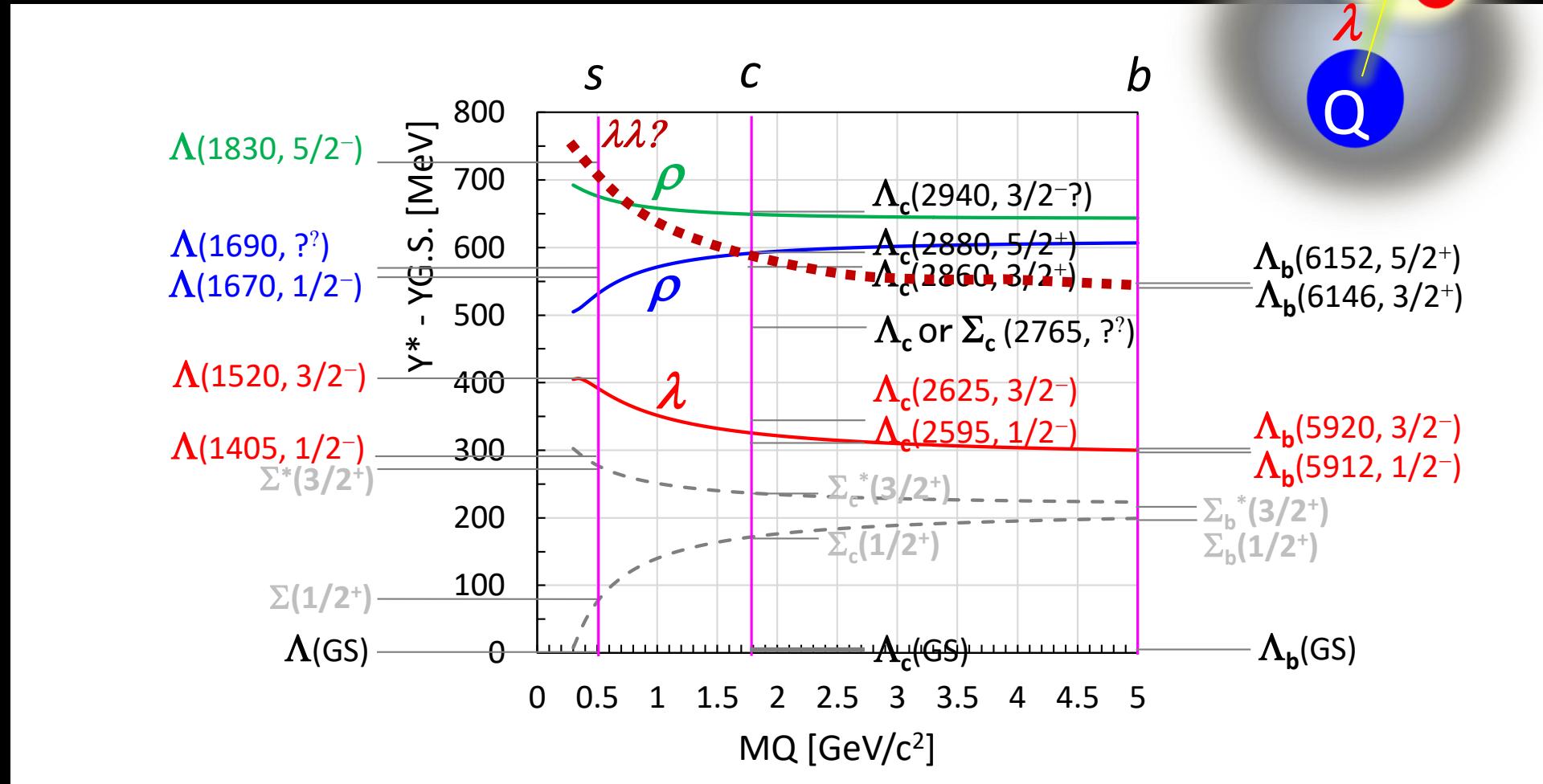
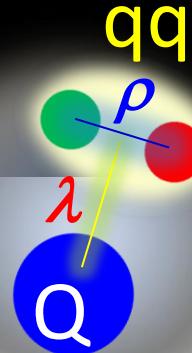
Lambda Baryons



non-rel. QM: $H = H_0 + V_{conf} + V_{SS} + V_{LS} + V_T$
 $\rho-\lambda$ mixing (cal. By T. Yoshida)

T. Yoshida et al.,
 Phys. Rev. D92, 114029(2015)

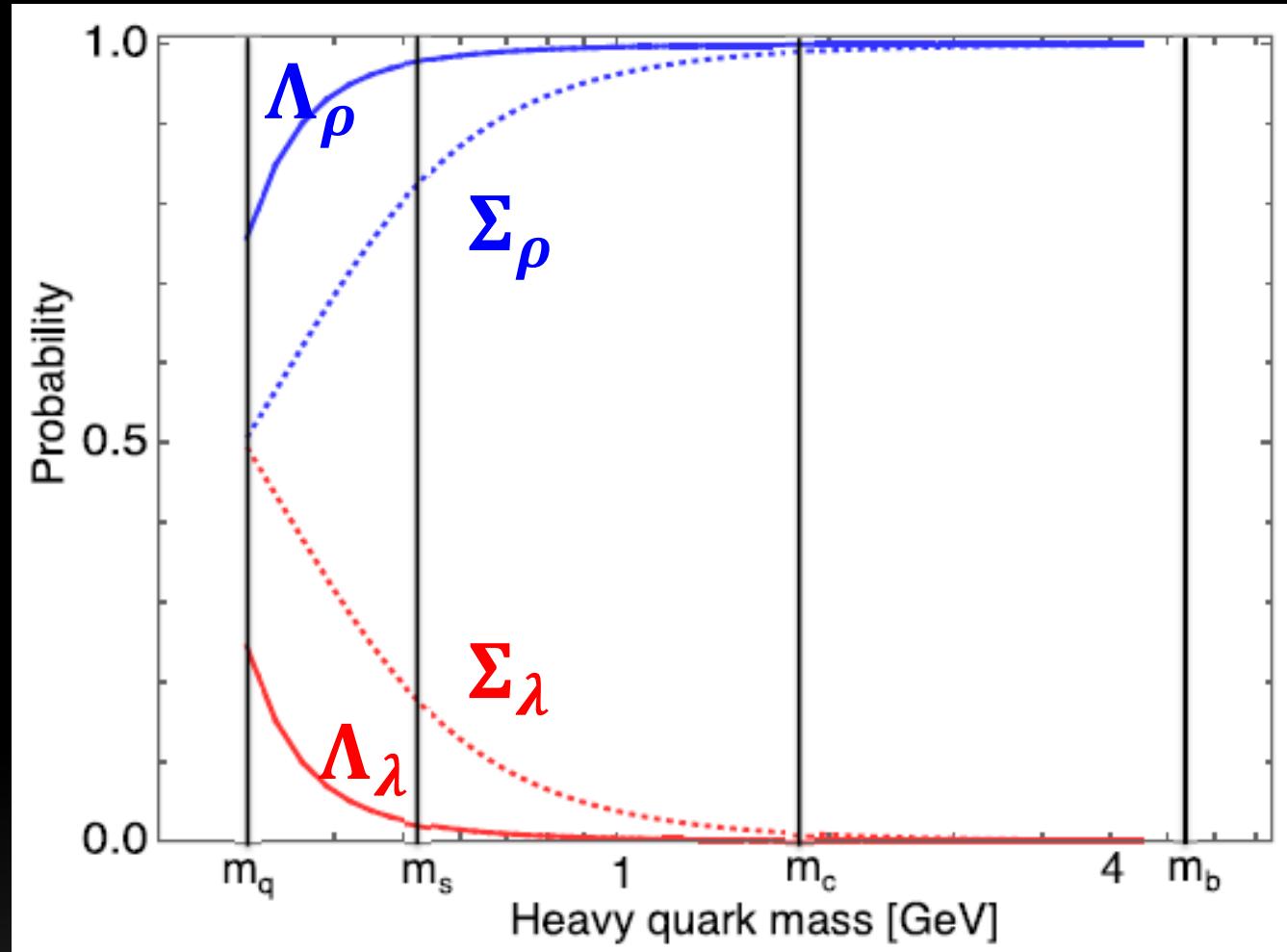
Lambda Baryons



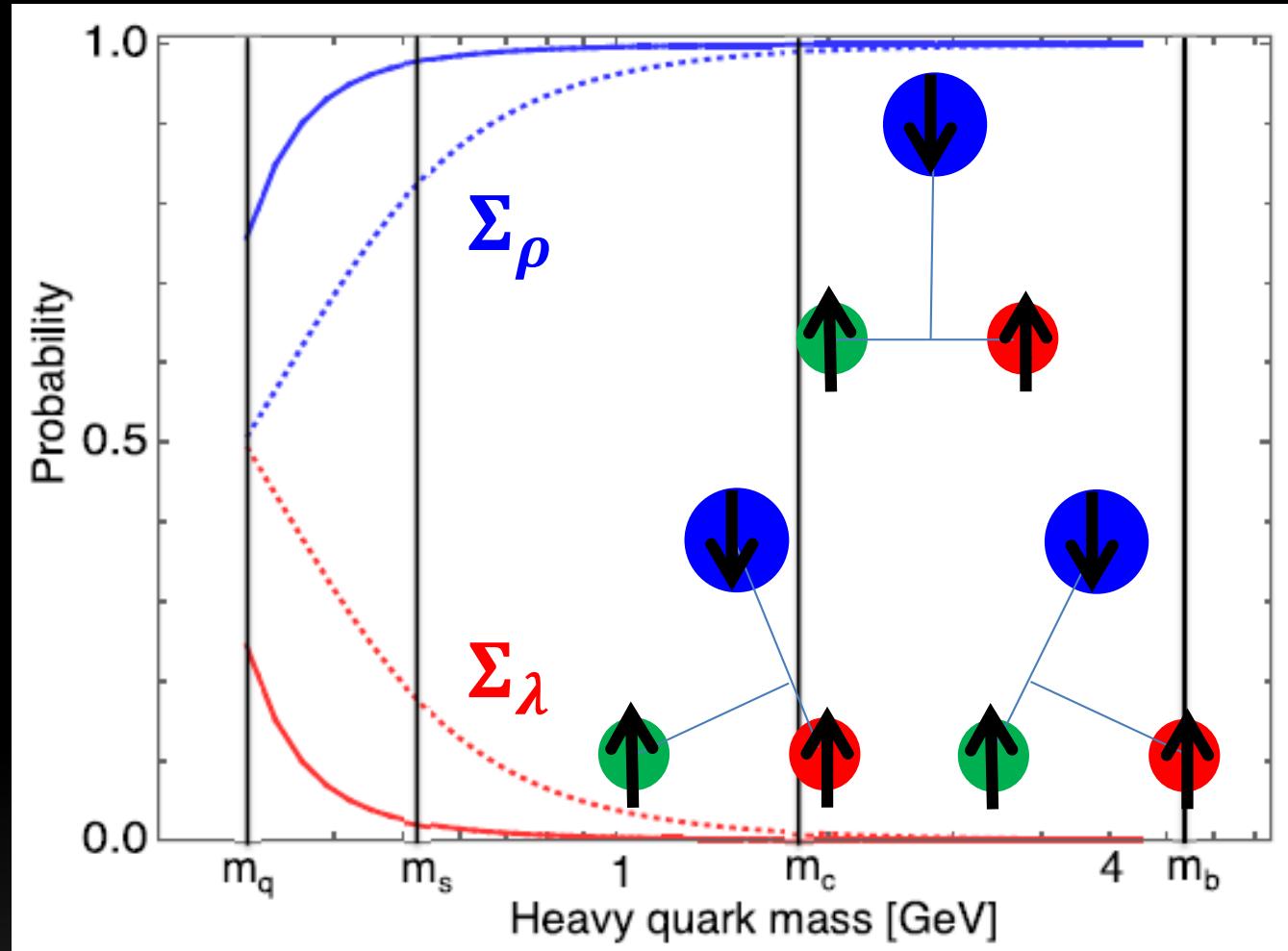
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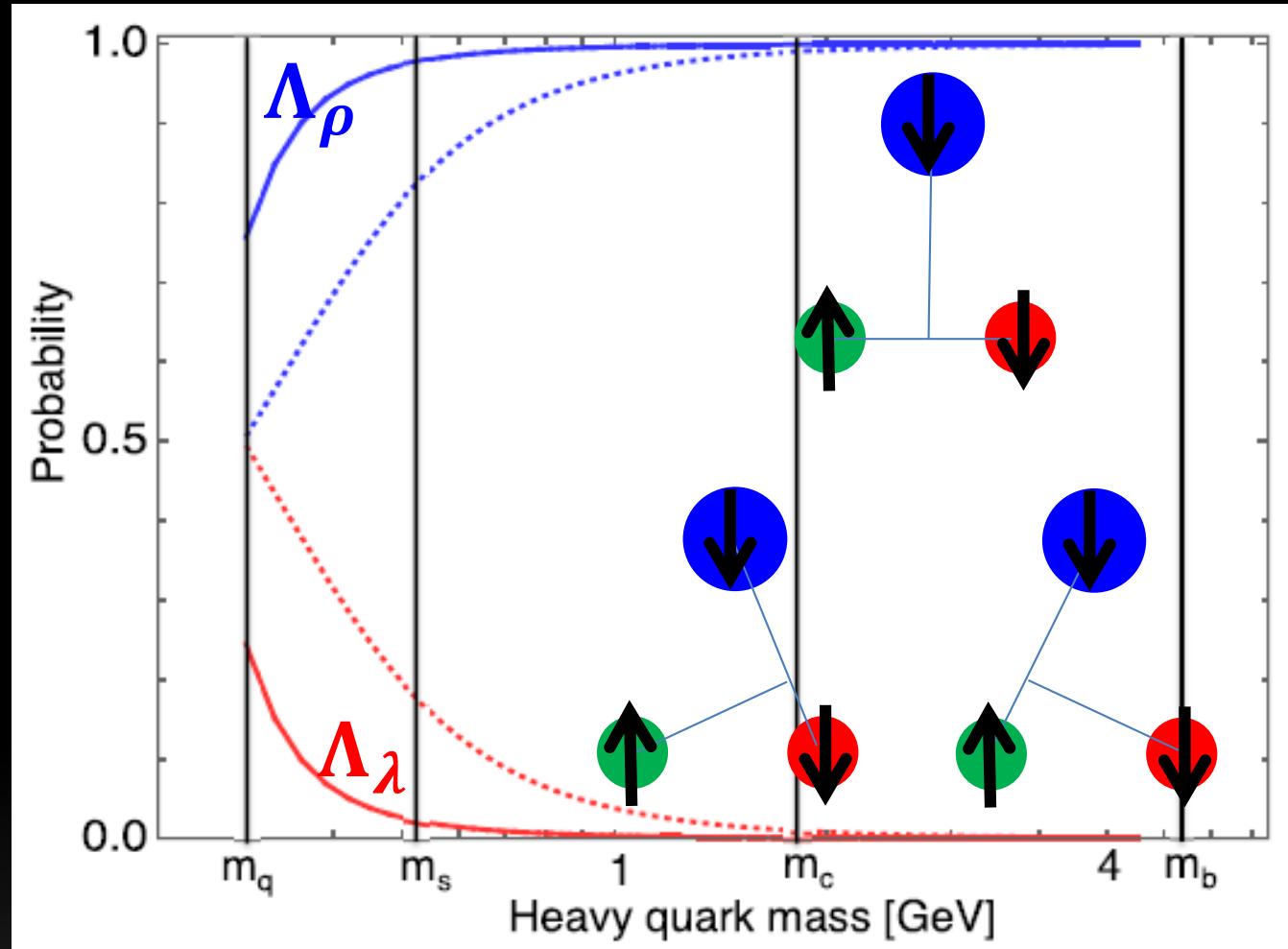
$\Lambda, \Sigma(P, \chi)$ (ρ , λ -mode mixing w/ V_{ss})



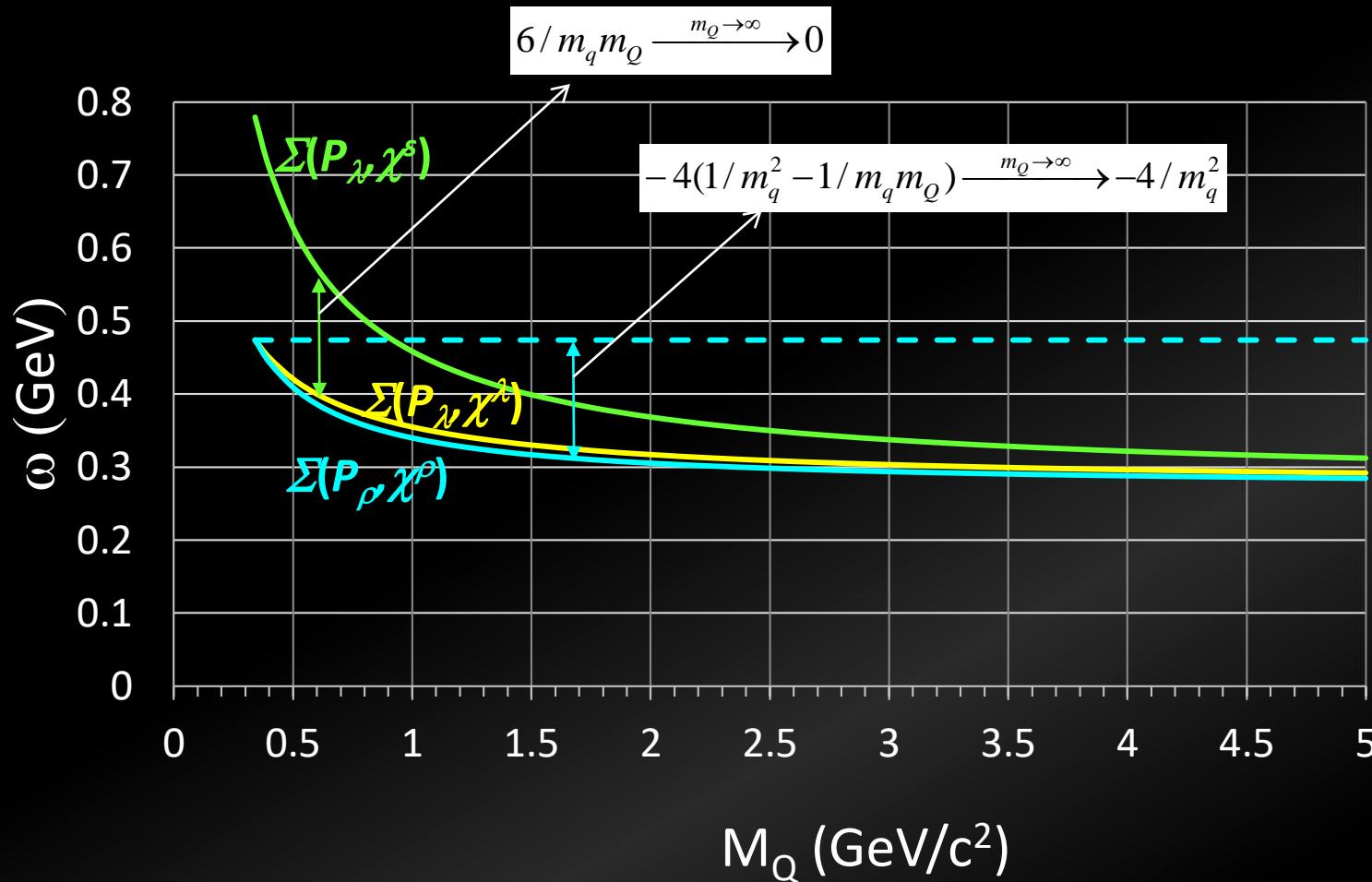
$\Lambda, \Sigma(P, \chi)$ (ρ, λ -mode mixing w/ V_{ss})



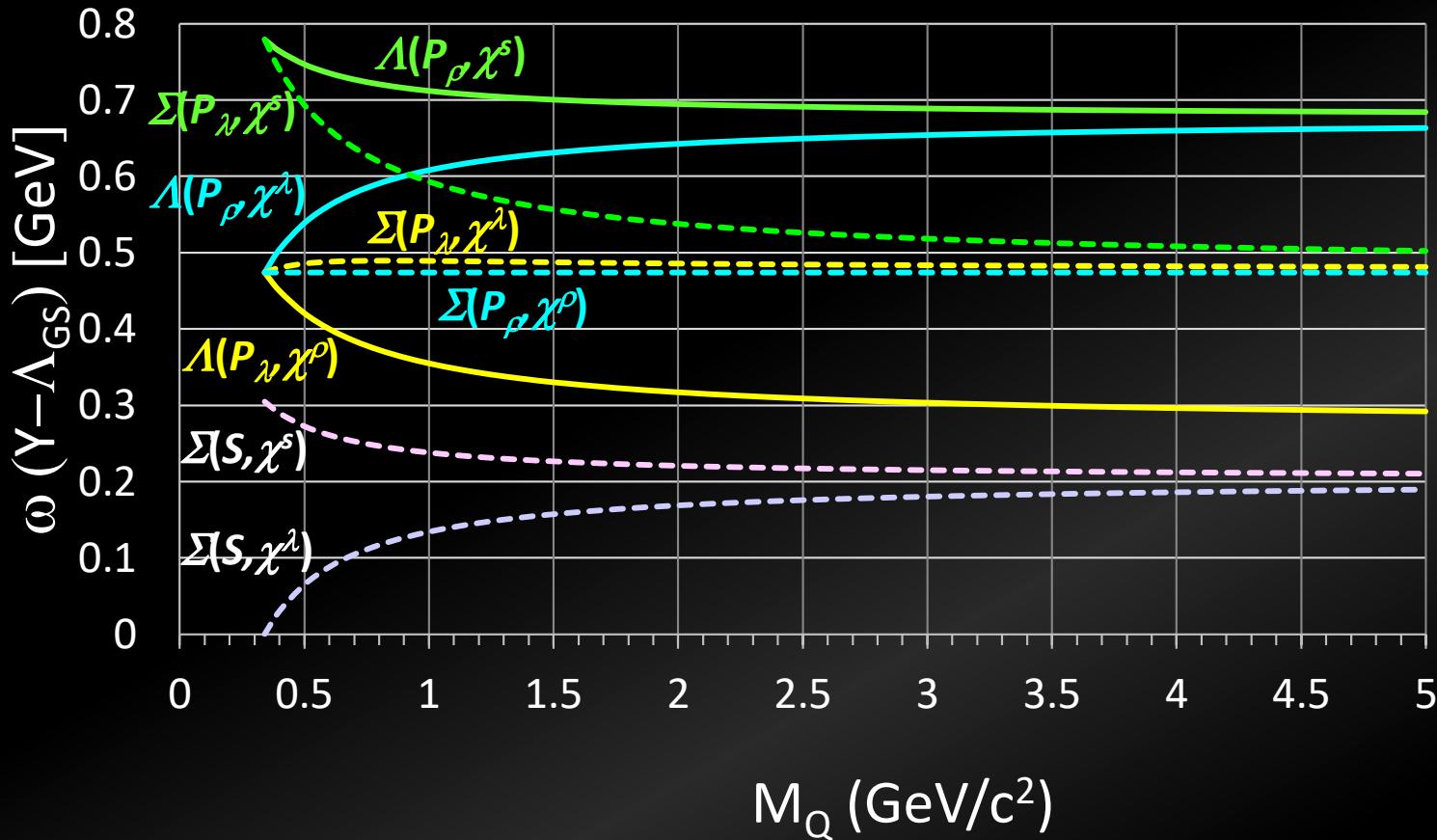
$\Lambda, \Sigma(P, \chi)$ (ρ , λ -mode mixing w/ V_{ss})



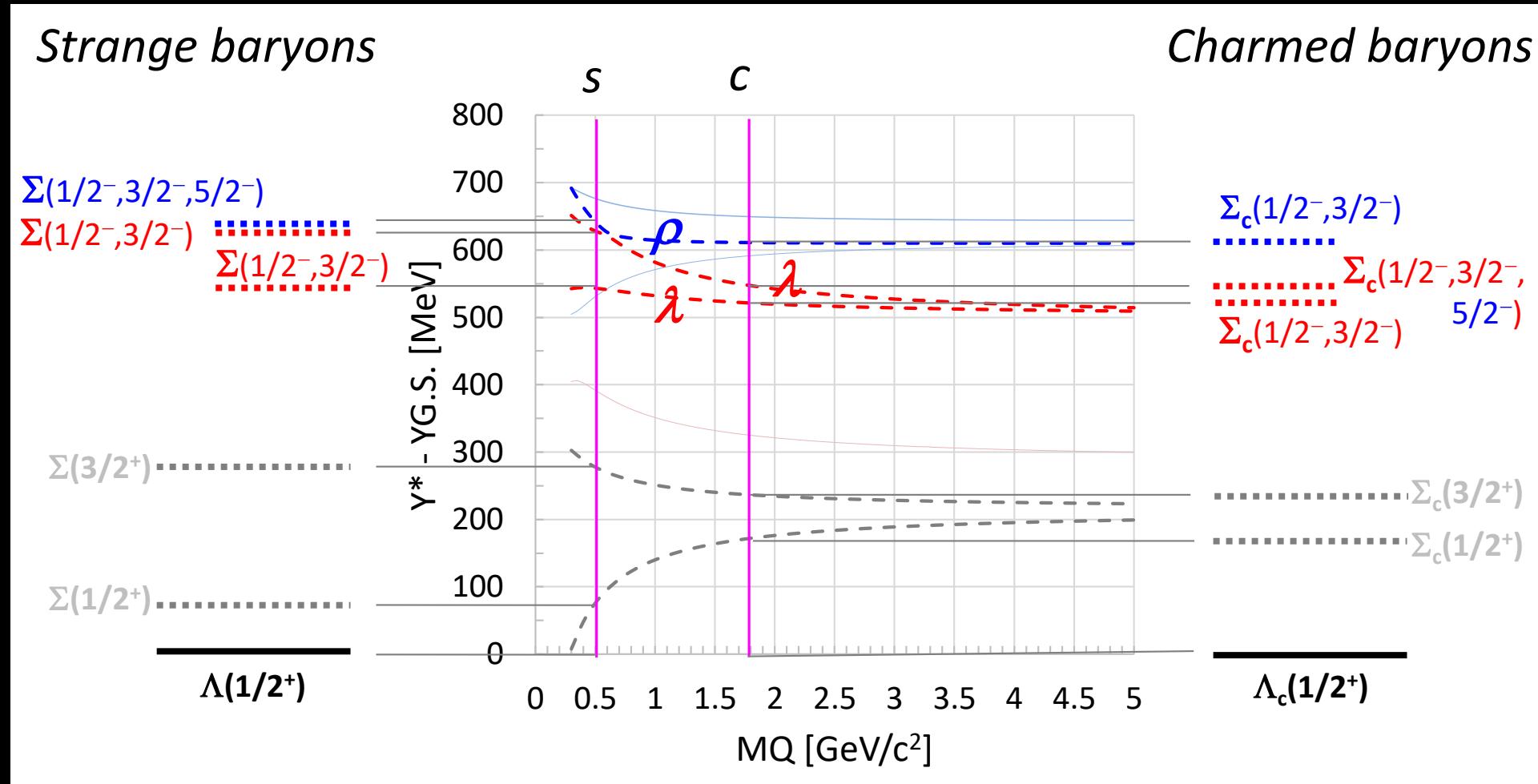
$\Sigma(P,\chi)$ (ρ, λ -mode excitations w/ V_{ss})



$\Lambda, \Sigma(1/2^-)$ (ρ, λ -mode excitations w/ V_{ss})



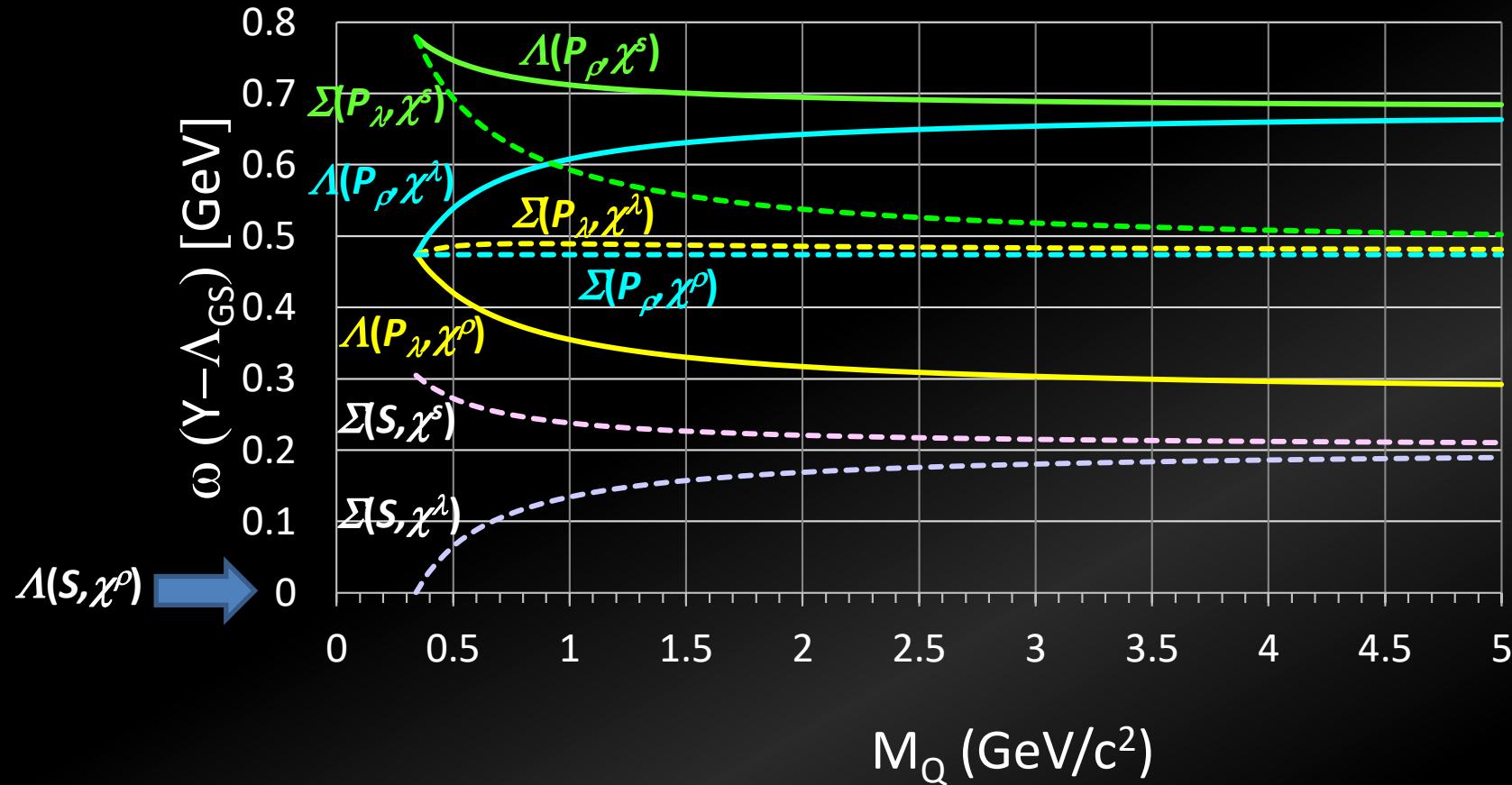
CQM calculation (P-wave Sigma)



non-rel. QM: $H = H_0 + V_{conf} + V_{SS} + V_{LS} + V_T$
 ρ–λ mixing (cal. By T. Yoshida)

T. Yoshida et al.,
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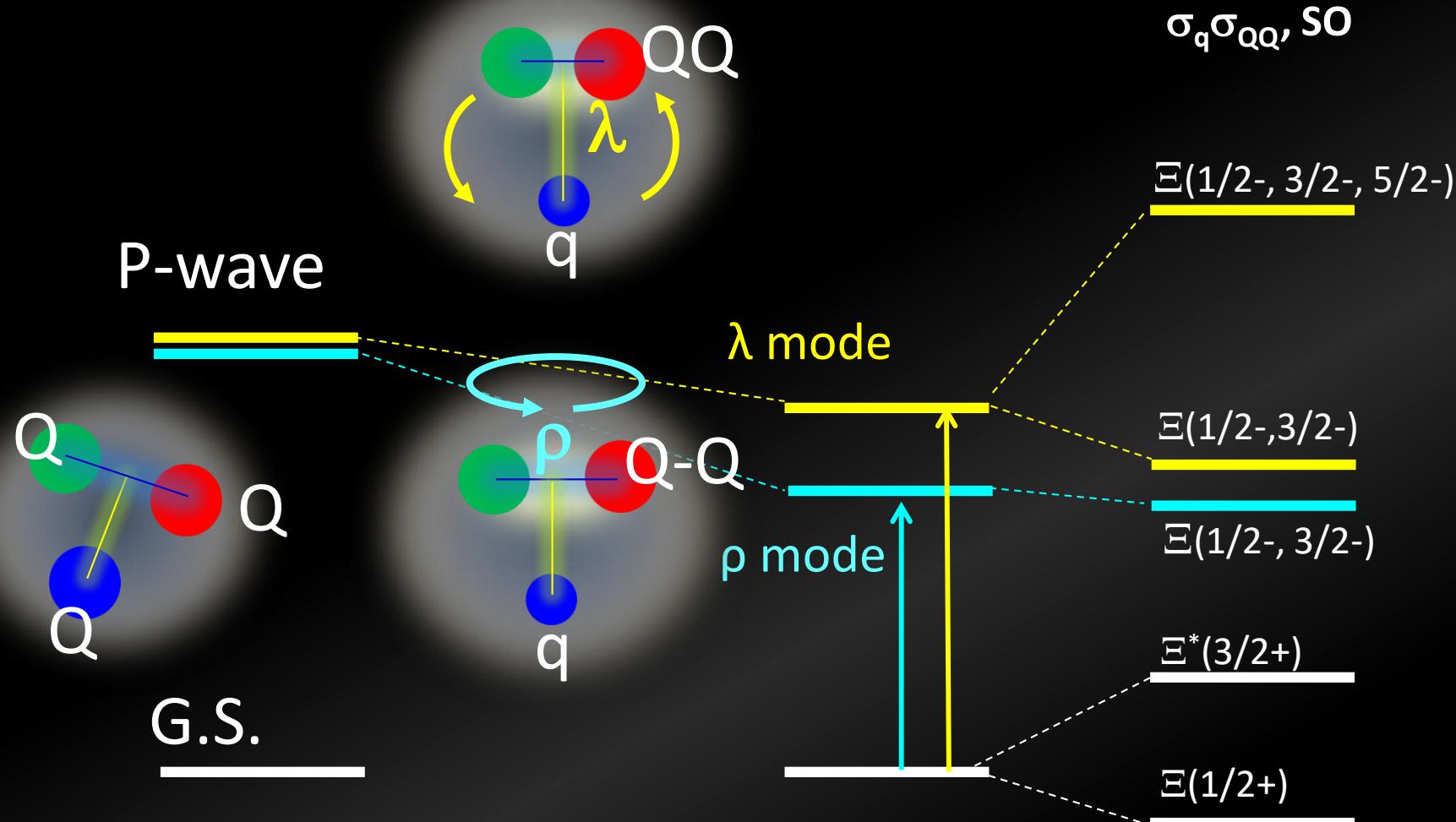
$\Lambda, \Sigma(1/2^-)$ (ρ, λ -mode excitations w/ V_{ss})



QQq

Level Structure of double-strange baryons

- λ and ρ mode excitations interchange



Confinement

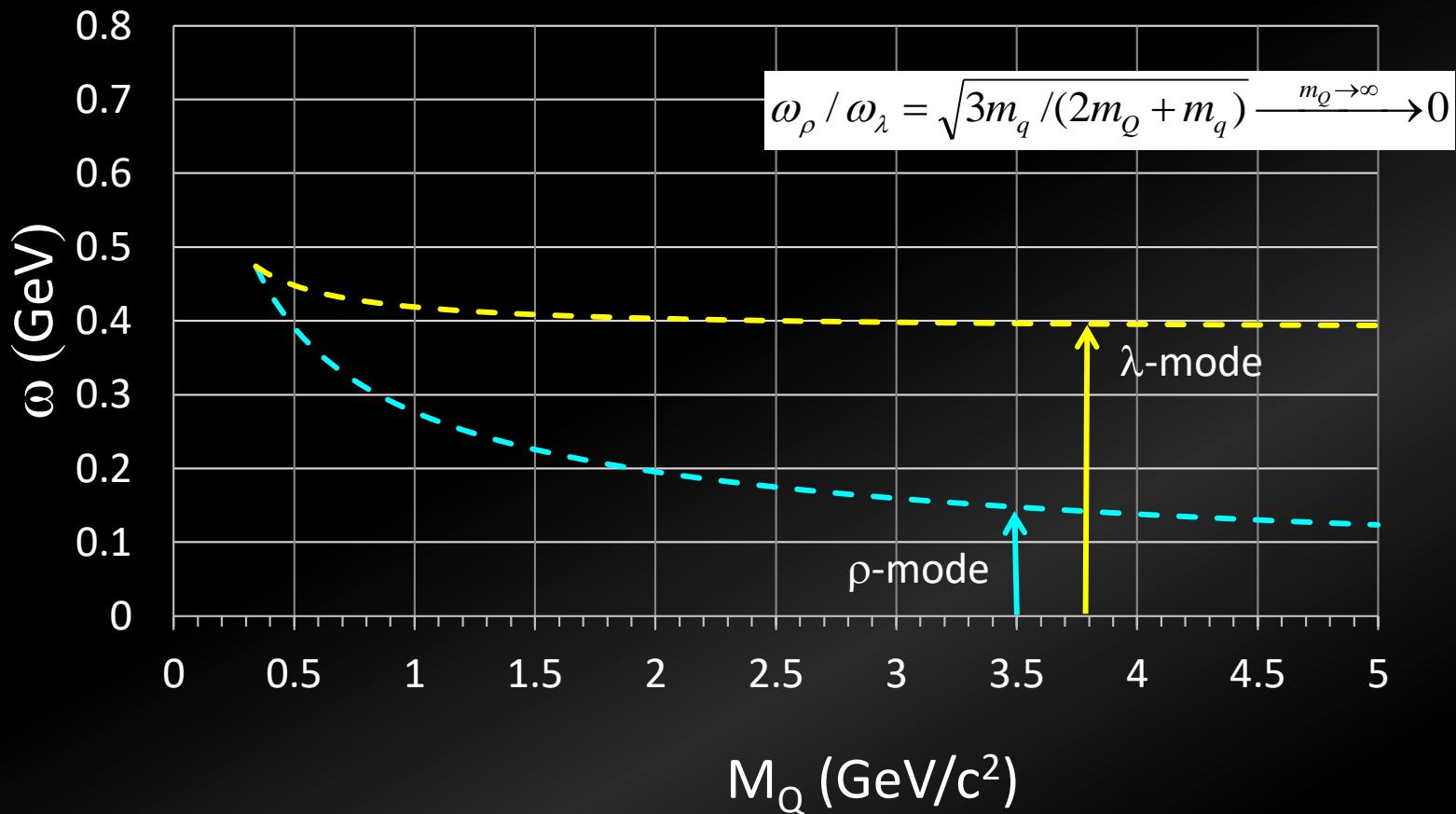
- $H = H_0 + V_c + V_{ss} + \dots$
 $\Psi \sim \psi_\ell \chi_S (\text{Isospin * color}) \rightarrow \text{symmetrize (anti-symm.)}$

- $V_c = k/2 \sum r_{ij}^2$

$$\omega_{\lambda,\rho} = \sqrt{3k/m_{\lambda,\rho}}, \quad \left(m_\lambda = \frac{3m_Q m_q}{2m_Q + m_q}, m_\rho = m_Q \right)$$

P-wave (ρ , λ -mode excitations)

isotope shift



Spin-spin Interaction

- $H = H_0 + V_c + V_{ss} + \dots$

$\Psi \sim \psi_\ell \chi_S$ (Isospin * color) \rightarrow symmetrize (anti-symm.)

- $V_c = k/2 \sum r_{ij}^2$

$$\omega_{\lambda,\rho} = \sqrt{3k/m_{\lambda,\rho}}, \quad \left(m_\lambda = \frac{3m_Q m_q}{2m_Q + m_q}, m_\rho = m_Q \right)$$

- $V_{ss} = c_s \sum \frac{\sigma_i \cdot \sigma_j}{m_i m_j} \delta(r_{ij})$ $\langle \chi_S | V_{ss} | \chi_S \rangle:$

$$\Xi \left(\frac{1}{2}^+\right) = \omega_0 + c_s \left(\frac{1}{m_Q^2} - \frac{4}{m_q m_Q} \right) \quad (S, \chi^\lambda) : \text{"QQ"-spin symm., [QQq]}^{1/2}$$

$$\Xi^* \left(\frac{3}{2}^+\right) = \omega_0 + c_s \left(\frac{1}{m_Q^2} + \frac{2}{m_q m_Q} \right) \quad (S, \chi^s) : \text{"QQq" spin symm.}$$

Spin-spin Interaction

- $H = H_0 + V_c + V_{ss} + \dots$

$\Psi \sim \psi_\ell \chi_S$ (Isospin * color) \rightarrow symmetrize (anti-symm.)

- $V_c = k/2 \sum r_{ij}^2$

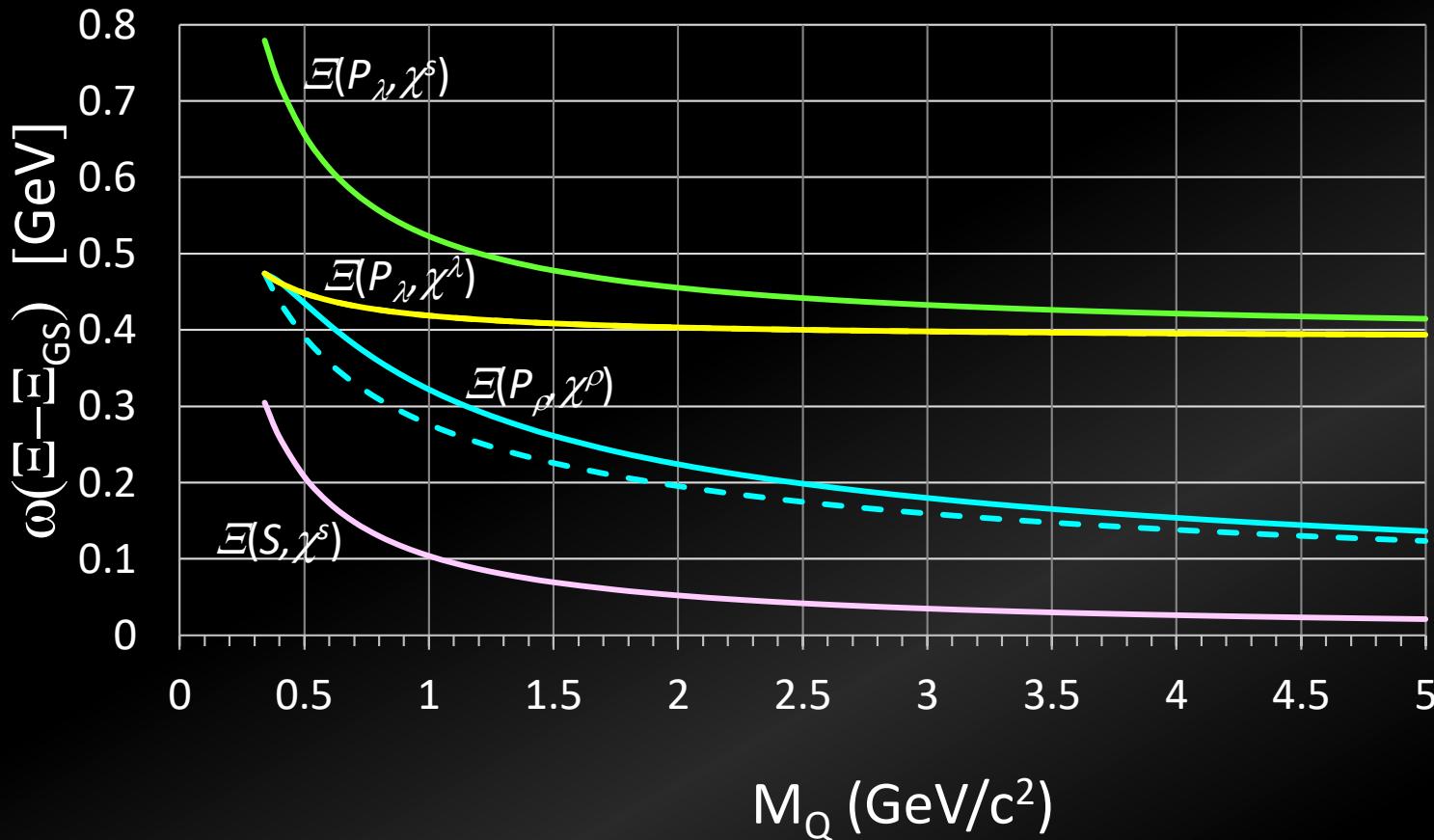
$$\omega_{\lambda,\rho} = \sqrt{3k/m_{\lambda,\rho}}, \quad \left(m_\lambda = \frac{3m_Q m_q}{2m_Q + m_q}, m_\rho = m_Q \right)$$

- $V_{ss} \sim c_s \sum \frac{\sigma_i \cdot \sigma_j}{m_i m_j} \delta(r_{ij})$ $\langle \chi_S | V_{ss} | \chi_S \rangle:$

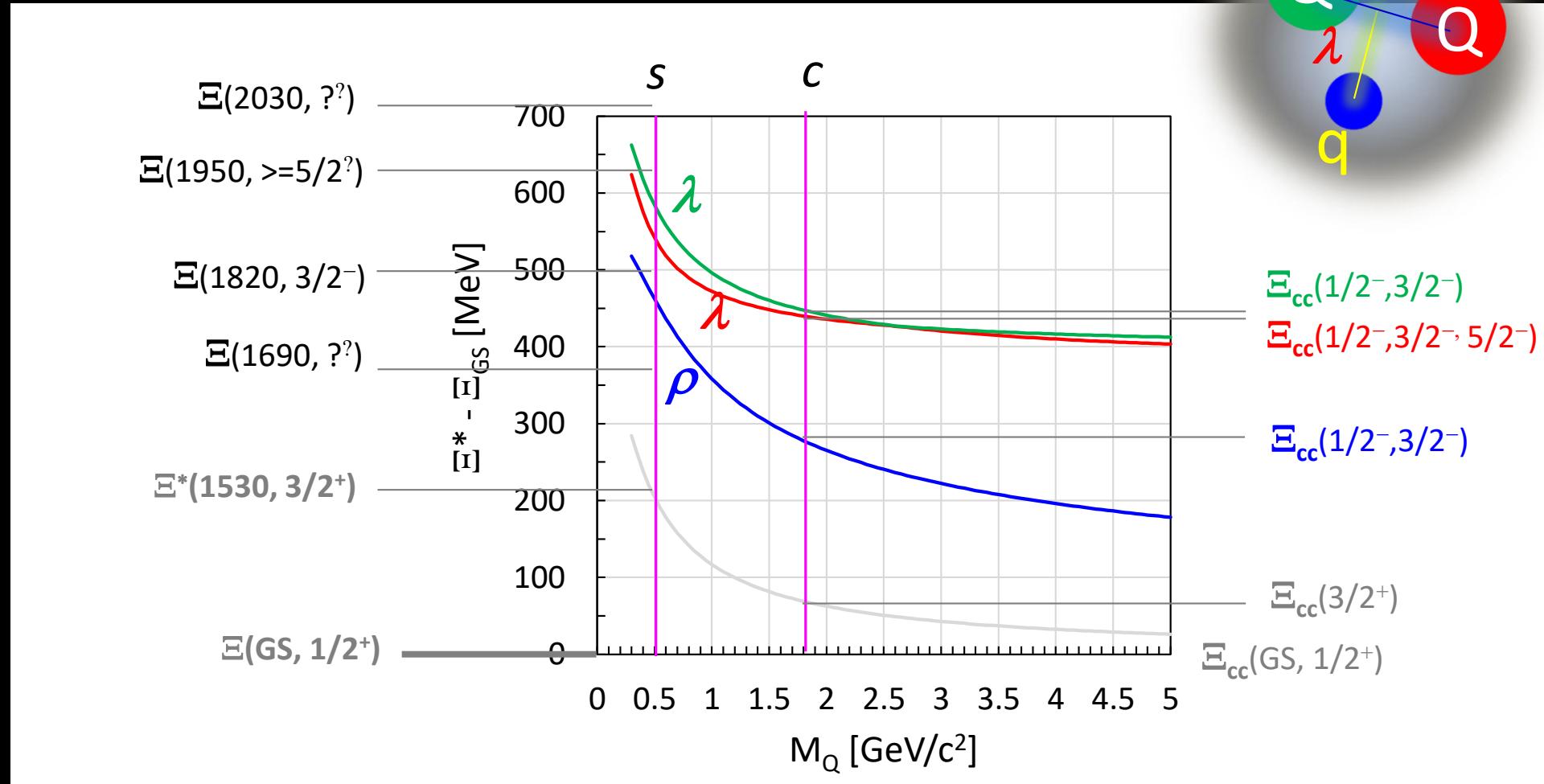
$$\Xi\left(\frac{1}{2}^-, \frac{3}{2}^-\right) = \begin{cases} \omega_\rho - 3c_s/m_Q^2 & (\ell_\rho = 1, \chi^\rho) \\ \omega_\lambda + c_s \left(\frac{1}{m_Q^2} - \frac{4}{m_q m_Q} \right) & (\ell_\lambda = 1, \chi^\lambda) \end{cases}$$

$$\Xi\left(\frac{1}{2}^-, \frac{3}{2}^-, \frac{5}{2}^-\right) = \omega_\lambda + c_s \left(\frac{1}{m_Q^2} + \frac{2}{m_q m_Q} \right) (\ell_\lambda = 1, \chi^s)$$

Ξ (ρ , λ -mode excitations w/ V_{ss})



Xi Baryons



non-rel. QM: $H = H_0 + V_{conf} + V_{SS} + V_{LS} + V_T$
 $\rho - \lambda$ mixing (cal. By T. Yoshida)

T. Yoshida et al.,
 Phys. Rev. D92, 114029(2015)

Measured Ξ (PDG)

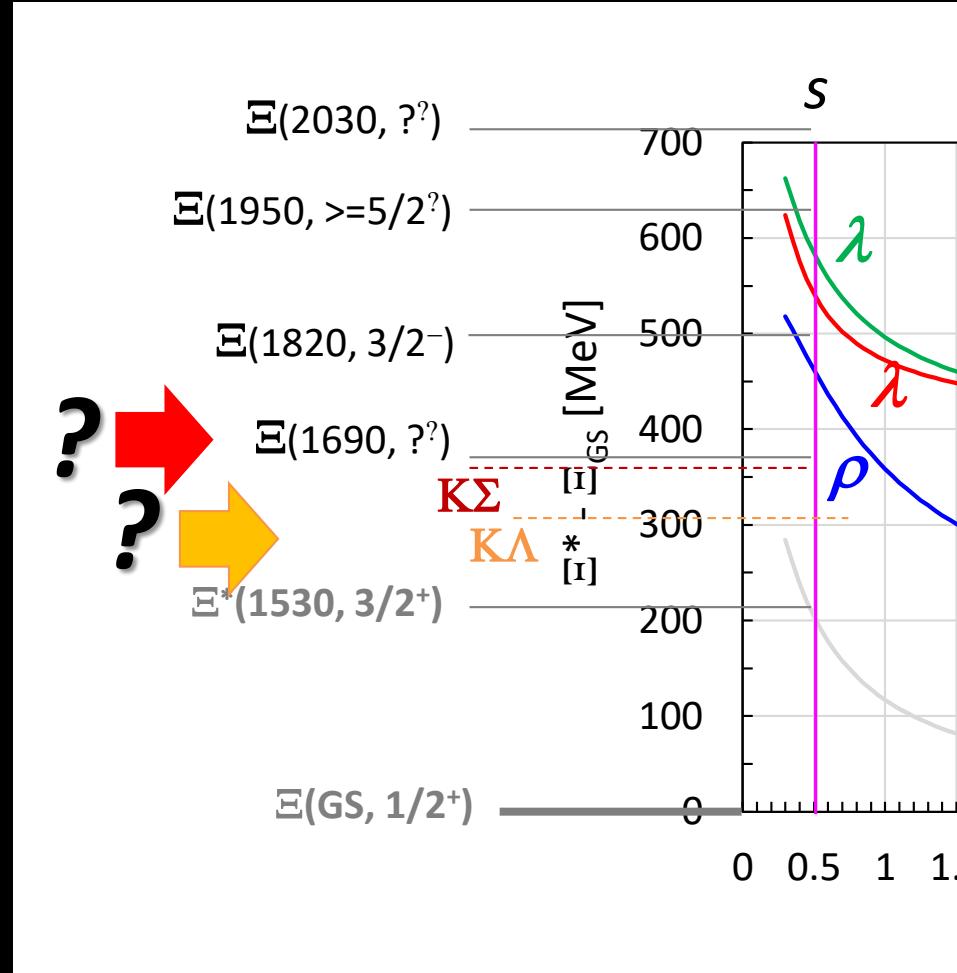
Threshold		JP	rating	Width [MeV]	$\rightarrow \Xi\pi$ [%]	$\rightarrow \Lambda K$ [%]	$\rightarrow \Sigma K$ [%]	
	$\Xi(2500)$??	1*	150?				
	$\Xi(2370)$??	2*	80?				$\Omega K \sim 9 \pm 4$
	$\Xi(2250)$??	2*	47+-27?				
	$\Xi(2120)$??	1* $\Sigma \bar{K}$	25?				
	$\Xi(2030)$	$>=5/2?$	3*	20^{+15}_{-5}	small	~ 20	~ 80	Why ΣK ?
	$\Xi(1950)$??	3*	60+-20	seen	seen		
	$\Xi(1820)$	3/2-	3*	24^{+15}_{-10}	small	Large	Small	
	$\Xi(1690)$??	3*	<30	seen	seen	seen	
	$\Xi(1620)$??	1*	20~40?				
	$\Xi(1530)$	3/2+	4*	19	100			

✓ Most of spins/parities have NOT been determined yet.

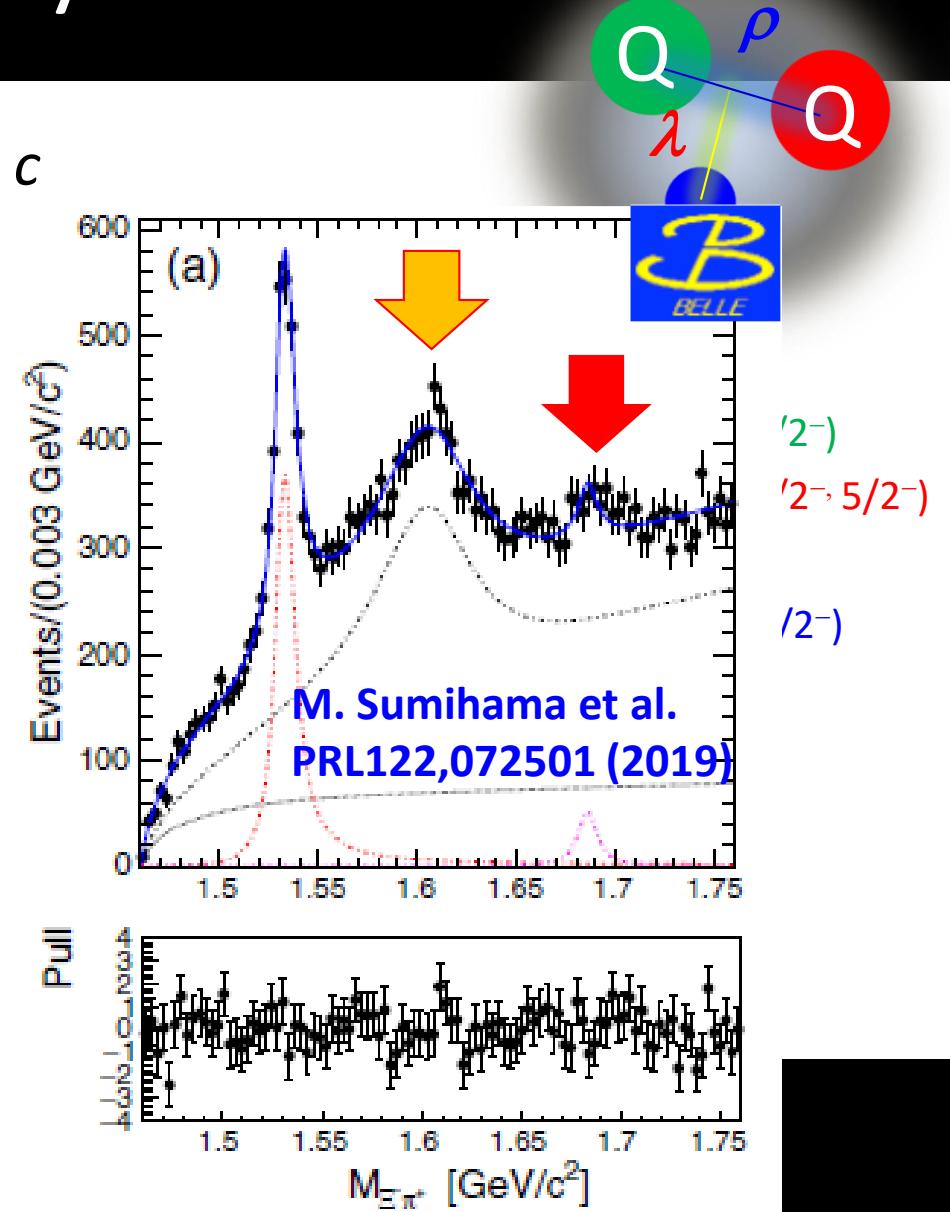
✓ Why the $\Xi^* \rightarrow \pi \Xi$ decay seems to be suppressed?

✓ expected to reflect QQq configuration.

Xi Baryons

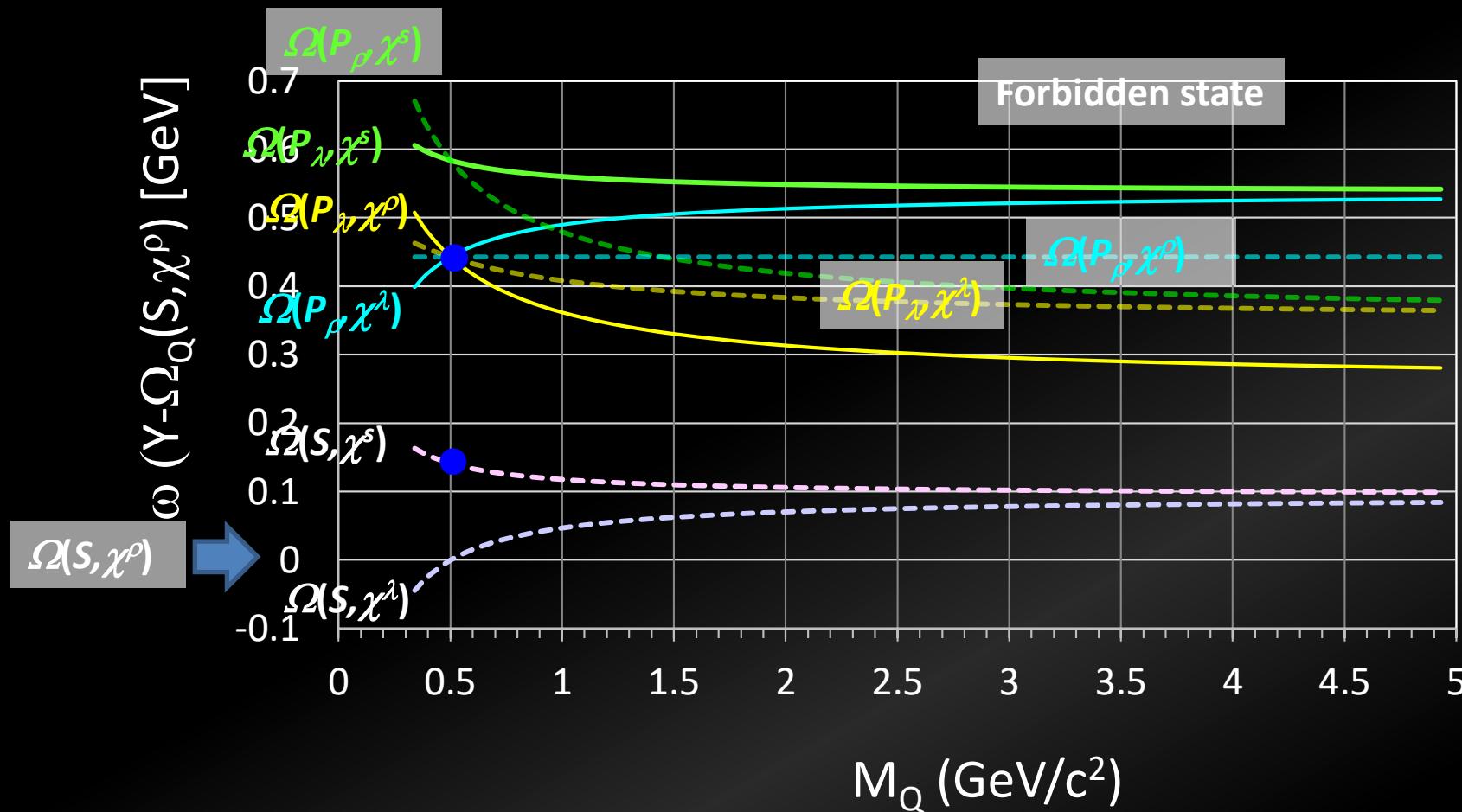


non-rel. QM: $H = H_0 + V_{conf} + V_{SS} + V_{LS} + V_T$
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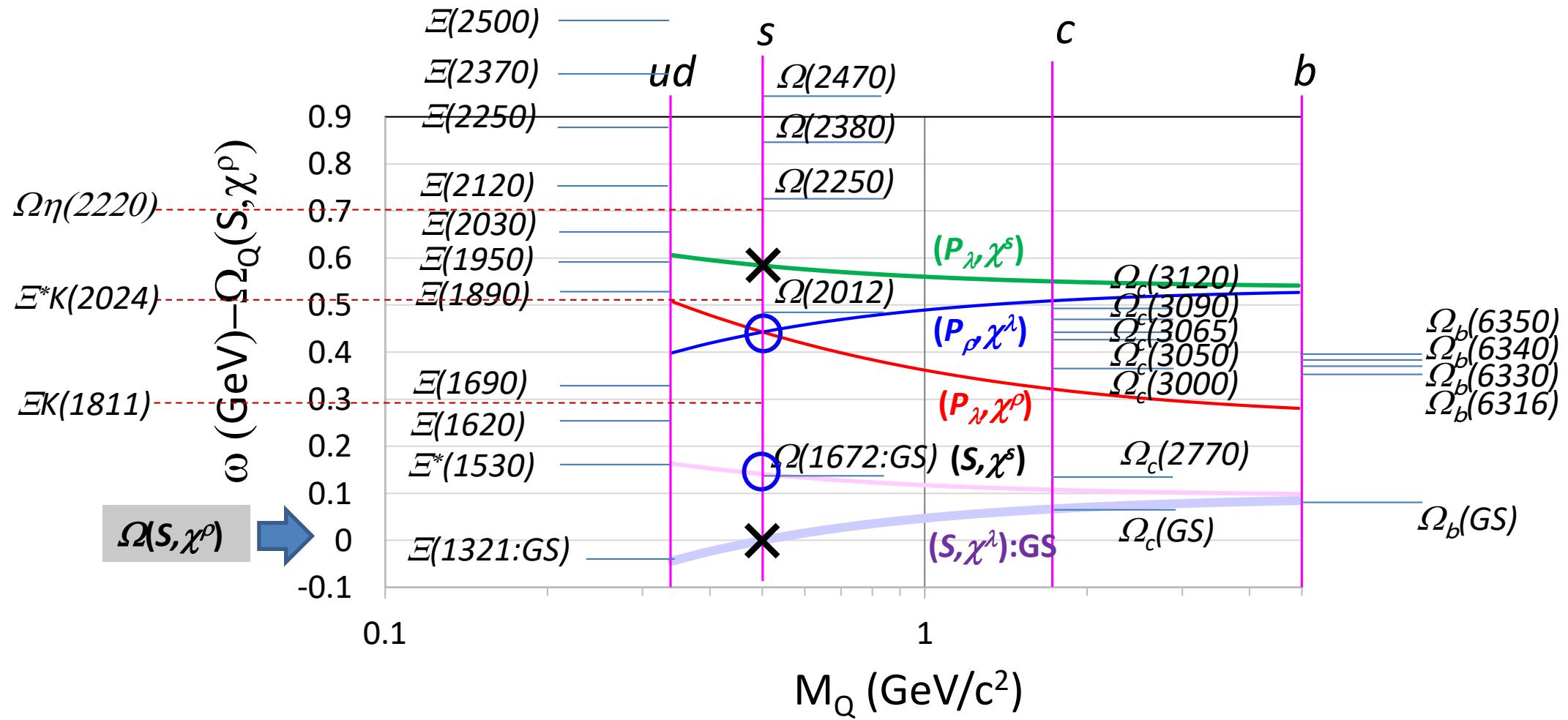


Qss

$\Omega(S,P)$ (ρ, λ -mode excitations w/ V_{ss}) when $m_q = m_s$



Omega Baryon



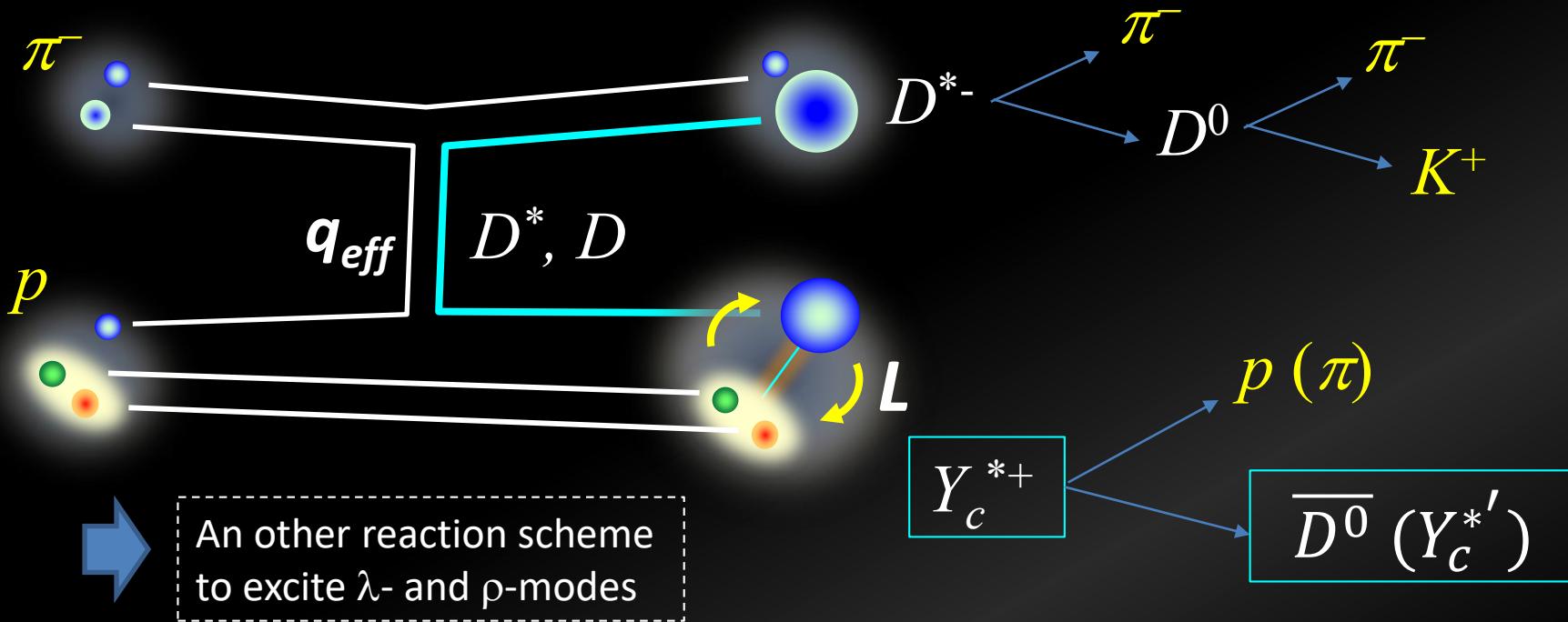
Measured Ω (PDG)

		JP	rating	Width [MeV]	$\rightarrow \Xi K$ (1)	$\rightarrow \Xi^* K$ (2)	$\rightarrow \Xi K^*$ (3)	$\rightarrow \Xi K\pi$ (4)	$\rightarrow \Omega\pi\pi$ (5)	
Threshold	$\Omega(2470)$??	2*	72+-33					seen	LASS (113MK-, 11GeV/c) (290+-90)/(5) nb
$\Xi^0 K^*- 2109$	$\Omega(2380)$??	2*	26+-23		<0.44 to (4)	0.5+-0.3 to (4)			Xi Beam
$\Xi^0 K^- 2024$	$\Omega(2250)$??	3*	55+-18		0.7+-0.2 to (4)		Seen		Xi Beam LASS (113MK-, 11GeV/c) (630+-180)/(2) nb
$\Xi^0 K-\pi^0 1956$	$\Omega(2012)$?-	3*	$6.4^{+2.5}_{-2.0}$ +-1.6	$1.2^{+0.3}_{-0.3}$ $(=\Xi^0/\Xi^-)$	<0.119 (1)				$\rightarrow \Xi^* K$ dominant if $\Xi^* K$ mol?
$\Omega\pi^0\pi^0 1942$	$\Omega(1672)$	3/2+	4*	-						
$\Xi^0 K^- 1811$ ($\Omega\pi^0 1807$)										

- ✓ Most of spins/parities/decay branches have yet to be determined.
- ✓ What the production $\Xi^* \rightarrow \Omega^* K$ and $\Omega^* \pi$ decay modes tell us Ω^* 's internal structure.

Charmed Baryon Spectroscopy

Charmed Baryon Spectroscopy Using Missing Mass Techniques



- ✓ Production and Decay reflect [qq] correlation in Excited Y_c^*
- ✓ C.S. DOES NOT go down at higher L when $q_{eff} > 1 \text{ GeV}/c.$

S.H. Kim, A. Hosaka, H.C. Kim, and HN, PTEP, (2014) 103D01,

S.H. Kim, A. Hosaka, H.C. Kim, and HN, Phys.Rev. D92 (2015) 094021

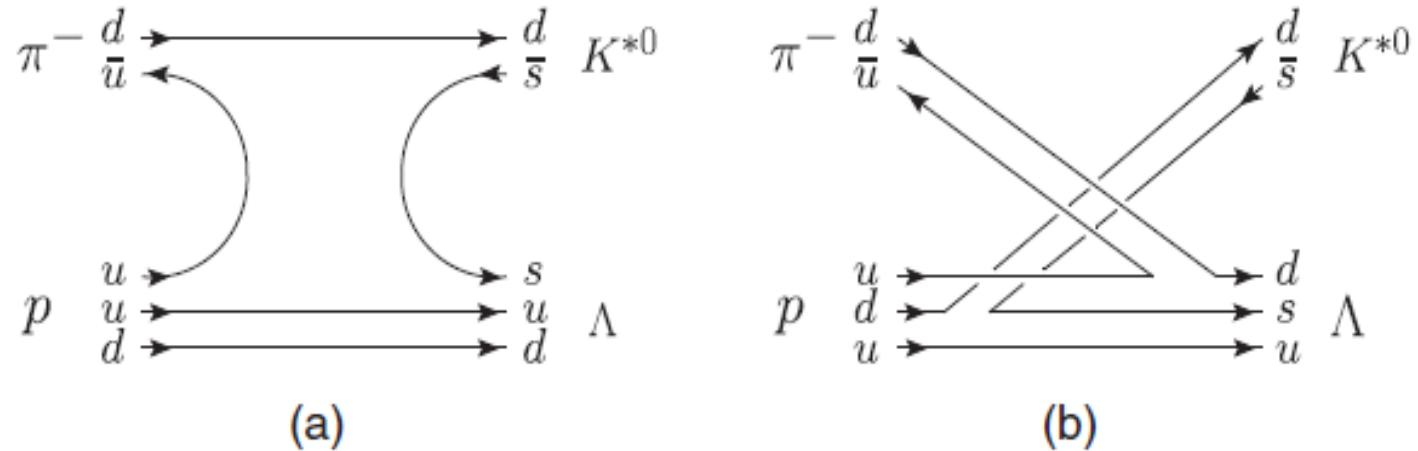
Revisit the Regge Theory

- shows the typical s-dependence of binary reaction cross sections at the large s region;

$$\frac{d\sigma}{dt} = \frac{1}{64\pi s(p_\pi^{cm})^2} |\langle f | T | i \rangle|^2 \quad \langle f | T | i \rangle = g_1 g_2 \Gamma(-\alpha(t)) (s/s_0)^{\alpha(t)}$$

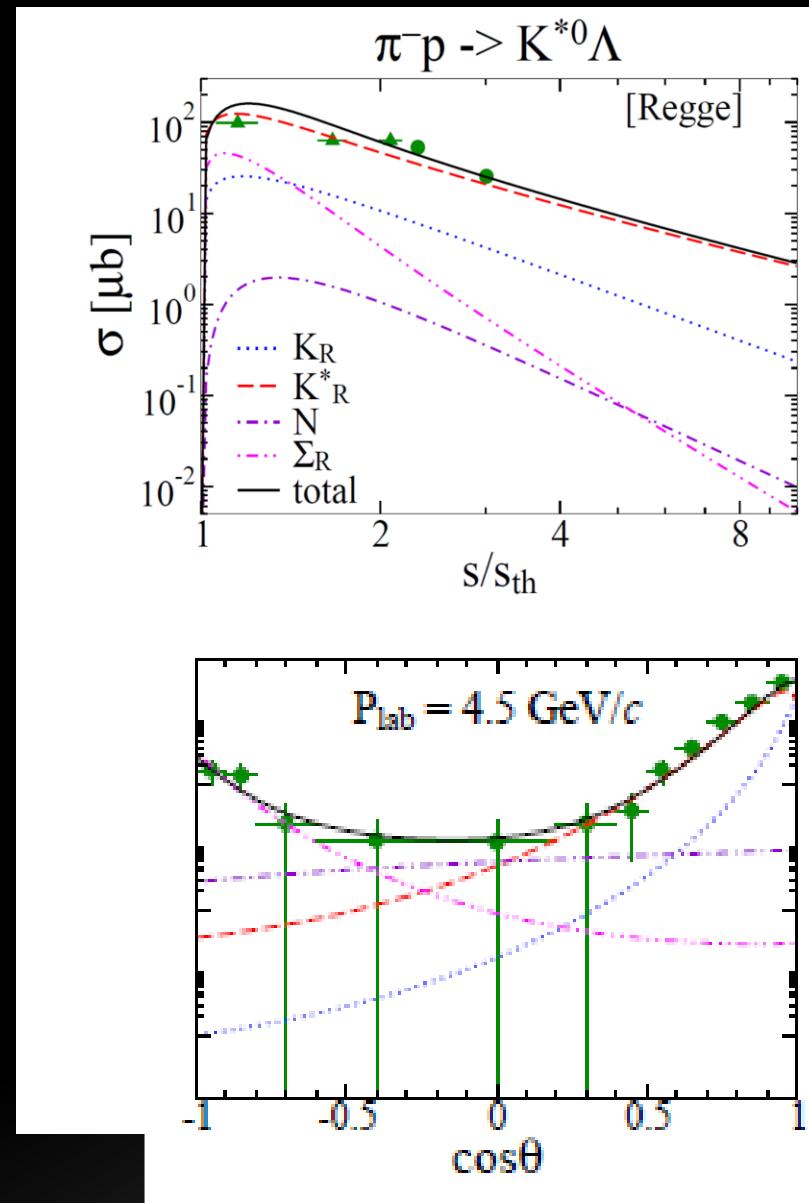
- Regge trajectory: $\alpha(t) = \alpha(0) - \gamma[\sqrt{T} - \sqrt{t-T}]$
- scale parameter s_0 :
s at the threshold energy of the reaction $AB \rightarrow CD$
(*In Kaidalov's Model: $s_0^{2(\alpha_D(0)-1)} = s_{CD}^{\alpha_P(0)-1} * s_{CD}^{\alpha_J/\psi(0)-1}$
 $s_{AB} = (\sum m_i)_A * (\sum m'_j)_B$, m_i :transversal masses of the constituent quark)

Amplitude

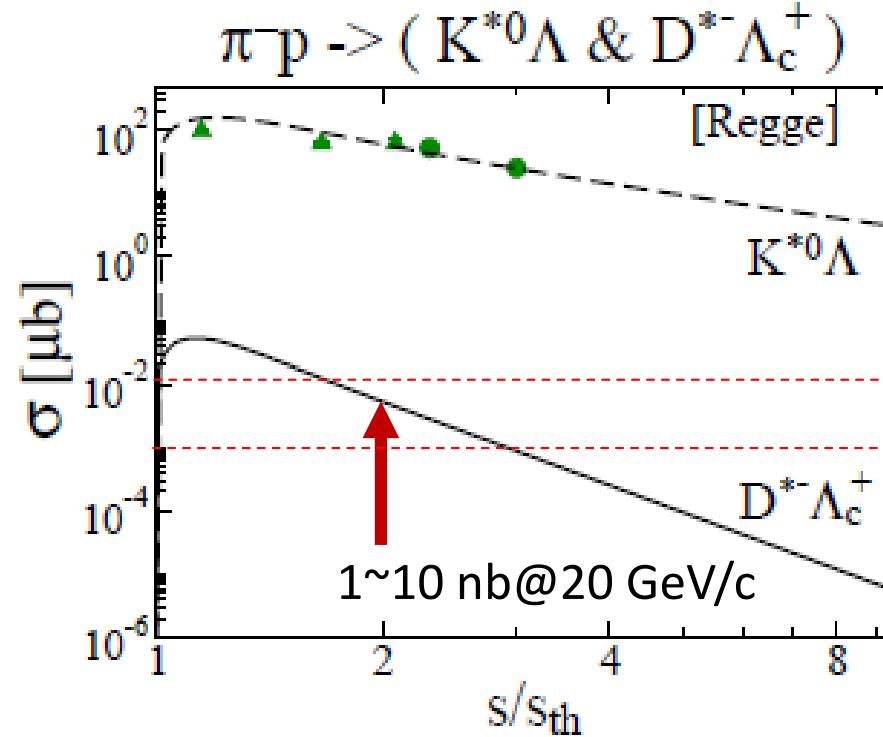


$$\begin{aligned}
 P_K^R(s, t) &= \left(\frac{1}{e^{-i\pi\alpha_K(t)}} \right) \left(\frac{s}{s_K} \right)^{\alpha_K(t)} \Gamma[-\alpha_K(t)] \alpha'_K, \\
 P_{K^*}^R(s, t) &= \left(\frac{1}{e^{-i\pi\alpha_{K^*}(t)}} \right) \left(\frac{s}{s_{K^*}} \right)^{\alpha_{K^*}(t)-1} \Gamma[1 - \alpha_{K^*}(t)] \alpha'_{K^*}, \\
 P_\Sigma^R(s, u) &= \left(\frac{1}{e^{-i\pi\alpha_\Sigma(u)}} \right) \left(\frac{s}{s_\Sigma} \right)^{\alpha_\Sigma(u)-\frac{1}{2}} \Gamma\left[\frac{1}{2} - \alpha_\Sigma(u)\right] \alpha'_\Sigma,
 \end{aligned}$$

Production Cross Section

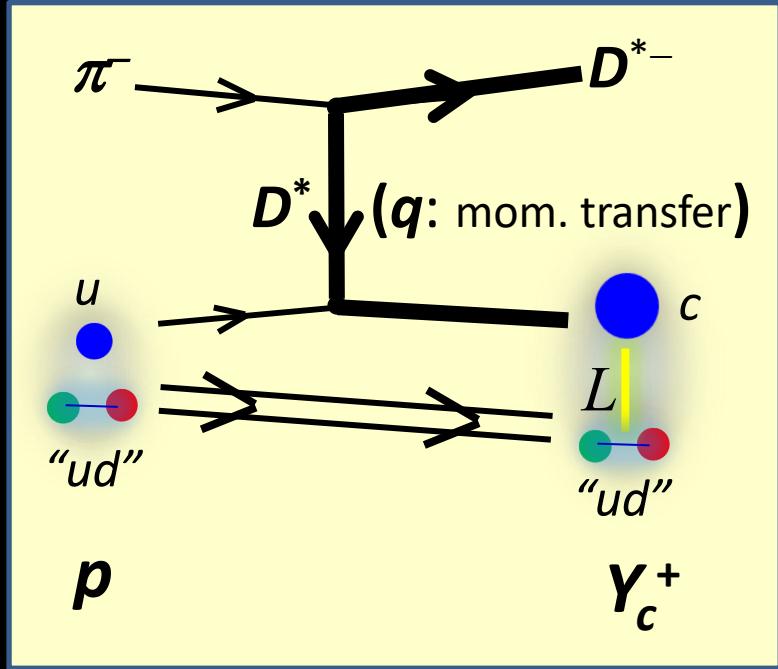


S.H. Kim, A. Hosaka, H.C. Kim, and HN
PRD92, 094021(2015)



Production Rate:

$$R \sim \left\langle \varphi_f \left| \sqrt{2} \sigma_- \exp(i \vec{q}_{eff} \cdot \vec{r}) \right| \varphi_i \right\rangle$$



- t-channel D^* **Reggeon** at a forward angle

S. H. Kim, et al.,
PTEP, 2014, 103D01(2014)

1. Momentum transfer (q_{eff})

$$I_L \sim (q_{eff}/A)^L \exp(-q_{eff}^2/2A^2)$$

$$q_{eff} \sim 1.4 \text{ GeV/c} \quad A \sim 0.4 \text{ GeV ([Baryon size])}^{-1}$$

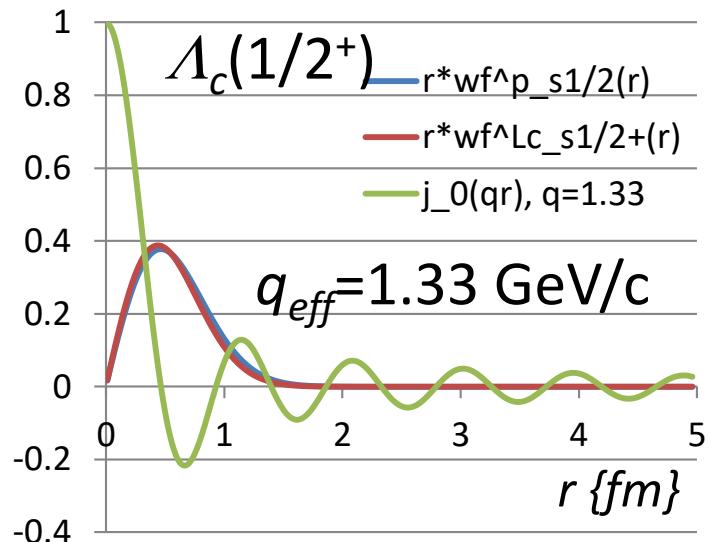
2. Population shared among HQ-spin multiplet

$$J_{BM} - s_{HQ} : J_{BM} + s_{HQ} = L : L + 1$$

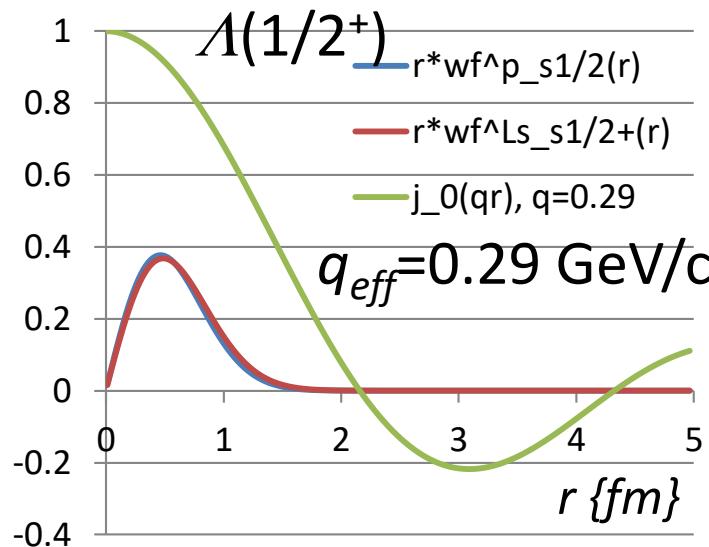
3. Spectroscopic Factor ("ud" configuration)

$$\gamma = 1/2 \text{ for } \Lambda_c \text{'s}, \quad = 1/6 \text{ for } \Sigma_c \text{'s}$$

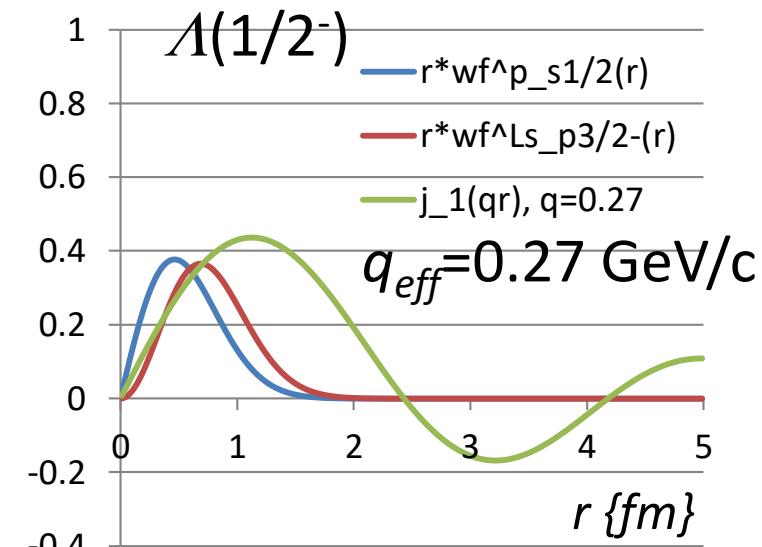
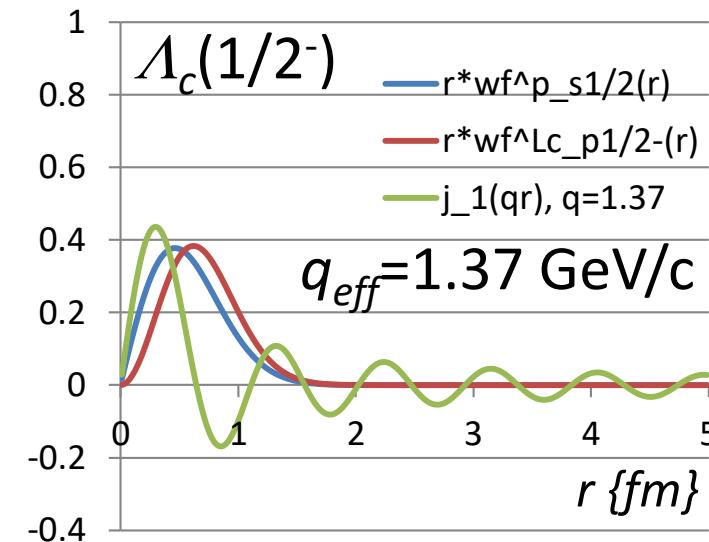
charm



strange



$L=0$



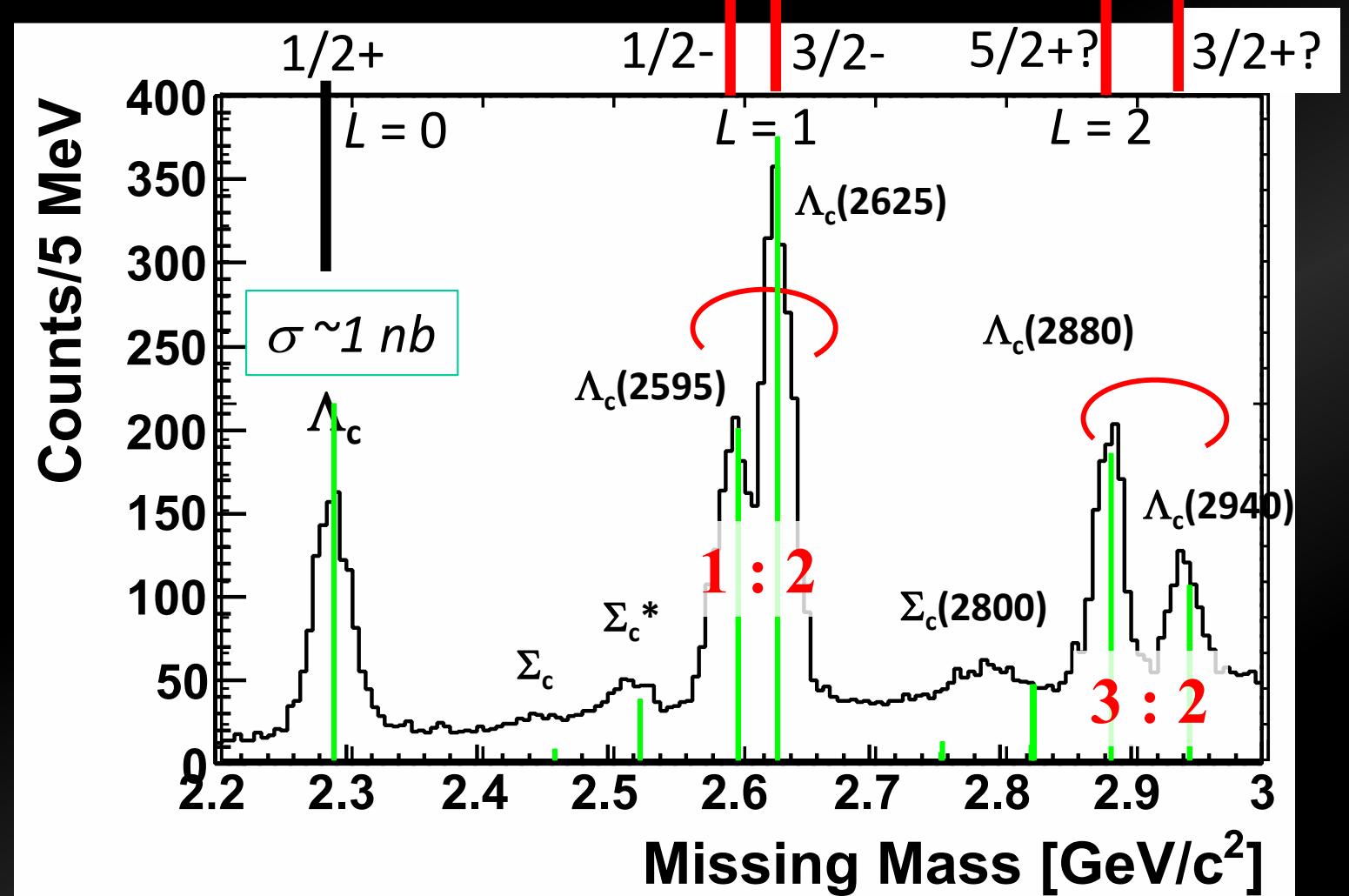
$L=1$

Missing Mass Spectrum (Sim.)

- $\sim 1000 Y_c^*/1 \text{ nb}/100 \text{ days}$
- Sensitivity: $\sigma \sim 0.1 \text{ nb}$ for Y_c^* w/ $\Gamma = 100 \text{ MeV}$

λ mode

$\lambda\lambda$ mode?



Missing Mass Spectrum (Sim.)

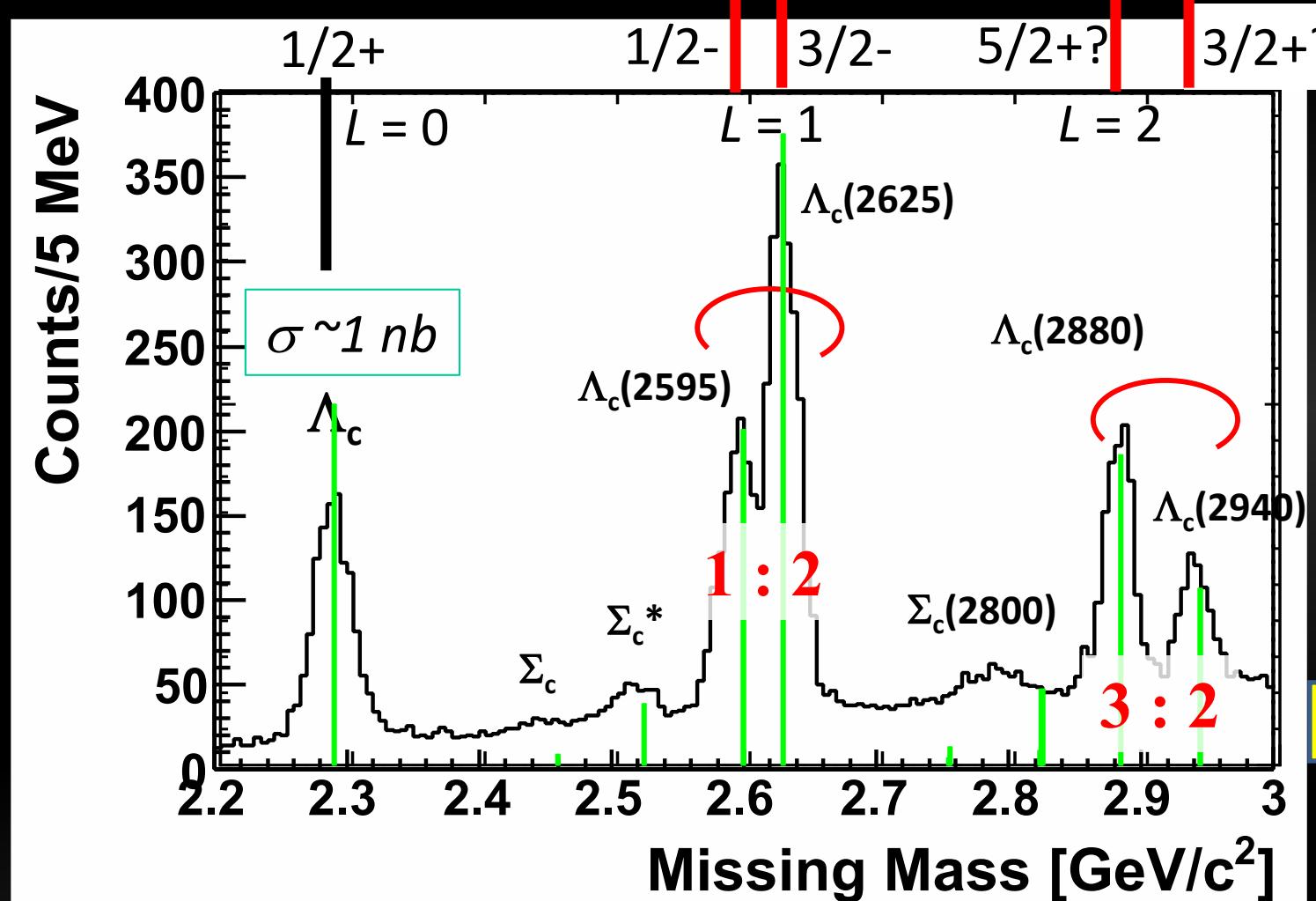
- $\sim 1000 Y_c^*/1 \text{ nb}/100 \text{ days}$
- Sensitivity: $\sigma \sim 0.1 \text{ nb}$ for Y_c^* w/ $\Gamma = 100 \text{ MeV}$

λ mode

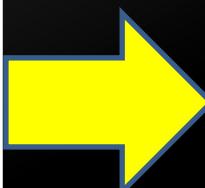
$\lambda\lambda$ mode?

LS partner
(HQS doublet)

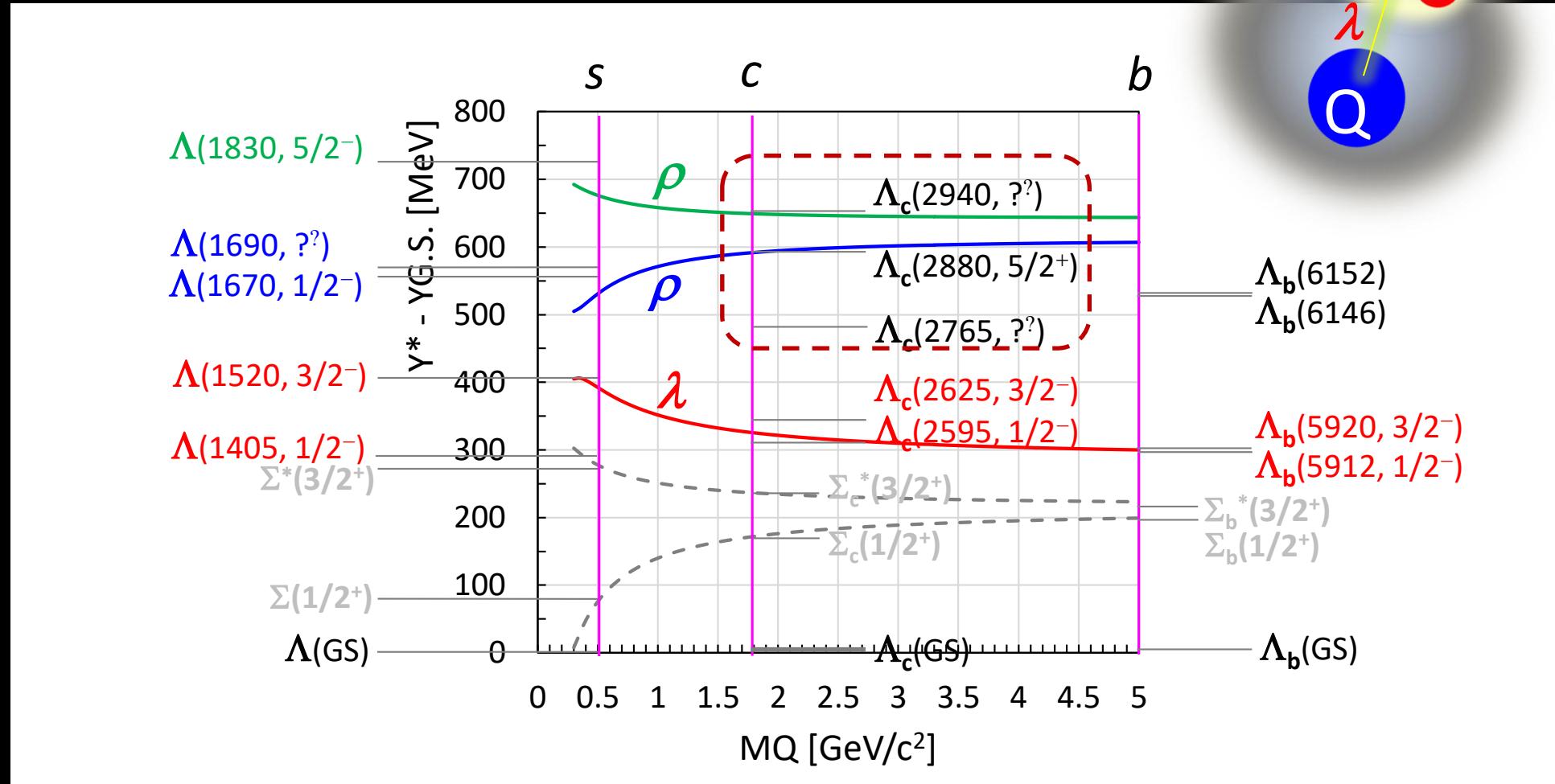
LS partner?
(HQS doublet?)



Are
There
More
Excited
States?



Lambda Baryons



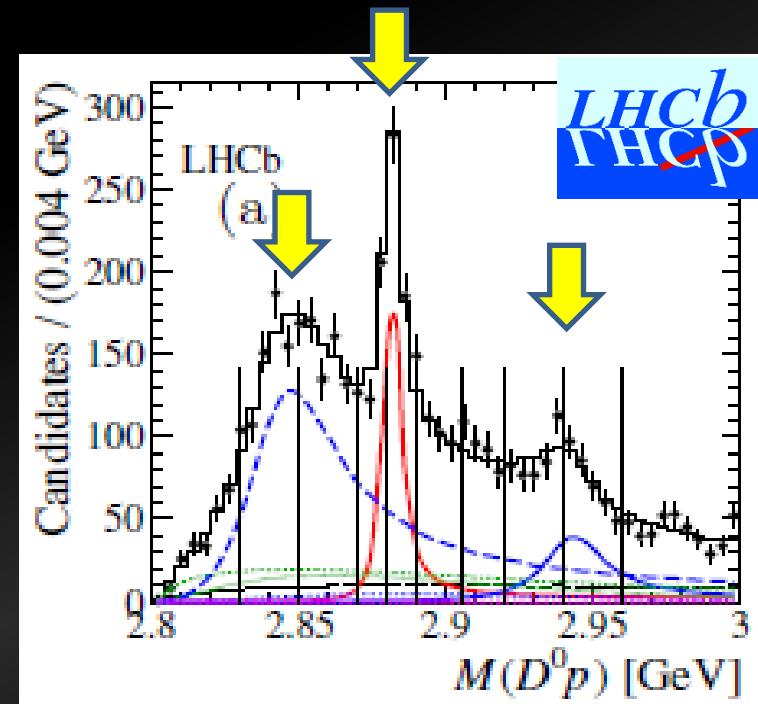
non-rel. QM: $H = H_0 + V_{conf} + V_{SS} + V_{LS} + V_T$
 ρ - λ mixing (cal. By T. Yoshida)

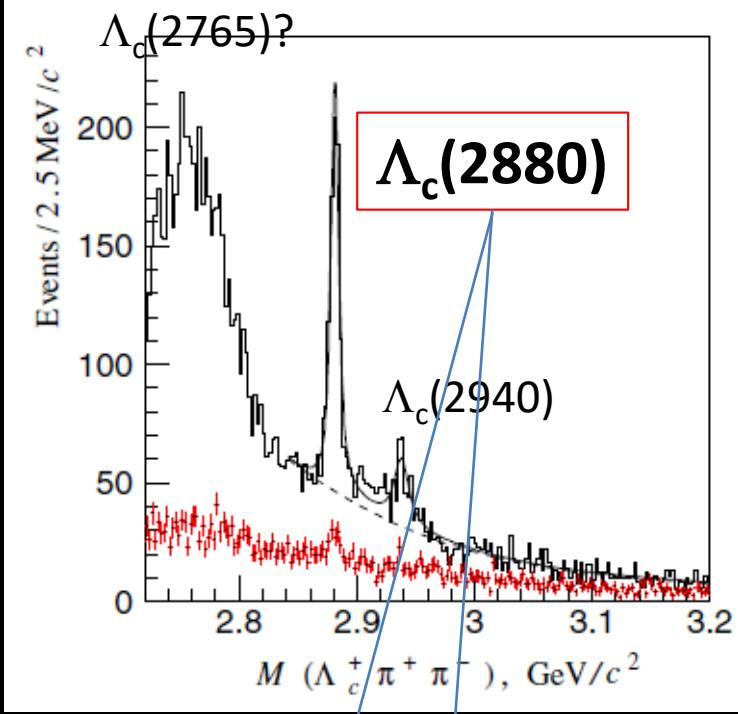
T. Yoshida et al.,
Phys. Rev. D92, 114029(2015)

LHCb data in Λ_c^*

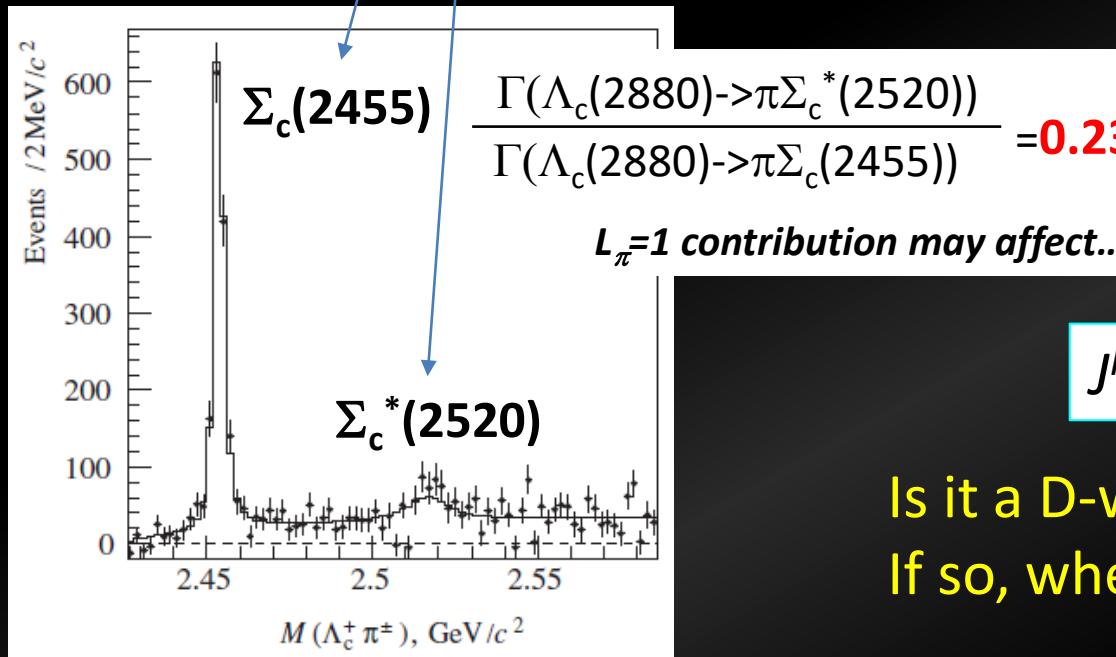
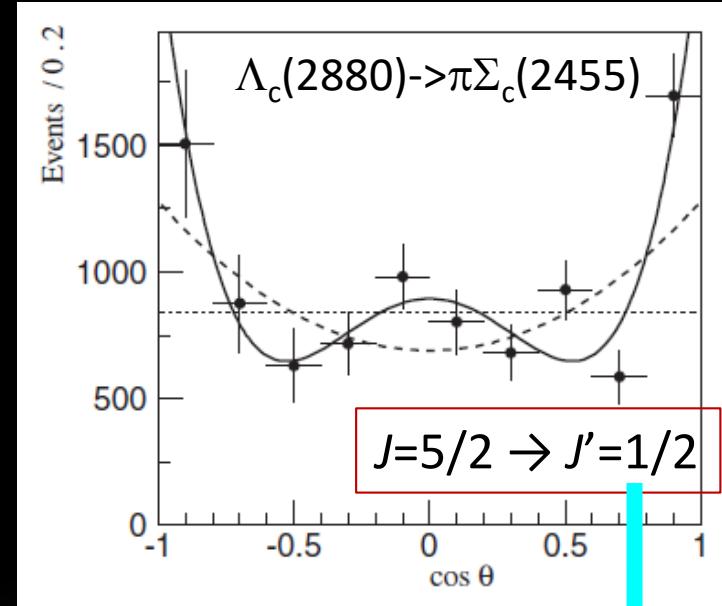
J. High Energ. Phys. (2017) 2017

- $D^0 p$ invariant mass in $\Lambda_b \rightarrow D^0 p \pi^-$
 - $\Lambda_c(2940)$: known
 - likely $3/2^-$, (acceptable $1/2, 7/2$)
 - $\Lambda_c(2880)$: known
 - $5/2^+$ confirmed
 - $\Lambda_c(2860)$: new
 - likely $3/2^+$, D-wave ($L=2$) resonance?
- Questions arise;
 - Is $\Lambda_c(2940)$ an $L=3$ state (λ mode)?
 - Are $\Lambda_c(2880)$ and $\Lambda_c(2860)$ LS partners of $L=2$ (λ modes)?
- Production rates in $p(\pi^-, D^{*-})\Lambda_c^*$ will give answer.





Lc(2880)Belle, PRL98, 262001('07)



Is it a D-wave Lambda-c Baryon?
If so, where is a spin partner?

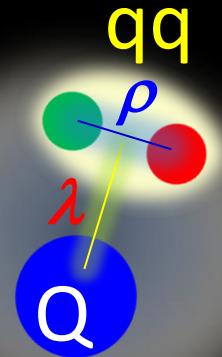
$\Lambda(2880)$ likely to be $\lambda\rho$ mode?

H. Nagahiro et al., PRD95 (2017) no.1, 014023

- P-wave transition seems to be suppressed in

$$\Lambda_c(2880)^{\frac{5}{2}+} \rightarrow \Sigma_c^*(2520)^{\frac{3}{2}+} + \pi(0^-).$$

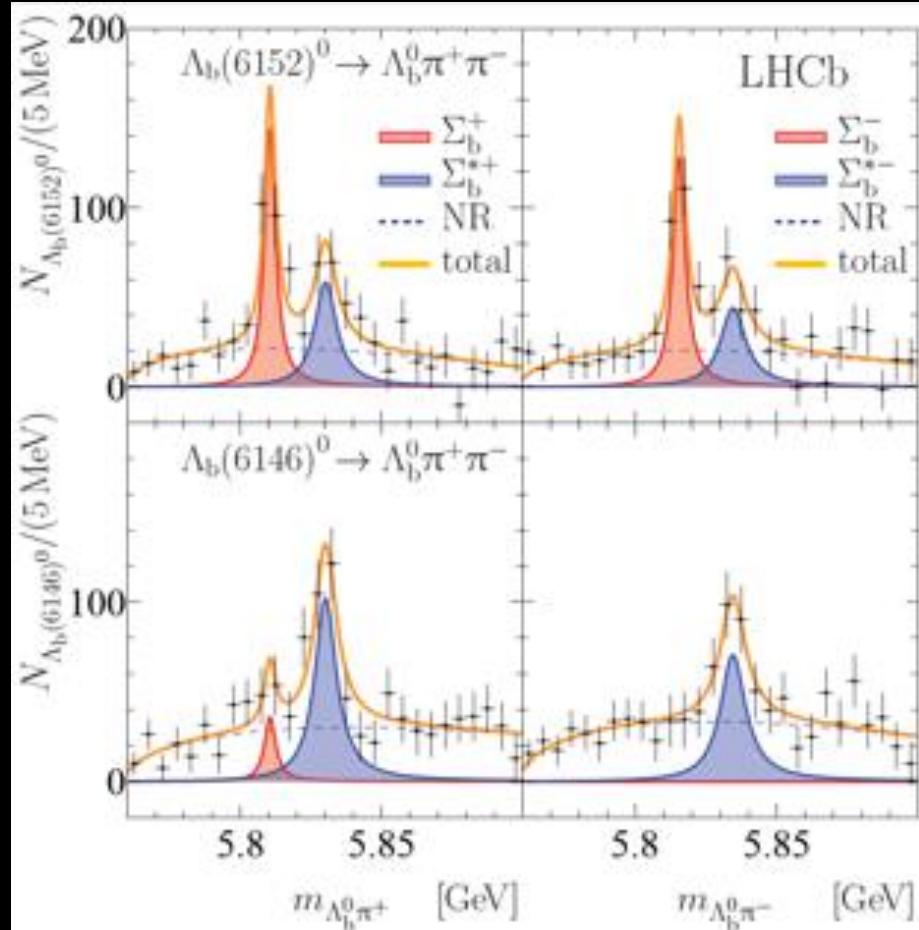
- It would be forbidden **only in the case of** $J_{BM}^P = 3^+$: “5/2-” state have large widths.
- $\Lambda_c(2880)^{\frac{5}{2}+}$ is likely to be a $\lambda\rho$ mode ($\lambda=1$, $\rho=1$) state.



$\Lambda_c(2880) 5/2+$	$\lambda\lambda$	$\lambda\rho$	$\rho\rho$	$\Sigma_c^*(2520) 3/2+$
color	Asymm.			
Isospin	Asymm. ($I=0$)			
Diquark spin	Asymm. 0	Symm. 1	Asymm. 0	Symm. 1
Diquark orbit	Symm. 0	Asymm. 1	Symm. 2	Symm. 0
Lambda orbit	2	1	0	0
J_{BM}^P	2+	1+, 2+, 3+	2+	1+

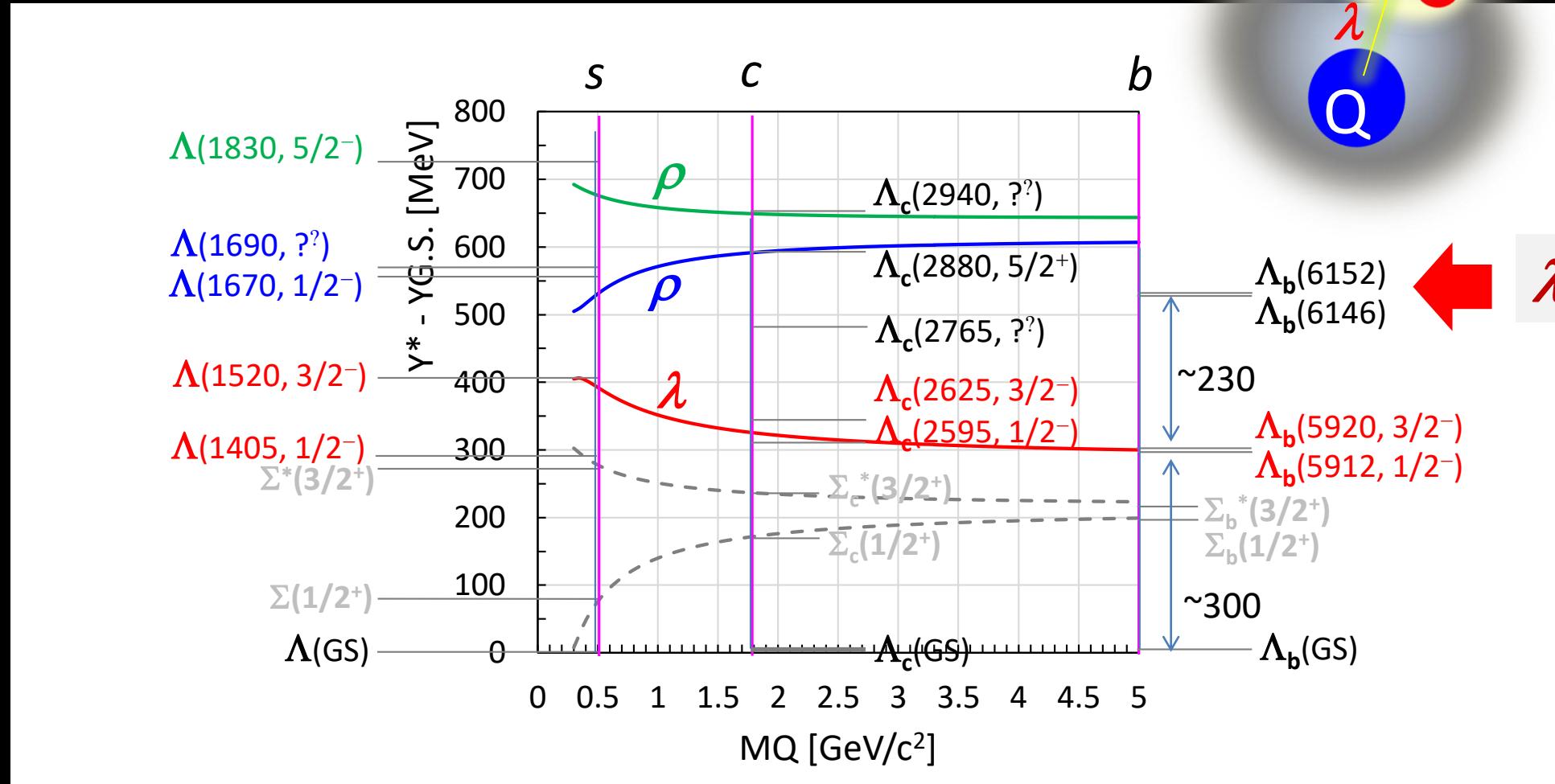
To be confirmed by measuring the prod. ratio

Lambda_b(6146/6152)



- A new doublet Λ_b^* states decaying into $\Lambda_b \pi^+ \pi^-$ have been observed.
 - $M_{\Lambda_b(6146)} = 6146.17 \pm 0.33 \pm 0.22 \pm 0.16 \text{ MeV}$
 - $M_{\Lambda_b(6152)} = 6152.51 \pm 0.26 \pm 0.22 \pm 0.16 \text{ MeV}$
 - $\Gamma_{\Lambda_b(6146)} = 2.9 \pm 1.3 \pm 0.3 \text{ MeV}$
 - $\Gamma_{\Lambda_b(6146)} = 2.1 \pm 0.8 \pm 0.3 \text{ MeV}$
- They are likely to be λ -mode with $L=2\dots$
- $\Lambda_b(6146)$ dominantly decays to Σ_b ?
 - Similar to the case of $\Lambda_c(2880, 5/2^+)$

Lambda Baryons



non-rel. QM: $H = H_0 + V_{conf} + V_{SS} + V_{LS} + V_T$
 ρ - λ mixing (cal. By T. Yoshida)

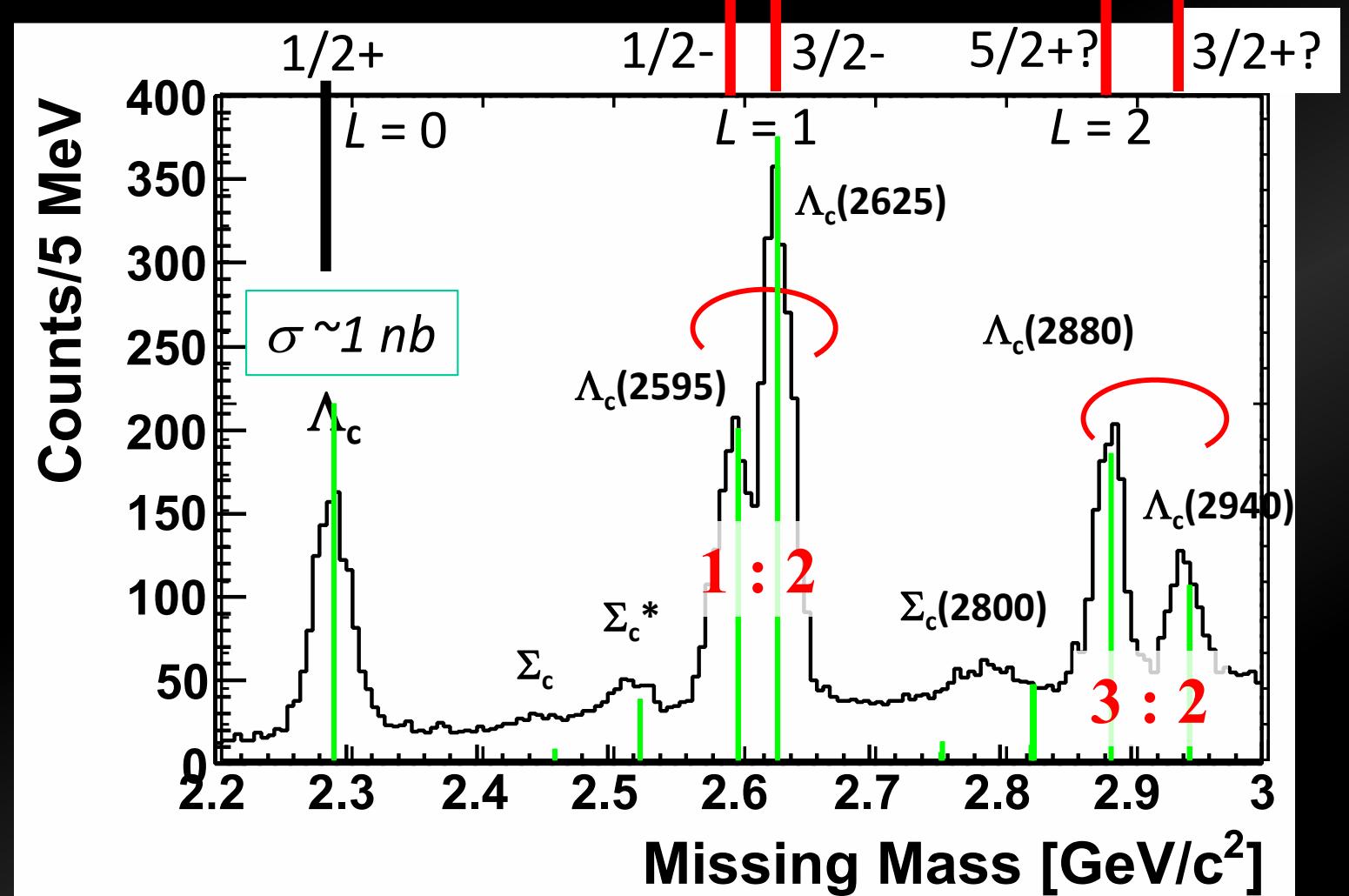
T. Yoshida et al.,
Phys. Rev. D92, 114029(2015)

Missing Mass Spectrum (Sim.)

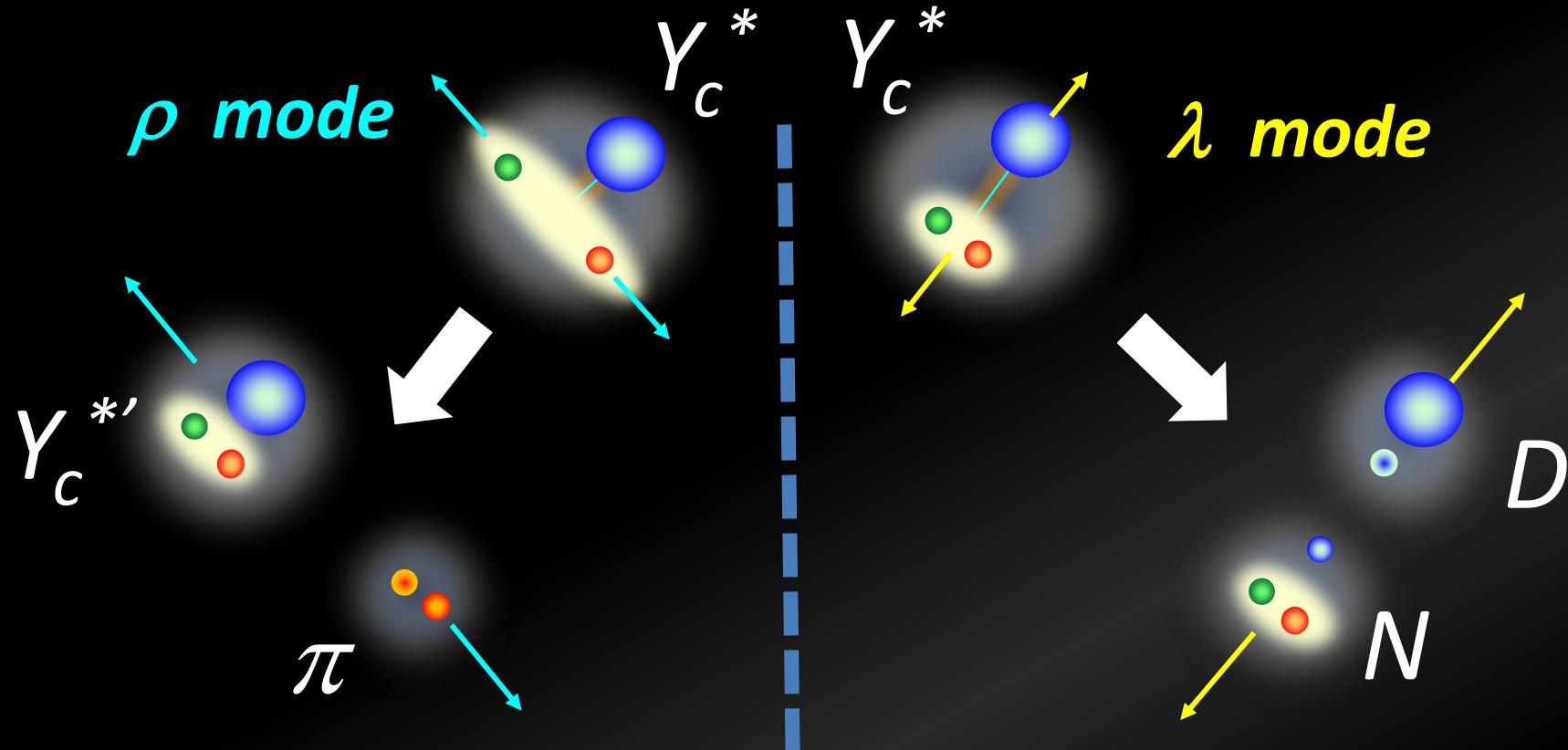
- $\sim 1000 Y_c^*/1 \text{ nb}/100 \text{ days}$
- Sensitivity: $\sigma \sim 0.1 \text{ nb}$ for Y_c^* w/ $\Gamma = 100 \text{ MeV}$

λ mode

$\lambda\lambda$ mode?



Y_c^* Decay Branching Ratio

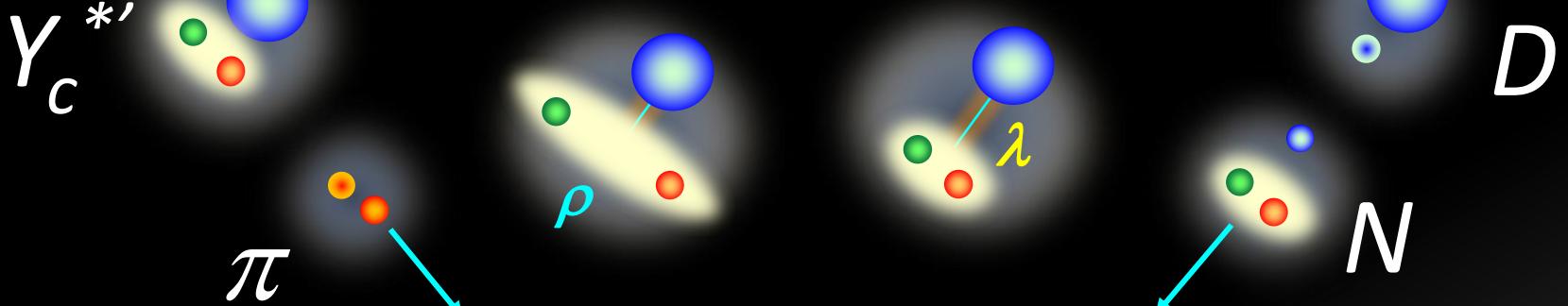


$$\Gamma(Y\pi) > \Gamma(DN)$$

$$\Gamma(DN) > \Gamma(Y\pi)$$

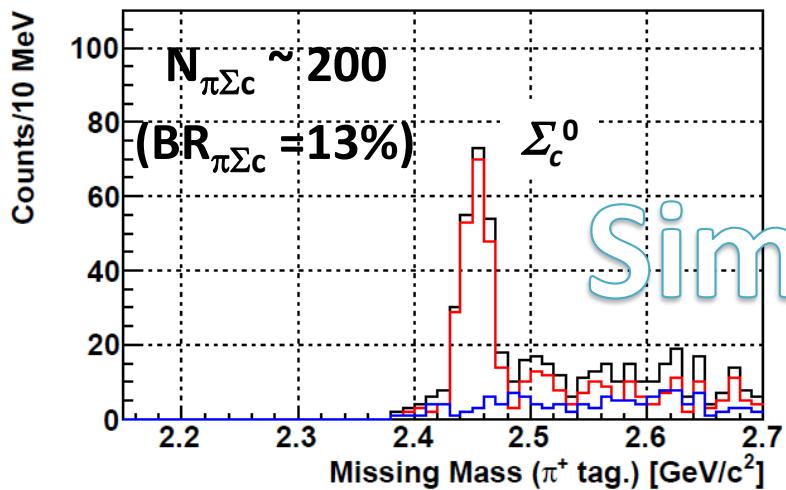
Λ_c^* Decays

$\Lambda_c^{*''}$

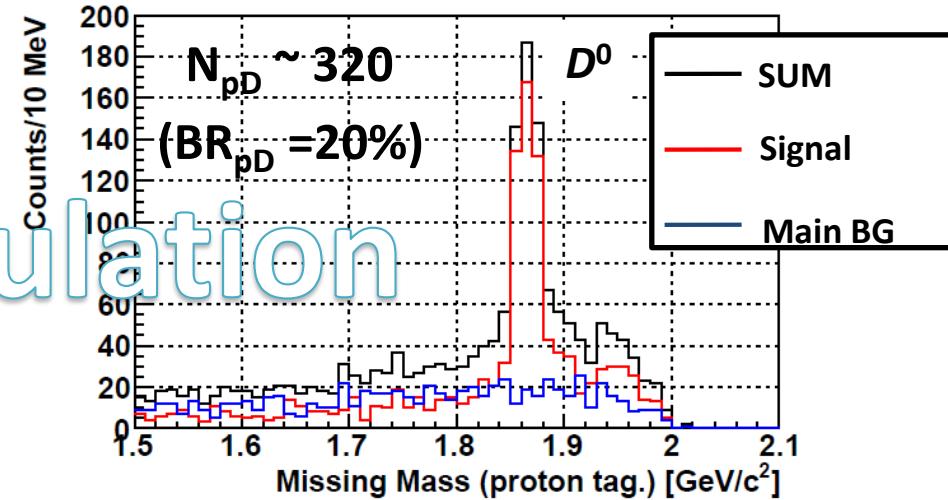


$$\Lambda_c(2940) \rightarrow \Sigma_c^0 \pi^+$$

with $\Lambda_c^+ \pi^+ \pi^-$ selected



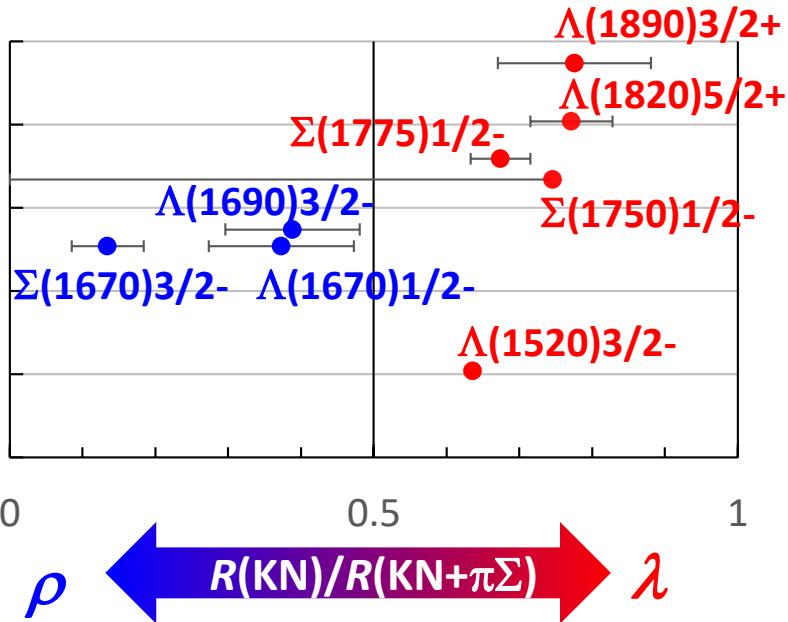
$$\Lambda_c(2940) \rightarrow p D^0$$



* Branching ratios: Diquark corr. affects $\Gamma(\Lambda_c^* \rightarrow pD)/\Gamma(\Lambda_c^* \rightarrow \Sigma_c \pi)$.

Hint in $R(NK)/R(\pi\Sigma)$

PDG Data



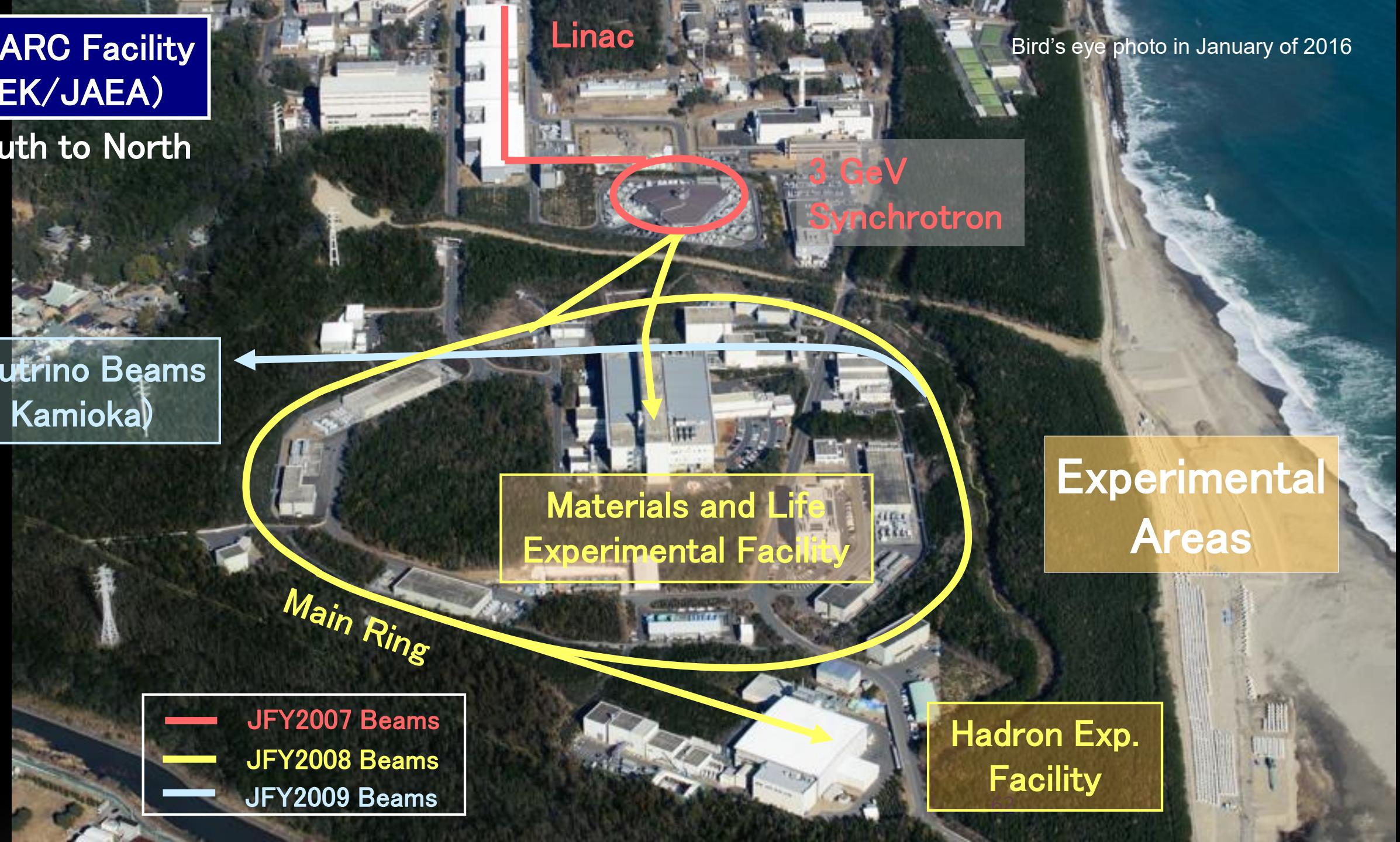
- Decay ratios in known hyperons SUGGEST the λ/ρ mode states
- λ/ρ mode ID by productions correlate w/ Decay Ratios
→ to be established

- Hyperon data indicate mode dependence
→ Errors should be improved.
- No data in charmed baryons

New Platform of Hadron Physics at J-PARC

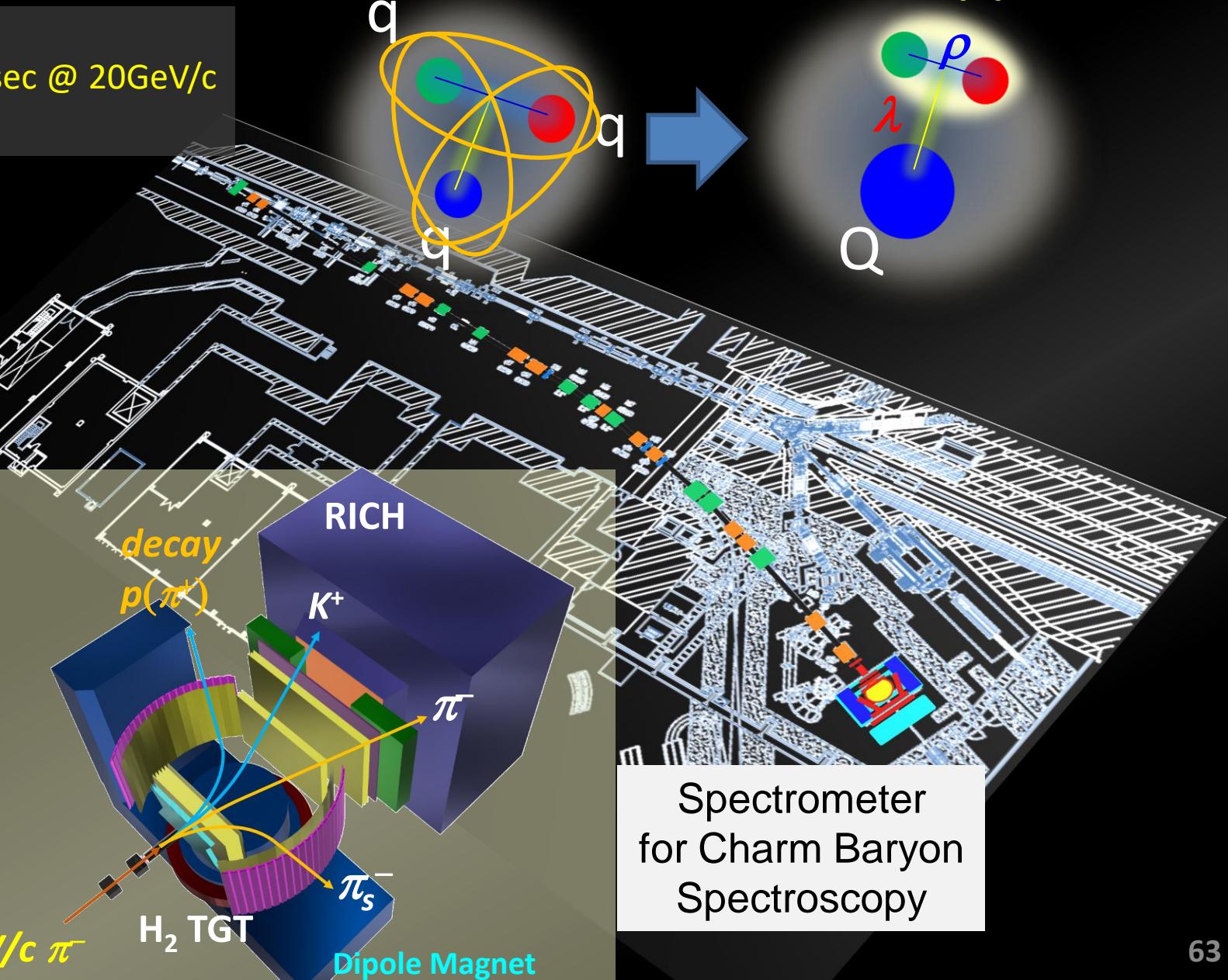
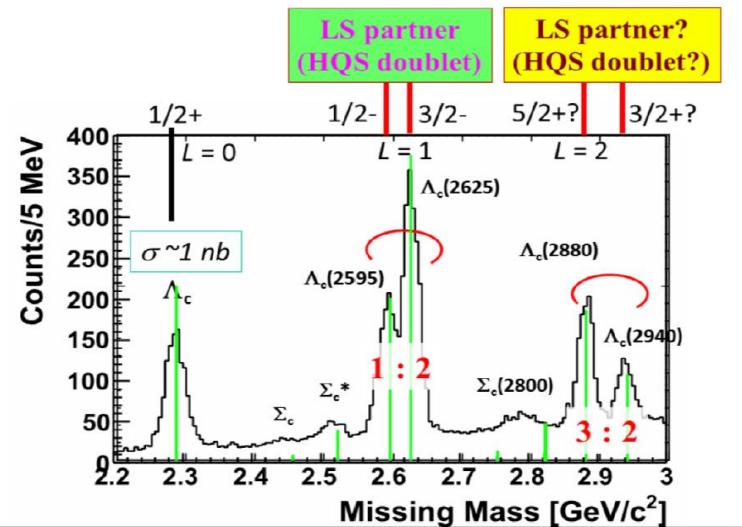
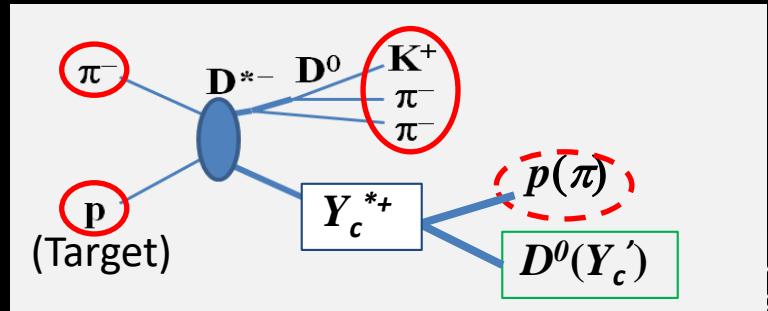
J-PARC Facility (KEK/JAEA)

South to North



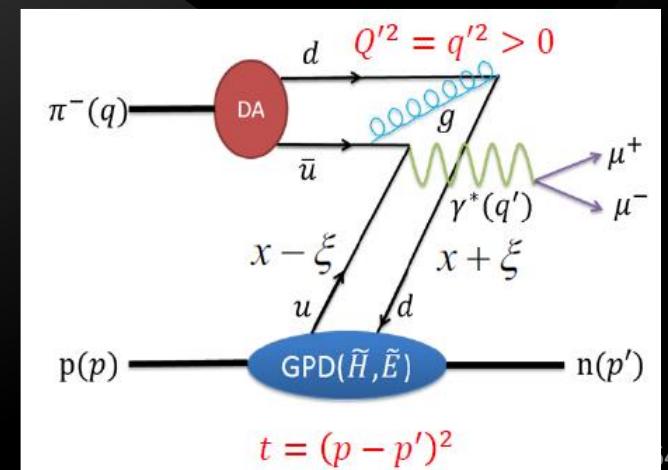
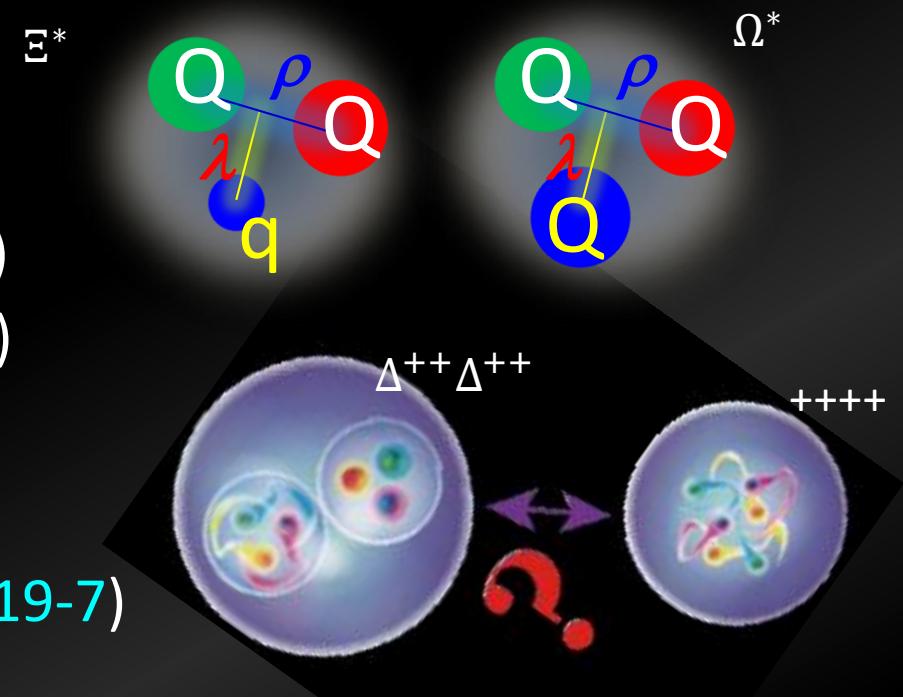
Hadron Physics at the High-p Beam Line

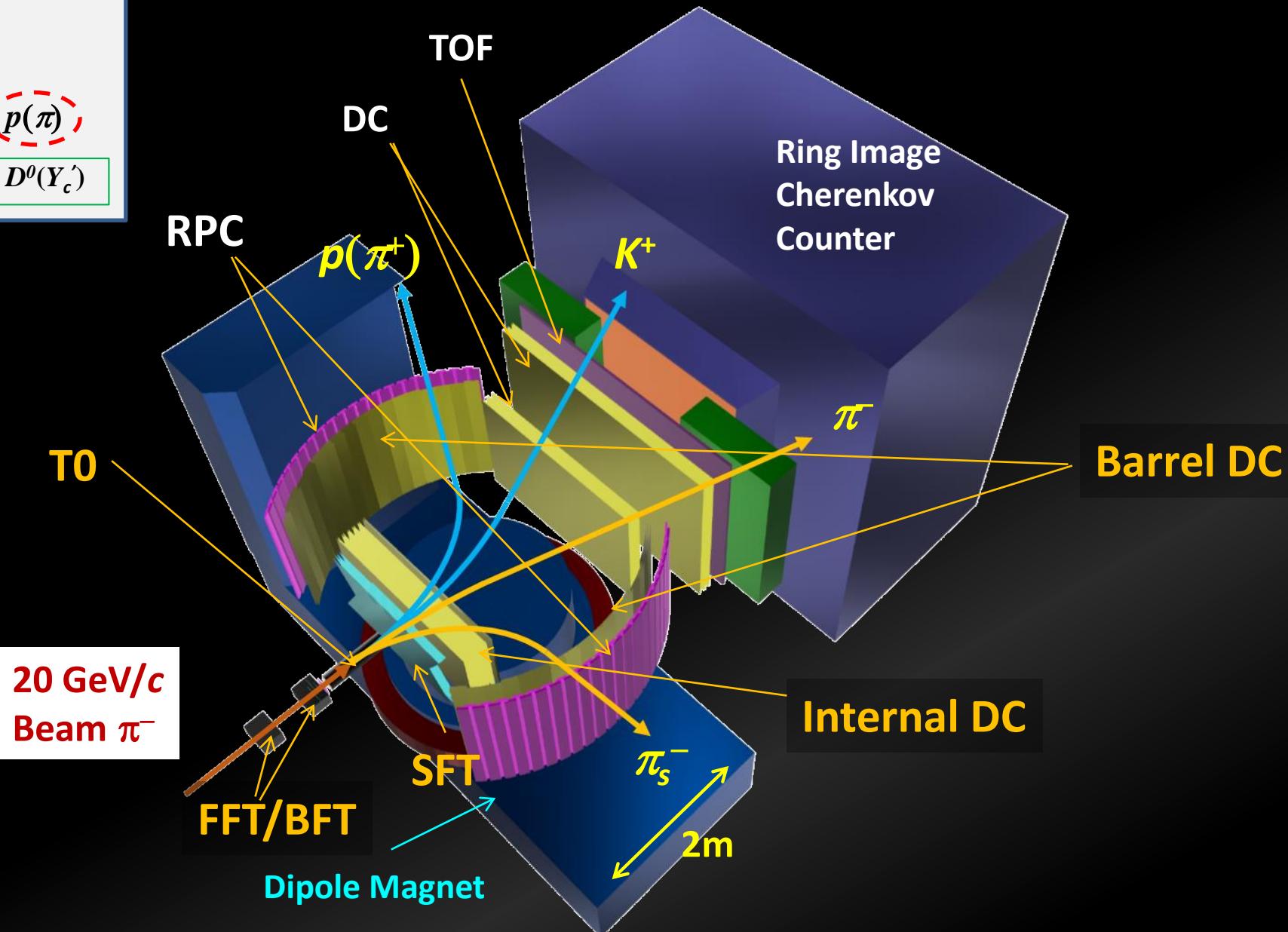
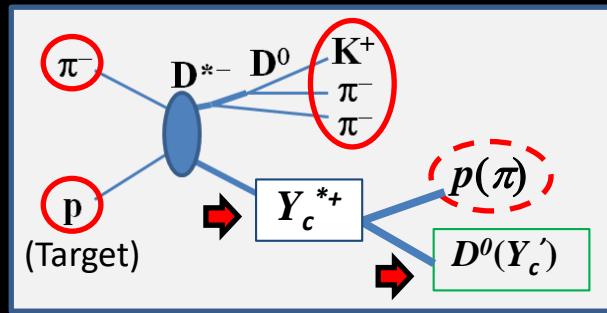
- $\pi 20$ beam line :
 - 1.0×10^7 pions/sec @ 20GeV/c
 - $\Delta p/p \sim 0.1\%$



Hadron Physics at High-p BL

- Baryon Spectroscopy
 - $p(\pi^-, D^{*-})\Lambda_c^*$ (E50)
 - $p(K^-, K^*)\Xi^*$, $p(K^-, K^+K^*)\Omega^*$ (Lol: KEK/J-PARC-PAC 2014-4)
 - Search for D_{30} Dibaryon State in $pp \rightarrow \pi^-\pi^-D_{30}$ (E79)
 - $p(\pi^-, K^*)\Lambda(1405)$ at large s, t (to be proposed)
- Hadron Tomography
 - Exclusive DY, $\pi^- p \rightarrow \mu^-\mu^+ n$ (Lol: KEK/J-PARC-PAC 2019-7)
- For Strangeness Nuclear Physics
 - Λp Scattering for the study of high-dense nuclear matter (Lol: KEK/J-PARC-PAC 2020-08)
- For Neutrino Physics
 - Hadron Production for neutrino beams



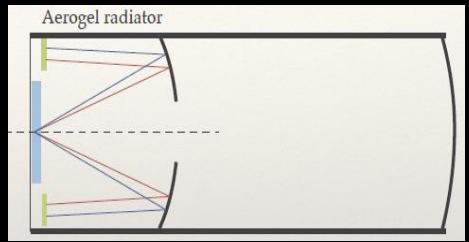
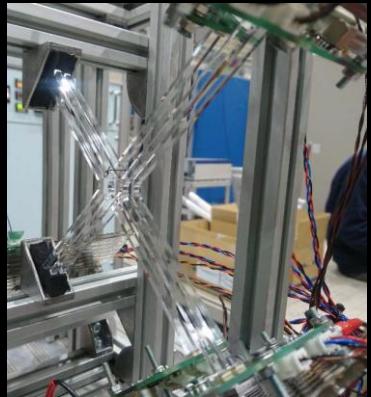
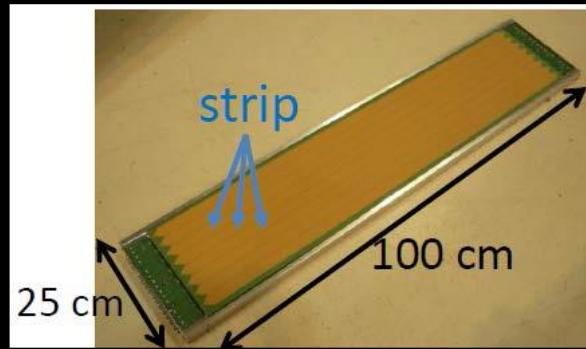


Acceptance: $\sim 60\%$ for D^* , $\sim 80\%$ for decay π^+

Resolution: $\Delta p/p \sim 0.2\%$ at ~ 5 GeV/c (Rigidity: ~ 2.1 Tm)

R&D Works

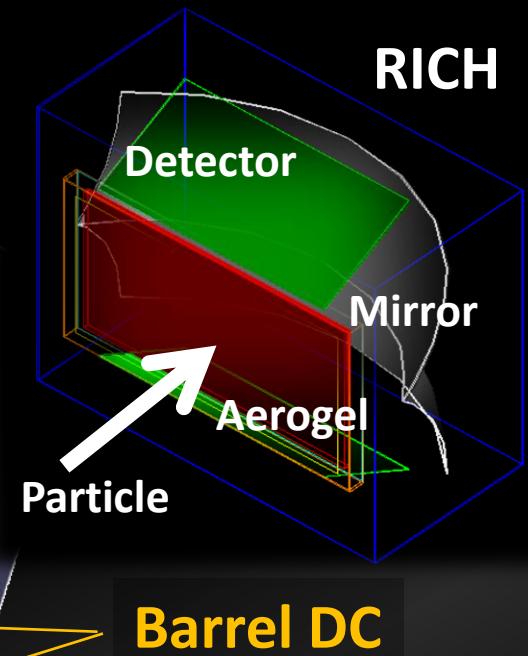
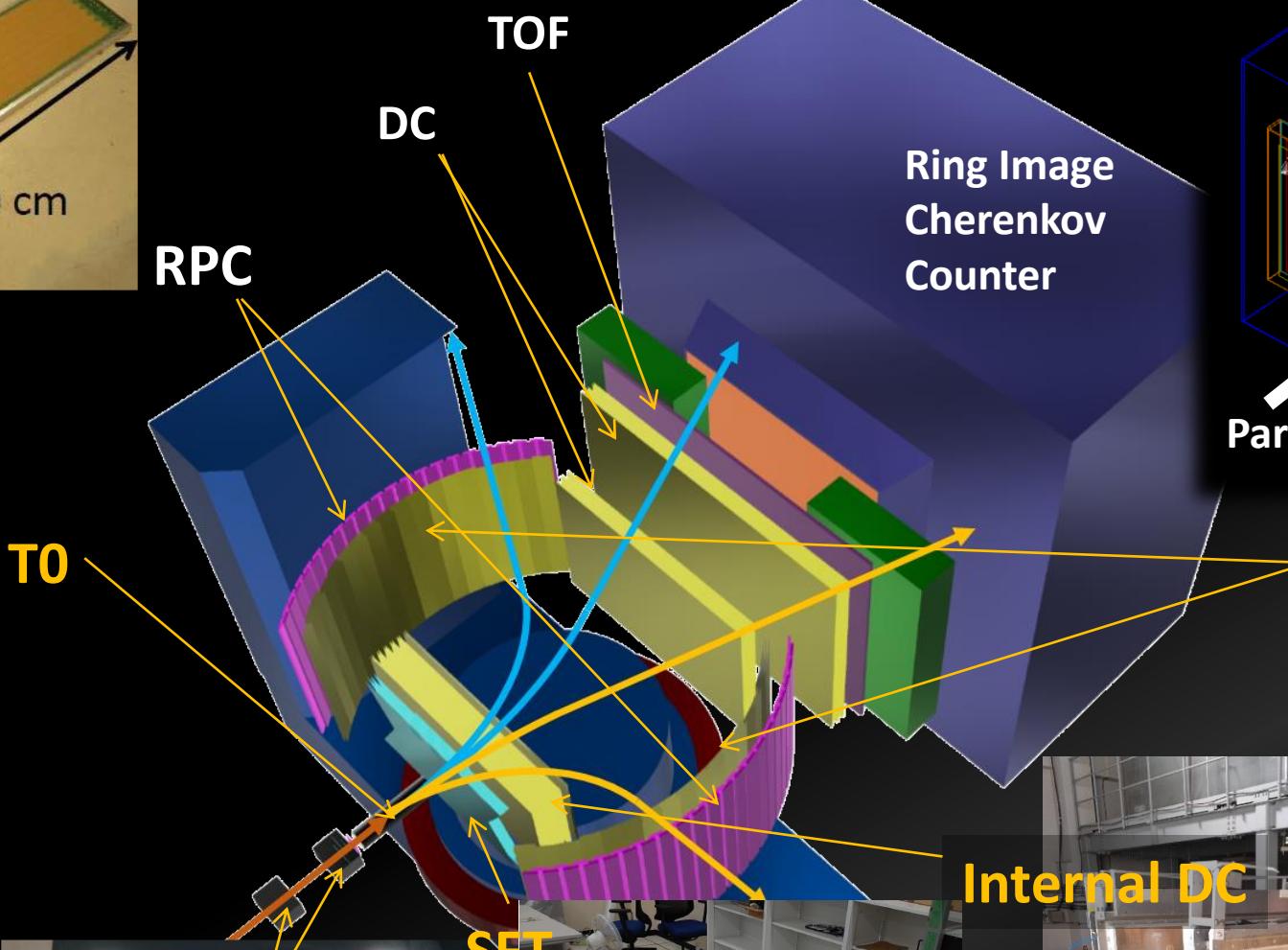
- Particle Identification (Osaka/Kyoto/Tohoku/RIKEN...)
 - Timing counters
 - T-Zero (Osaka): Cherenkov type ~50 ps
 - Resistive Plate Chamber (RCNP/ELPH/Taiwan/Kyoto/JAEA/Tsukuba): Large Size~60 ps
 - Ring Image Cherenkov Detector
 - BeamRICH/RICH (Tohoku/KEK/Osaka/Kyoto/RIKEN/...)
 - Muon ID (Academia Sinica)
- Trackers (Tohoku/RCNP/RIKEN...)
 - SciFi Tracker (Focal Plane/Beam/Scattered particle)
 - DC (Forward, Barrel)
- High-speed DAQ system (RCNP/Tohoku/Taiwan/KEK/RIKEN...)
 - PC cluster-based DAQ scheme
 - Flexible “trigger”: not only (π^- ,D*) but also (K^- , K^*),...



Beam RICH



Active area



Internal DC



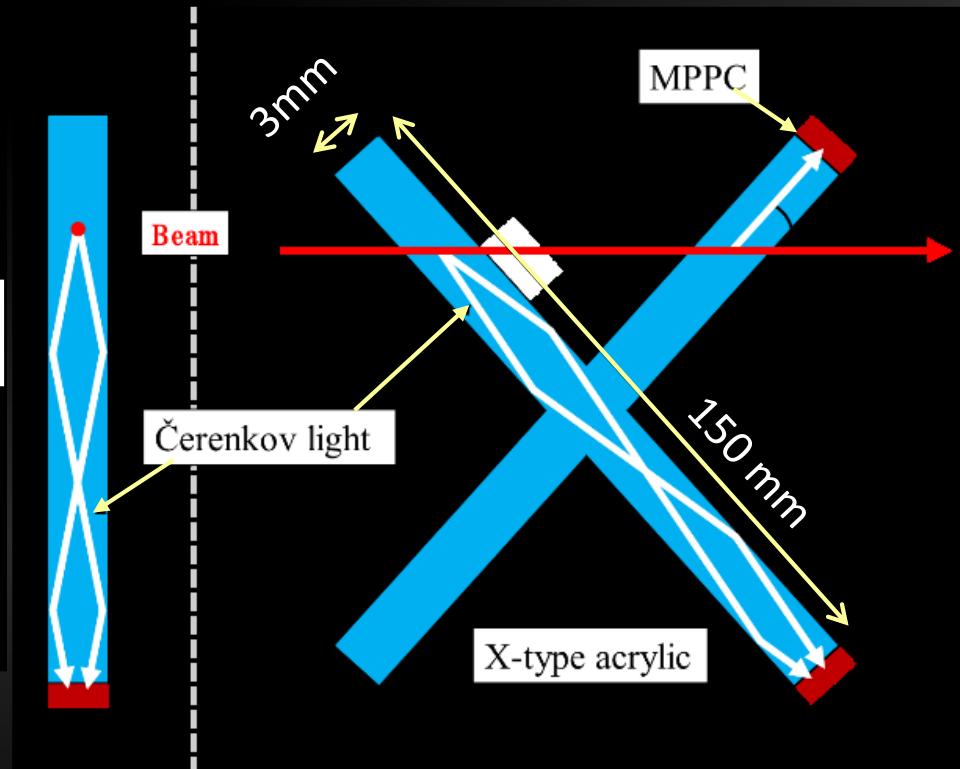
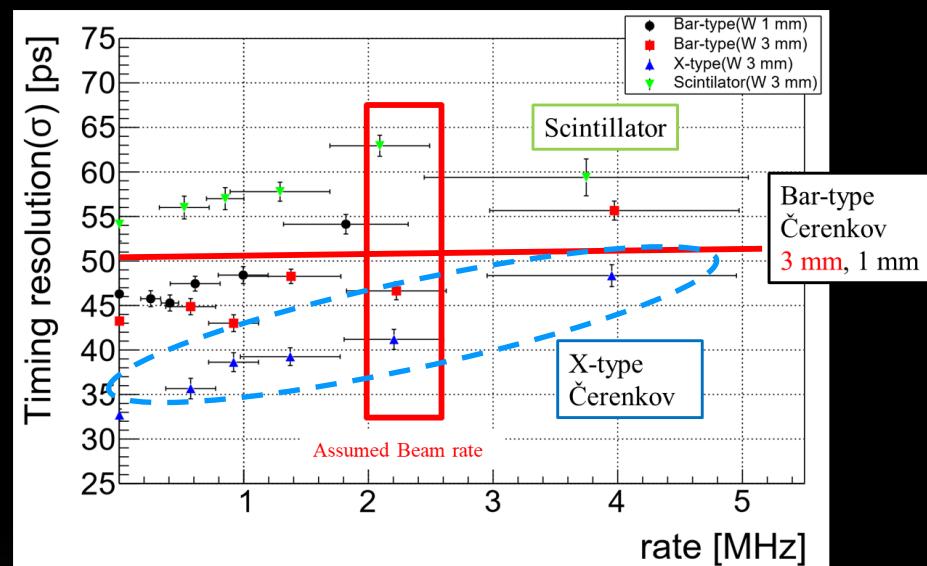
Time Zero Counter

- Hodoscope w/ Cherenkov Radiator for a Beam Rate: 60 M/spill (30 MHz)
 - X-shape to cancel position dependence by taking mean time
 - $\sigma < 50$ ps at 3-5 MHz

By T. Akaishi, K. Shirotori et al.

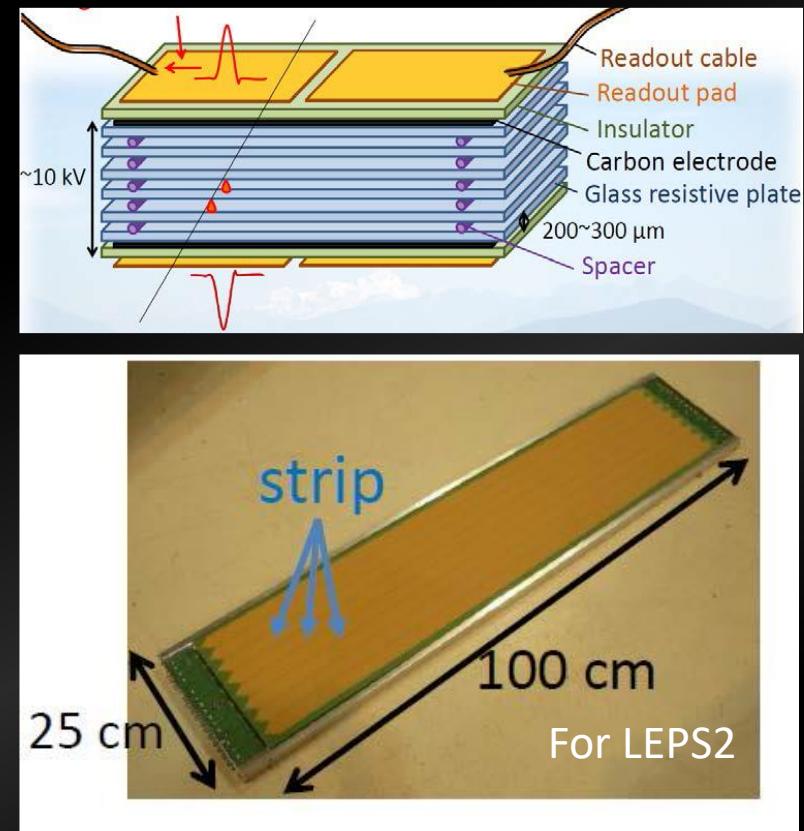
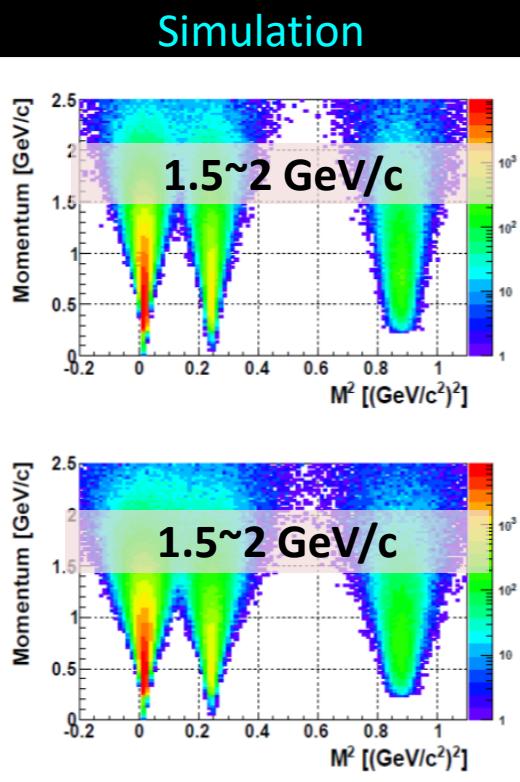
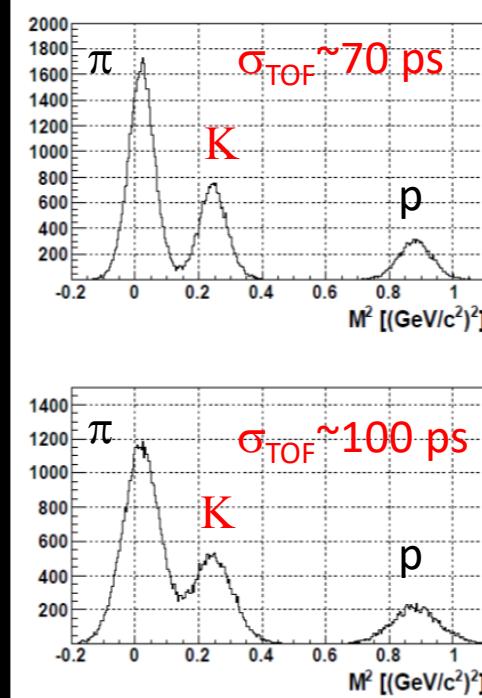


Measured Performance for MIPs



Resistive Plate Chamber

- TOF meas. for Scattered Particles
 - Developed in LEPS
 - $\sigma \sim 60$ ps

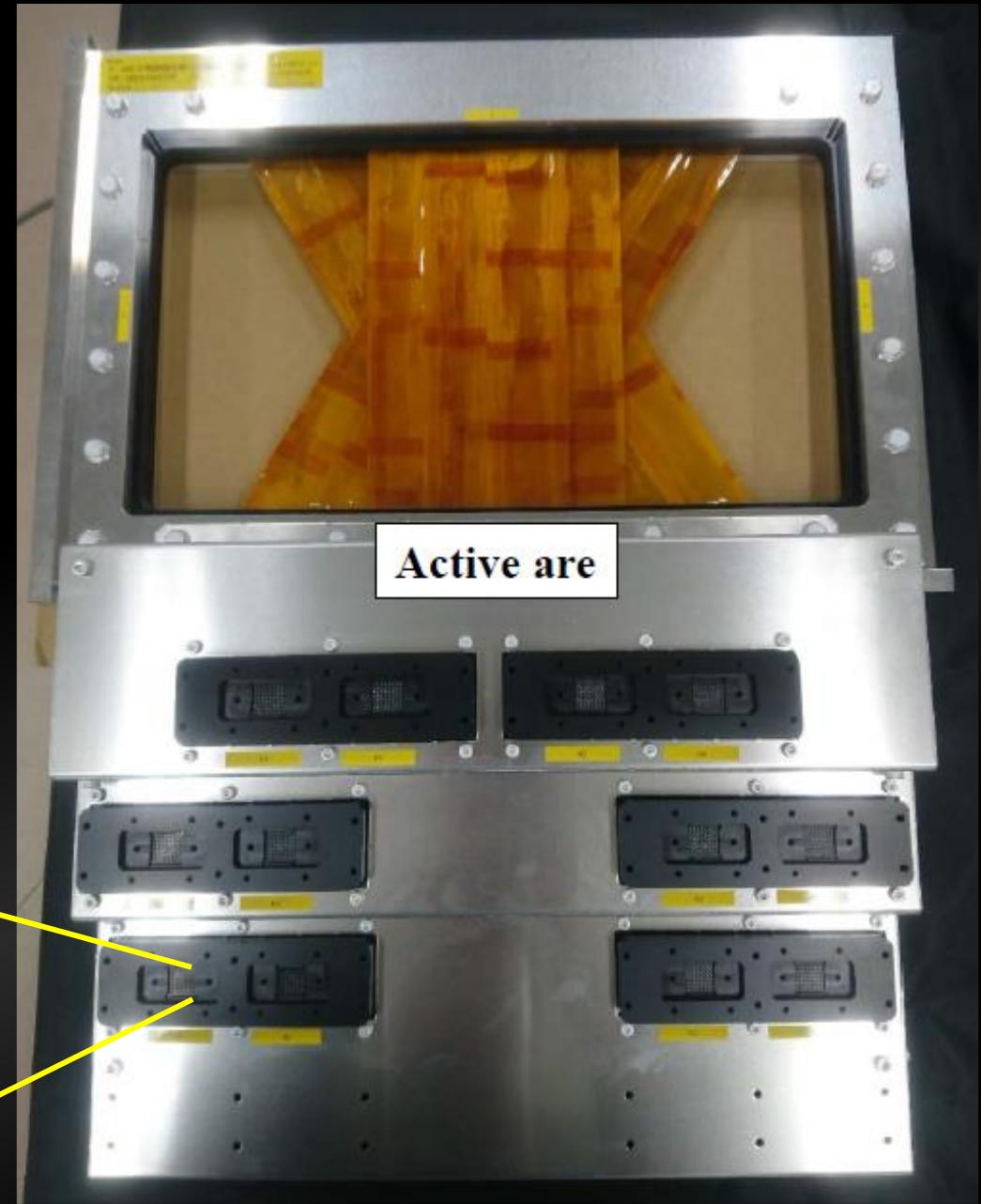
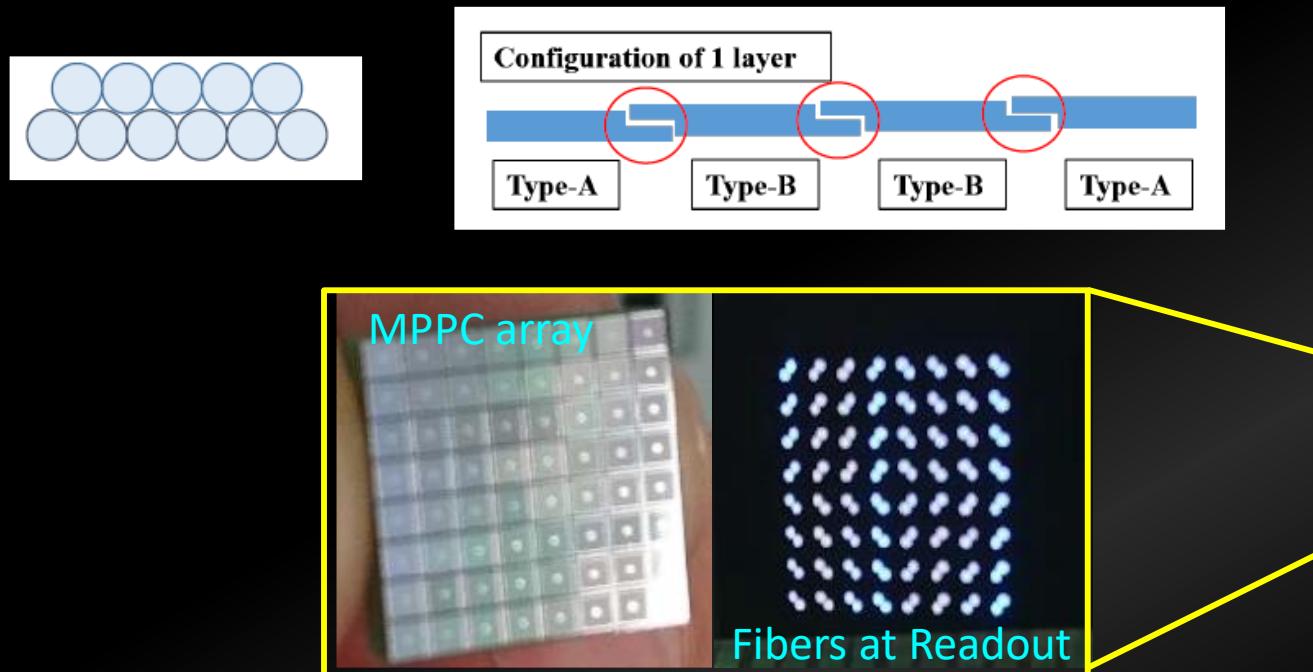


Expected PID performance w/ TOF vs Mom.

Fiber Tracker

By K. Shirotori et al.

- Faster Responding Trackers are needed for a Beam Rate: 60 M/spill (30 MHz)
 - Focal plane: XUV 1 set w/ $\phi 1\text{mm}$ fiber
 - Beam Trackers: XUV 2 sets w/ $\phi 0.5\text{mm}$ fiber
 - Scattered Particle Trackers : **in Fabrication**



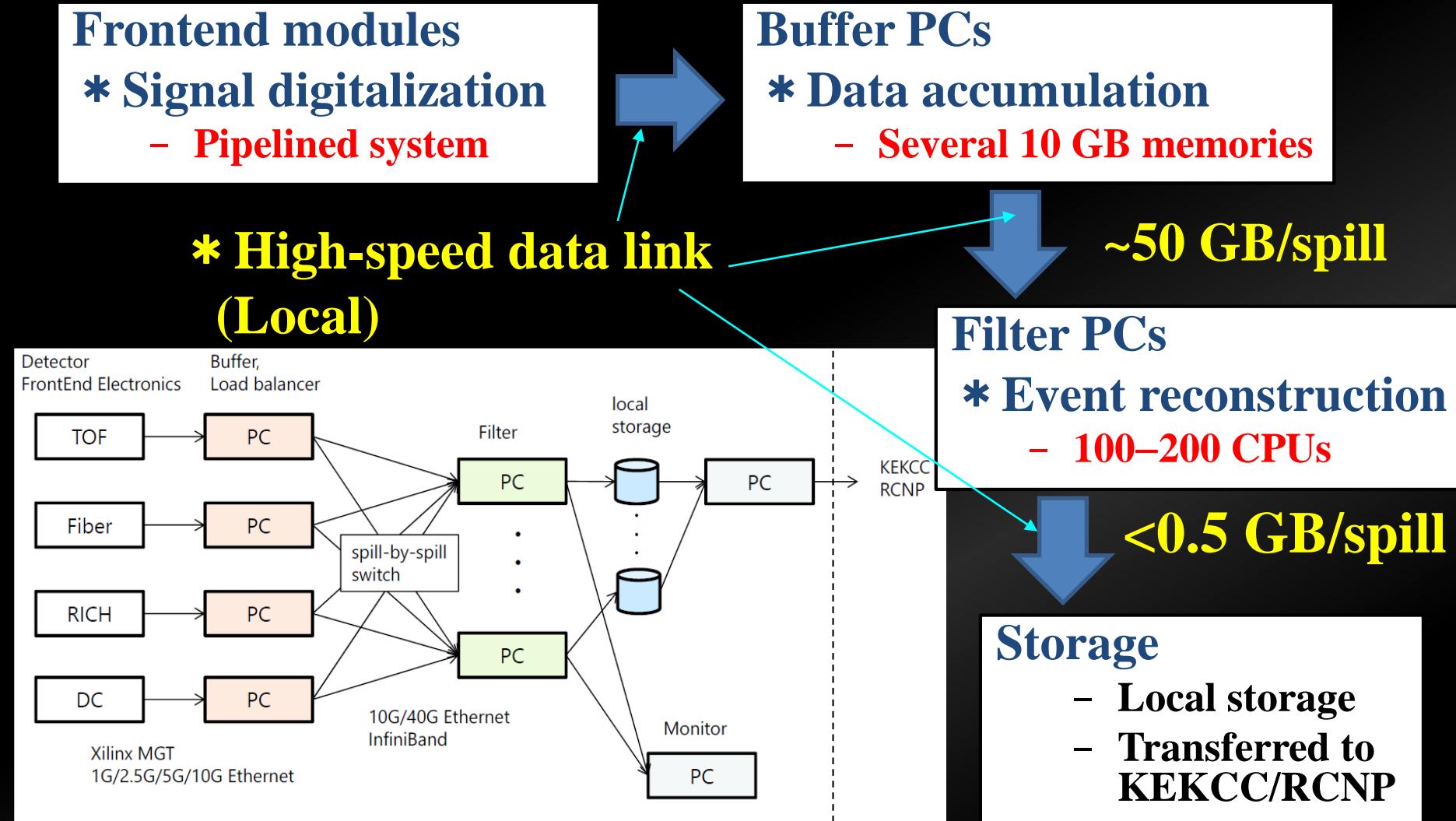
Drift Chamber

- Barrel DC (Side DC) for backward-emitted, low mom. particles
 - Two arms are ready and **waiting for FEEs.**
- Front/Forward DC for Forward-emitted particles
 - To be prepared
 - **still missing pieces for better redundancy**

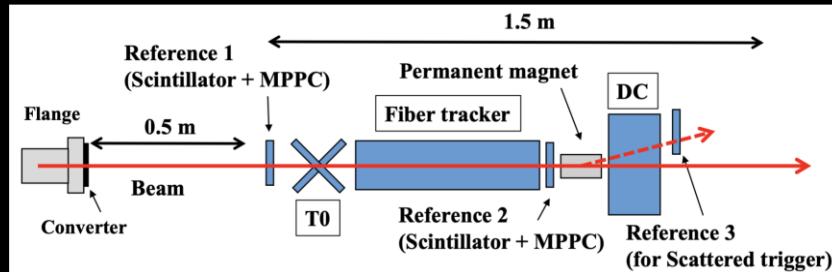


High-speed DAQ system

Streaming DAQ(~50 GB/spill)

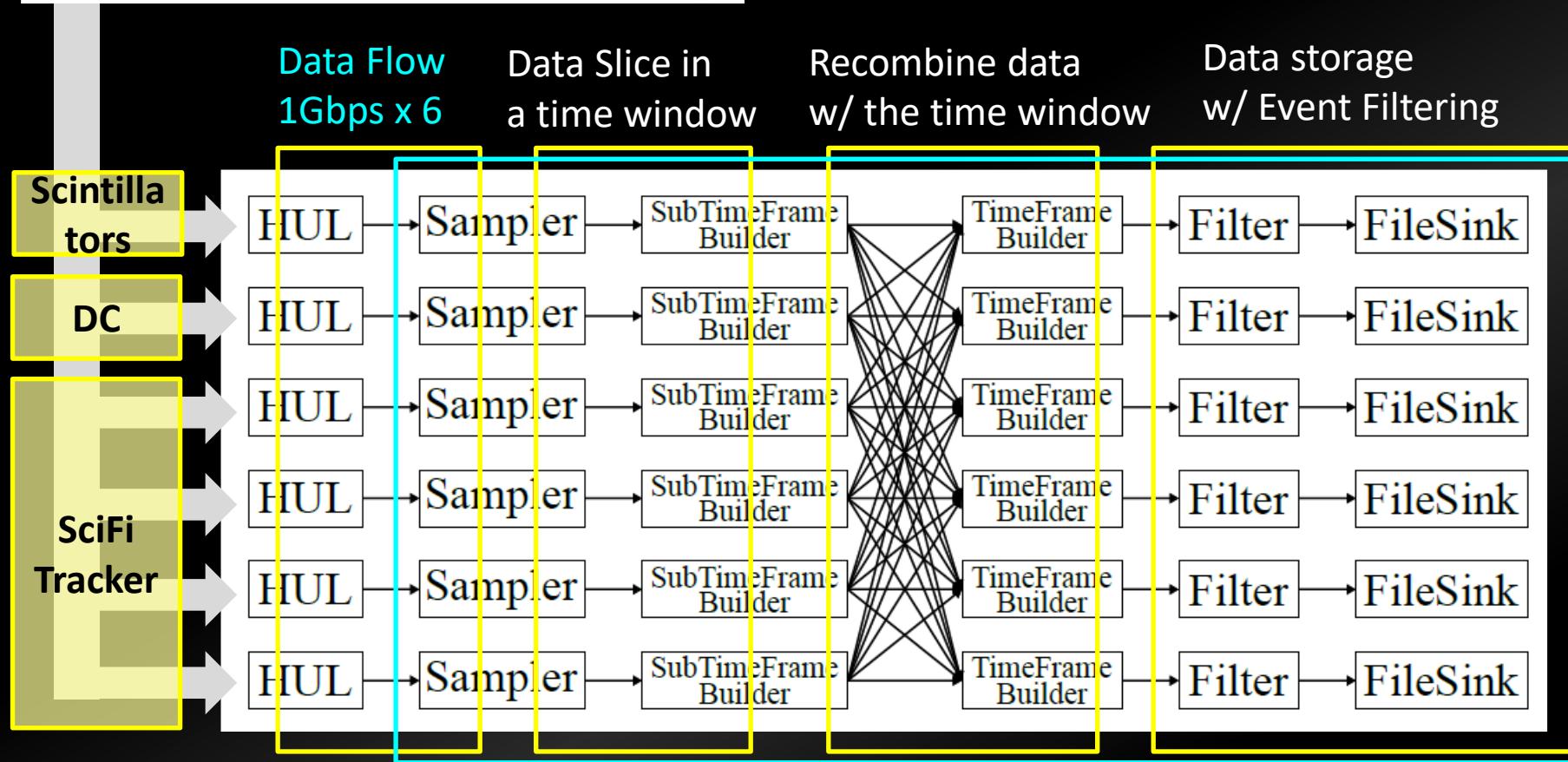


Demonstration of High-speed DAQ



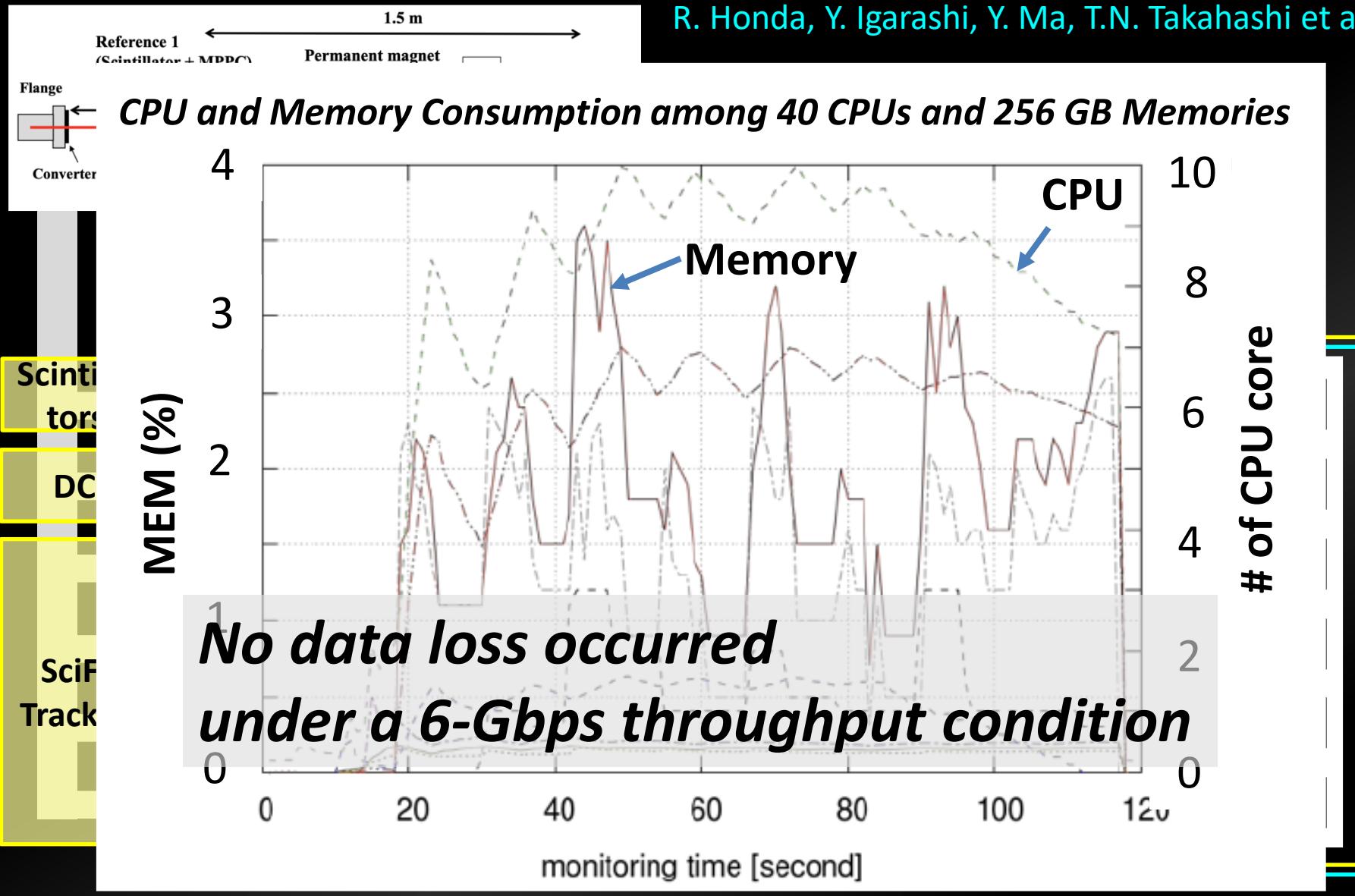
under the Highest Throughput (SiTCP: GbE)

R. Honda, Y. Igarashi, Y. Ma, T.N. Takahashi
et al. (paper in preparation)



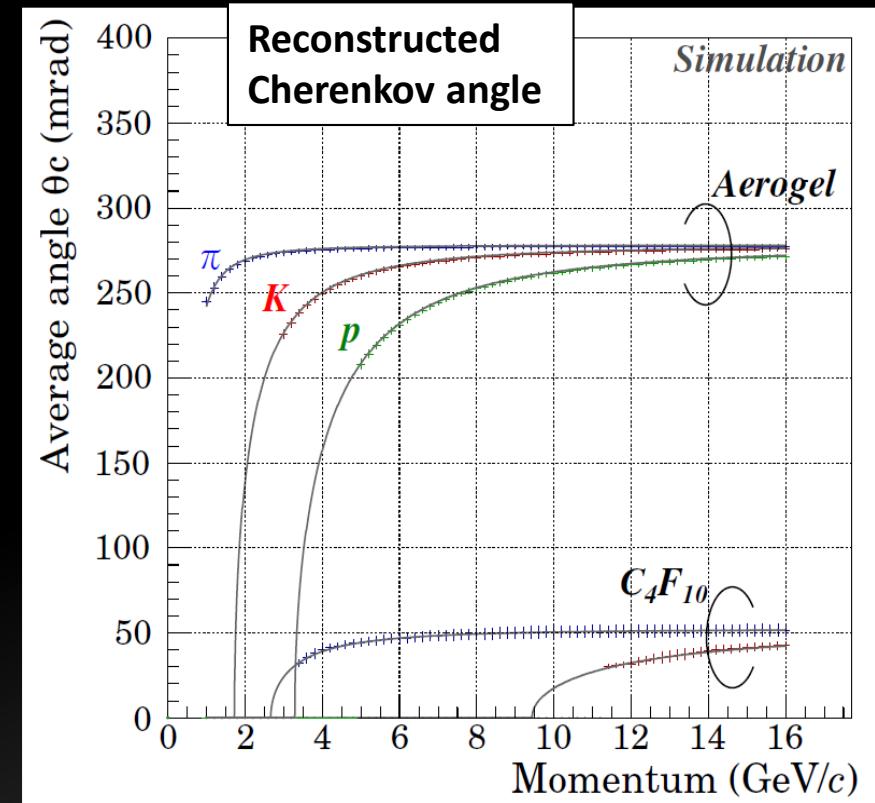
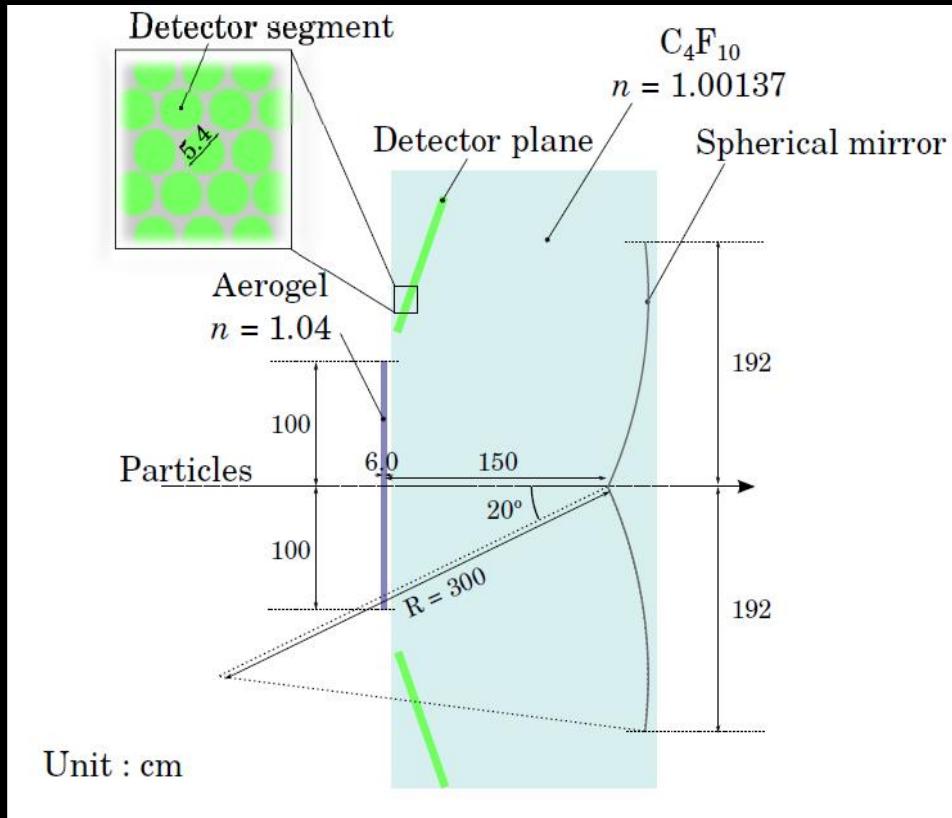
Demonstration of High-speed DAQ

R. Honda, Y. Igarashi, Y. Ma, T.N. Takahashi et al.



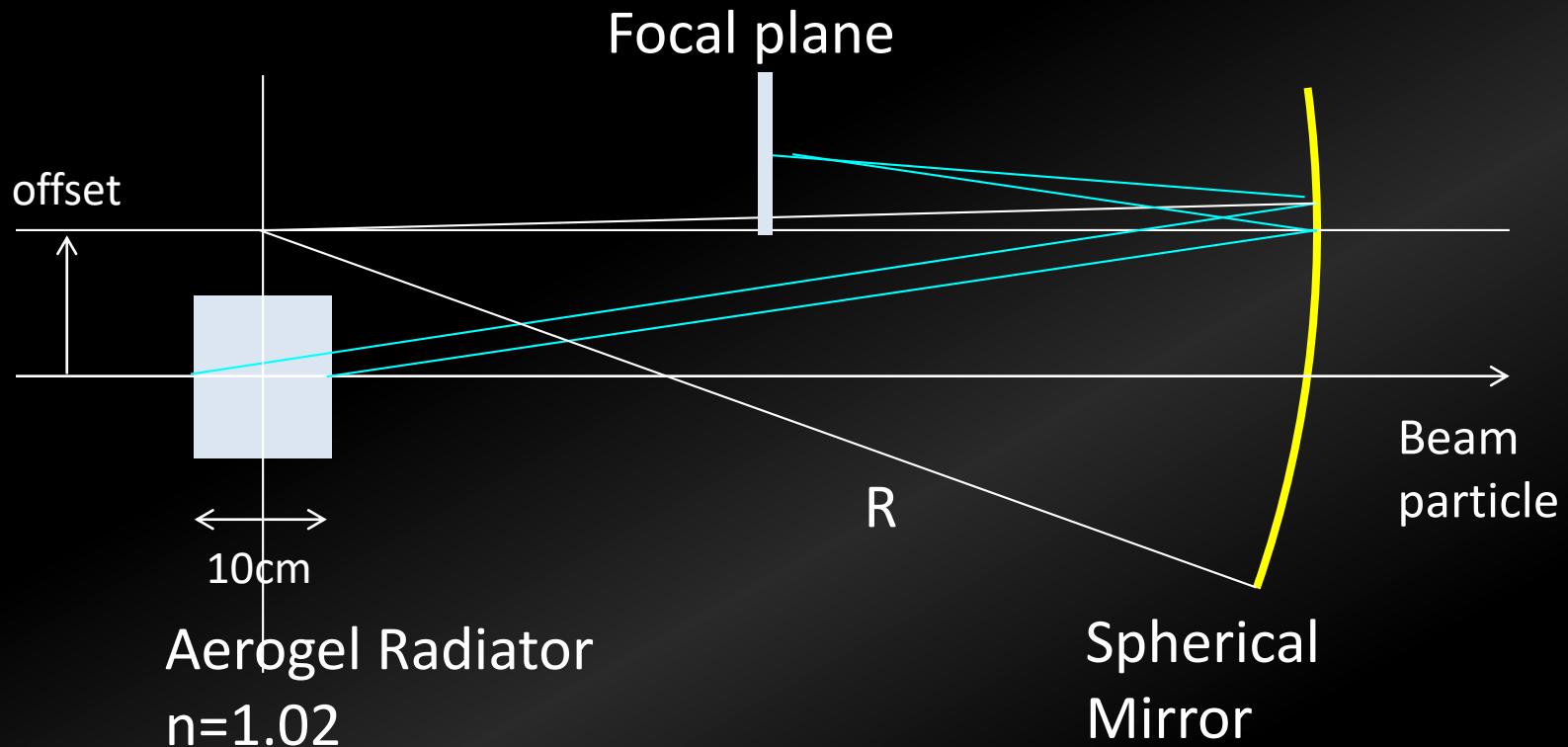
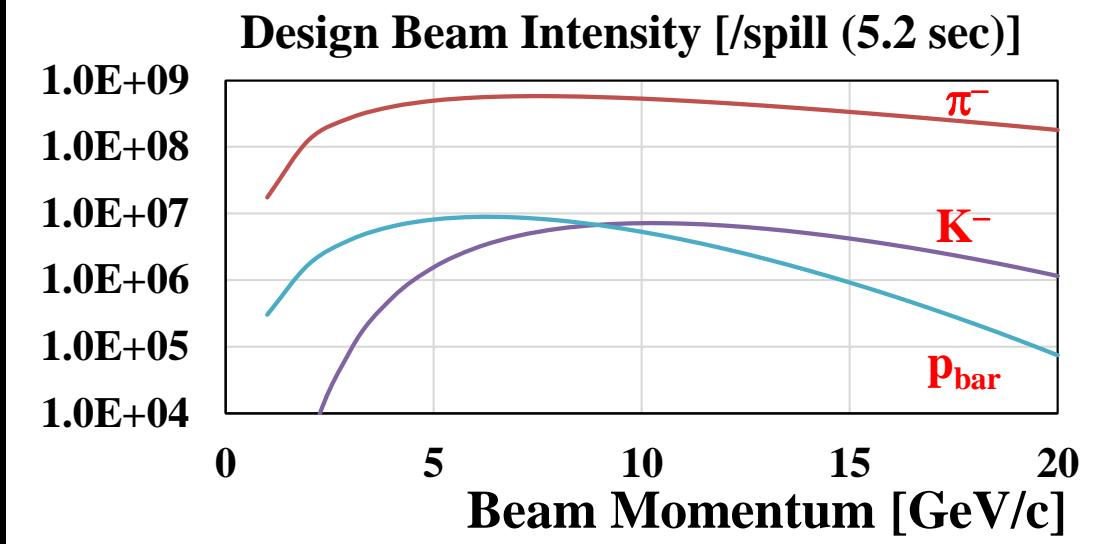
Ring Image Cherenkov detector

By T. Yamaga

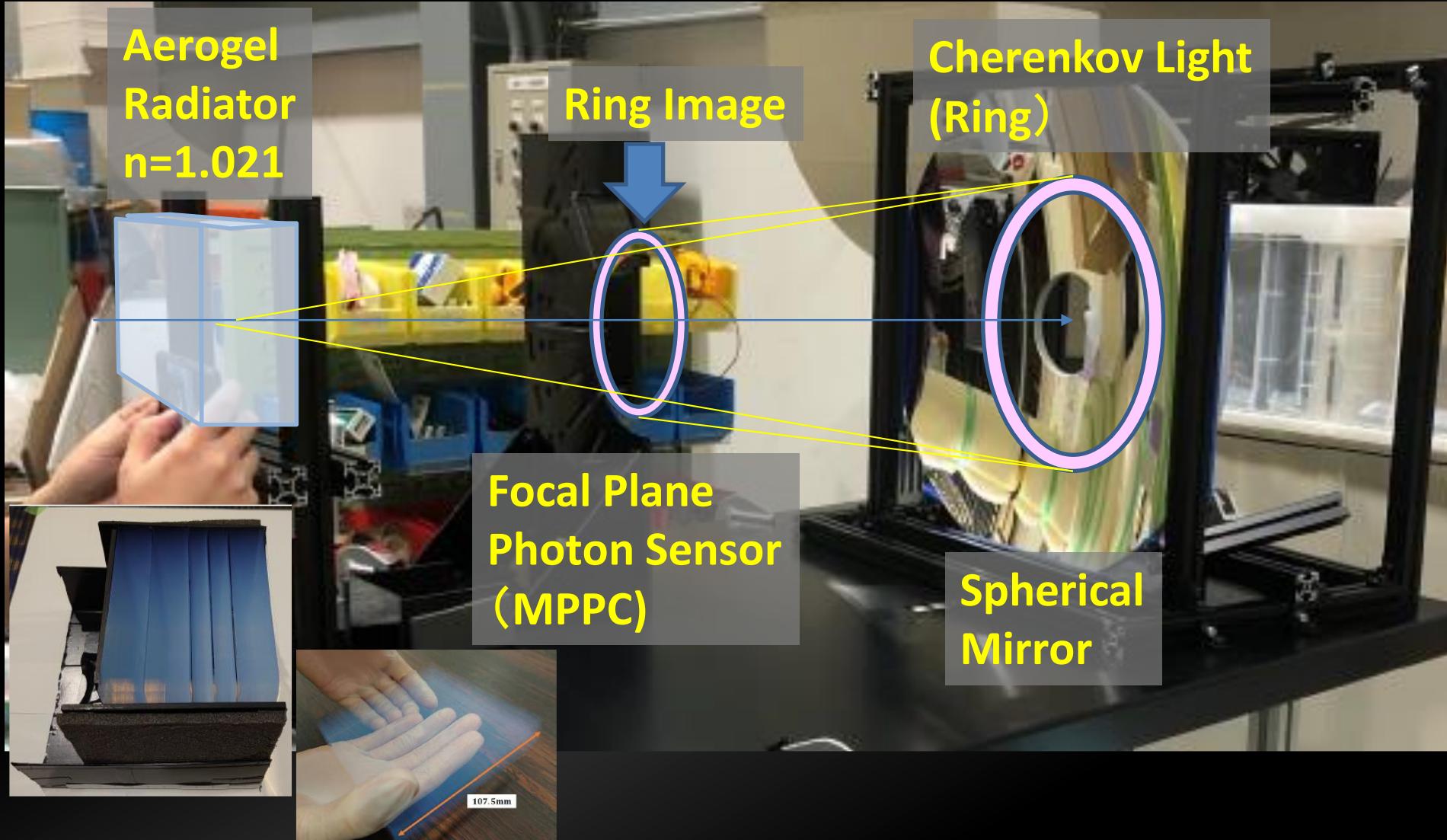


- High-momentum PID
 - Mom. range: 2-16 GeV/c
- ⇒ Hybrid RICH
 - Aerogel + C_4F_{10} gas
 - MPPC detector plane
 - Spherical mirror
- Design Performances
 - Efficiency: ~99%
 - Wrong PID: 0.10~0.14%
- ⇒ Background level × 1.05

S. Kajikawa
Master Thesis
2021 Tohoku Univ.

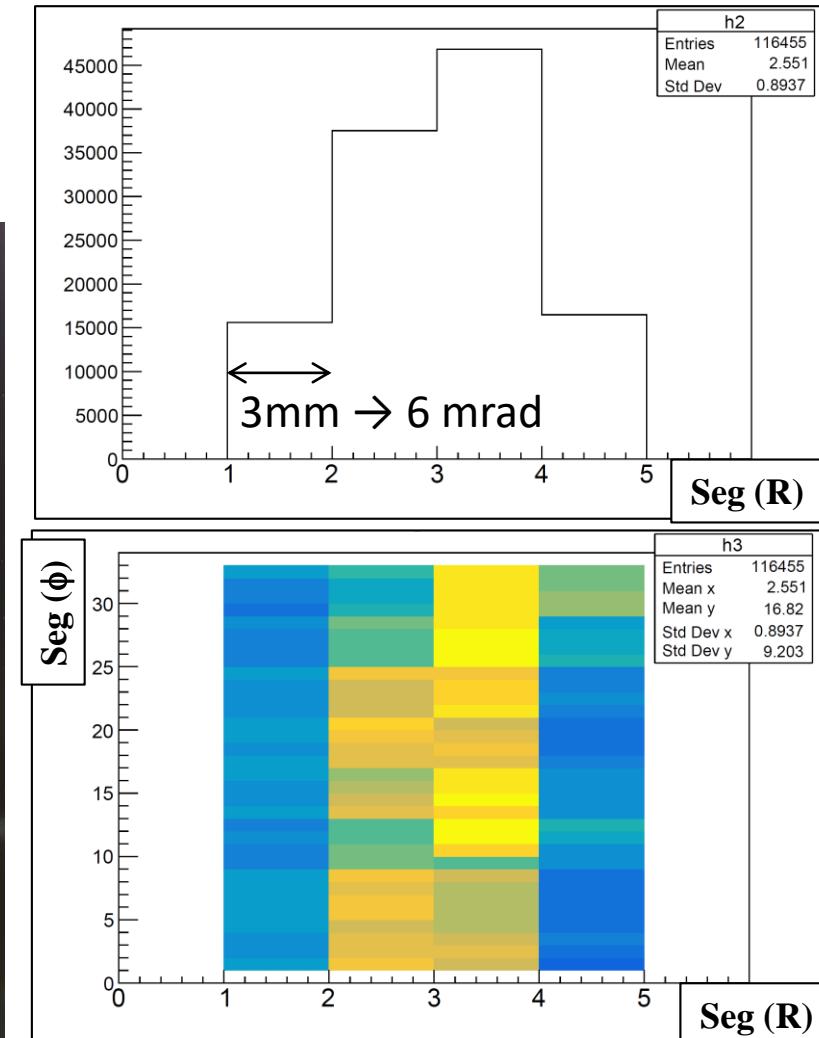
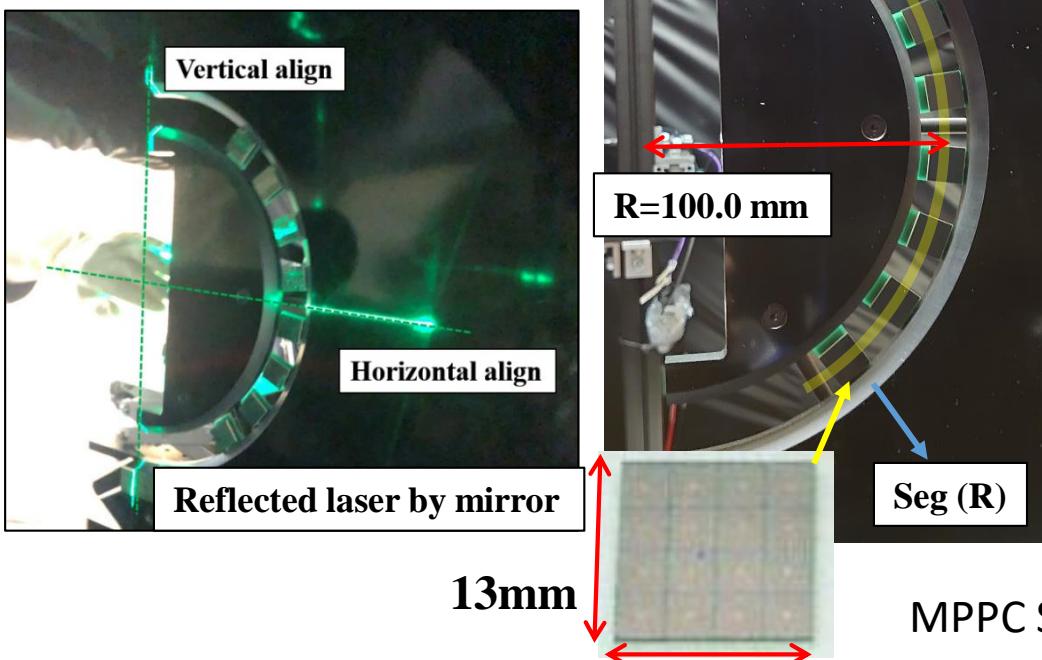


Beam RICH: test w/ electron@LEPS



Hit pattern on detector plane

- MPPC segment configuration
(3-mm 4×4 array, 8 sets)
 - Seg(R): 1–4 seg
 - Seg(ϕ): 1–32 seg
(half a full space)
- Clear hit pattern
 - Alignment: OK

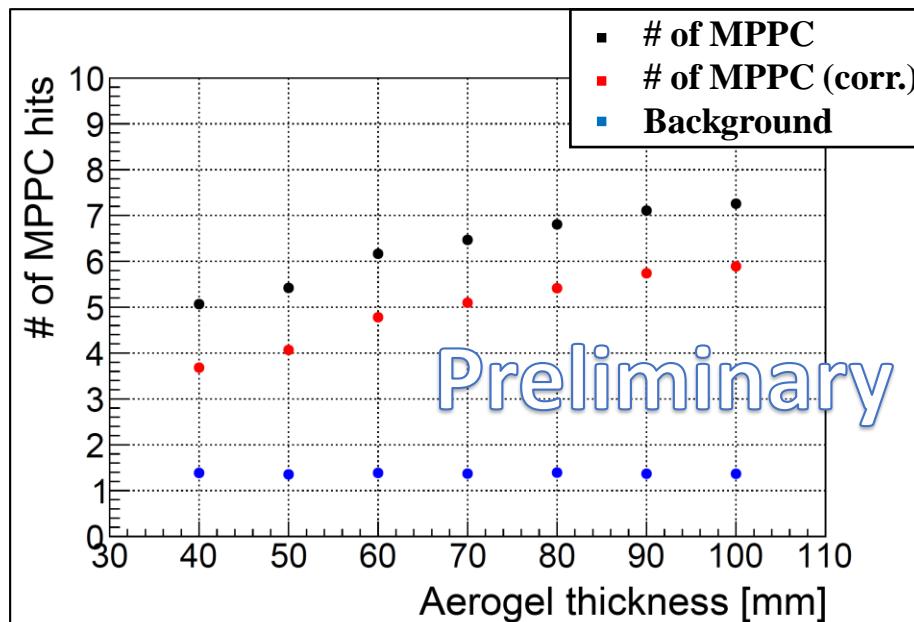


MPPC S13361-3050: 3-mm 4x4 array

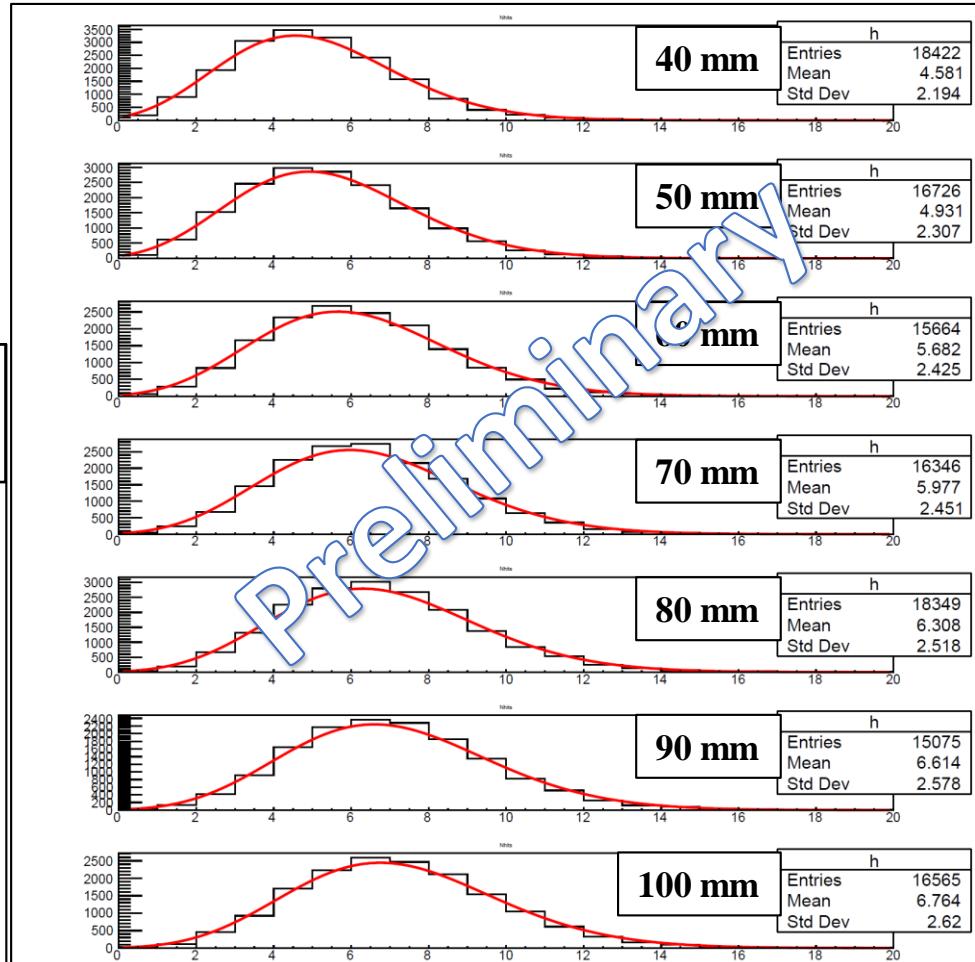
Number of MPPC hits (# of photons) : Aerogel thickness dependence

- 1 MPPC hit = 1 p.e.
- Poisson fitting
- $V_{ov} = 54.0 \text{ V}$ (+3.0V)
- Background: Side band of TOF
 - 30 ns rough time gate

※ 1 p.e. Vth operation of MPPC has been confirmed.



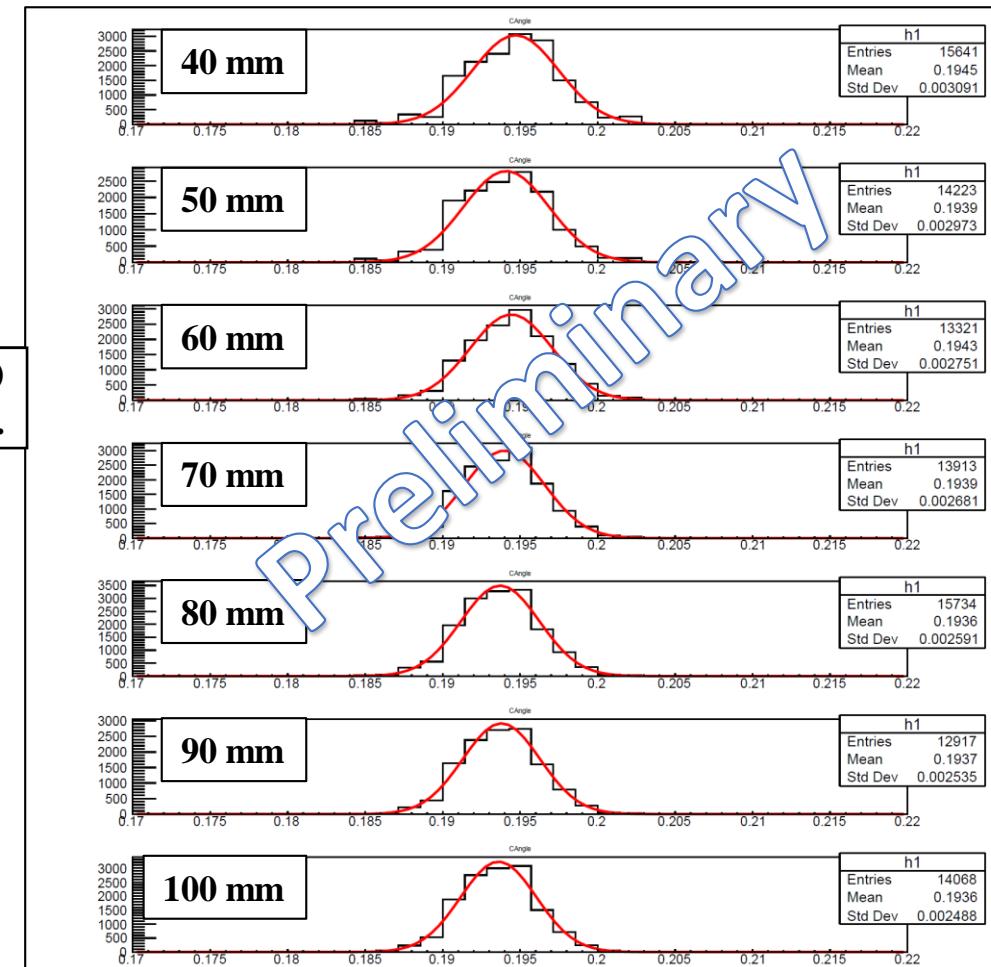
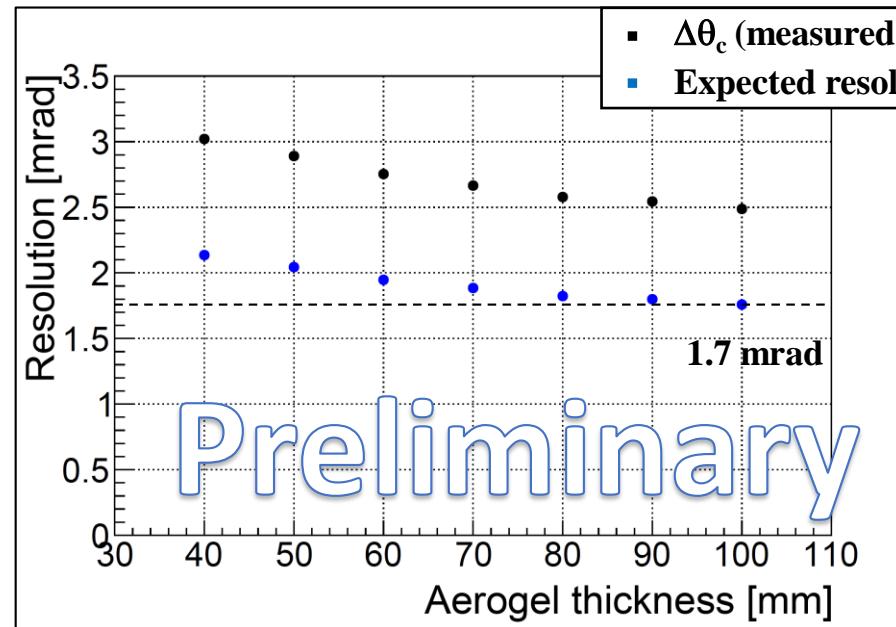
Estimation: ~5 p.e. (100 mm)



Cherenkov angle resolution: Aerogel thickness dependence

- $V_{ov} = 54.0 \text{ V (+3V)}$
- Gaussian fitting resolution:
- Expected performance achieved:
 $\sigma_{\theta_c} \sim 2.5 \text{ mrad}$
 $\Rightarrow \# \text{ of MPPCs can be double}$
 $2.5/\sqrt{2} \sim 1.7 \text{ mrad}$

※ Aerogel is sufficiently transparent up to 100 mm.



Summary

- A heavy quark plays an inert particle in a hadron and is quite helpful to investigate internal motions and/or correlations of quarks.
 - Excitation Energy, Production Rate, and Decay Branching Ratio
- We conduct charmed baryon spectroscopy by means of missing mass technique at the J-PARC high-momentum beam line, where the intense pion beams up to 20 GeV/c will be delivered.
 - New platform of hadron physics will be covered owing to the general purpose spectrometer