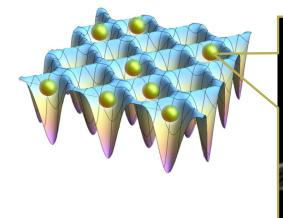
Fundamental Physics with Cold Radioactive Heavy Elements

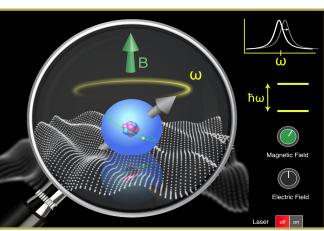
Yasuhiro SAKEMI (UTokyo)

Quantum sensing with the crystal of the radioactive atoms

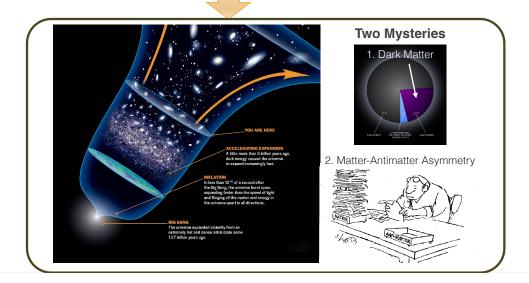
- Introduction
- Electric Dipole Moment (EDM) and Heavy Element
- Status of the experiment
- Co-magnetometer
- Future plan

Anti-matter disappearance mechanism in the universe





Quantum sensing with RI crystal of the cold atoms to search for the new particles



Sakharov's Conditions: Need CP-Violation



VIOLATION OF CP INVARIANCE, C ASYMMETRY, AND BARYON ASYMMETRY OF THE UNIVERSE

A. D. Sakharov Submitted 23 September 1966 ZhETF Pis'ma <u>5</u>, No. 1, 32-35, 1 January 1967

The theory of the expanding Universe, which presupposes a superdense initial state of matter, apparently excludes the possibility of macroscopic separation of matter from antimatter; it must therefore be assumed that there are no antimatter bodies in nature, i.e., the Universe is asymmetrical with respect to the number of particles and antiparticles (C asymmetry). In particular, the absence of antibaryons and the proposed absence of baryonic neutrinos implies a non-zero baryon charge (baryonic asymmetry). We wish to point out a possible explanation of C asymmetry in the hot model of the expanding Universe (see [1]) by making use of effects of CP invariance violation (see [2]). To explain baryon asymmetry, we propose in addition an approximate character for the baryon conservation law.

The Nobel Foundation

Matter dominated universe :

- 1. A baryon number violating interaction exists.
- 2. Departure from thermal equilibrium.
- 3. CP(Charge conjugation and Parity)-symmetry must be violated.

Standard Model CP-Violation : Not Enough

$$\eta = \frac{(\text{matter}) - (\text{antimatter})}{\text{relic photons}} \propto \sin(\delta)$$
$$\eta_{\text{exp}} \approx 10^{-9} \qquad \text{PDG2022}$$
$$\eta_{\text{CKM}} \approx 10^{-26} \qquad \text{Huet \& Sother PRD 51:379 (1995)}$$

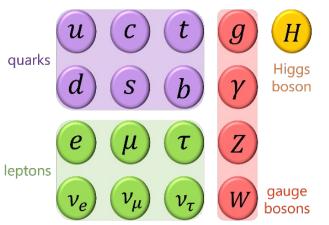
 $V = \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}\exp(-i\delta) \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}\exp(+i\delta) & +c_{12}c_{23} - s_{12}s_{23}s_{13}\exp(+i\delta) & s_{23}c_{13} \\ +s_{12}s_{23} - c_{12}c_{23}s_{13}\exp(+i\delta) & -c_{12}s_{23} - s_{12}c_{23}s_{13}\exp(+i\delta) & c_{23}c_{13} \end{bmatrix}$

 $\delta = CP$ -violating "phase"

Search for beyond-standard-model (BSM)

Big issues in particle physics now

- Is the discovered Higgs particle the SM one ?
- Where is BSM ? TeV scale or higher energy scale ?



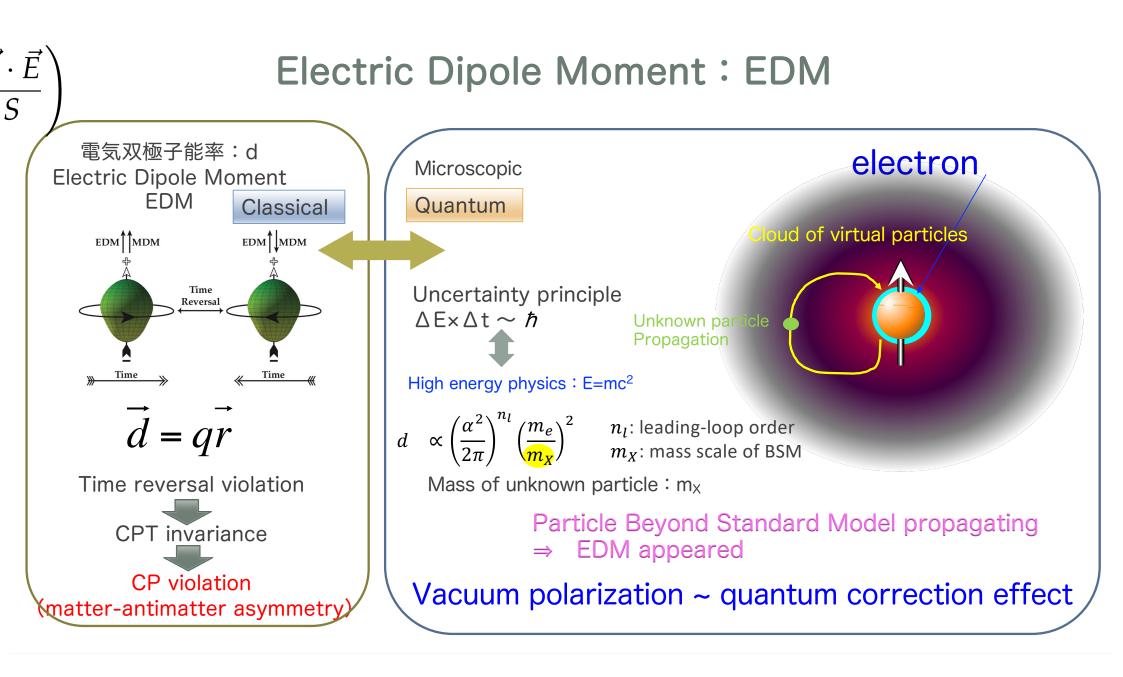
Remaining mysteries

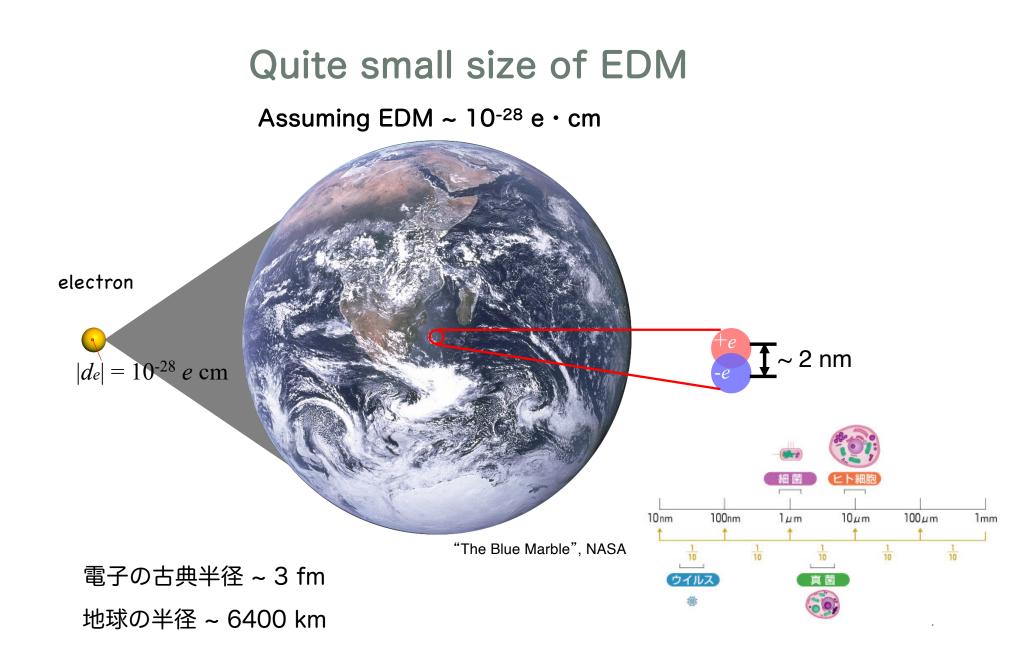
- What is dark matter?
- How is neutrino mass generated?
- Why is the Higgs boson so light?
- Why do the proton and electron have same magnitude of charge?
- Why is antimatter so scarce in the Universe compared to matter?

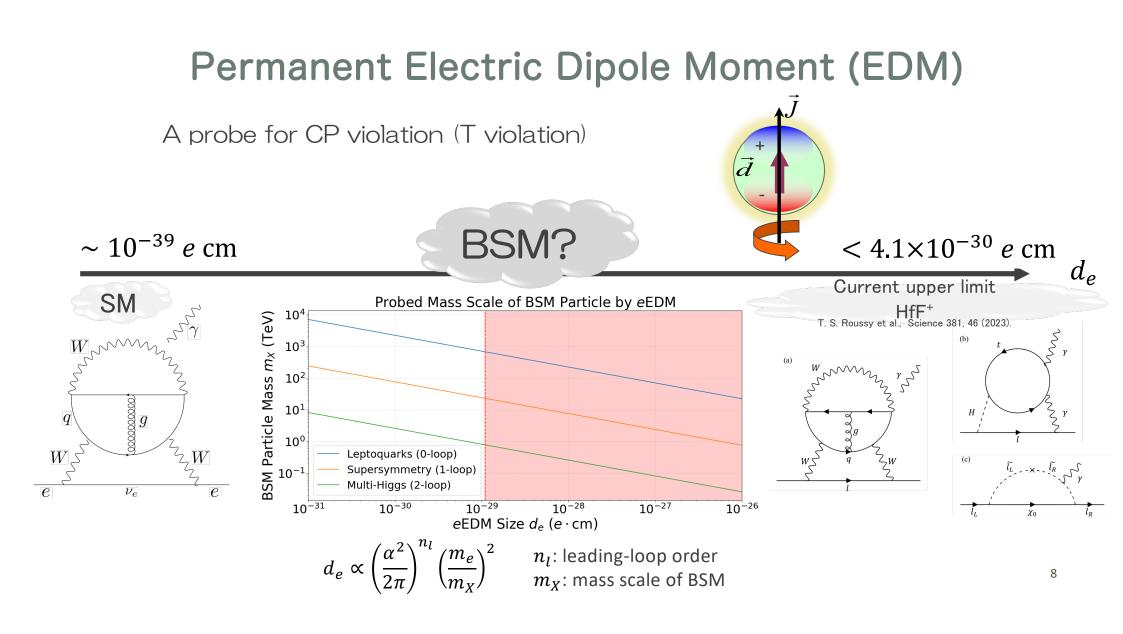
• •••

Where do we look for more CP violation ?

Decays of B-mesons
D-coefficient in nuclear beta-decay
Double polarized neutron transmission
Angular decay correlations of positronium
Neutrino oscillations
Neutrino less Double Beta Decay 0ν β β
Nuclear magnetic quadrupole moments
Electric Dipole Moment (EDM)



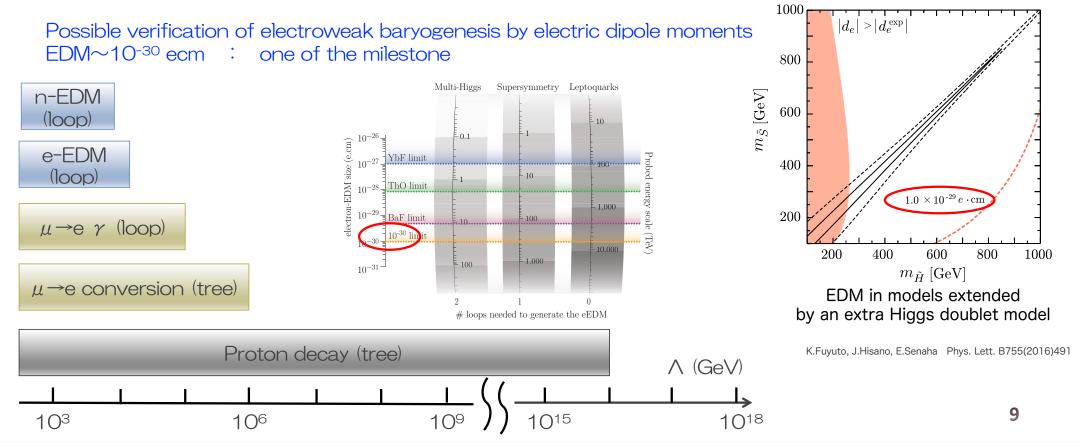




Search for symmetry breaking

Sensitivities of current experimental bounds on new physics scale (Λ) . Only one loop factors are included for the loop processes. Small symmetry breaking parameters suppress the sensitivities.

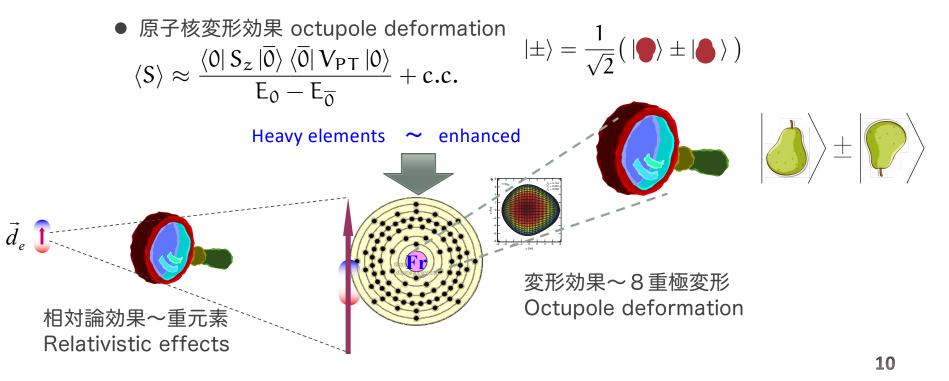
EDMs and muon LFV are important to probe new physics at and beyond TeV scale



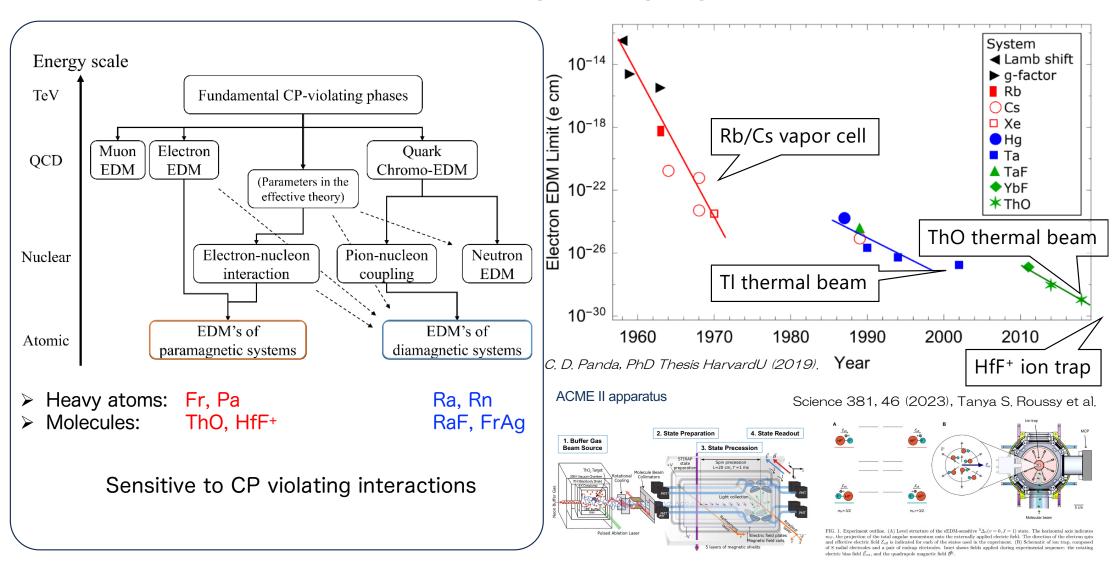
Heavy Elements (Radioisotope) ~ microscope to detect the CP violation~ 重元素~CP対称性の破れ (EDM)を拡大する顕微鏡

● 相対論効果 relativistic effects

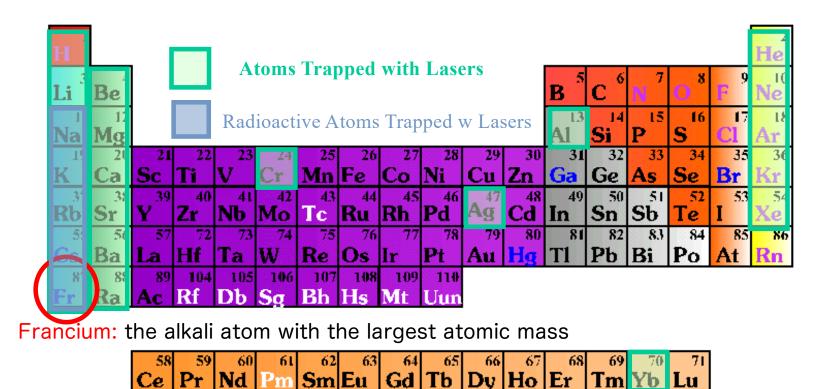
$$K \sim \frac{d_{atom}}{d_e} \sim \frac{Z^3 \alpha^2}{2} \sim \left| \psi_s(0) \right|^2 V Z^5 \alpha^2 \frac{e}{a_0^2}$$



EDM of many-body system



Candidate of the EDM search



• 7th period heavy elements

Τh

Pa

• Super Heavy Element

From ChemiCool.com

Marguerite Perey announcing the discovery of Francium 1939

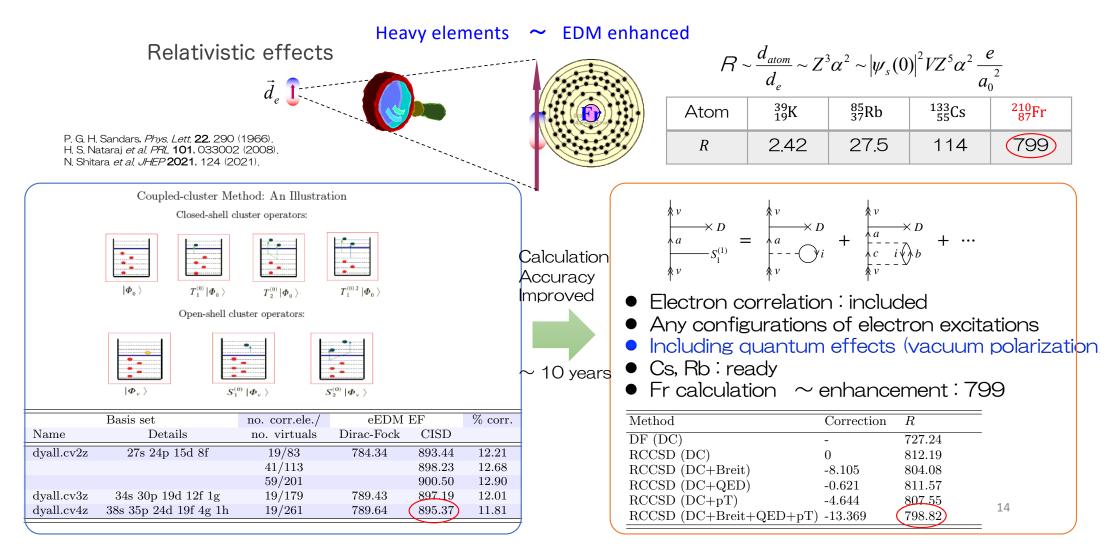


Courtesy of the Curie Institute, Paris

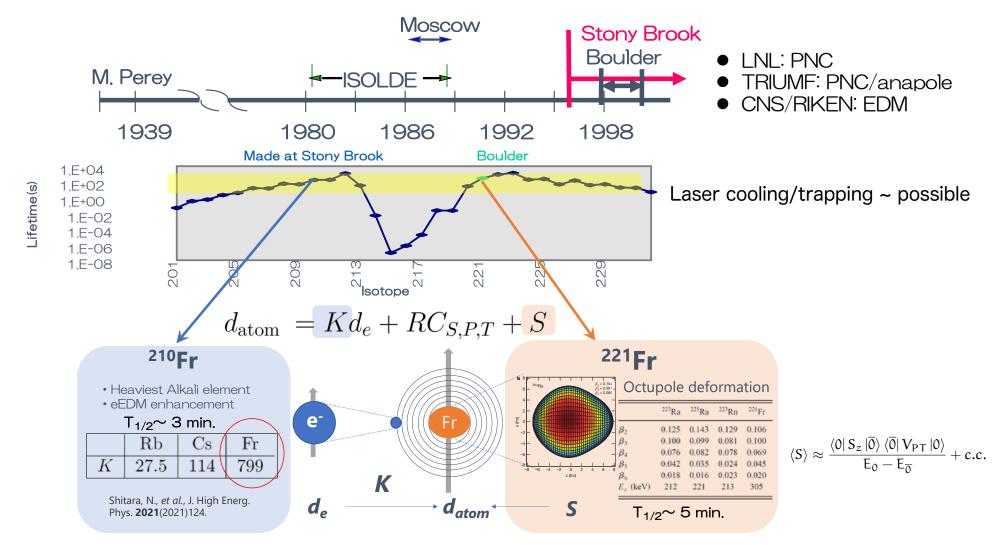
e-EDM enhancement in $a_{\Sigma}^{\dagger \nu}$ $m_{i}c_{(b)}s_{V_{a}}^{\dagger \nu}$ $s_{L}em_{\Sigma}^{\dagger \nu}$ ~Relativistic Coupled Cluster model ~

 $S^{(1)}$

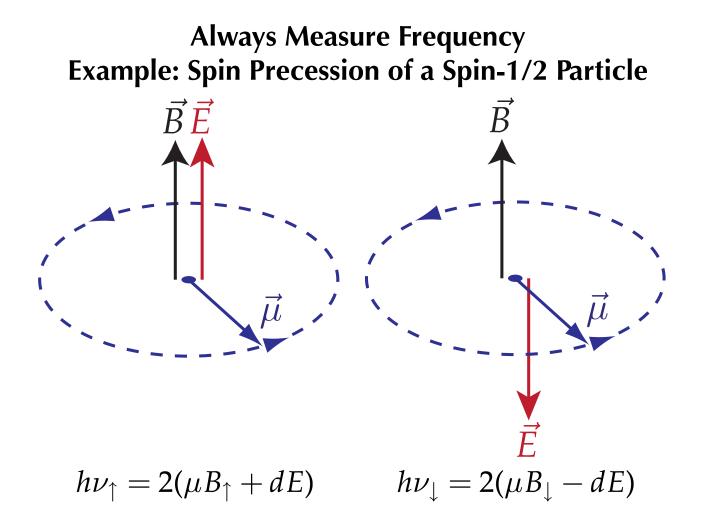
 $D \times$



Francium – MANY ISOTOPES : New laboratory for fundamental physics



Measurement of the EDM

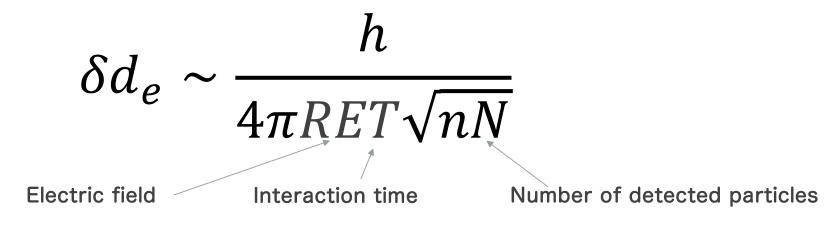


EDM sensitivity

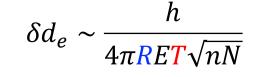
Statistics & Systematics

$$\Delta \nu = \nu_{\uparrow} - \nu_{\downarrow} = \frac{4dE}{h} + \frac{2\mu(B_{\uparrow} - B_{\downarrow})}{h}$$

Quantum Projection Noise:

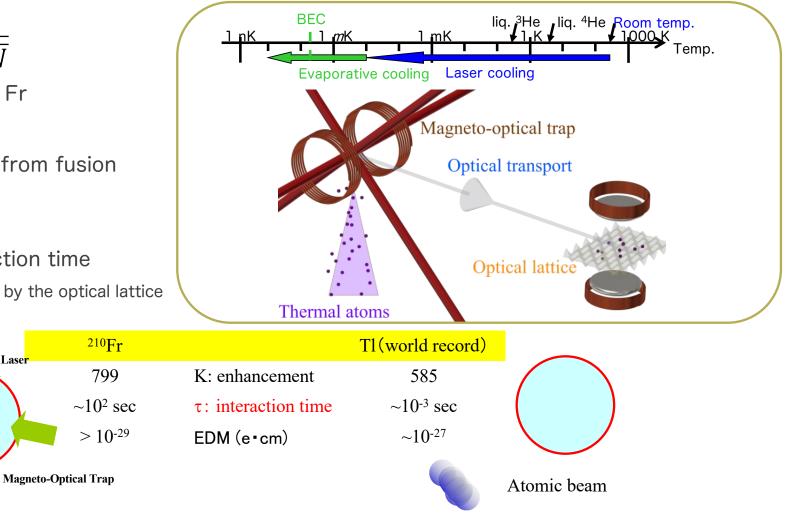


Electron EDM search using laser-cooled Fr atoms

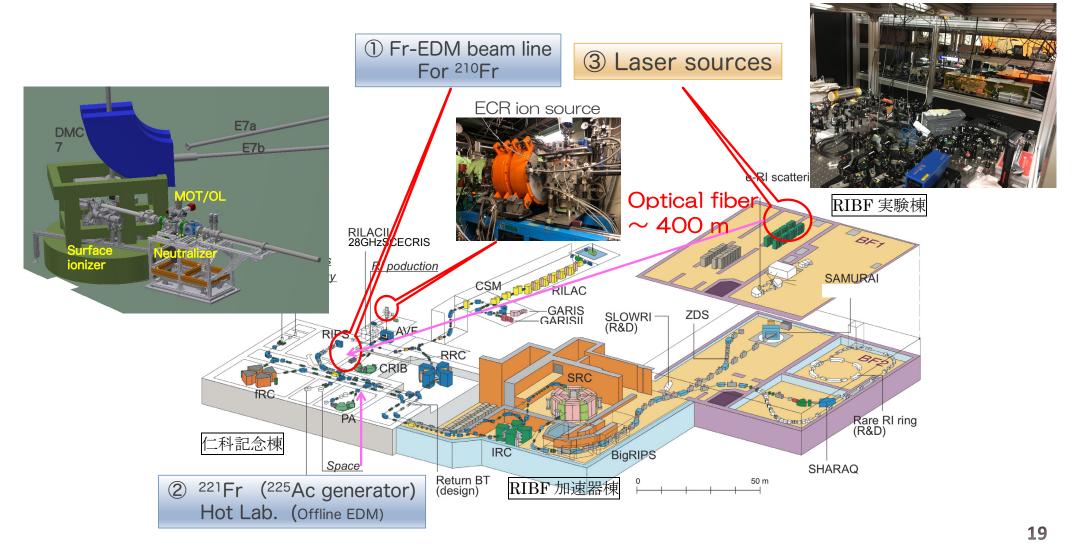


- Large R (\sim 799) : for Fr
 - > EDM enhancement factor
- Large nN (~10⁷ Fr⁺/s from fusion reaction)
- Long T (> 10 s) : interaction time
 - > Collision-suppressed trapping by the optical lattice

Laser



RIKEN RI BEAM FACTORY and EDM apparatus





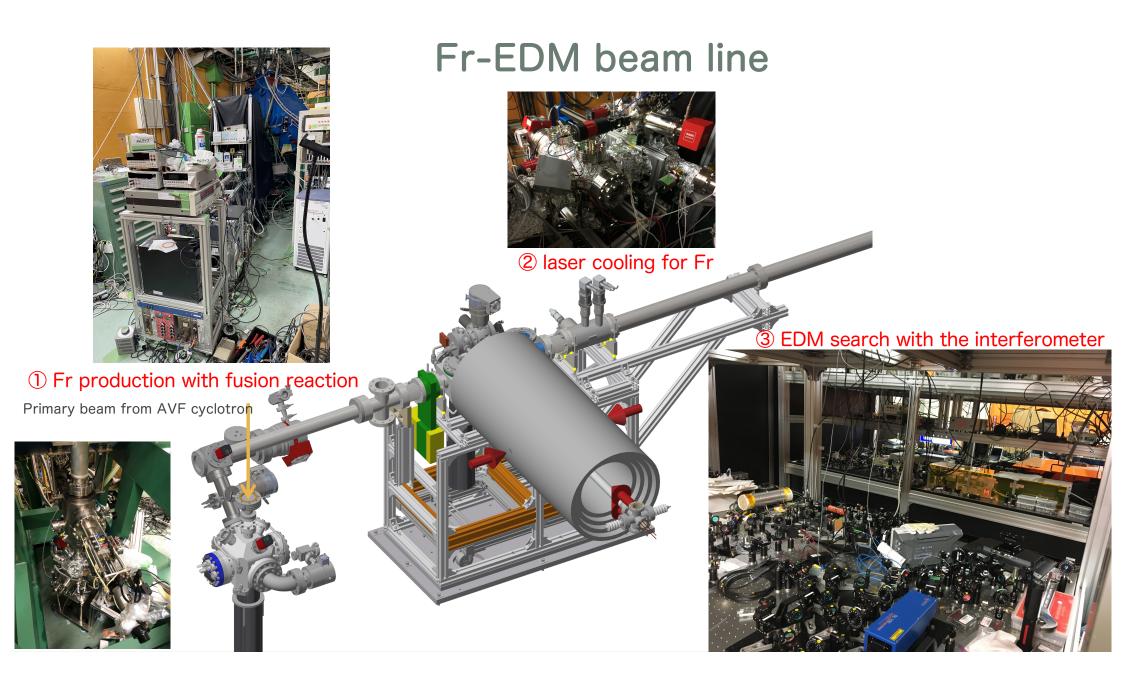
Electron EDM enhancement

Calculated by the relativistic coupled-cluster theory to an accuracy of about 3%

	Rb	Cs	Fr
K	27.5	114	799

Shitara, N., *et al*., J. High Energ. Phys. **2021**(2021)124.

- Radioactive $T_{1/2}$ ~3min
- Make it by nuclear fusion reaction



The striking imbalance between matter and antinatter observed in the universe is known to be one of ne biggest mysteries in modern physics. For decades, hysicists have been attempting to investigate CP (C: narge conjugation, P: parity transformation) violaons for different physical systems, which could lead to ne discovery of the fundamental reasons for this mysery. One of these intriguing systems is the electric ipole moment of an electron (eEDM). The existence f an absolute value of the EDM implies that SP symptons that the same time, this would ${f n}$ so reveal the existence of a new theory beyond the andard model of particle physics.

Since the eEDM is expected to be small,¹⁾ it is need ssary to adopt a system hat onhances the eEDM to sufficiently large value. For this purpose, we have nosen francium. In Indua has an eEDM enhancenent factor of approximately 10^3 , which is known to e the largest among any ground-state atom.²⁾ T leal environment for **R**amsey spectroscopy to **A** re the eEDM using fr<mark>a</mark>ncium atom is one in **x** rge amount of francium atoms are trapped in a tiny olume where the interactions with other particles ighly suppressed. To this end, we are planning evelop a three-dimensional optical lattice to confine ser-cooled francium a<mark>to</mark>ms. In this study, we tins<mark>tructual</mark> a surface ionizer to cr

eflection electrodes ons will experience e guided at an ang<mark>l</mark>

te a high-intensity **franci**um ion beam. The surf mizer consists of a vacuum chambe eflection electroder and an infrared radiation heat s shown in Fig. 1. An ¹⁸O⁶⁺ beam (6.28 MeV/a) pro ided by the AVF cycl<mark>o</mark>tron in RIKEN RIBF will be radiated on the gold target to induce the h on reactions inside the gold: ¹⁹⁷Au(180, xn)^{215-x}Fr. y heating up the target using the infrared heater, the nermal diffusion process of the franciums will occur, nd some fractions will reach the surface. Most of the anciums at the surface will be ionized according to ne Saha-Langmuir equation³⁾ and, subsequently, therally released into the vacuum. Since high voltages re applied to the target $(\sim 1000 \text{ V})$ as well as to the 30 V), the thermal francium electric-field gradient and be cansported as a secondary beam ($\sim 1 \, \text{keV}$), which will 45° wiRregeter¹⁸

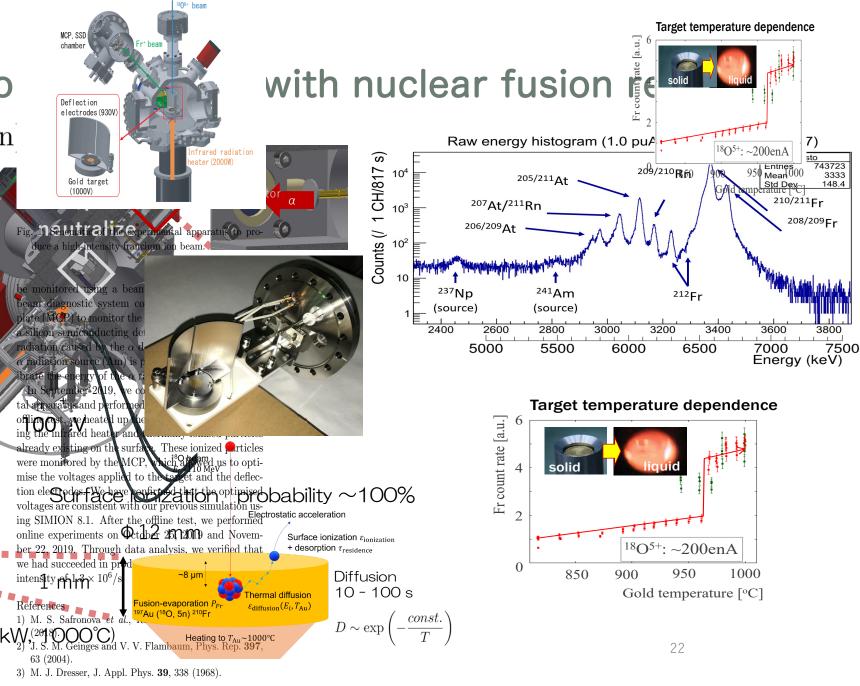
Center for Nuclear Study, the University of Tokyo

eam. The francium ion beam generated will finally

RIKEN Nishina Center

Waseda University Honjo Senior High School

Cyclotron and Radioisotope Center, Tohoku University



Fr production yield

- ²¹⁰Fr yield analysis
 - 1. Fit energy spectrum with multi-skewed Gaussian
 - $f(E) \ni \frac{1}{2\tau} \exp\left[\pm \frac{E-\mu}{\tau} + \frac{\sigma^2}{2\tau^2}\right] \operatorname{erfc}\left[\frac{1}{\sqrt{2}}\left(\pm \frac{E-\mu}{\sigma} + \frac{\sigma}{\tau}\right)\right]$
 - 2. Fit decay curve at each peak $e^{-\frac{t-t_1}{2}}$

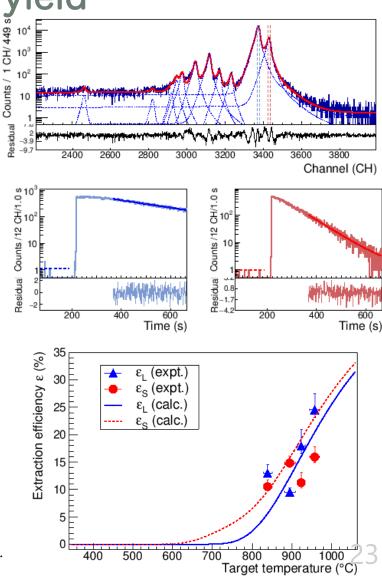
$$S(t) = S_L e^{-\tau_L}$$

3. Calculate ion beam flux

$$f_{210} = \frac{P_{210}}{b_{210}P_{210} + b_{211}P_{211}} \frac{\tilde{S}_L}{\Omega_{\text{SSD}}(e^{t_1/\tau_L} - e^{t_0/\tau_L})} = 6.7^{+0.8}_{-0.6} \times 10^6 \text{ s}^{-1}$$

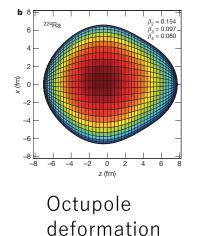
- Extraction efficiency: (Faraday cup) / (Target) $\varepsilon_L = \frac{1}{b_{210}P_{210} + b_{211}P_{211}} \frac{\tilde{S}_L}{\Omega_{\text{SSD}}(e^{t_1/\tau_L} - e^{t_0/\tau_L})} = 24.5^{+3.0}_{-1.4} \%$
 - ✓ Efficient heating of the target✓ Efficient electrostatic transportation

N. Ozawa et al., RSI 94 (2023) 023306.



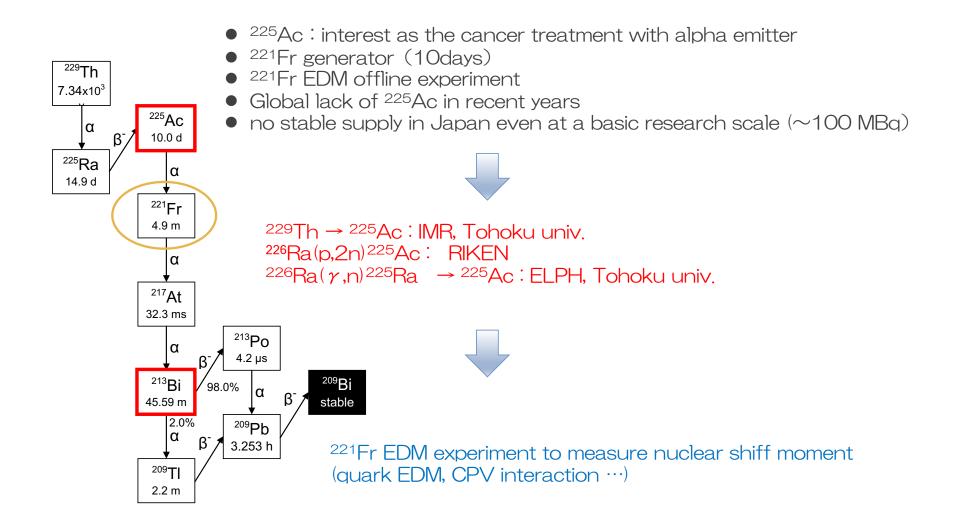


Searching for the nuclear EDM

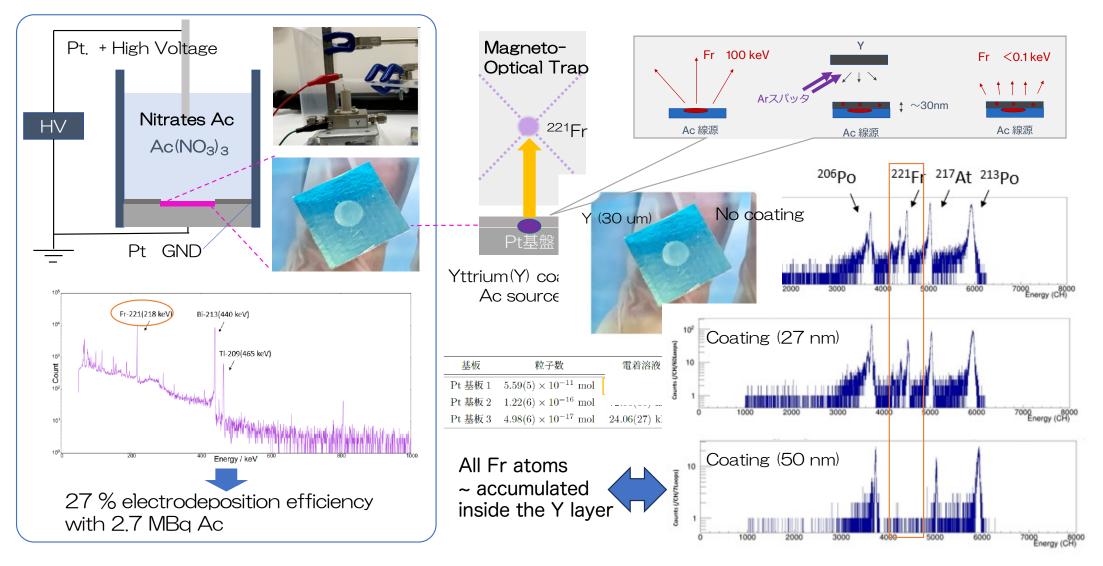


- Radioactive $T_{1/2}$ ~5min
- Make it by generator ²²⁵Ac
- Large Octupole deformation \rightarrow Schiff moment nuclear EDM enhancement

²²⁵Ac as the ²²¹Fr generator

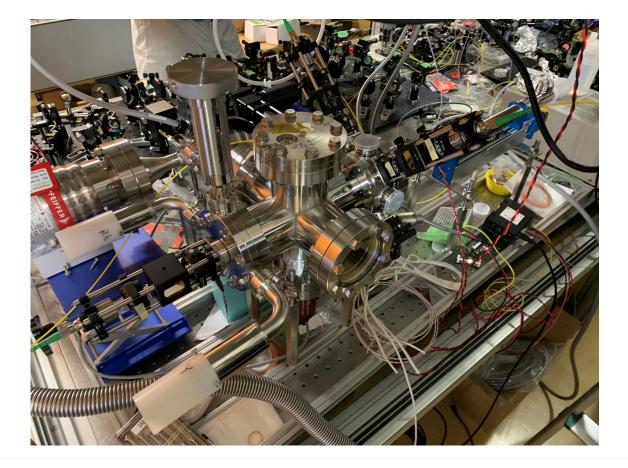


Ac source with molecular plating method



Offline EDM with ²²¹Fr generator

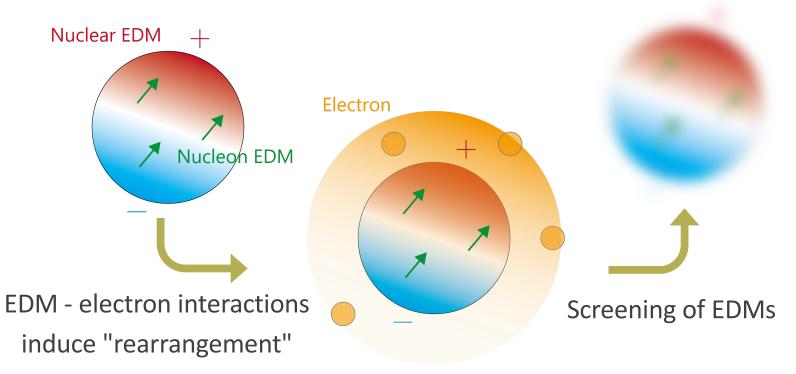
- 1st challenge of MOT with $^{225}Ac/^{221}Fr$ source \sim in progress
- $\bullet\,$ All the apparatus \sim installed in the glove box at the hot lab. In RIKEN





Si detector

Nuclear EDM ~ Shiff moment of ²²¹Fr



2nd & 3rd order nuclear Schiff moment (NSM) operators

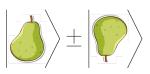
$$S_{2,k} = \frac{1}{6} \sum_{a=1}^{A} d_{a,k} \left(r_a^2 - \left\langle r^2 \right\rangle_{ch} \right) + \frac{2}{15} \sum_{a=1}^{A} d_{a,j} \left(Q_{a,jk} - \left\langle Q_{jk} \right\rangle_{ch} \right)$$
$$S_{3,k} = \frac{e}{10} \sum_{a=1}^{Z} \left[r_a^2 r_{a,k} - \frac{5}{3} r_{a,k} \left\langle r^2 \right\rangle_{ch} - \frac{4}{3} r_{a,j} \left\langle Q_{jk} \right\rangle_{ch} \right]$$
By Dr. K. Yanase

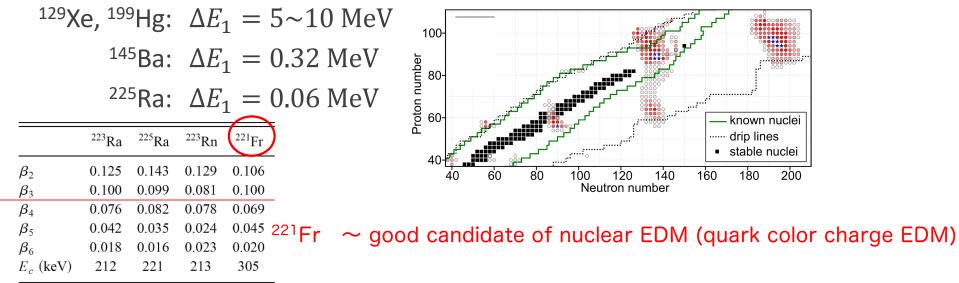
Enhancement of NSM due to octupole correlation

3rd order NSM contains the energy denominator

$$\left\langle \widetilde{\psi}_{g.s.}^{(N)} \big| S_{3z} \big| \widetilde{\psi}_{g.s.}^{(N)} \right\rangle = \sum_{n} \frac{\left\langle \psi_{g.s.}^{(N)} \big| S_{3z} \big| \psi_{n}^{(N)} \right\rangle \left\langle \psi_{n}^{(N)} \big| \widetilde{V} \big| \psi_{g.s.}^{(N)} \right\rangle}{\left[E_{g.s.}^{(N)} - E_{n}^{(N)} \right]} + c.c.$$

Parity doublet caused by nuclear octupole correlations enhances the Schiff moment by orders of magnitude

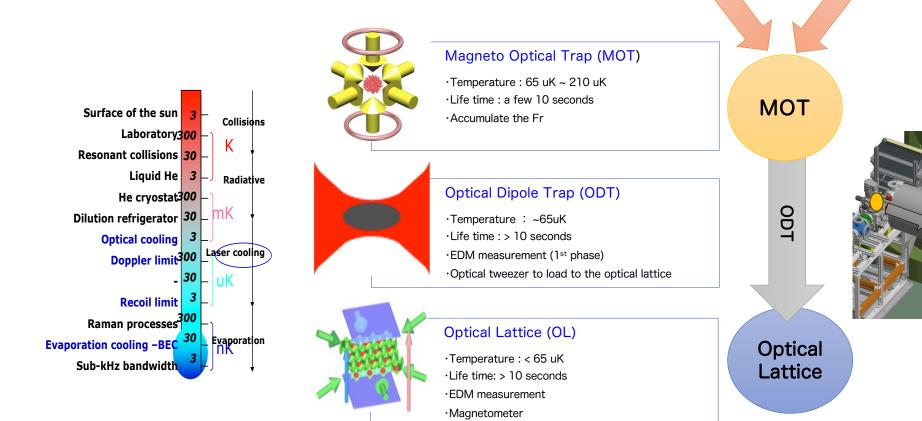




Spevak, V., N. Auerbach, and V. V. Flambaum.. *Physical Review C* 56.3 (1997): 1357.

Cooling Procedure

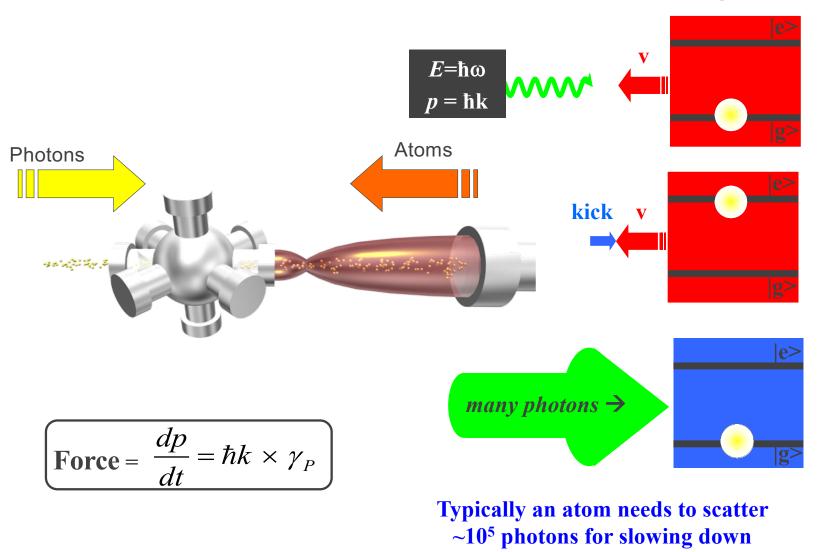




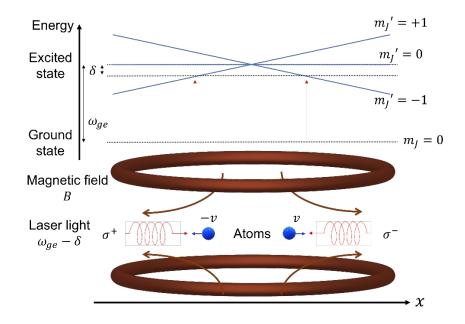
²²¹Fr

Radiation chemistry

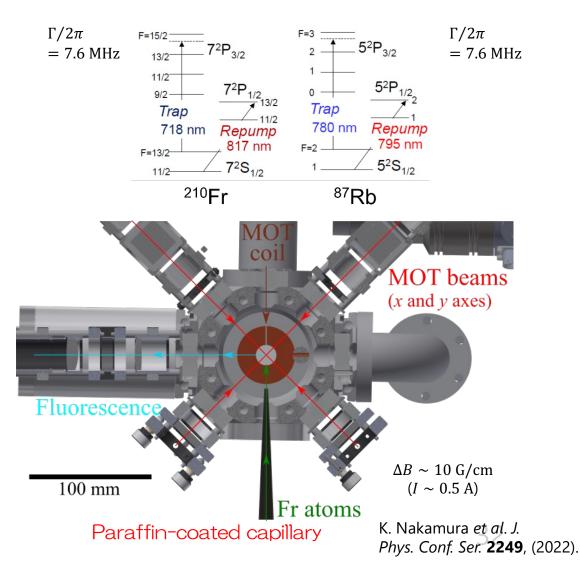
Neutral atom trap \sim Laser cooling



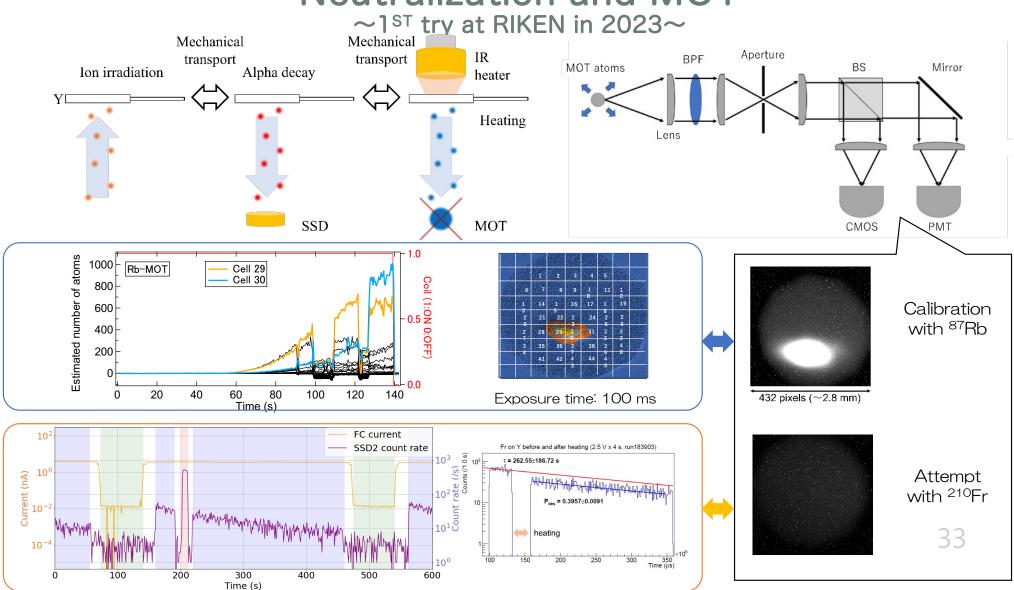
Laser cooling: MAGNETO-OPTICAL TRAP (MOT)



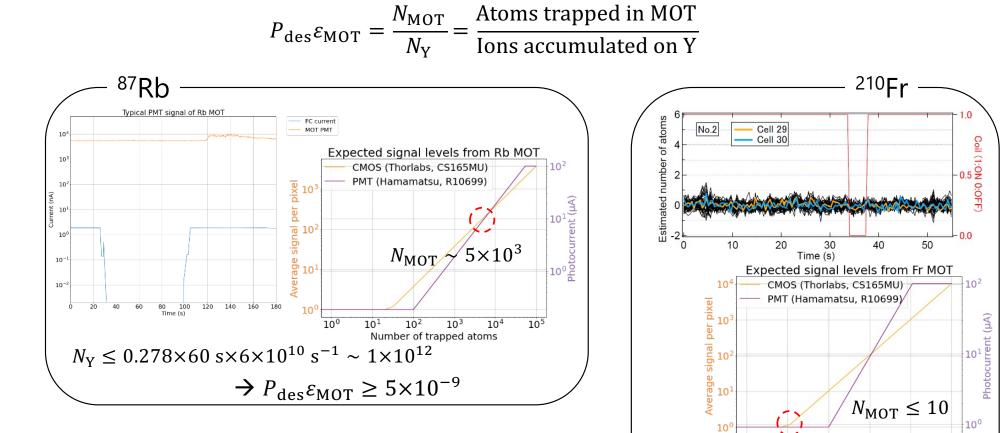
Effective restorative force exerted on atoms by repeated absorption/emission of photons



Neutralization and MOT



Trapping efficiency



10¹

10⁰

10²

 $N_{\rm Y} \sim 275 \, \rm s \times 10^6 \, \rm s^{-1} \times 0.46 \sim 1 \times 10^8$ $\rightarrow P_{\rm des} \varepsilon_{\rm MOT} \leq 8 \times 10^{-8}$

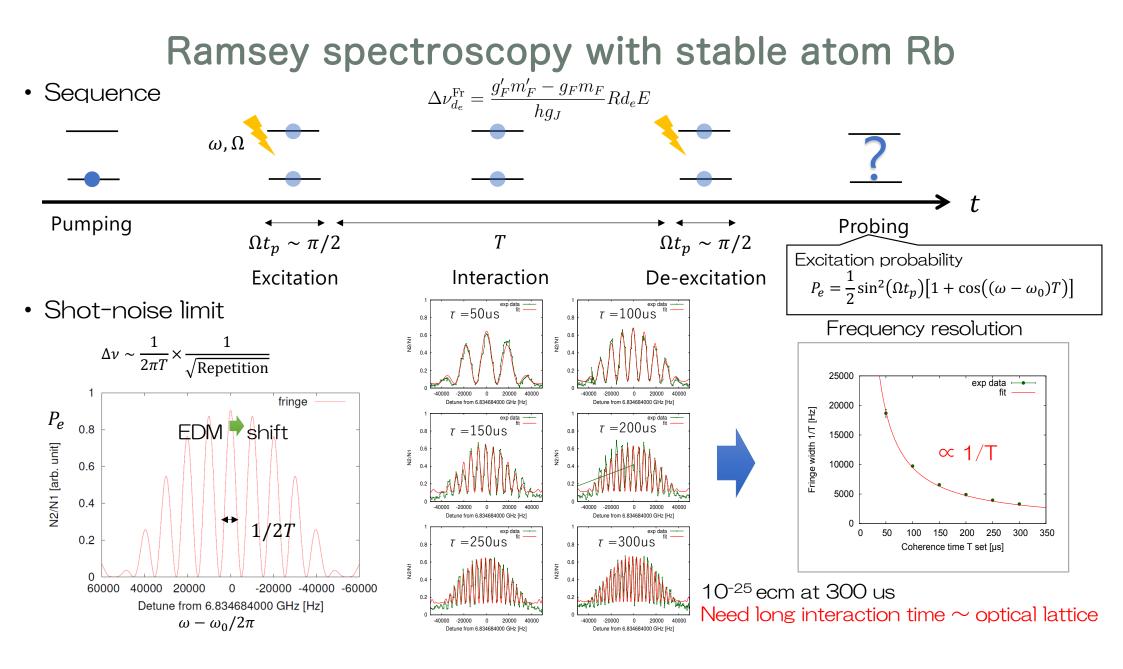
Number of trapped atoms

10³

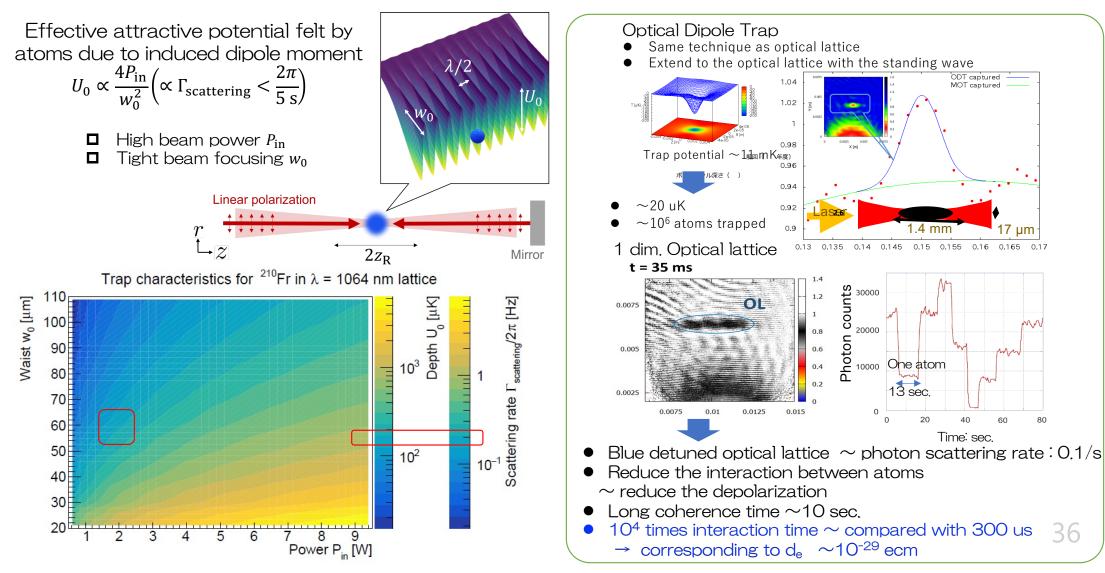
104

105

✓ Proof-of-principle for each component \sim Efficiency still low



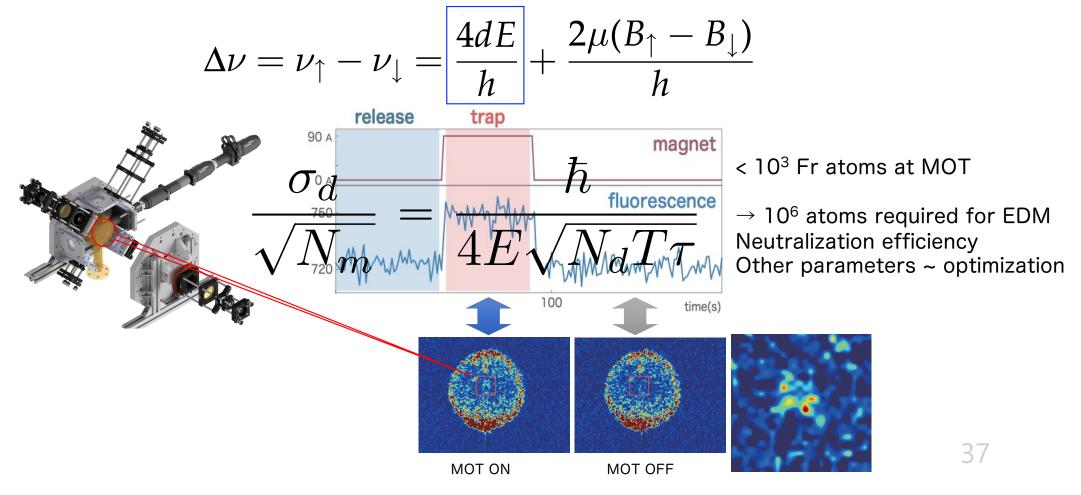
Laser trapping: Optical Lattice



Cold Fr source

Laser cooled Fr apparatus ~ ready

- Laser cooling/trapping techniques for radioactive elements ~ established
- Optimization to increase the number of trapped atoms ~ in progress



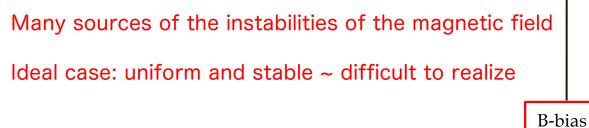
Systematic error and Magnetometer

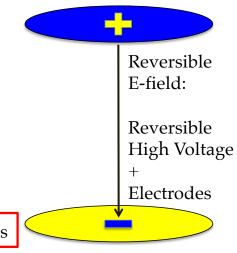
Magnetic Field Instabilities: Annoying

$$\Delta
u =
u_{\uparrow} -
u_{\downarrow} = rac{4dE}{h} + rac{2\mu(B_{\uparrow} - B_{\downarrow})}{h}$$

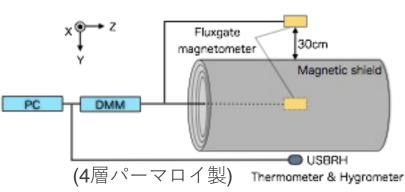
• Quite large contribution EDM false signal

Instabilities adds noise & limits the statistical precision.

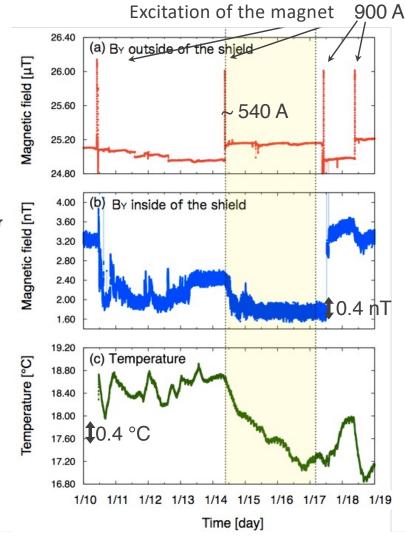


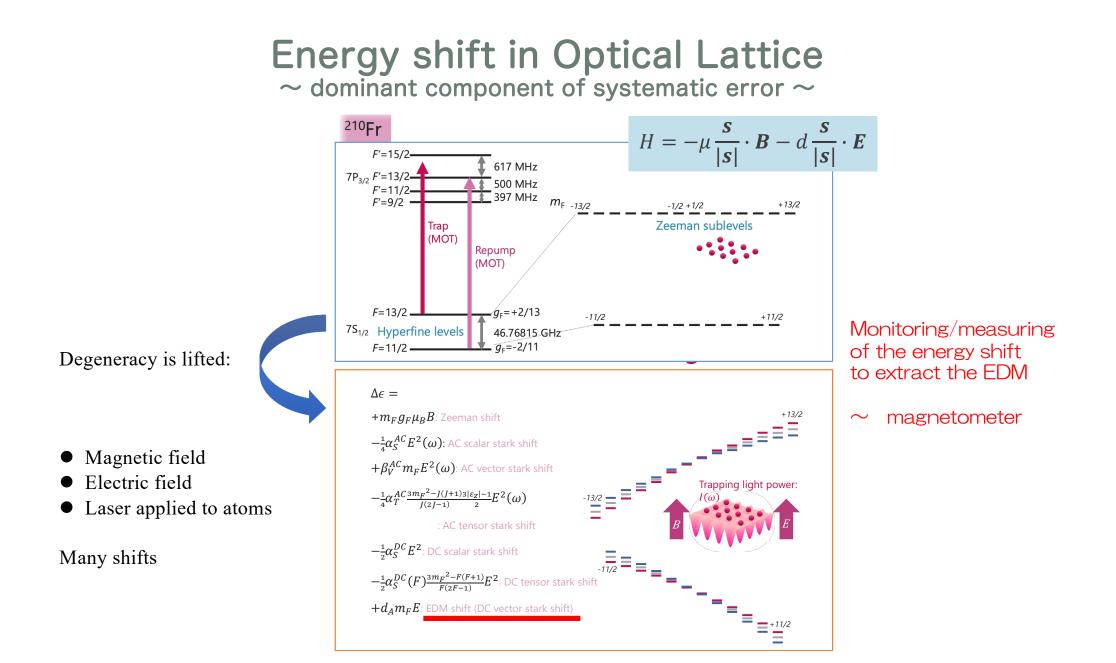


Environmental magnetic field measurement

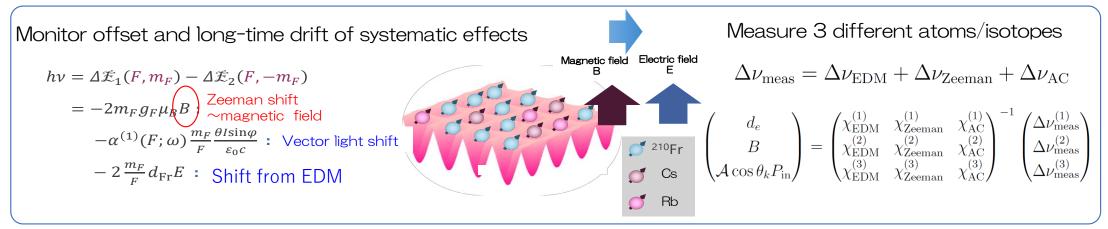


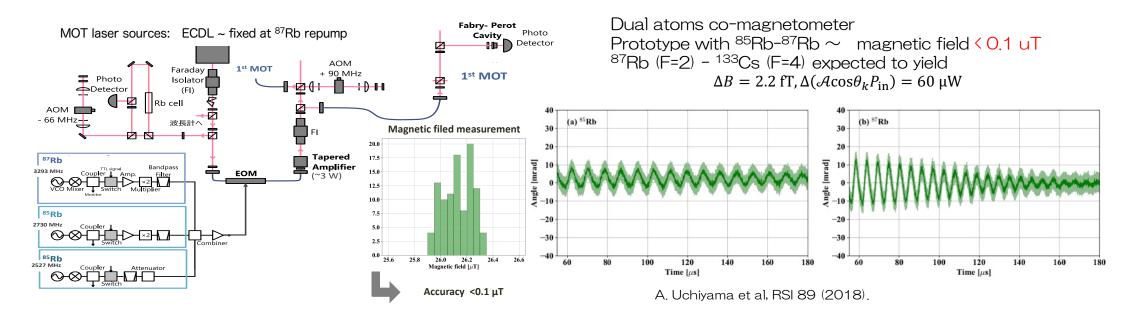
- Magnetic field shielding : 25 uT \rightarrow 2nT
- Magnetic field change : clearly observed
 magnet excitation, metro/car movement…
- Inside the shield : still change \sim 0.4 nT



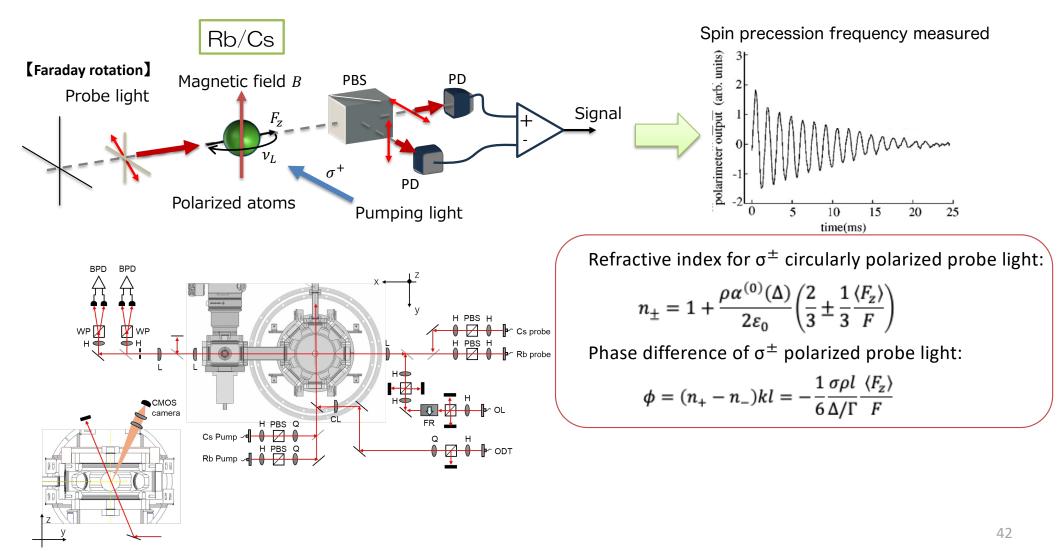


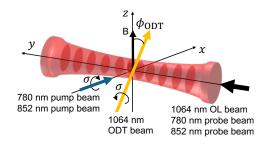
Dual atoms co-magnetometry \sim new idea

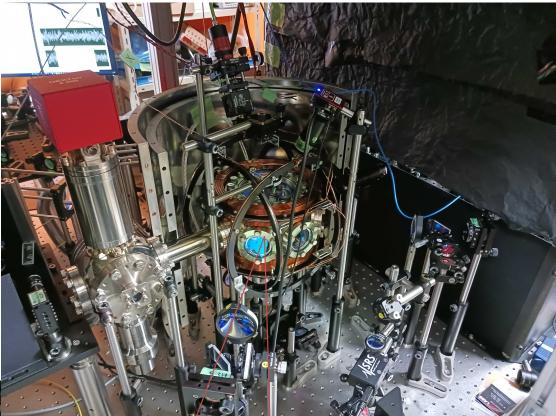


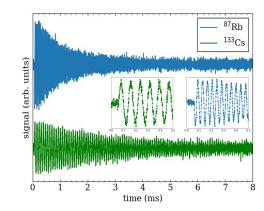


Co-magnetometer setup



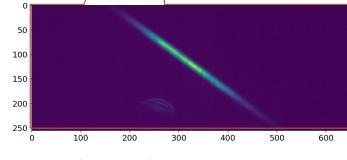




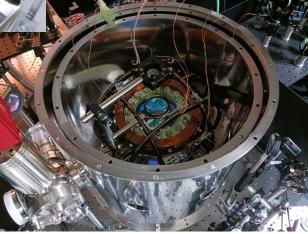


Developed by Dr. Nagase@UTokyo

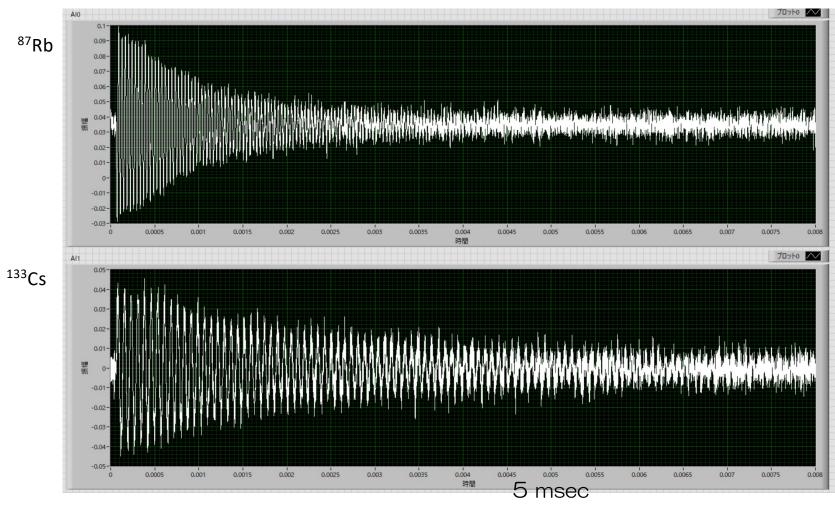




Cold Rb/Cs in the optical lattice

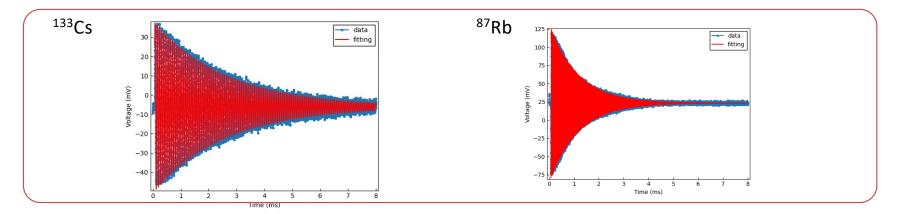


Signal of Faraday rotation

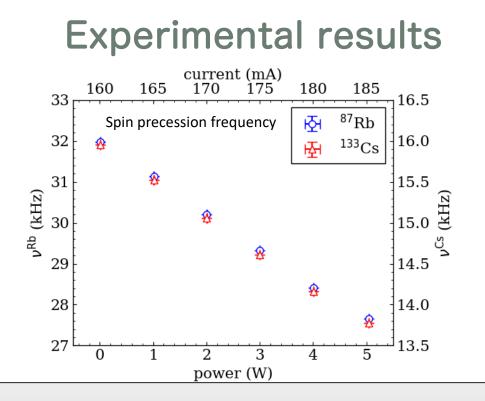


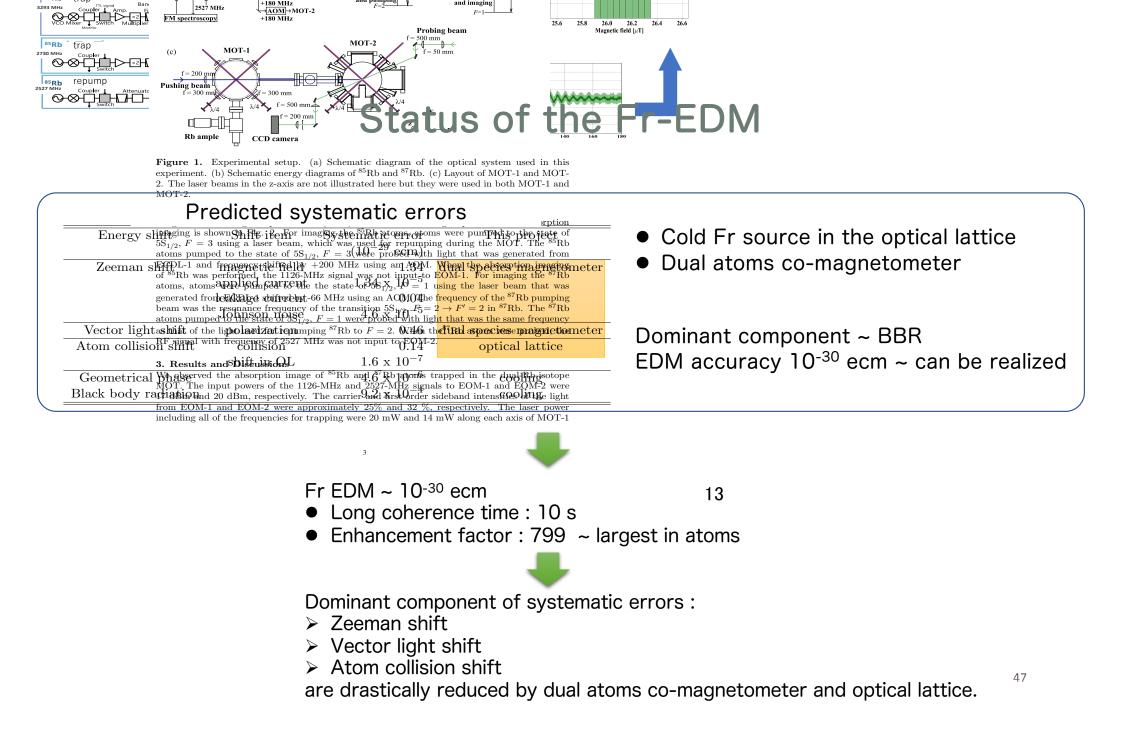
Spin precession \sim observed clearly

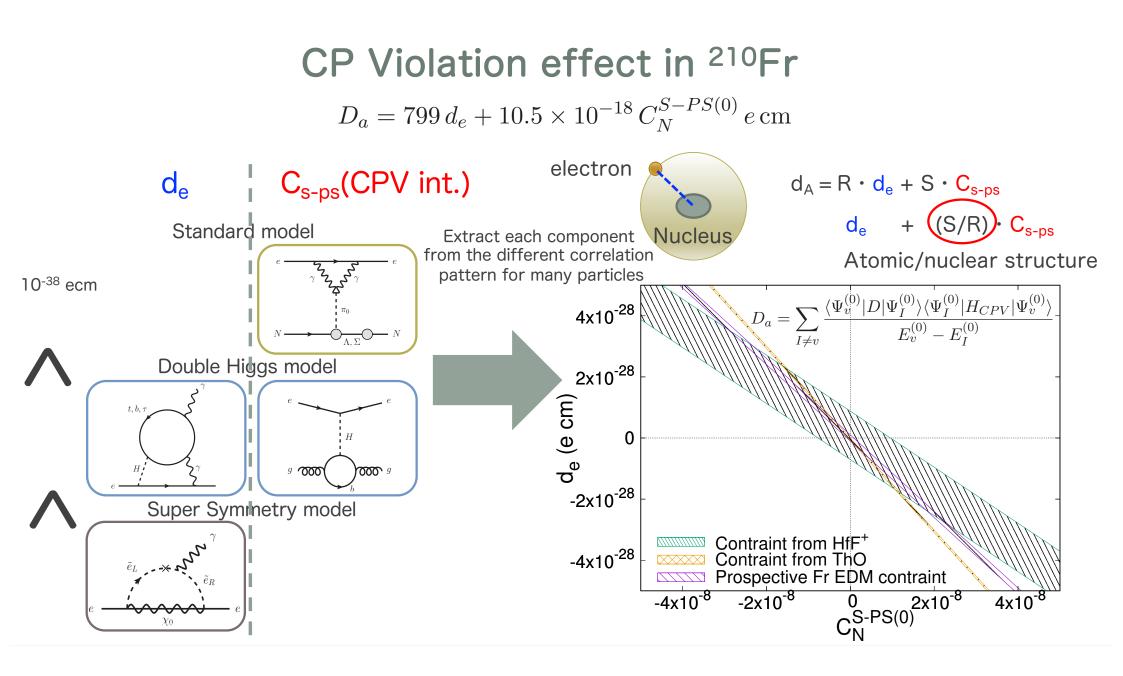
Spin precession frequencies for Rb/Cs



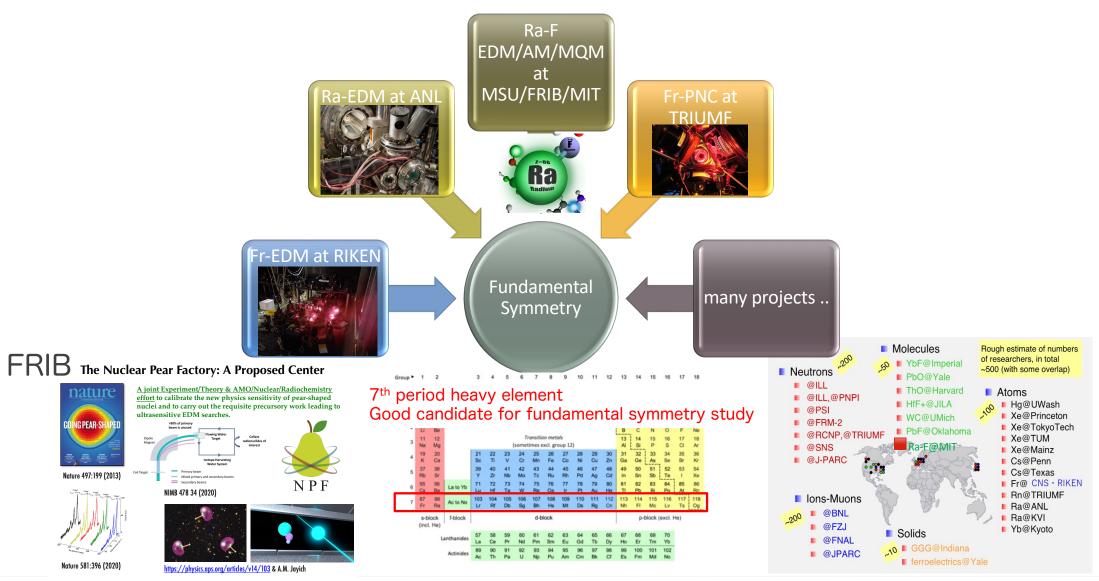
$$\begin{aligned} \nu^{\mathrm{Rb}} &= g_F^{\mathrm{Rb}} \mu_B B / h + \alpha_{\mathrm{Rb}}^{(1)} (F = 2; \Delta) \frac{1}{h\epsilon_0 c} I \varepsilon_{cp} \cos \phi \\ \nu^{\mathrm{Cs}} &= g_F^{\mathrm{Cs}} \mu_B B / h + \alpha_{\mathrm{Cs}}^{(1)} (F = 4; \Delta) \frac{1}{2h\epsilon_0 c} I \varepsilon_{cp} \cos \phi \\ & \swarrow \\ B &= \frac{\Gamma_{\mathrm{Cs}} \nu^{\mathrm{Rb}} - \Gamma_{\mathrm{Rb}} \nu^{\mathrm{Cs}}}{\gamma_{\mathrm{Rb}} \Gamma_{\mathrm{Cs}} - \gamma_{\mathrm{Cs}} \Gamma_{\mathrm{Rb}}}, \qquad \delta B = \frac{\sqrt{\Gamma_{\mathrm{Cs}}^2 (\delta \nu^{\mathrm{Rb}})^2 + \Gamma_{\mathrm{Rb}}^2 (\delta \nu^{\mathrm{Cs}})^2}}{|\gamma_{\mathrm{Rb}} \Gamma_{\mathrm{Cs}} - \gamma_{\mathrm{Cs}} \Gamma_{\mathrm{Rb}}|} \\ I \varepsilon_{\mathrm{cp}} \cos \phi &= \frac{-\gamma_{\mathrm{Cs}} \nu^{\mathrm{Rb}} + \gamma_{\mathrm{Rb}} \nu^{\mathrm{Cs}}}{\gamma_{\mathrm{Rb}} \Gamma_{\mathrm{Cs}} - \gamma_{\mathrm{Cs}} \Gamma_{\mathrm{Rb}}}, \qquad \delta (I \varepsilon_{\mathrm{cp}} \cos \phi) = \frac{\sqrt{\gamma_{\mathrm{Cs}}^2 (\delta \nu^{\mathrm{Rb}})^2 + \gamma_{\mathrm{Rb}}^2 (\delta \nu^{\mathrm{Cs}})^2}}{|\gamma_{\mathrm{Rb}} \Gamma_{\mathrm{Cs}} - \gamma_{\mathrm{Cs}} \Gamma_{\mathrm{Rb}}|} \end{aligned}$$

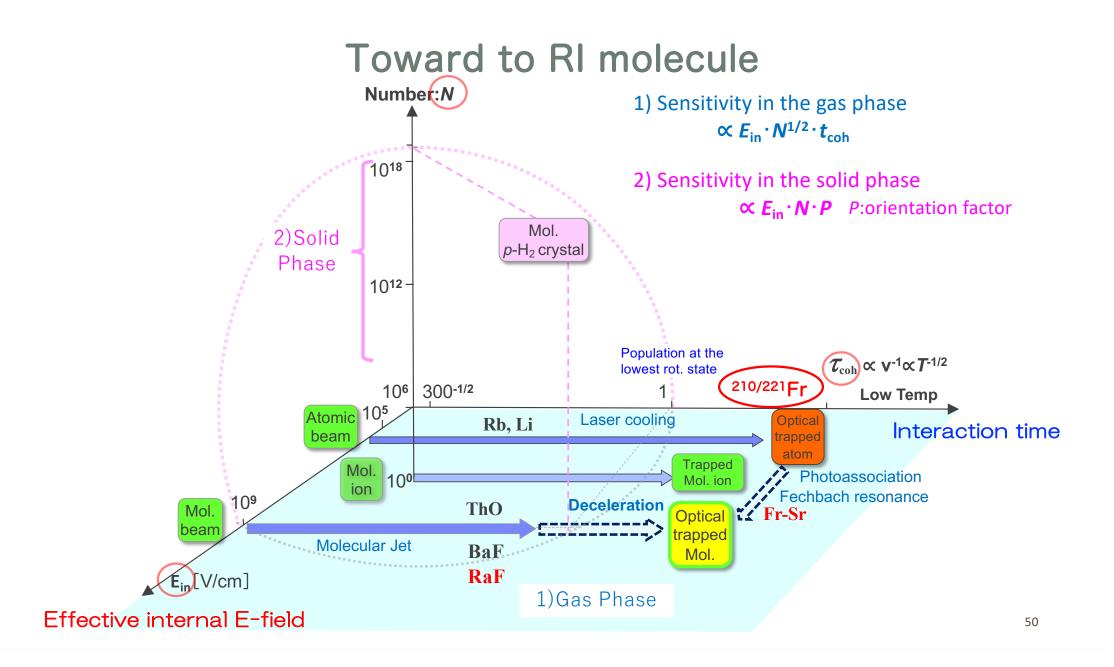






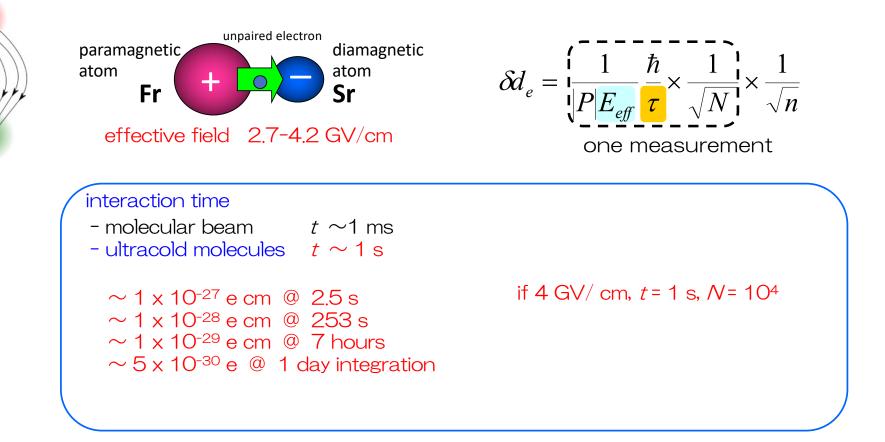
Fundamental physics with radioactive elements



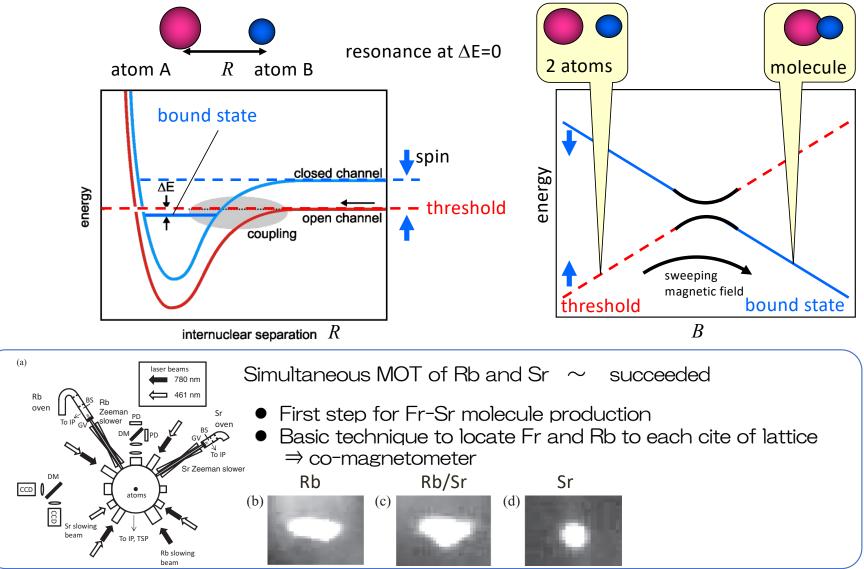


Fr-Sr molecule

Ultracold trapped molecules associated by Feshbach resonance from laser cooled Fr and Sr atoms.



Association of molecules near Feshbach resonance



Possibility of superheavy molecules: LrX

PHYSICAL REVIEW A 104, 062801 (2021)

Towards CP-violation studies on superheavy molecules: Theoretical and experimental perspectives

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Molecules containing superheavy atoms can be artificially created to serve as sensitive probes to study symmetry-violating phenomena. Here, we provide detailed theoretical studies of quantities relevant to the electron electric dipole moment (eEDM) and nucleus-electron scalar-pseudoscalar interactions in diatomic molecules containing superheavy lawrencium nuclei. The sensitivity to parity and time (or, equivalently, CP) reversal violating properties is studied for different neutral and ionic molecules. The effective electric fields in these systems are found to be about 3-4 times larger than other known molecules on which eEDM experiments are being performed. Similarly, these superheavy molecules exhibit an enhancement of more than 3 times for CP-violating scalar-pseudoscalar nucleus-electron interactions. Our preliminary analysis using the Woods-Saxon nuclear model also demonstrates that these results are sensitive to the diffuse surface interactions inside the Lt nucleus. We also briefly comment on some experimental aspects by discussing the production of these systems.



Molecule	$\mathcal{E}_{\mathrm{eff}}$ (GV/cm)	
LrO	258.92 250.21	Large
LrF ⁺ LrH ⁺	246.5 246.31 343.38	Effective Field
ThO	87	
HgF	115.42	
HfF^+	22.5	
YbF	23.2	53

Lr molecules production yield

	Reaction	σ _{prod} / nb	beam current	Transport +Mass sep. eff.	Beam intensity after Mass sep.(ion/min)
@JAEA	$^{249}Cf + {}^{11}B$	122	200 pnA	~10 %	2
@RIKEN GARIS	²⁰⁹ Bi + ⁴⁸ Ca	60	2000 pnA	~30%	30
@RIKEN ISOL*	²⁰⁹ Bi + ⁴⁸ Ca	60	2000 pnA	~10%	10

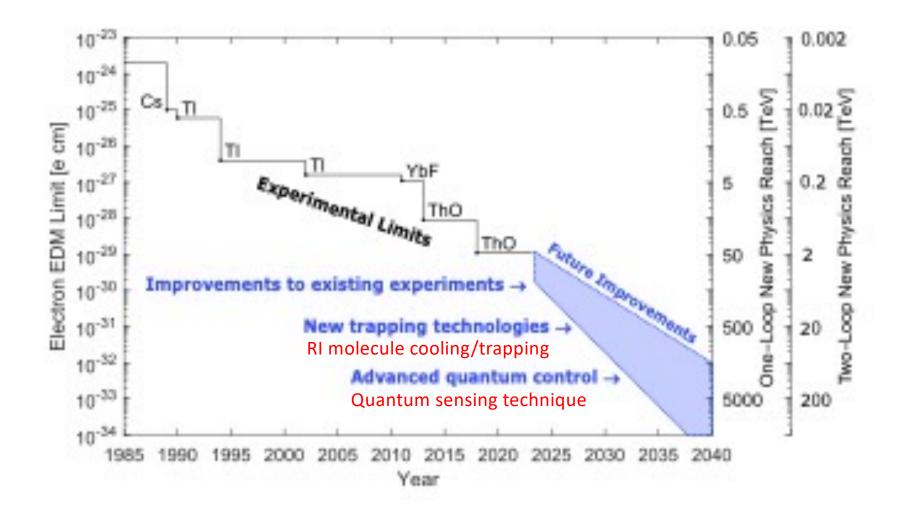
* Assuming a gas-jet ISOL (similar to JAEA) is available at RIKEN RILAC

Production of Lr molecules

- LrO+ : Lr + trace amount O_2 (contaminants in a carrier gas or additional O_2)
- $LrF+: LnF_x$ ion beam has been produced @ CERN-ISOLDE in a molten metal target. In the case of Lr, fluorination gas (ex. CF_4) needs to be introduced.

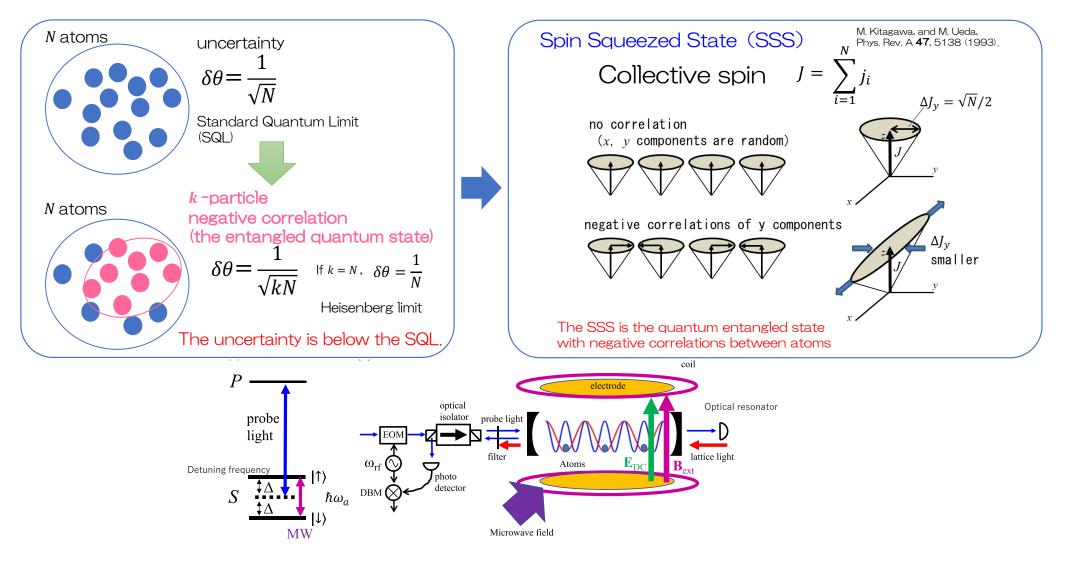
New idea of the quantum sensing is required for the limited number of the particles .

Future strategy



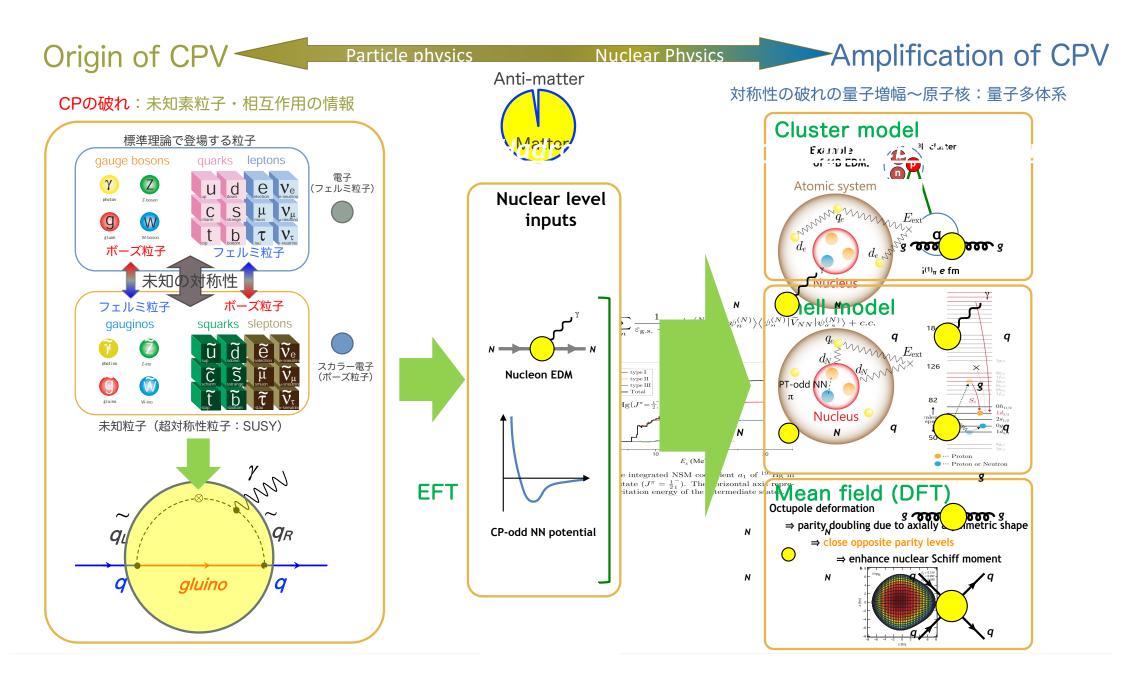
Quantum sensing of EDM with entangled Fr atoms

 \sim Uncertainty of the phase of Ramsey resonance with Squeezed Spin State \sim

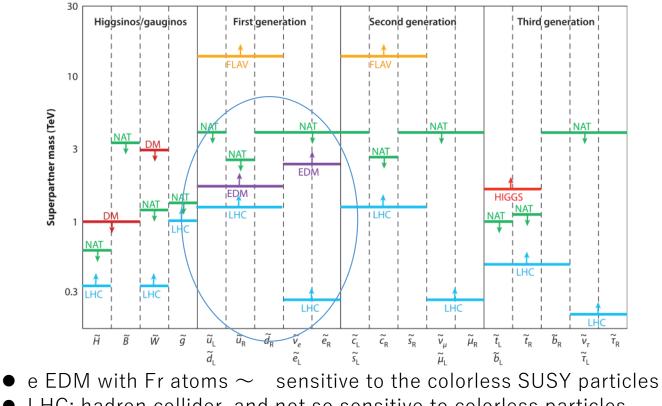


EDM detection scheme and sensitivity with SSS $F = \frac{13}{2}$ $M_F = -\frac{11}{r}$ Ramsey resonance Quantum sensing $7P_{3/2}$ probe ΔN light $M_F = +\frac{1}{2}$ $\Delta \theta$ $F = \frac{11}{2}$ uantum 11 11 squeezing $\frac{11}{2}M_F$ detection sensing F = 13 2τ 2τ $\frac{\pi}{2}$ $\frac{\pi}{\pi}$ $\frac{\pi}{2}$ $7S_{1/2}$ MW π π MW v V х х probe F = 11light $\frac{11}{2}$ $\frac{11}{2}$ 10-27 Quantum limit below SQL 10-28 ·g/2π` $\delta d_{\rm e} = 3.3 \times 10^{-29} \, e {\rm cm}$ $\delta d_{\rm e}^{\rm SSS} = \frac{13\hbar}{22RE_{\rm DC}} \frac{\xi_{\rm R}}{\sqrt{NTt_{\rm total}}}$ $\delta d_{e} (e \text{cm})$ 10^{3} - SQ 10⁻²⁹ 10^{4} $E = 100 \text{ kV cm}^{-1}$ T = 1 s 10^{5} *t*_{total} = 24h = 86400 s *N* = 2.5x10⁵ 10⁻³⁰ 106 $g/2p = 10^5 \,\text{Hz}$ *x*_B = 1 / 44.2 10⁻³¹ 10⁵ 10^{4} 106 108 107 $\delta d_{\rm e} = 7.5 \times 10^{-31} \, e {\rm cm}$ Ν

Quantum Sci. Technol. 6, 044008 (2021).



Explored energy range

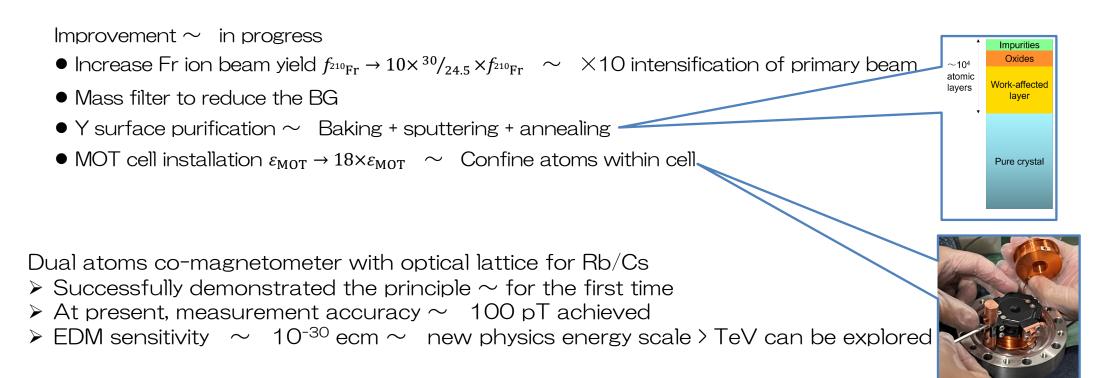


- LHC: hadron collider, and not so sensitive to colorless particles
- Fr EDM \sim can explore the mass scale > TeV region : 10^{-30} ecm

Summary

Optical lattice (OL) atomic interferometer development \sim in progress to search for e-EDM \checkmark Fr production \sim ready : extraction efficiency close to maximum

 \checkmark OL and co-magnetometer : proof-of-principle with Rb/Cs atoms \sim ready



Thank you very much !

Collaborators

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