

AMD triple-Sによる水素と ヘリウムアイソトープの分析

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AMD triple-S: AMD Superposition of Selected Snapshot

主な課題

**クラスターの取り扱いによる $A=20$ 領域
までの軽い原子核の系統的分析**

$A=20$ 領域までにある原子核を全て調べる (つもり)。

クラスター (広義) : α 、 t 、 h 、 d 、 $2n$ 、 $2p$

**軽い領域の不安定核において定量性を
持った系統的理解 (ハロー、共鳴、クラ
スター構造等) を目指す。**

「核力と核構造」の理解

Present Method (AMD triple-S)

In order to treat the spatial extension, we use

AMD+GCM

Enyo *et al.*

N. Itagaki and S. Aoyama, Phys. Rev. C61, 024303 (2000)

In order to treat the large model space, we use

SVM

Kukulin

K. Varga, Y. Suzuki, Y. Ohbayashi, Phys. Rev. C50, 189 (1994)

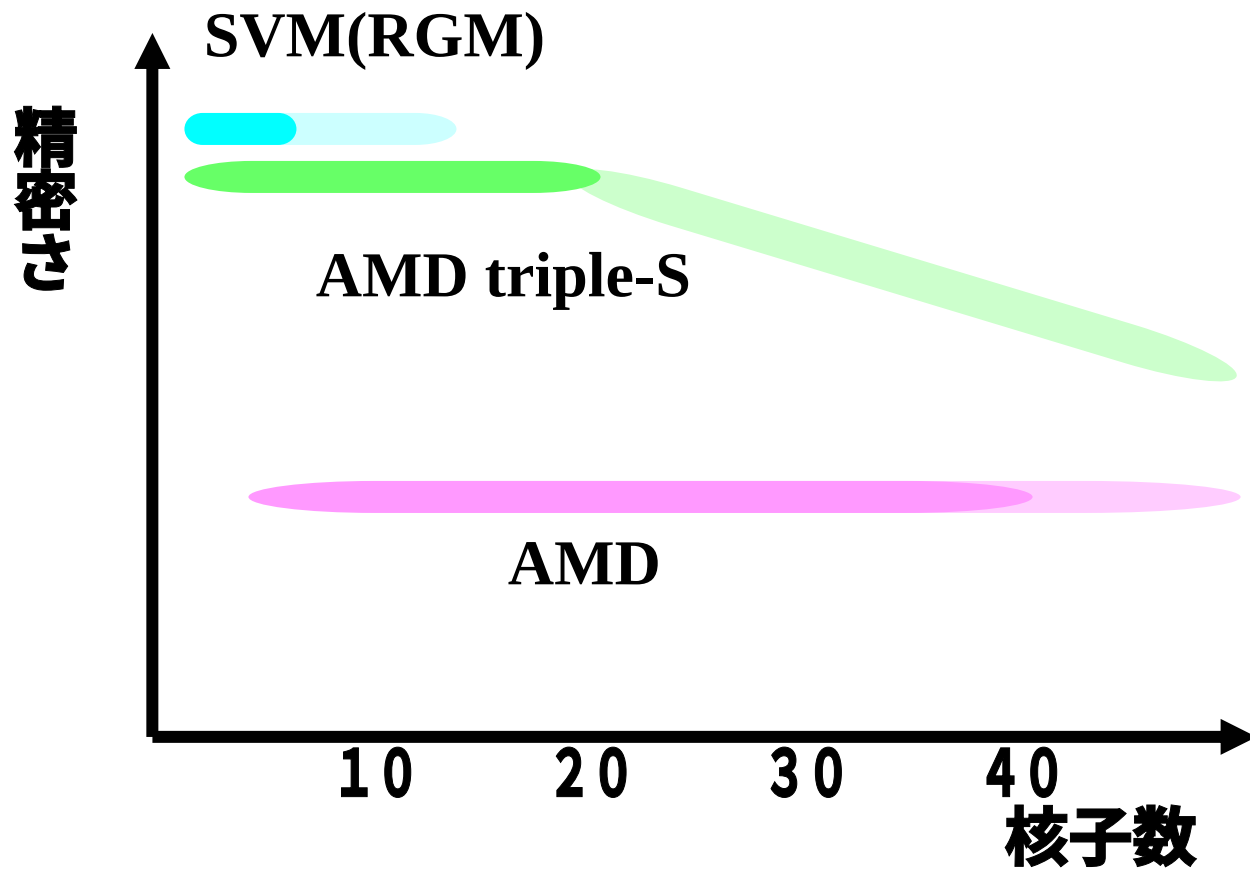
In order to solve the many body resonance, we use

ACCC

Kukulin

S. Aoyama, Phys. Rev. Lett. 89, 052501 (2002)

AMD+GCM+SVM (+ACCC)



1. The Gaussian center ($\mathbf{z}$) of the AMD w.f. is **randomly generated.**

AMD w.f.

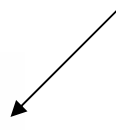
$$\Psi_k = \mathcal{A}[(\psi_1 \chi_1)(\psi_2 \chi_2) \cdots]_k.$$
$$\psi_i = \left(\frac{2\nu}{\pi}\right)^{\frac{3}{4}} \exp[-\nu(\vec{r} - \vec{z}_i/\sqrt{\nu})^2],$$

2. We solve **the frictional cooling equation only for imaginary part.**

$$\frac{dE}{d\tau} = \sum_i^A \frac{\partial E}{\partial \vec{z}_i} \cdot \frac{d\vec{z}_i}{d\tau} + \sum_i^A \frac{\partial E}{\partial \vec{z}_i^*} \cdot \frac{d\vec{z}_i^*}{d\tau},$$

$$\frac{d\vec{z}_i}{d\tau} = -Im\left[\frac{\partial E}{\partial \vec{z}_i^*}\right], \quad \frac{d\vec{z}_i^*}{d\tau} = -Im\left[\frac{\partial E}{\partial \vec{z}_i}\right].$$

3. We regard it as a basis function for the **GCM.**

$$\Phi = \sum_k c_k P_{MK}^J \Psi_k$$


4. If the obtained energy **decrease, we adopt it.**

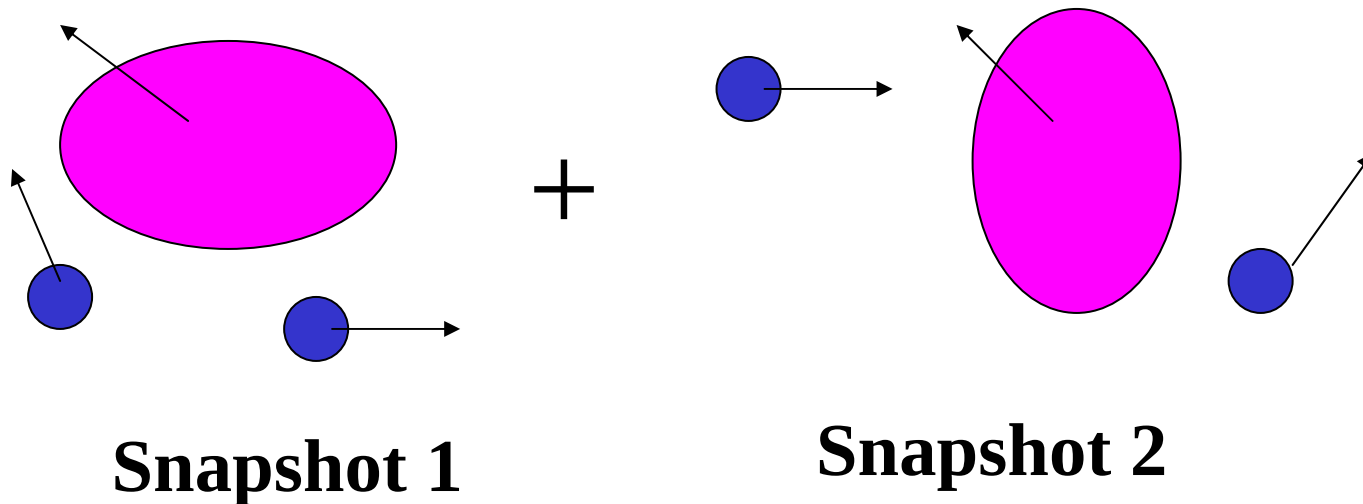
$$\sum_{k=1}^N |E^k(j) - E^k(j-1)| > \epsilon,$$

e.g. N=3, ϵ =0.05 MeV

5. We return 1, if the energy does not converge.

AMD triple-S基底関数の解釈

核子 (クラスター) の運動量期待値 $= \langle p_i \rangle \neq 0$



AMD triple-S = AMD Superposition Selected Snapshot

Hamiltonian

$$\hat{H} = \sum_{i=1}^A \hat{t}_i - \hat{T}_{c.m.} + \sum_{i>j}^A \hat{v}_{ij},$$

central Volkov No.2

$$V(r) = (W - MP^\sigma P^\tau + BP^\sigma - HP^\tau) \\ \times (V_1 \exp(-r^2/c_1^2) + V_2 \exp(-r^2/c_2^2)),$$

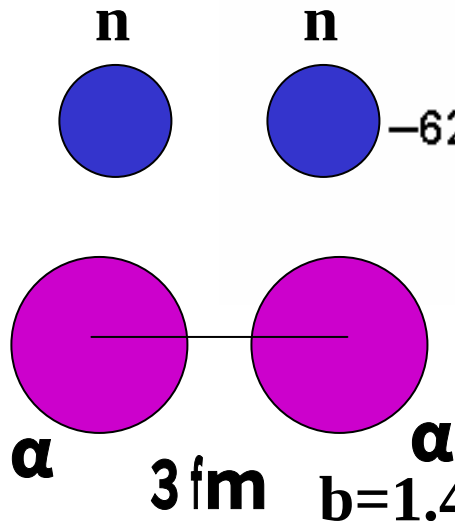
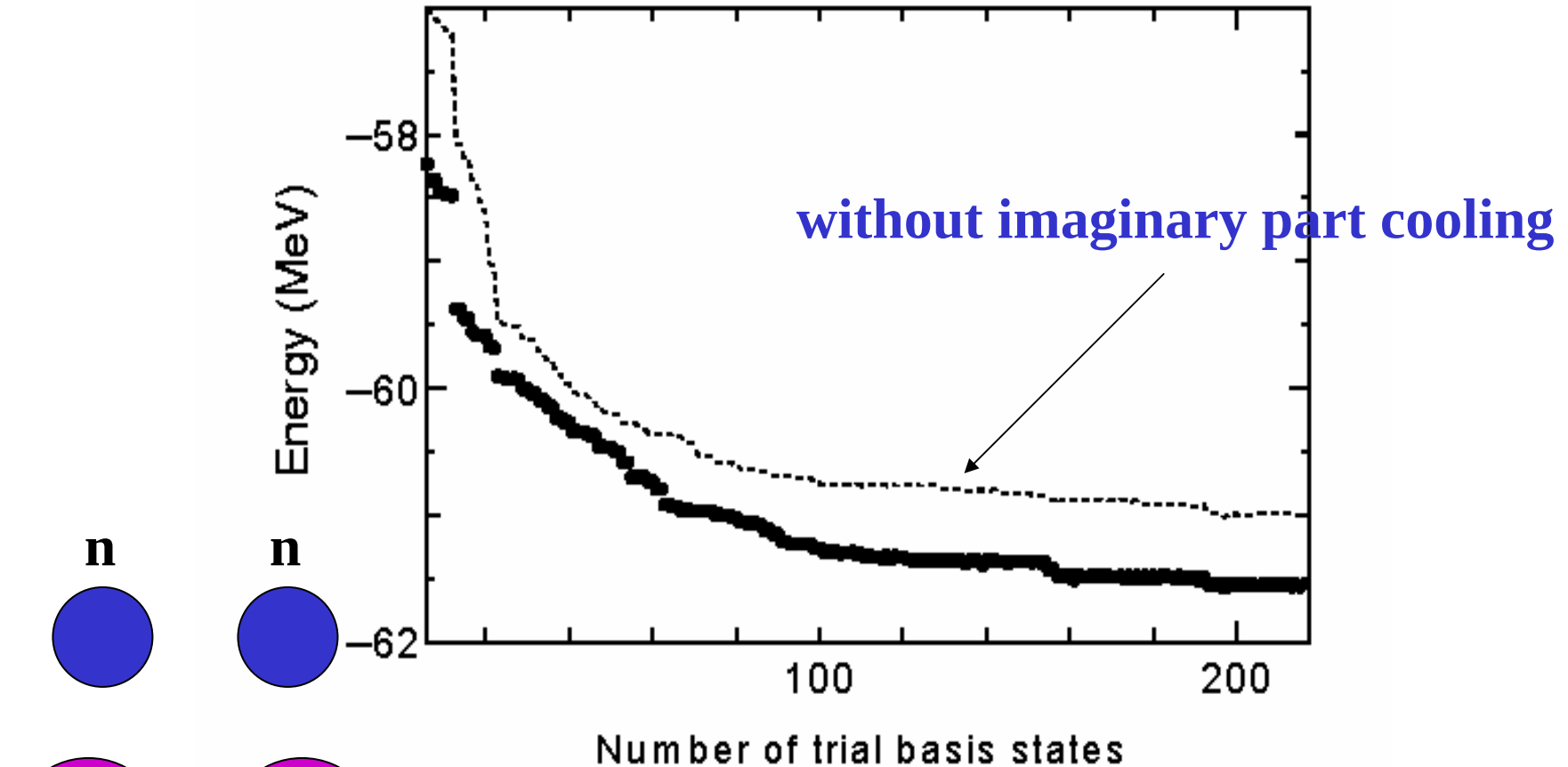
$$W = 1 - M, \quad M = 0.60 \text{ and } B = H = 0.125$$

Spin-orbit G3RS

$$V_{ls} = V_0 \{e^{-d_1 r^2} - e^{-d_2 r^2}\} P(^3O) \vec{L} \cdot \vec{S},$$

$$d_1 = 5.0 \text{ fm}^{-2}, \quad d_2 = 2.778 \text{ fm}^{-2}, \quad V_0 = 2000 \text{ MeV}.$$

The energy convergence of the ground state(0+) of ^{10}Be



Volkov No.2($M=0.6$, $B=H=0.125$) + G3RS

⁶He

The energy of single AMD calculation is 6.74 MeV higher than present one.

	S_{2n} (MeV)	$E_x(2_1^+)$ (MeV)	r.m.s. radius (fm)
single AMD	-5.75	1.95	2.02
AMD+GCM	-0.73	2.45	2.28
AMD triple-S	+0.99	1.89	2.37
Exp.	+0.98	1.80 ^a	2.33 ± 0.04 ^b
			2.48 ± 0.03 ^c
			2.57 ± 0.1 ^d

We can not describe the halo structure.

A-body Hamiltonian

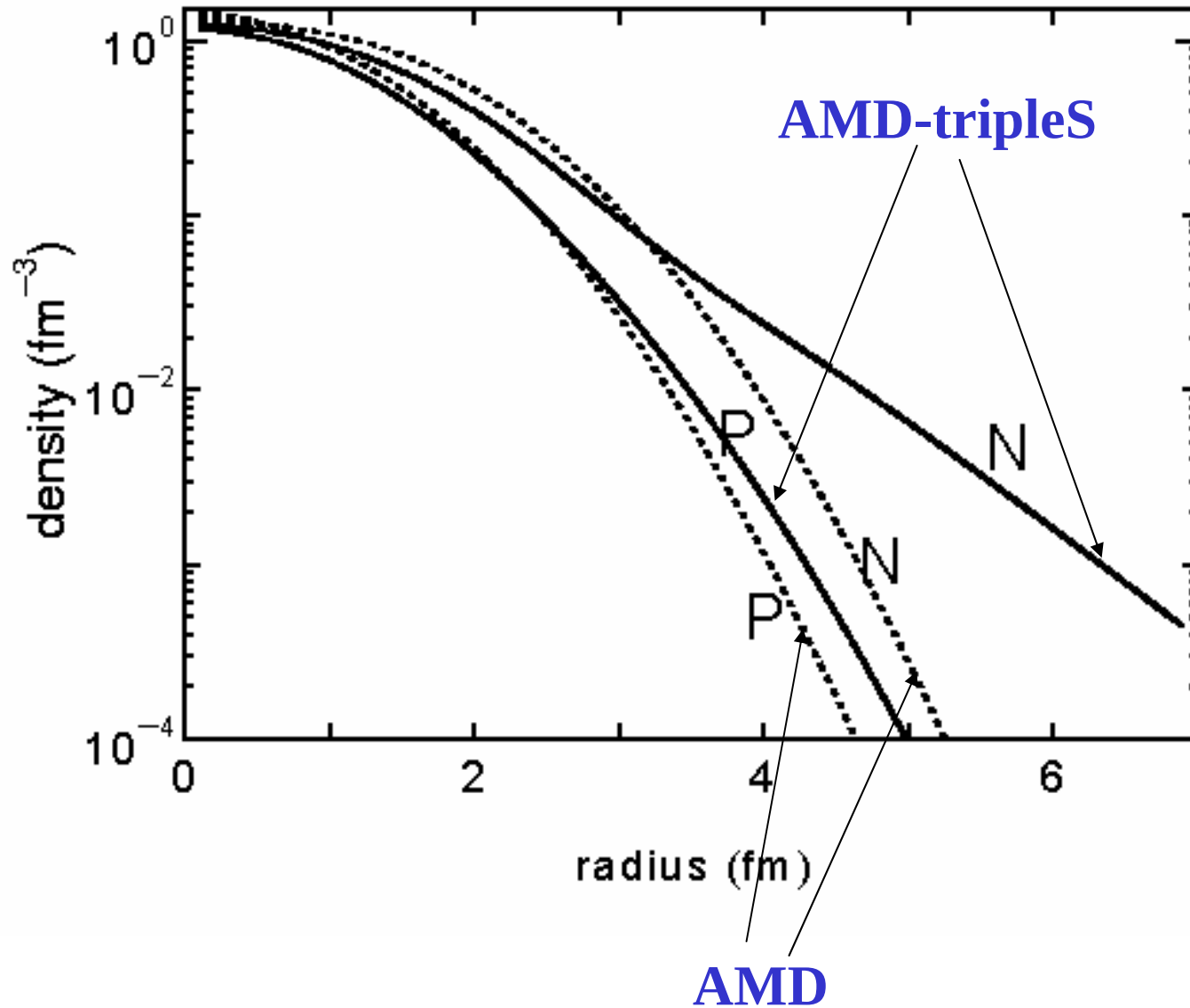
$$\hat{H} = \sum_{i=1}^A \hat{t}_i - \hat{T}_{c.m.} + \sum_{i>j}^A \hat{v}_{ij},$$

Volkov No.2 (M=0.6, B=H=0.125)

+G3RS (*ls* potential)

b=1.46fm

Neutron Tail of ${}^6\text{He}$



Check for absolute values of the binding energy

AMD
triple-S

RGM

${}^6\text{He}$ $E = -28.56\text{MeV}$ $E = -28.34\text{MeV}$

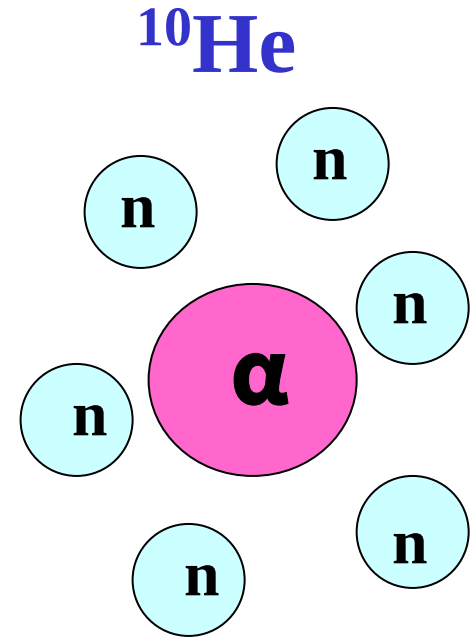
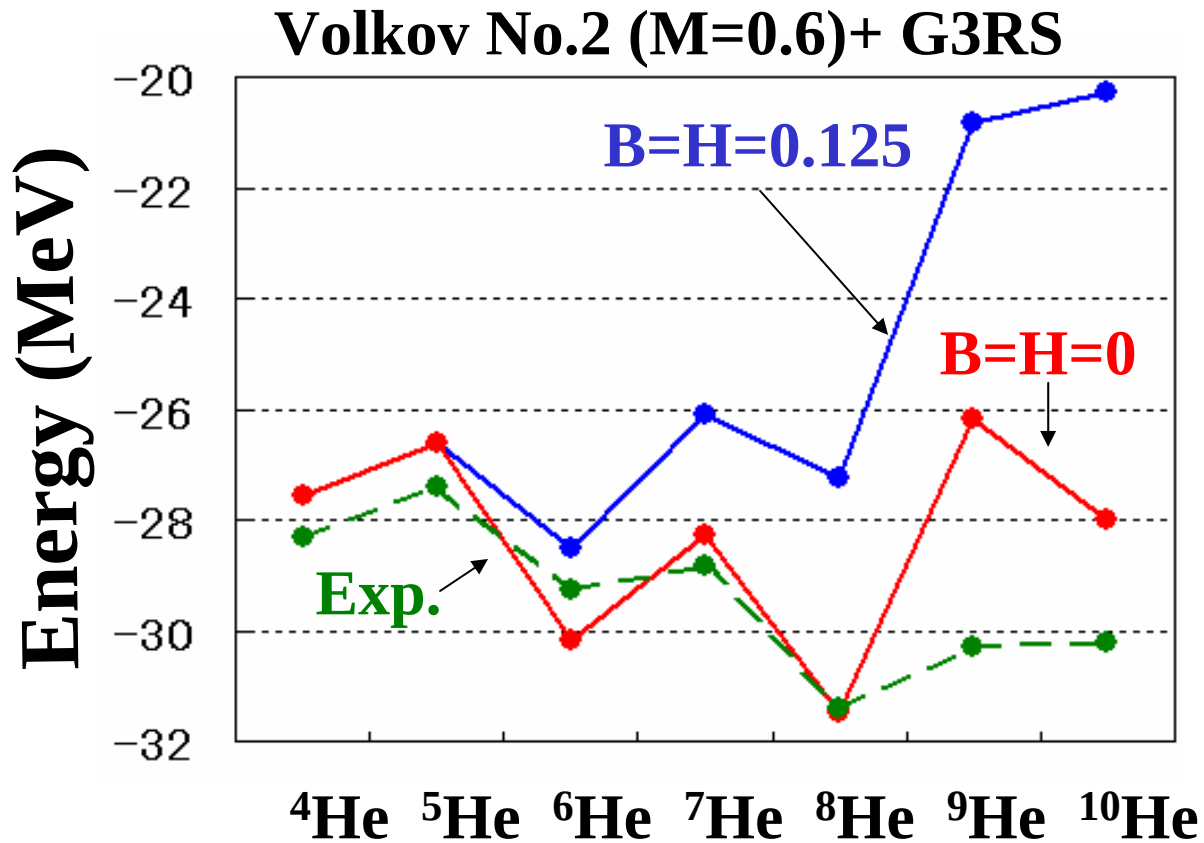
K. Arai

6.26 MeV



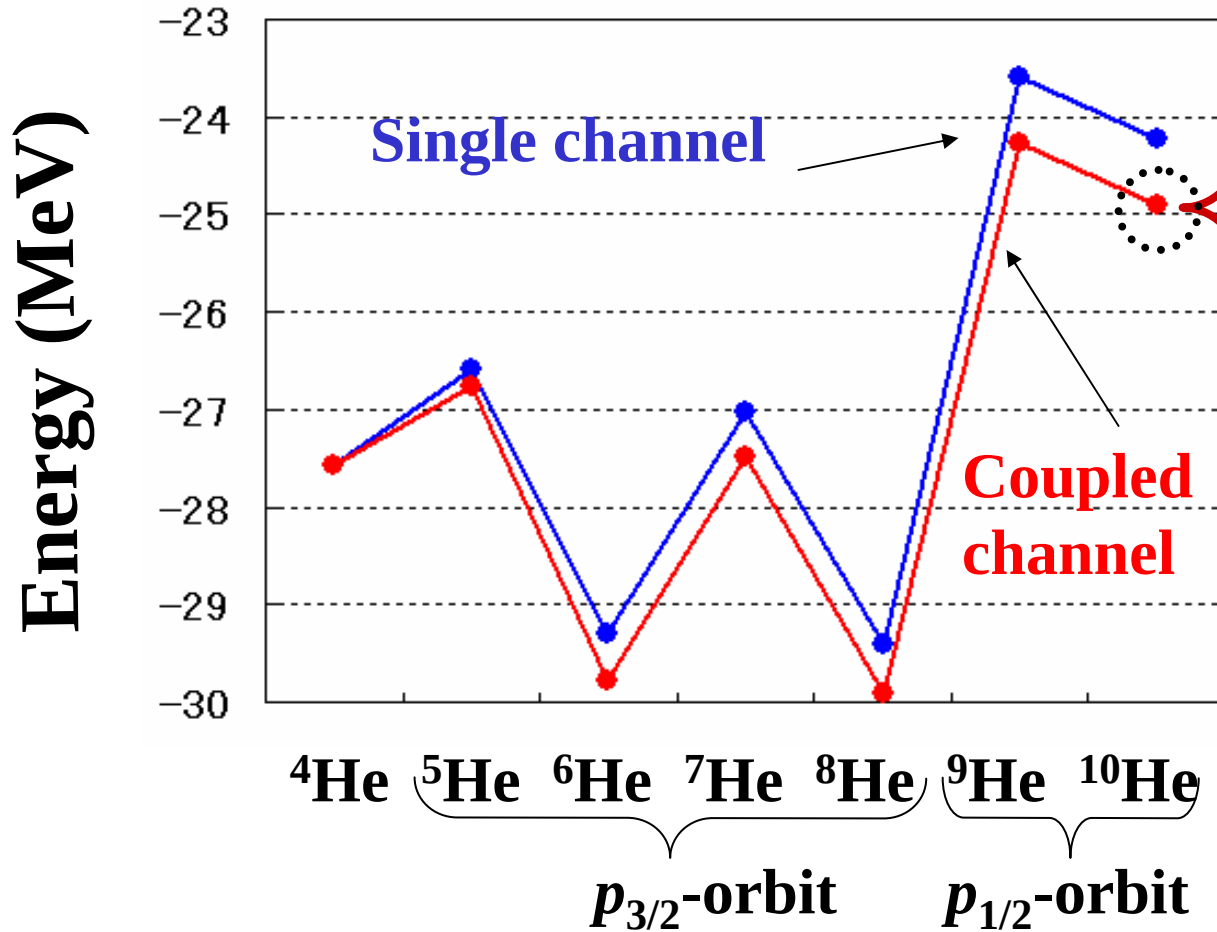
AMD: $E = \underline{-21.82\text{ MeV}}$
(Single Slater, VBP)

Calculated Energies for He-isotopes



t-t contribution in He-isotopes

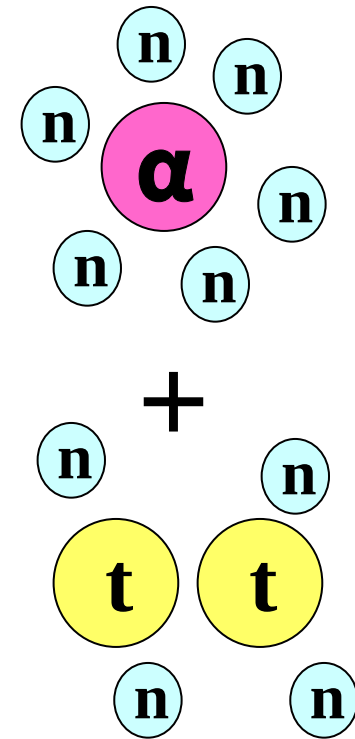
Volkov No.2 (M=0.6, B=H=0.06)+ G3RS



Single channel: α -core model

Coupled channel: α -core + (t+t)-core model

^{10}He



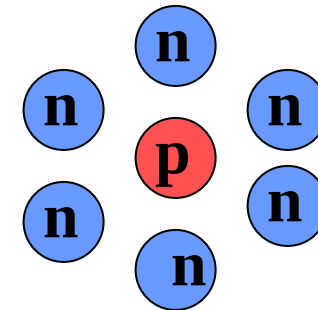
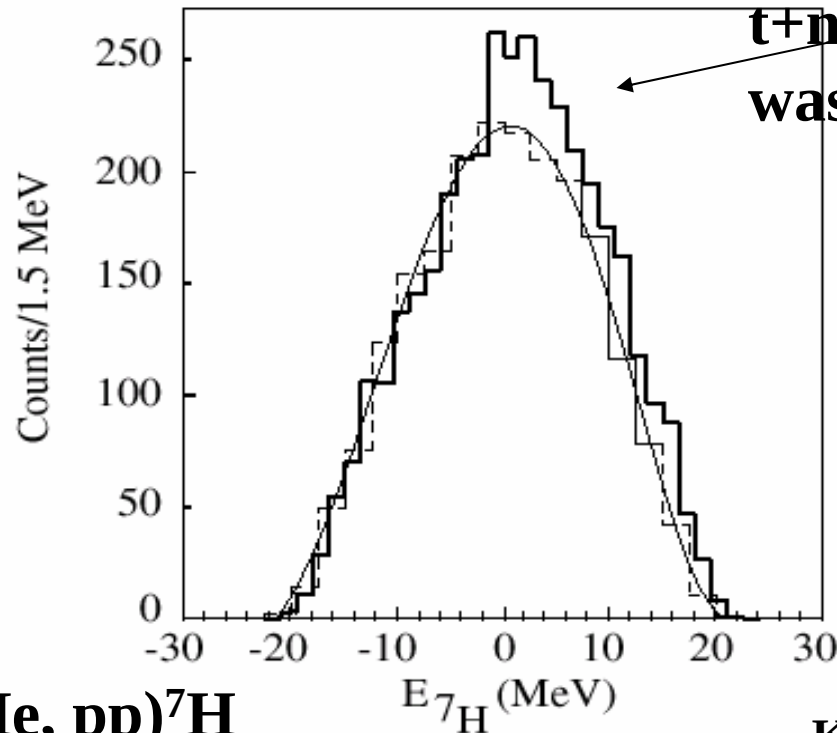
Coupled channel
problem for 7-body
system

Very recently, ${}^7\text{H}$ is observed in

DIKENI

Enhancement above

$t+n+n+n+n$ threshold
was observed



$N/Z=6$ very large value
 $\sim$ neutron star

$p({}^8\text{He}, pp){}^7\text{H}$

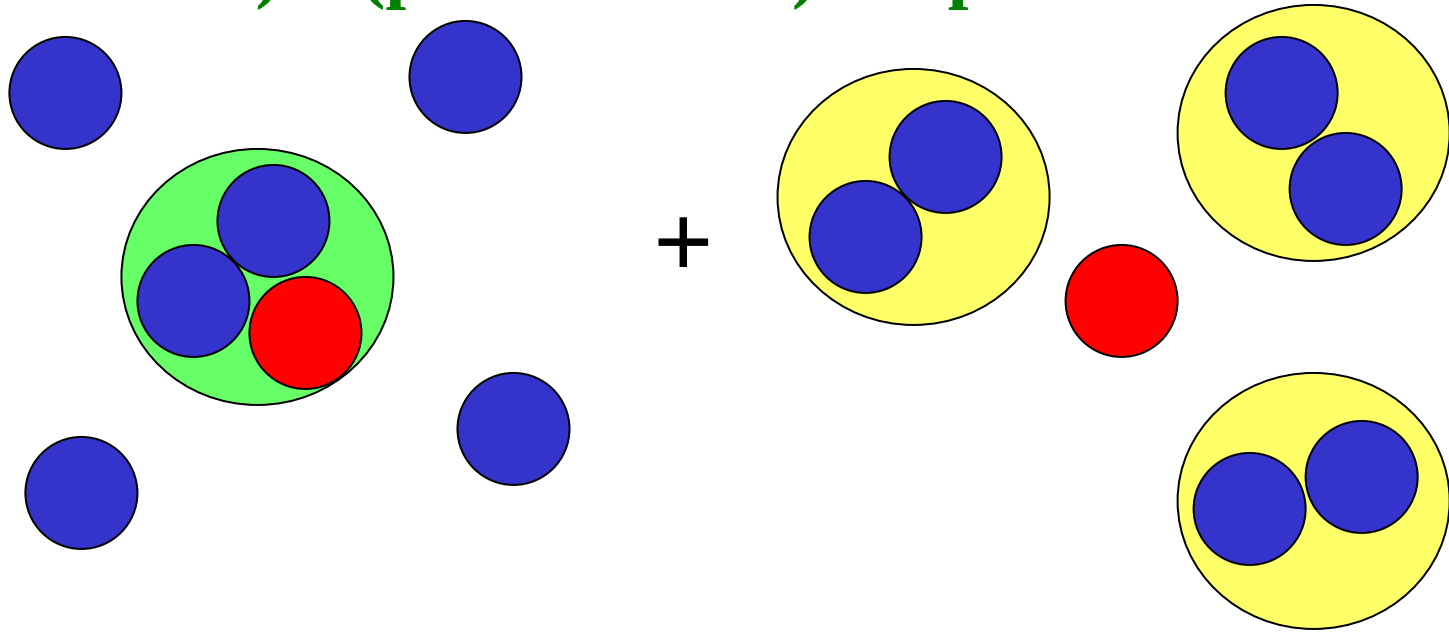
Korshenninnikov et al., PRL90(2003)

FIG. 3. Spectrum of ${}^7\text{H}$ from the reaction $p({}^8\text{He}, pp){}^7\text{H}$. The solid histogram was obtained with the proton target. The dashed histogram shows the empty-target background.

One proton in the six-neutron system

Ikeda and Horiuchi's advise

$(t+n+n+n+n) + (p+2n+2n+2n)$ coupled channel model



In this study, we investigate H-isotopes by using **AMD triple-S**.
(**A**ntisymmetrized **M**olecular **D**ynamics **S**uperposition **S**electd **S**napshot)



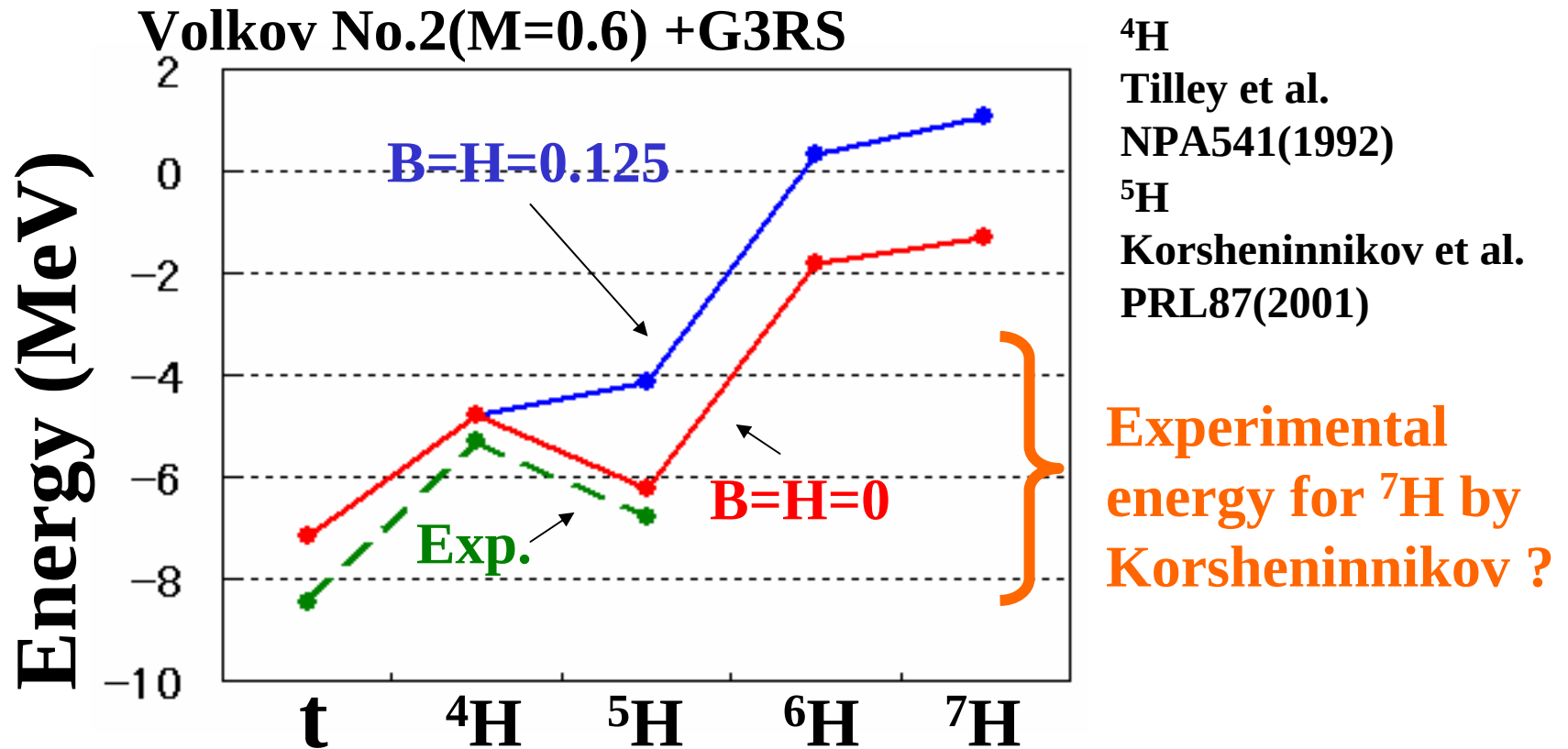
AMD+GCM is combined with **SVM**.

(**S**tochastic **V**ariational **M**ethod)

In order to solve the many body resonance, we will use **ACCC**.
(**A**nalytical **C**ontinuation in the **C**oupling **C**onstant)

V.I.Kukulin et al., JPA10 (1977)

Calculated Energies for H-isotopes



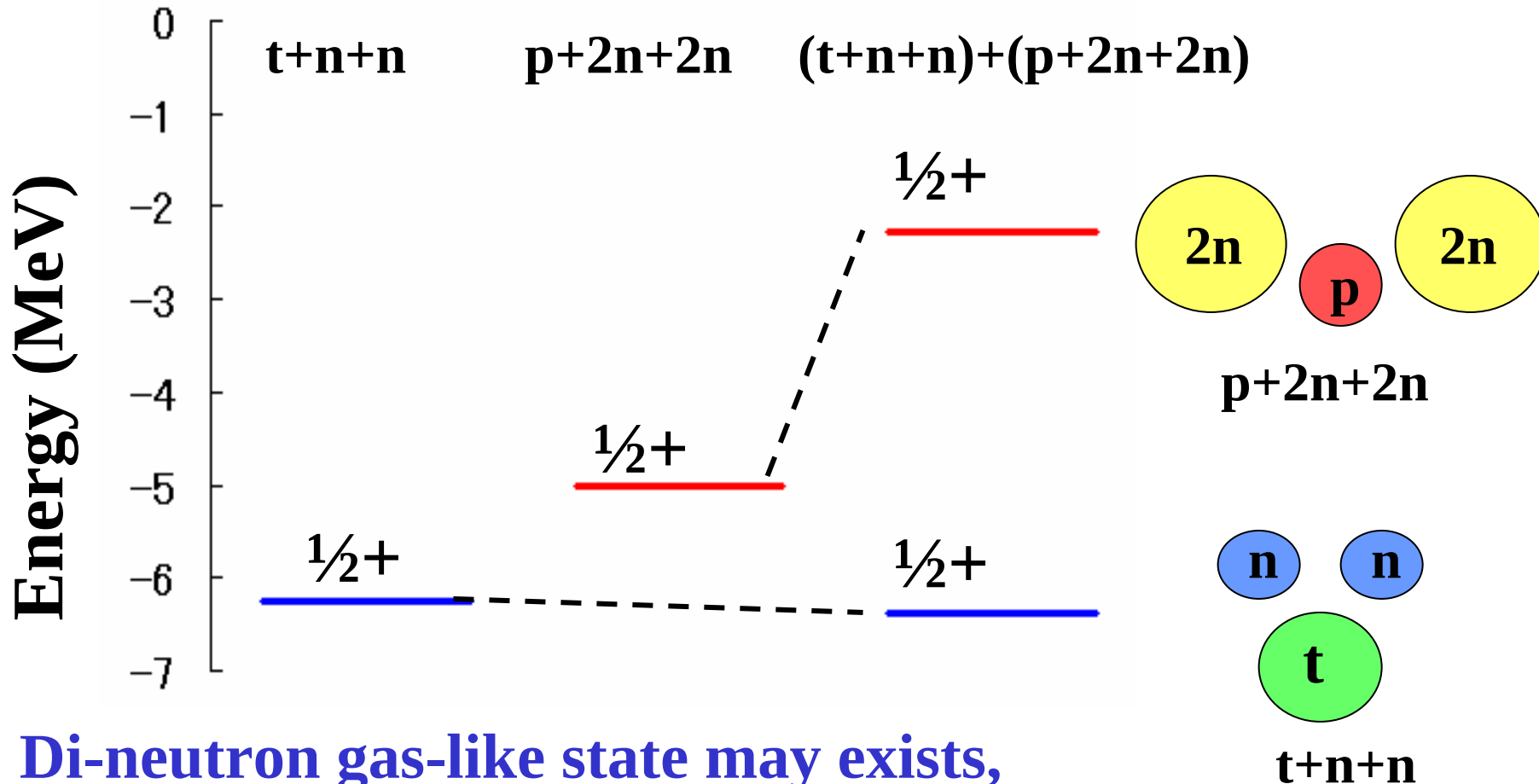
The present effective interaction does not reproduce tendency of the experimental binding energy for ${}^7\text{H}$.



Search for other effective interactions and correlations

Gas-like State in ^5H

Volkov No.2(M=0.6, B=H=0) +G3RS



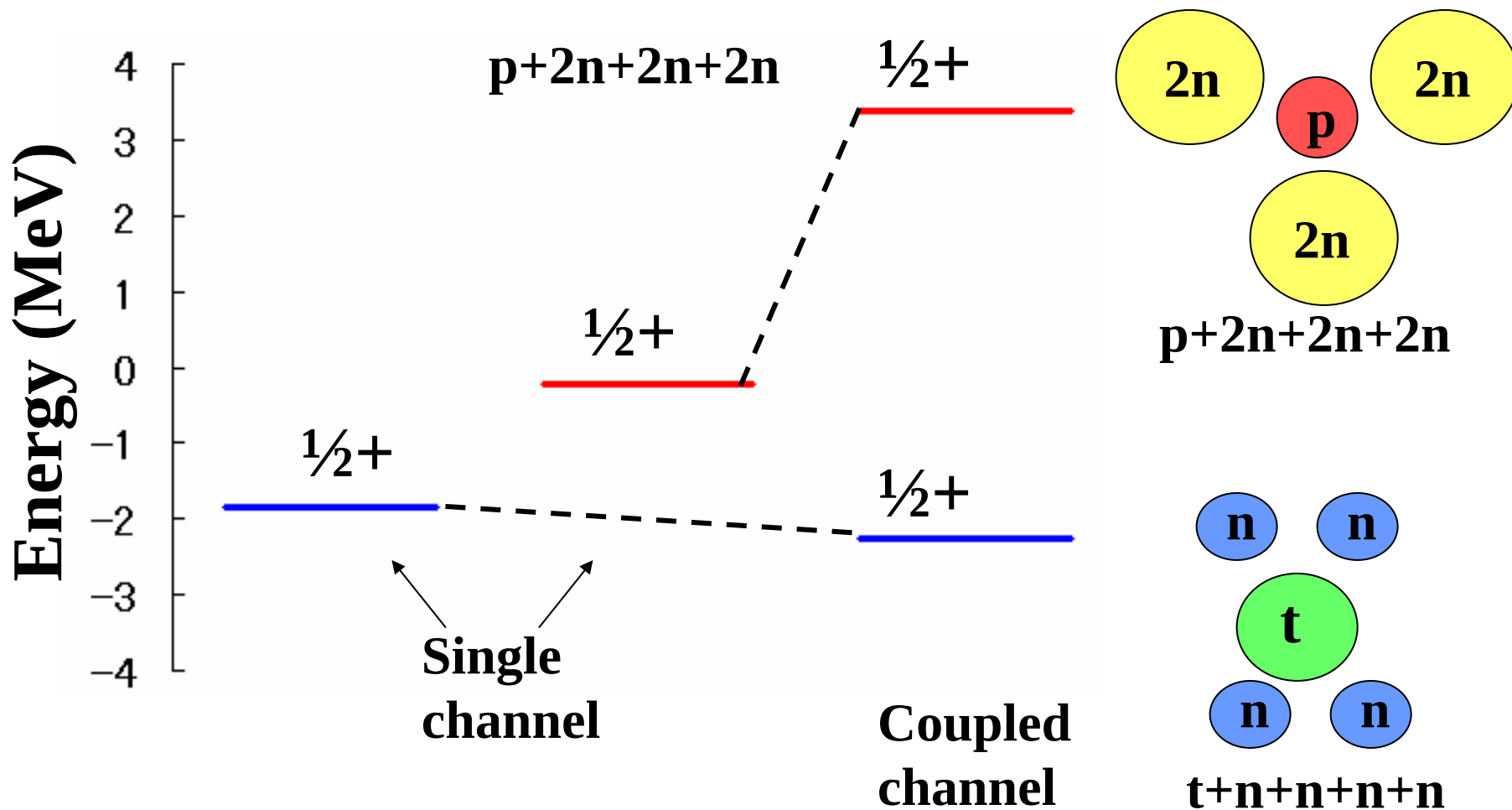
Di-neutron gas-like state may exist, although we should solve it as a resonance.

Gas-like State in ${}^7\text{H}$

Volkov No.2(M=0.6, B=H=0) +G3RS

$t+n+n+n+n$

$(t+n+n+n+n)+(p+2n+2n+2n)$



ACCC (Analytical Continuation in the Coupling Constant)

V.I.Kukulin et al., J. Phys. A10 (1977)

Application of ${}^9\text{Be}$ ($\alpha + \alpha + n$)

N.Tanaka, Y.Suzuki, K.Varga, G.Lovas, Phys. Rev. C59 (1999)

$$H = H_0 + \delta V \qquad H = T + V + \delta V = T + \lambda V$$

coupling constant
potential

Pade approximation

$$k_l^{MN}(\delta) = i \frac{c_0 + c_1 x + c_2 x^2 + \dots + c^M x^M}{d_0 + d_1 x + d_2 x^2 + \dots + c^N x^N}$$

$$x = (\delta - \delta_0)^{1/2}$$

coupling
constant

coupling constant at threshold
(except for s- wave)

Energy Levels of ${}^6\text{He}$

We can solve the resonance by using

ACCC ! $E_r=0.825$

2^+

$\Gamma=0.113$

$E_r=0.875$

$\Gamma=0.115$

Bound state approximation

$E_r=0.90 \text{ MeV}$

0^+

$E= -0.98$

$E= -0.99$

in MeV

Exp.

Cal.

N. Itagaki, A. Kobayakawa, S. Aoyama, PRC68, 055302(2003)

$$H=H_0 + \delta V$$

Pade approximation

$$k_l^{MN}(\delta) = i \frac{c_0 + c_1 x + c_2 x^2 + \cdots + c^M x^M}{d_0 + d_1 x + d_2 x^2 + \cdots + c^N x^N}$$

$$x = (\delta - \delta_0)^{1/2}$$

Summary

We develop a new method: AMD triple-S.

- We understand that the AMD triple-S is useful for the analyses of extremely neutron-rich nuclei .
- t-t component is also important for extremely neutron-rich He-isotopes.
- Gas-like configurations may exist in H-isotopes.
- The present effective interaction does not reproduce tendency of the experimental binding energy of H- and He-isotopes.
- We can solve three-body resonance: As a first step, 2^+ state of ${}^6\text{He}$ is solved by using AMD triple-S combined with ACCC.

Next

1. We are going to investigate other effective interactions and correlations.
2. We solve it as a resonance with ACCC.