

Studies on excited states of Be isotopes

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Contents

1. Introduction

2. Our model

3. Results of ^{10}Be

4. Effect of Spin-orbit interaction

5. Discussion on breakup reaction of ^xBe

Introduction

1. Studies on nuclear structure

Model Space $\leftrightarrow$ Eff. Interaction (**Close connection**)

↓
should depend on the subjects of studies

2. Studies on Be isotopes ($\alpha + \alpha + n + n + \dots$)

- Low-lying states $\rightarrow$ Molecular orbitals (π^- , σ^+)
- High-lying states $\rightarrow$ Resonances (${}^6\text{He}$, ${}^8\text{He}$??)

They should be studied in a unified way !!

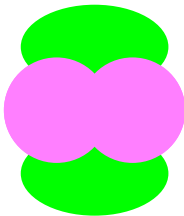
3. Appropriate model space

We should include high $\hbar\omega$ state around two α cores.

Small α - α distance

Molecular orbitals

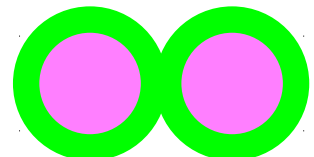
Low $\hbar\omega$



Large α - α distance

Valence bonding

High $\hbar\omega$



4. Present approach

Valence neutron's orbital $\rightarrow$ Simple LCAO (0p-shell)

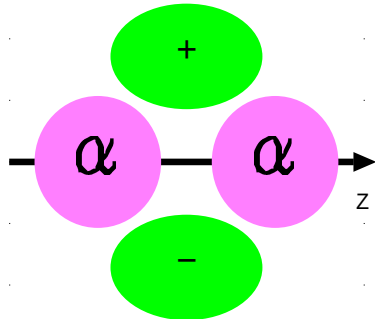
Full configuration of LCAO (0~2 $\hbar\omega$ excitation)

NN force $\rightarrow$ reproduce the threshold (Parameters)

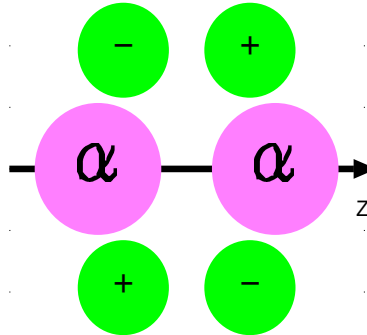
Molecular orbitals in LCAO

1. Linear Comb. of Atomic Orbitals (LCAO)

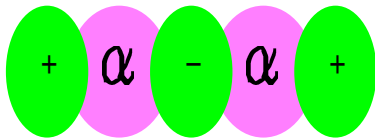
π^- ($0\hbar\omega$, 0p-shell)



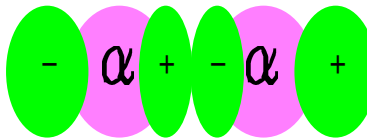
π^+ ($1\hbar\omega$, sd-shell)



σ^+ ($1\hbar\omega$, sd-shell)

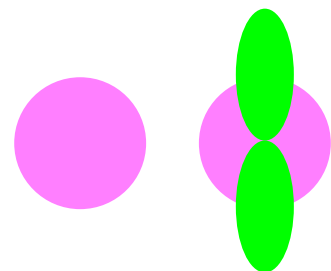


σ^- ($2\hbar\omega$, pf-shell)

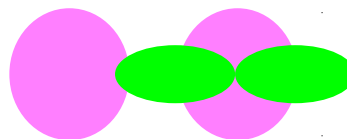


2. Projection to the atomic orbitals

$$\pi^- + \pi^+ = 0\hbar\omega + 1\hbar\omega$$



$$\sigma^+ + \sigma^- = 1\hbar\omega + 2\hbar\omega$$



3. Formation of $^{10}\text{Be} = \alpha + {}^6\text{He}(0_1^+)$

+ parity $\rightarrow 0, 2, 4 \hbar\omega : (\pi^-)^2, (\pi^+)^2, (\sigma^+)^2, (\sigma^-)^2, \dots$

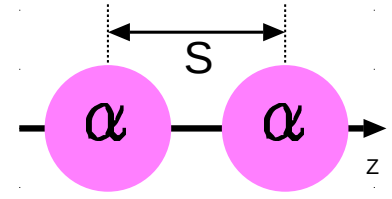
- parity $\rightarrow 1, 3 \hbar\omega : (\pi^- \sigma^+), (\sigma^+ \sigma^-), \dots$

Generalized Two-center Cluster Model (GTCM)

1. Construction of the basis functions

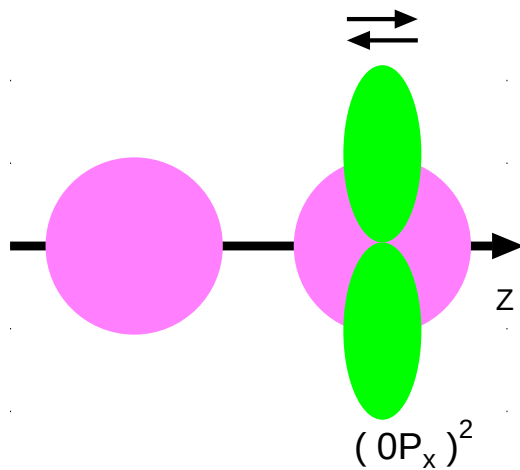
System : $^{10}\text{Be} = \alpha + \alpha + n + n$ $\alpha : (0s)^4$

Neutrons are arranged around one of α cores with $0p_x, 0p_y, 0p_z$



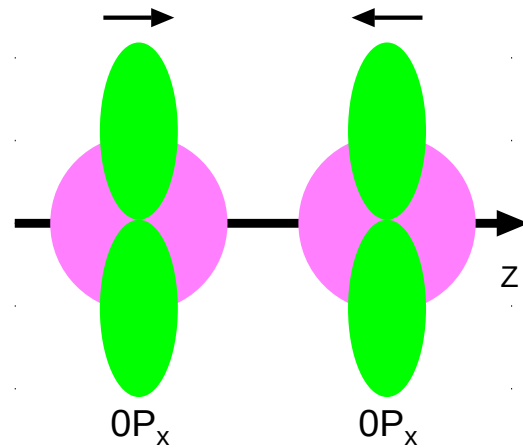
$\alpha + {}^6\text{He}$ Partition

$\alpha + (\alpha + n + n)$



${}^5\text{He} + {}^5\text{He}$ Partition

$(\alpha + n) + (\alpha + n)$



$$\Phi^{J^\pi K}(S, n_1, n_2) = \mathcal{P}^\pi P_{MK}^J \mathcal{A} \{ [\varphi_L(\alpha) \varphi_R(\alpha) \phi(n_1, n_2)]_{\mathbf{S}} \}$$

n_i : (Left or Right, $0p_x, 0p_y, 0p_z$, $\uparrow$ or $\downarrow$), $\mathbf{b}=1.46$ fm

2. Total wave function

$$\Psi^{J^\pi} = \int dS \sum_{K, \beta} f(K, S, \beta) \Phi^{J^\pi K}(S, \beta) \quad \beta = (n_1, n_2)$$

3. Full eigenvalue problem

$$\Psi^{J^\pi} = \int dS \sum_{K, \beta} f(K, S, \beta) \Phi^{J^\pi K}(S, \beta)$$

$$\equiv \int dS \chi^{J^\pi}(S)$$

$$\chi^{J^\pi}(S) \equiv \sum_{K, \beta} f(K, S, \beta) \Phi^{J^\pi K}(S, \beta)$$

$$< \Phi^{JK'\pi}(S', \beta') | H - E | \int dS \chi^{J^\pi}(S) > = 0$$

4. Adiabatic energy surfaces

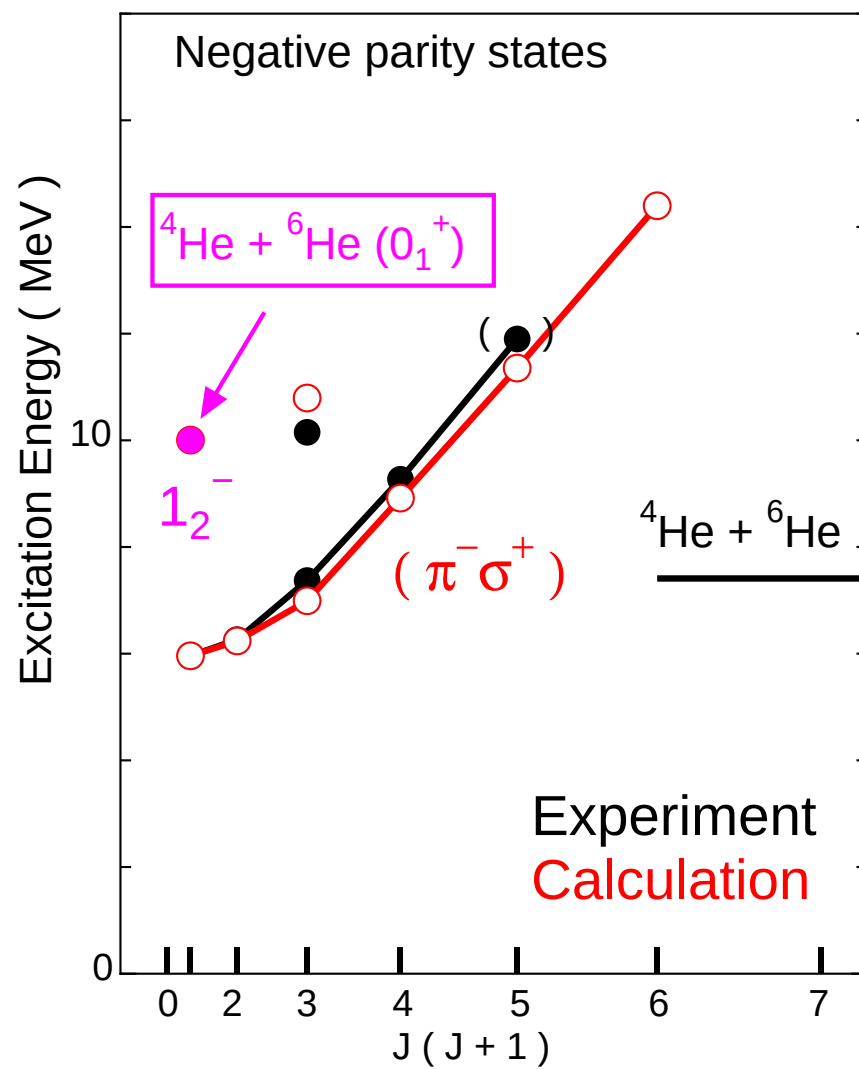
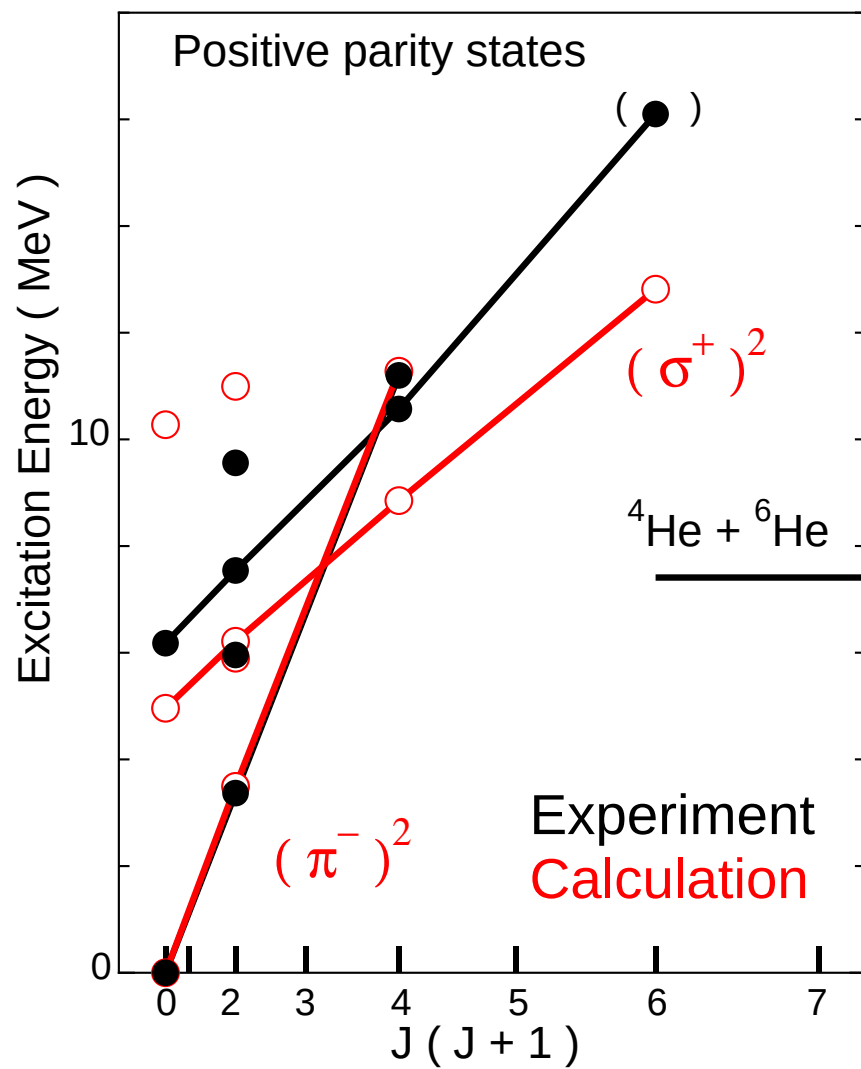
We can solve eigenvalue equations at a fixed distance S .

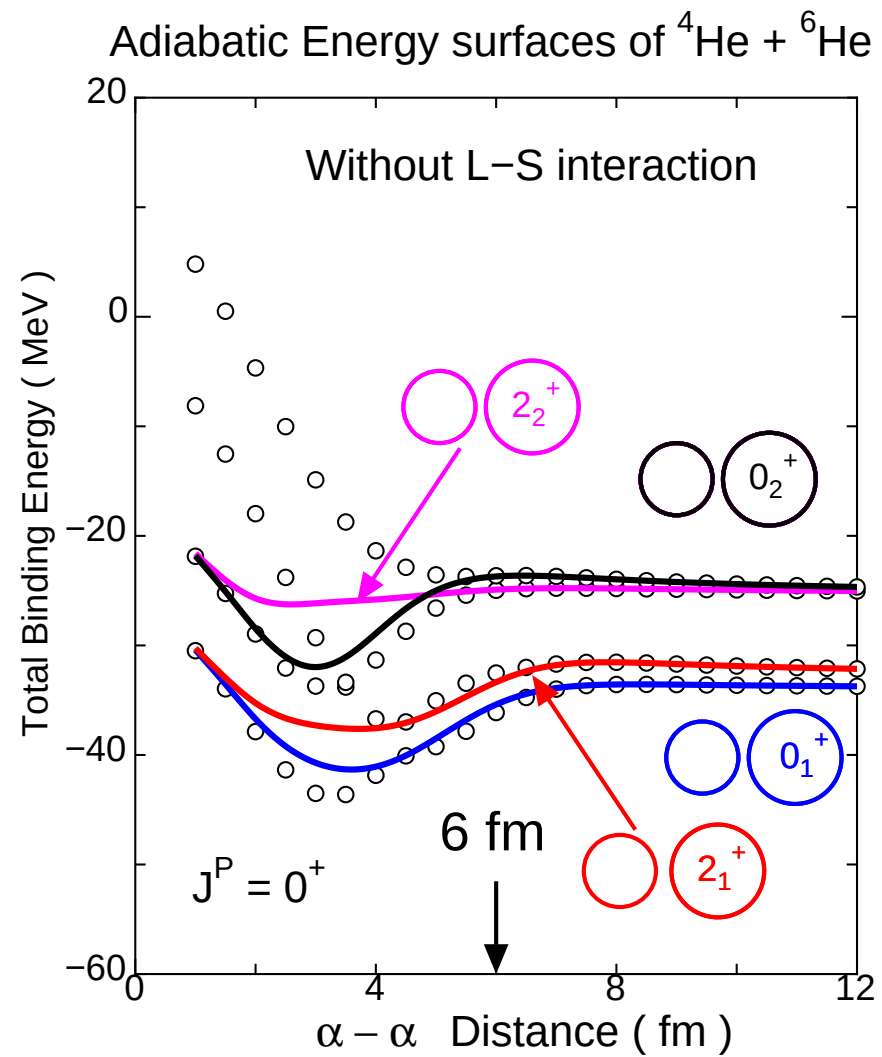
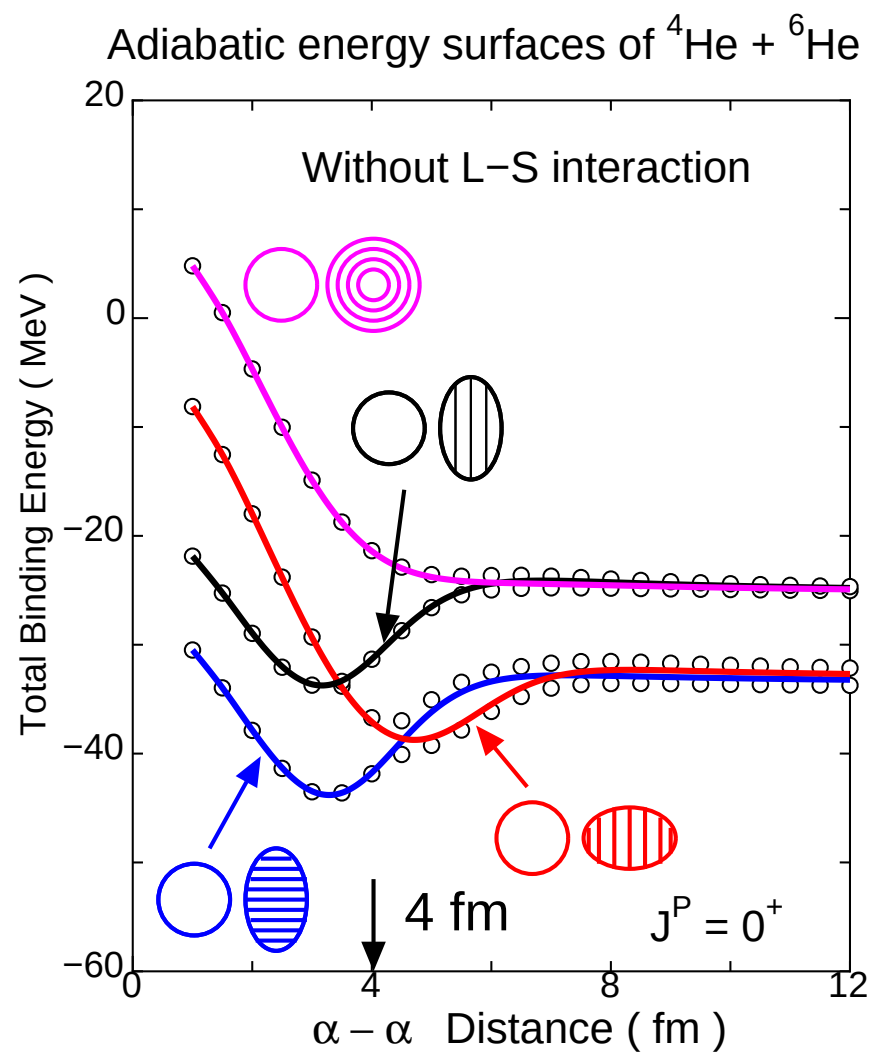
$$< \Phi^{JK'\pi}(S, \beta') | H - E(S) | \chi^{J^\pi}(S) > = 0$$

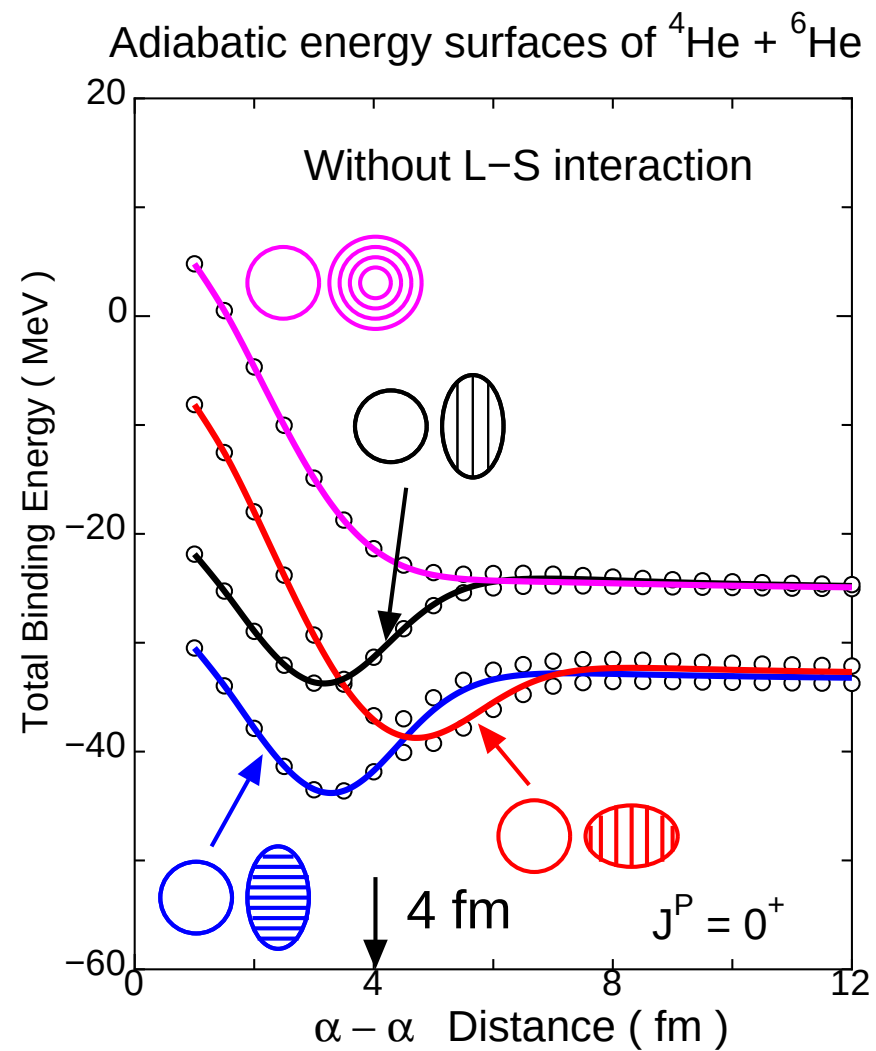
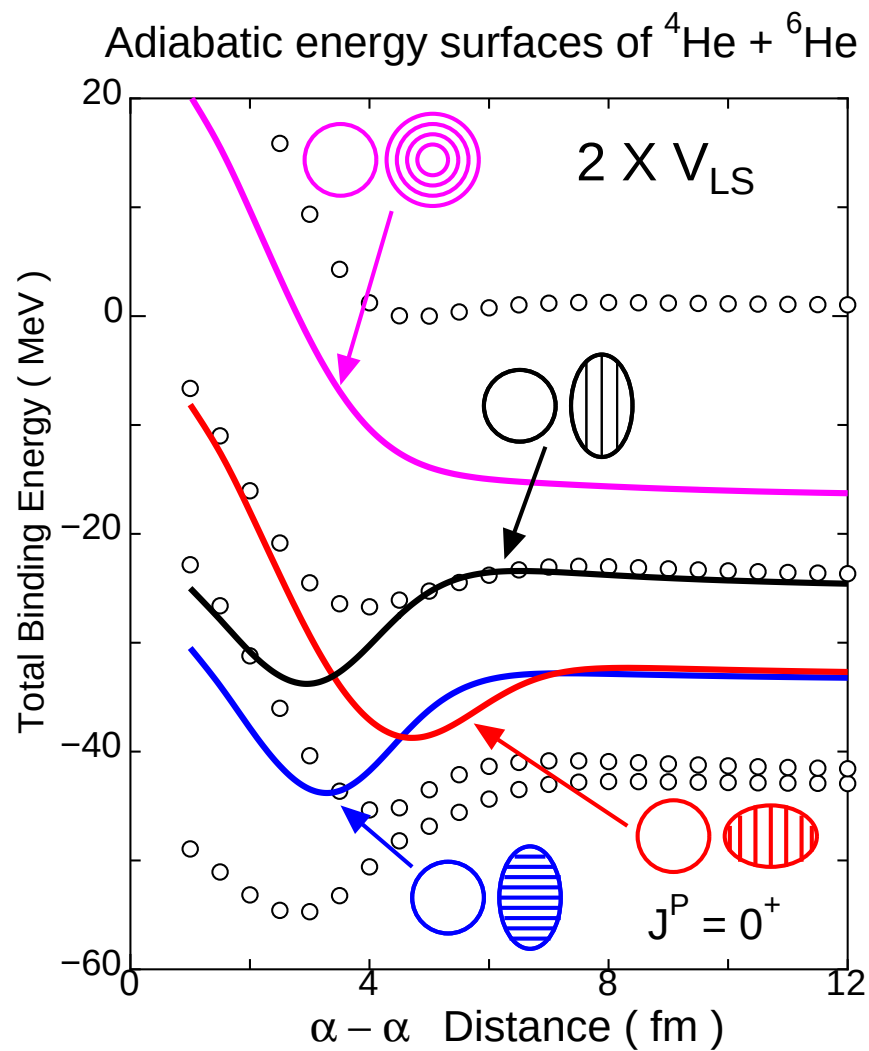
5. nucleon–nucleon interaction

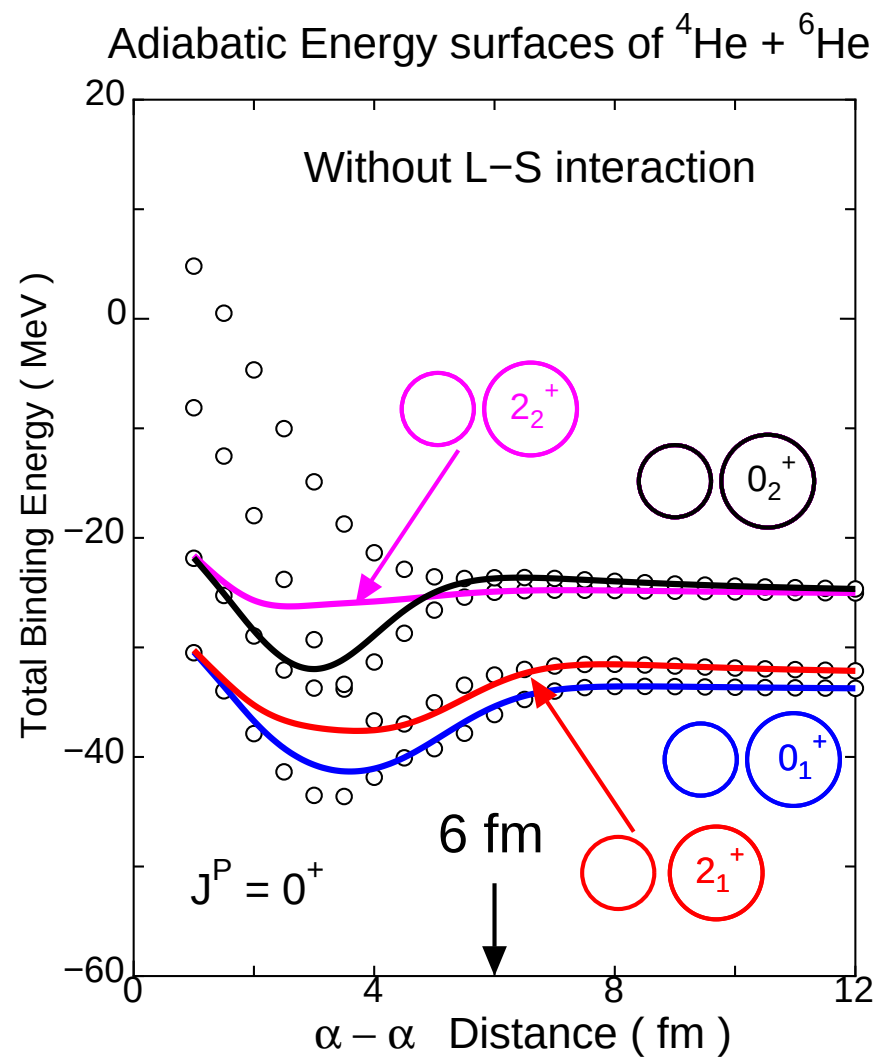
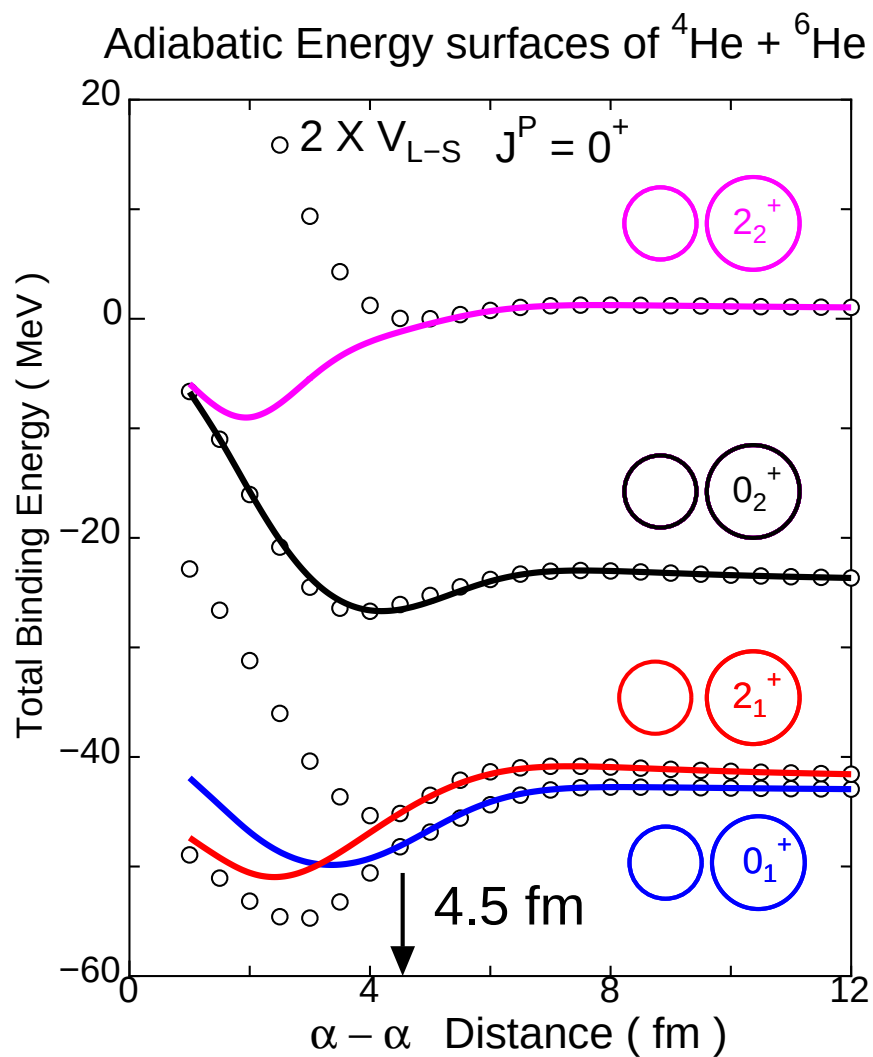
- **Cent. : Volkov (2) $\rightarrow$ M=0.643, B=−H=0.125**
- **L–S : G3RS $\rightarrow$ Rep. +3000, Att. −2000 (MeV)**
- **Size b=1.46 fm**

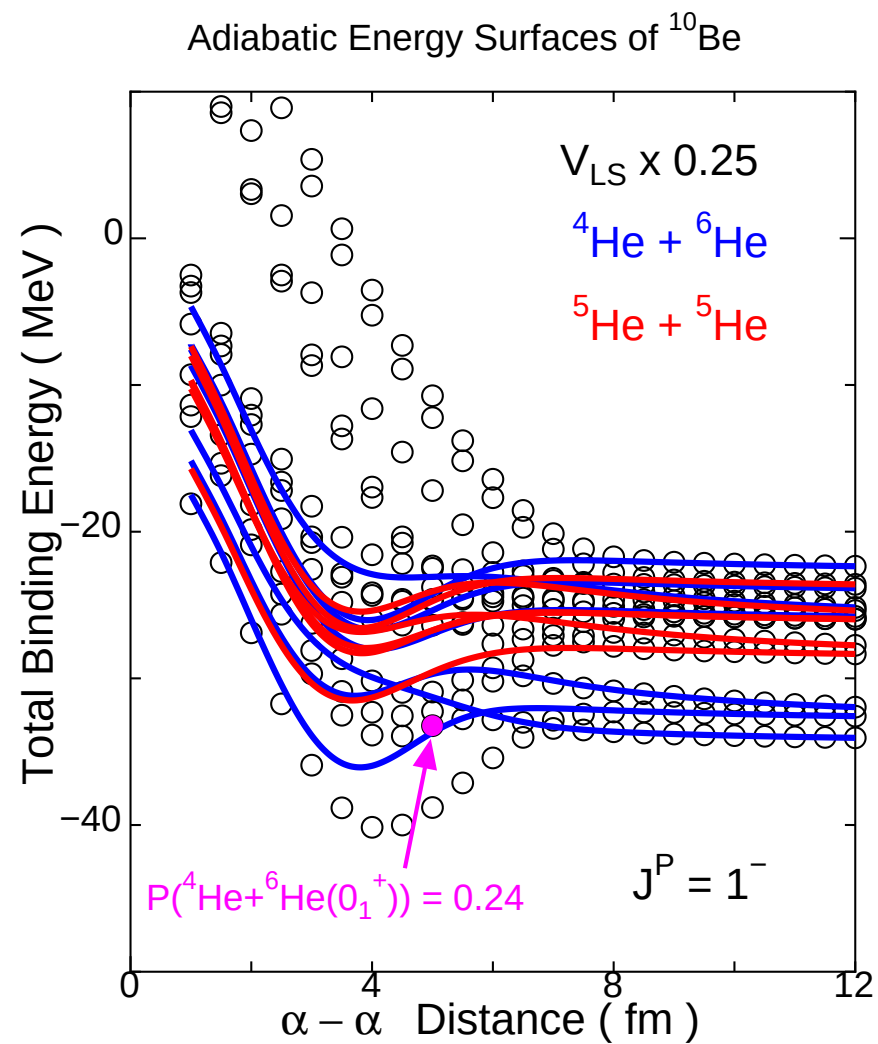
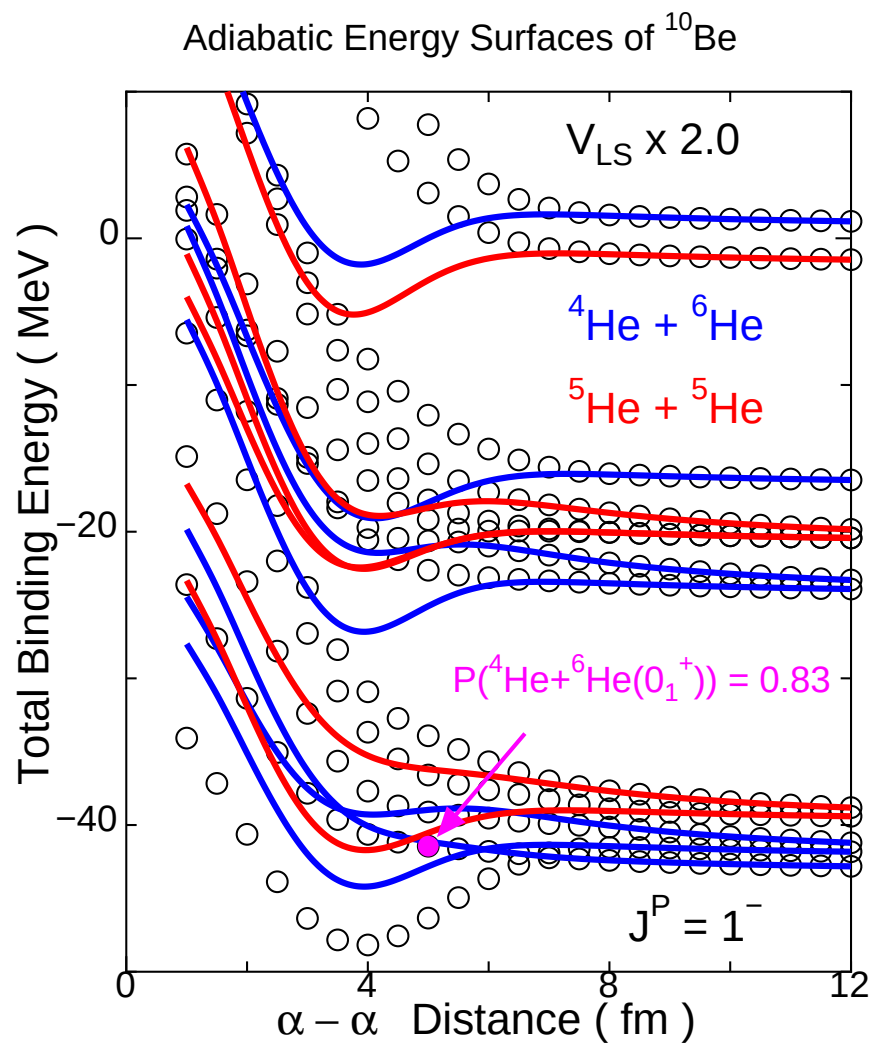
	Cal. (MeV)	Exp. (MeV)
$E_{th}(\alpha + {}^6\text{He})$	7.41	7.41
${}^6\text{He}(2_1^+)$	1.9	1.8
${}^5\text{He}$ L–S splitting	6.0	1.47











Observation of the $\alpha + {}^6\text{He}_{g.s.}$ ($J^\pi = 1_2^-$)

E1 Coulomb Excitations (Breakup)

$${}^{10}\text{Be}_{g.s.} \rightarrow {}^{10}\text{Be}_{ex.} = | \alpha + {}^6\text{He} \rangle + | {}^5\text{He} + {}^5\text{He} \rangle$$

$${}^{10}\text{Be}_{g.s.} \rightarrow \alpha + {}^6\text{He} : \text{E1 transition is possible.}$$

$${}^{10}\text{Be}_{g.s.} \rightarrow {}^5\text{He} + {}^5\text{He} : \text{E1 transition is forbidden !!}$$

The transition ${}^{10}\text{Be}$ (g.s. $\rightarrow 1_2^-$) will be enhanced in **E1 Coulomb breakup !!**

${}^{10}\text{Be}$ + Target inelastic scattering

${}^{12}\text{C}$ target : Nuclear dominance

${}^{208}\text{Pb}$ target : Coulomb dominance

$\rightarrow {}^{10}\text{Be}(\text{g.s.} \rightarrow 1_2^-)$ will be enhanced !!

e. g. ${}^{12}\text{Be}_{g.s.} \rightarrow {}^{12}\text{Be} (1^-)$ by Iwasaki *et al.*

Systematics in Be isotopes

${}^x\text{Be}$	Occupation	J^π	Cluster	Probe
${}^{10}\text{Be}$	$(\pi_{3/2}^- \pi_{3/2}^+)$	1^-	${}^4\text{He} + {}^6\text{He}$	Coulomb
${}^{12}\text{Be}$	$(\pi_{3/2}^-)^2 (\pi_{3/2}^+)^2$	0^+	${}^6\text{He} + {}^6\text{He}$	Nuclear
${}^{14}\text{Be}$	$(\pi_{3/2}^-)^2 (\pi_{3/2}^+)^2 (\sigma_{1/2}^+ \sigma_{1/2}^-)$	1^-	${}^6\text{He} + {}^8\text{He}$	Coulomb

Summary

We applied GTCM to ^{10}Be and discussed the effect of the L–S interaction for the weak-coupling formation of $\alpha + {}^6\text{He}$.

Conclusion

1. Model space

The high $\hbar\omega$ states : (0p) $\sim$ (1p 0f) shell

$0 \sim 2\hbar\omega$ excitation in a single particle state

2. L–S interaction

Break the geometrical configuration

Enhance the weak-coupling configuration

3. Observation of ${}^6\text{He}$ or ${}^8\text{He}$ cluster states

Coulomb breakup $\rightarrow \alpha + {}^6\text{He}$, ${}^6\text{He} + {}^8\text{He}$

Nuclear breakup $\rightarrow {}^6\text{He} + {}^6\text{He}$

Future plans

1. Application of Complex Scaling Method

Identification of resonance poles and their width

2. Transition density of ${}^x\text{Be}$

Coupled channels of ${}^x\text{Be} + \text{Target}$