Right-handed weak currents in neutrinoless double-beta decays and ton scale double-beta detectors part II

Toru Sato

RCNP, Osaka University

arXiv:2501.0345v1 H. Ejiri, T. Fukuyama, TS

Objective:

- ${\rm \bullet}\,$ transparent derivation of the effective nuclear transition operators for $0^+ \rightarrow 2^+$ DBD
- meson-exchange current of DBD

In this talk:

• revisit sensitivity of η - and λ -terms in $0^+ \rightarrow 0^+/2^+ \ 0\nu$ DBD

$$[t_{1/2}^{0\nu}]^{-1} = C_{\lambda\lambda}^J < \lambda >^2 + C_{\eta\eta}^J < \eta >^2 + \ldots$$

$$C_{\eta\eta}^{(0)}/C_{\lambda\lambda}^{(0)} \sim 10^3 - 10^4 \quad C_{\eta\eta}^{(2)}/C_{\lambda\lambda}^{(2)} \sim 10 - 10^2$$

J. Suhonen, O.Civitarese, Phys.Rep. 300(1998)123, D Fang, A. Faessler PRC 107(2023) 015501

$$0^+
ightarrow 0^+$$
 case: $\chi^2_R G_{09} \sim g_A \mu_V C^2_{ss}$

• $N \leftrightarrow \Delta$ DBD (neutrino exchange inside baryon) in $0^+ \rightarrow 2^+$ DBD $N\Delta \ 0\nu$ DBD Interference between L-handed and R-handed terms \rightarrow p-wave neutrino propagator

$$P_L(\gamma_0\omega - \boldsymbol{\gamma}\cdot\boldsymbol{k} + m_i)P_R \to \boldsymbol{\gamma}\cdot\boldsymbol{r}$$

possibility to have positive parity nuclear effective operator for $0^+ \rightarrow 0^+/2^+$ transition

(1) One of the electrons is p-wave \rightarrow Fermi(V - V) and GT A - A

$$0^+ \to 0^+ : (s_{1/2}p_{1/2}) \qquad 0^+ \to 2^+ : (s_{1/2}p_{3/2})$$

(2) Both electrons are s-wave $\rightarrow V - A$ interference term

$$0^+ \to 0^+ : (s_{1/2}s_{1/2}) \qquad 0^+ \to 2^+ : X$$

(2) is the source of large sensitivity to η term, which does not exist for $0^+ - 2^+$ transition,

Hadron current

$$<\eta>J_LJ_L+<\lambda>J_LJ_R\rightarrow<\eta>(V-A)(V-A)+<\lambda>(V-A)(V+A)$$

transition amplitude of DBD

$$\mathcal{M}_{fi} \sim \int d\boldsymbol{x} d\boldsymbol{y} \frac{1}{|\boldsymbol{x} - \boldsymbol{y}|} \bar{e}_2(\boldsymbol{y}) \mathcal{O} e_1^c(\boldsymbol{x})$$

 \mathcal{O} : Difference between η and λ terms.

$$<\eta>[(\bigvee(y)(r\cdot\gamma)\bigvee(x)+\bigwedge(y)(r\cdot\gamma)\bigwedge(x))-(\bigvee(y)(r\cdot\gamma)\bigwedge(x)+\bigwedge(y)(r\cdot\gamma)\bigvee(x))]\\+<\lambda>[(\bigvee(y)(r\cdot\gamma)\bigvee(x)-\bigwedge(y)(r\cdot\gamma)\bigwedge(x))-(\bigvee(y)(r\cdot\gamma)\bigwedge(x)-\bigwedge(y)(r\cdot\gamma)\bigvee(x))\gamma_5].$$

• $g_A \mu_V$ does not contribute to λ

•
$$s^2$$
, $VA \leftrightarrow sp$, $V_0^2 A^2$: $\frac{4.7 \times q_\nu}{2M_N} \leftrightarrow m_e R, V_C R$

• $0^+ - 0^+$ transition is sensitive to η : Right-handed current and V - A structure



SU(6)/static quark model

$$\Delta \sim \chi^S \eta^S ~~ N \sim (\chi^{MS} \eta^{MS} + \chi^{MA} \eta^{MA})/\sqrt{2}$$

• among 2^+ transition operators, spin-spin interaction contributes

$$<\Delta^{++}||\sum_{i,j=1}^{3}\frac{2}{3}[\boldsymbol{\sigma}_{i}\otimes\boldsymbol{\sigma}_{j}]^{(2)}h(r_{ij})\tau_{i}^{+}\tau_{j}^{+}||n>=-\frac{8\sqrt{10}}{3}<\frac{1}{r}>$$

6/9

 Δ component was calculated by π and ρ exchange model (T. Tomoda, NPA484 (1988) 635)



- Meson-exchange current for 0ν DBD
- contributes to $0^+ 2^+$ transition

Possible estimation of MEC: ratio of the matrix elements of IA and MEC

$$\sum_{m,n=1}^{A} \tau_m^+ \tau_n^+ \boldsymbol{\sigma}_m \cdot \boldsymbol{\sigma}_n [\hat{r}_{mn} \otimes \hat{r}_{mn}]^{(2)} h(r_{mn}).$$

h(r): 1/r for IA, $Y_2(m_\pi r)$ for MEC

From table of Tomoda(NPA484 635(1988) DBD of ⁷⁶Ge)

$$\frac{MEC}{IA} \sim 0.1 \propto \frac{f_{\pi N\Delta} f_{\pi NN}}{(m_{\Delta} - m_N)g_A^2} < \frac{1}{r} >$$

- $0^+ 0^+$ transition is sensitive to η -term than λ -term due to R-handed current and V A/V + A structure of weak currents.
- $\sigma \cdot \sigma Y_2$ type operator can be useful measure of MEC.