# Supersymmetry after Higgs discovery

DBD16 Workshop : 11/8/2016 Masahiro Ibe (ICRR&IPMU)

# <u>What have we learned from LHC results ?</u>

Higgs exists

✓ Higgs is more like an elementary scalar

Higgs can be fit by the simplest implementation

 $V = -m_{higgs}^2/2 h^+h + \lambda/4 (h^+h)^2$ 

 $m_{higgs} = \lambda^{1/2} v [v = 174.1 GeV]$ 

 $m_{higgs} \sim 125 \text{GeV} \longrightarrow \lambda \sim 0.5$ 



(We knew v=174.1GeV before the discovery of Higgs)



# <u>What have we learned from LHC results ?</u>

- No New Physics beyond the SM has been confirmed ...
   [Some tentative hints are/were reported though]
- Wide range of alternatives to the Standard Model have been excluded...
  - [e.g. Higgsless models]

# Could the Standard Model Be the Ultimate Theory?

# <u>Answer is NO</u>

There should be new physics which addresses at least

- 🗸 Quantum Gravity
- 🗸 Neutrino Masses
- 🗸 Dark Matter

Baryon Asymmetry of the Universe

- Inflation
- Strong CP problem
- Anomalies ( muon g-2, B-physics etc )

# Are some of these entries accessible at the LHC?

### Possibly, but not sure...

Most of the above entries can be solved in decoupled (ultra-violet) theories...

## Why New Physics at the LHC?

### We want to know why the weak scale $\langle M_{PL}$ .

Let us suppose that the Higgs couples to X in the UV theories.

$$\Delta L = \lambda_X |h|^2 |X|^2$$

$$\rightarrow \Delta m_h^2 = \frac{\lambda_X}{16\pi^2} \left[ \Lambda_{\rm UV}^2 - m_X^2 \left( \log \left( \frac{\Lambda_{\rm UV}^2}{m_X^2} \right) + (\text{finite}) \right) \right]$$

 $\rightarrow$  Even aside from the quadratic divergence, the Higgs boson mass is sensitive to the UV scale  $m_X$ !

If  $M_W << M_{PL}$  is explainable, there should be new physics close to  $M_W$  which renders the SM being UV insensitive.

 $\rightarrow$  New Physics Accessible at the LHC !

[No verifiable conflicts happen even if no new physics exist...]

# <u>Supersymmetric Standard Model</u>

We just enlarge spacetime symmetry to supersymmetry !



Higgs mass term can be protected !

Higgs mass term = Higgsino mass term Higgsino mass term can be protected by chiral symmetries! Hierarchy problem is solved if SUSY breaking is around M<sub>w</sub>.

['79 Maiani, '81 Witten, '81 Kaul ]

### **Supersymmetric Standard Model**

Big Bonus !

Just by introducing the superpartners at around *TeV*, gauge coupling unification becomes more precise!



Supersymmetric standard model is perfectly consistent with GUT!

### **Supersymmetric Standard Model**

Searchable at the LHC !

Supersymmetry = Spacetime symmetry

All the Standard Model particles are accompanied by superpartners.

→ Lots of colored new particles (*gluino*, *squarks*)





### <u>The Higgs boson ~ 125GeV requires multi-TeV SUSY ?</u>

In the MSSM, the tree-level Higgs boson mass is given by the gauge coupling constants.



The predicted Higgs boson mass is around Z-boson mass,

 $m_H = \lambda^{1/2} v \sim m_Z \cos 2\beta$ 

at the tree-level.



### <u>The Higgs boson ~ 125GeV requires multi-TeV SUSY ?</u>

The radiative corrections to the Higgs boson mass logarithmically depend on the stop masses.



['91 Okada, Yamaguchi, Yanagida, '91 Haber, Hempfling, '91 Ellis, Ridolfi, Zwirner ]

# The Higgs boson mass larger than m<sub>z</sub> can be obtained for larger SUSY breaking effects!



### <u>The Higgs boson ~ 125GeV requires multi-TeV SUSY ?</u>

The simplest interpretation:  $m_H \sim 125 \text{ GeV}$  suggests that the sfermion (stop) masses are above O(10-1000)TeV!



['12, MI, Matsumoto, Yanagida (µ<sub>H</sub>=O(M<sub>susy</sub>))]

### <u>m<sub>H</sub> ~ 125 GeV with O(1)TeV SUSY?</u>

- ✓ A large *stop A-term* (trilinear coupling)
- ✓ Singlet extension (NMSSM)
- Extra vector-like matter
- ✓ U(1) extension

Model building becomes complicated though.

Direct SUSY search at the LHC.

Proton stability

Baryon asymmetry of the universe

We assume *R*-parity (or *B-L* parity)

$$P_R = (-1)^{3(B-L)+2s}$$

The lightest superparticle (LSP) is *stable* (at least inside detectors).

The neutral LSP (neutralino, gravitino):

We look for missing transverse momentum



We look for heavy charged track.





# Direct SUSY search at the LHC.

squark mass	SUSY spectrum	Higgs	EW Naturalness	GUT	10 years from now?	
≫ <b>PeV</b>	DM ?? wino<3TeV Higgsino <2TeV	m <sub>Higgs</sub> ≫125GeV	We need to rethink	Unification is OK as long as M <sub>gaugino</sub> < O(TeV)	Stopping Gluinos ?	
10TeV- PeV	One-loop suppressed Gaugino mass	т <sub>Higgs</sub> ~125GeV in the MSSM	We need to rethink	Unification is OK as long as M <sub>gaugino</sub> < O(TeV)	Gaugino searches. DM Cosmic ray ? EDM	
TeV	Getting excluded by the LHC	Large A-term NMSSM Other extensions	sub% Tuning ( low scale mediation is preferred ? )	Unification is OK	The LHC will discover the SUSY (M <sub>SUSY</sub> < 3TeV)! Or merged into high-scale SUSY	
Light Stop& Higgsino	Stop <tev DM : Higgsino</tev 	Enhancement is required!	Better than O(1)% if M <sub>gulino</sub> < TeV	Most models have tensions with the GUT	<b>Searchable!</b> Gluino/Stop search will be discovered !	
Compressed Spectrum	Loophole in the LHC searches!	Enhancement is required!	Can be better than the simple TeV scenarios	GUT relation in the gaugino masses are not favored (exception: Mirage Mediation)	Searchable! ISR jet ? ISR photon ? soft lepton ?	
This table is an extended version of						

Unnatural Natural

http://www2.yukawa.kyoto-u.ac.jp/~fumihiro.takayama/YITP\_ws2012March\_BSM/talks/SUSY\_sub.pdf

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missing  $p_T$  search (the neutralino LSP)



 $M_{gluino} \sim M_{squark} > 1.8 \text{ TeV} (Run1)$  $M_{gluino} > 1.8 \text{ TeV}$  for  $M_{squark} \gg \text{TeV} (13 \text{TeV}, 13.3 \text{ fb}^{-1})$ 

[Future discovery reaches:  $M_{gluino} \sim M_{squark} \sim 2.8 \text{ TeV} (14 \text{TeV}, 300 \text{ fb}^{-1})$   $M_{gluino} \sim M_{squark} \sim 3.2 \text{ TeV} (14 \text{TeV}, 3000 \text{ fb}^{-1})$   $M_{gluino} \sim M_{squark} \sim 6.8 \text{ TeV} (33 \text{TeV}, 3000 \text{ fb}^{-1})$ ] [arXiv:1310.0077]

### missing $p_T$ search (the gravitino LSP with $m_{3/2} << O(1) keV$ )



✓ stable stau search (GMSB with  $m_{3/2} > O(100) keV$ )



via gluino,squark production Mstau > 480 GeV (GMSB Nm = 3, Mgluino > 3.3TeV) via Drell-Yan production Mstau > 340 GeV (Run1) (Mgluino,squark >> TeV)

# Direct SUSY searches at the LHC.

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### Supersymmetry with light stop and light Higgsino

Why light stop/ light Higgsino ?

Higgs mass terms in the MSSM

 $V = (m_{Hu}^{2} + |\mu_{H}|^{2}) |H_{u}|^{2} + (m_{Hd}^{2} + |\mu_{H}|^{2}) |H_{d}|^{2} + B\mu_{H} H_{u} H_{d} + h.c.$   $\mu_{H}: \text{Higgsino mass}$  $m_{Hu}^{2}, m_{Hd}^{2}, B\mu_{H}: \text{soft SUSY breaking Higgs mass}$ 

Soft SUSY breaking mass receives large contributions from the stop mass

$$\Delta m_{H_u}^2 = -\frac{12y_t^2}{16\pi^2} m_{\tilde{t}}^2 \log\left(\frac{M_{\rm med}}{m_{\tilde{t}}}\right)$$

 $\rightarrow m_{stop}$ ,  $\mu_H \sim M_{weak}$  is favorable for naturalness

[Model building with consistent Higgs boson mass is very difficult though...]

### Supersymmetry with light stop and light Higgsino



#### Current limit :

*M*<sub>stop</sub> < 800 GeV has been excluded

(when stop  $\rightarrow$  top +  $\chi^o$  mode is open )

Prospect :

Discovery could be possible for  $M_{stop} = 1 - 1.2 \text{ TeV} (14 \text{ TeV})$ .

# Direct SUSY searches at the LHC.

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- ✓ The simplest interpretation:  $m_H \sim 125 \text{ GeV}$  suggests that the sfermion (stop) masses are above O(10-1000)TeV!
- Model building is ridiculously simple !
- Consistent with cosmology (good DM candidate, no Polonyi problem)
- Gaugino Masses in the TeV range can be naturally obtained.
- Fine-tuning problem between O(10-1000)TeV and O(100)GeV...



Gaugino Masses = Loop-suppressed (AMSB + Heavy Higgs Mediation)





 $m_{gluino} = 2.5 \times 10^{-2} m_{3/2}$   $m_{wino} = 3.0 \times 10^{-3} (m_{3/2} + L)$   $m_{bino} = 9.6 \times 10^{-3} (m_{3/2} + L/11)$ for  $m_{3/2} = O(100) TeV$ .

 $L = O(m_{3/2})$ : Higgsino mediation effect

- The wino is the LSP in the most parameter space.
- ✓ The gluino can be lighter than the prediction in AMSB for  $L/m_{3/2} = O(1)$ .

Current limits via gluino production



 Multi-jets + Missing *p*<sub>T</sub> search (conventional SUSY search)

> $m_{gluino} > 1.8 \text{ TeV} \text{ or } m_{wino} > 600 \text{GeV}$  $m_{gluino} > 1 \text{TeV} \text{ or } m_{wino} > 800 \text{GeV}$ [@95%CL: ATLAS-CONF-2015-062 13TeV, 3.2fb<sup>-1</sup>]

For *gluino*  $\rightarrow$  *tt*+*wino* or *bb*+*wino*, the constraints get a little more stringent.

Prospects ?

M<sub>gluino</sub> ~ 2.3 TeV (14TeV,300fb<sup>-1</sup>) M<sub>gluino</sub> ~ 2.7 TeV (14TeV,3000fb<sup>-1</sup>) M<sub>gluino</sub> ~5.8 TeV (33TeV, 3000fb<sup>-1</sup>)

[arXiv:1310.0077]

🖊 DM is Pure Wino Dark Matter !

Indirect search by gamma-ray from dwarf Spheroidal galaxies are promising !



Fermi-LAT 6 years data excluded the triplet dark matter in

*m*<sub>triplet</sub> < 400 GeV (classical dSphs)

[For recent J-factor estimation '16 Hayashi, Ichikawa, Matsumoto, MI, Ishigaki, Sugai]

Future prospect at CTA

Dwarf looks better target than the galactic center by taking the DM profile of the galactic center into account!

['16 Lefranca, Moulina, Panci, Sala, Silk]

# Direct SUSY searches at the LHC.

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### <u>Split Supersymmetry : Squark Mass >> PeV</u>

['04 Arkani-Hamed, Dimopoulos, Giudice, Romanino]

### *MSSM fermions* = *O*(*TeV*) *« MSSM scalars*

- Bino, Wino, Higgsinos are good candidate for dark matter
- Coupling Unification is OK
- Severe fine-tuning problem...
- ✓ Higgs mass tends to be larger than 125GeV.
- ✓ Gaugino Masses are generically 10<sup>-2</sup> x MSSM scalars...



Gluino may have a long lifetime !

 $\tau_{gluino} = 5 \times 10^{-9} \sec \theta$ 

 $x (TeV/m_{gluino})^5 x (m_{squark}/10^4 TeV)^4$ 

Current limit from heavy charged track search :  $M_{gluino} > 1.4-1.6 TeV.$ 

# ✓ Direct SUSY searches at the LHC.

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# Summary

The Standard Model is in very good shape...

(Higgs can be fit by the simple elementary doublet.)

- The Higgs boson mass at 125GeV disfavors very natural and simple SUSY models.
- In view of 125GeV Higgs boson mass, the null SUSY results so far are not very surprising/discouraging.
- Although naturalness arguments become arguable than before, still SUSY is the most attractive model beyond the Standard Model.

[SUSY can be consistent with the SM Yukawa intereactions. SUSY can be consistent with cosmology including inflation and baryogengesis such as leptogenesis.]

We still have a lot of chances to see *TeV* SUSY at the LHC ! Other channels such as DM searches are getting more important !

## <u>Although we have no strong hints at this point, we must</u> <u>keep seeking/thinking beyond the Standard Model !</u>



picture from <a href="http://www.vecteezy.com">http://www.vecteezy.com</a>

Don't give up, we don't know how near we are with new physics!

# **Back Up Slides**



#### Neutralino mixing mass

$$M_{\chi} = \begin{pmatrix} M_{1} & 0 & -\frac{1}{2}g'v\cos\beta & \frac{1}{2}g'v\sin\beta \\ 0 & M_{2} & \frac{1}{2}gv\cos\beta & -\frac{1}{2}gv\sin\beta \\ -\frac{1}{2}g'v\cos\beta & \frac{1}{2}gv\cos\beta & 0 & -\mu \\ \frac{1}{2}g'v\sin\beta & -\frac{1}{2}g'v\cos\beta & -\mu & 0. \end{pmatrix}$$
bino  
wino  
Higgino1  
Higgino2

 $\rightarrow$  lightest Neutralino is DM !

#### Main component of the LSP

bino / wino / Higgsino DM



Thermal Relic Dark Matter

Pure Bino LSP : too small cross section to be WIMP Pure Wino LSP : WIMP cross section at M<sub>wino</sub> ~ 3TeV

Pure Higgsino LSP: WIMP cross section at M<sub>Higgsino</sub> ~ 1TeV

For WIMP with  $M_{\chi}$  < TeV, we need appropriate mixing !

- $\rightarrow$  couplings to Higgs and Z tend to be unsuppressed
- $\rightarrow$  Direct detection cross sections are rather unsuppressed.

['12 Cheung, Hall, Pinner, Ruderman]



#### Neutralino DM interaction with Nucleus

$$\mathcal{L}_{\text{int}} = \frac{c_{h\chi\chi}}{2} h(\chi\chi + \chi^{\dagger}\chi^{\dagger}) \rightarrow \mathcal{L}_{\text{int}} \propto \text{DM}^2 \times \bar{\psi}_n \psi_n \rightarrow \sigma_{\text{SI}} = 8 \times 10^{-45} \text{ cm}^2 \left(\frac{c_{h\chi\chi}}{0.1}\right)^2$$
$$\mathcal{L}_{\text{int}} = c_{Z\chi\chi} \chi^{\dagger} \bar{\sigma}^{\mu} \chi Z_{\mu} \rightarrow \mathcal{L}_{\text{int}} \propto (\text{DM}^2)_{\mu} \times \bar{\psi}_n \gamma_5 \gamma^{\mu} \psi_n \rightarrow \sigma_{\text{SD}} = 3 \times 10^{-39} \text{ cm}^2 \left(\frac{c_{Z\chi\chi}}{0.1}\right)^2$$

Bino/Higgsino DM

**Bino/Wino DM** 

$$c_{h\chi\chi}, c_{Z\chi\chi} \propto \frac{(\sin\beta \pm \cos\beta)\sin\theta_W}{\sqrt{2}} \left(\frac{M_Z}{\Delta M}\right),$$

$$c_{h\chi\chi}, c_{Z\chi\chi} \propto \frac{\sin 2\beta \sin 2\theta_W}{2} \left(\frac{M_Z^2}{\mu(M_2 - M_1)}\right)$$

#### Neutralino Dark Matter



On the brown lines, the dark matter abundance is consistent with observation !

✓ Direct detection searches give complemental information to the LHC searches and the indirect searches ( $<\sigma v > ~ 10^{-9}GeV^{-2}$ ).

Pure Wino Dark Matter

Wino Dark Matter Search (direct detections,  $\chi N \rightarrow \chi N$ )



One-loop diagrams which contribute to the Wino-nucleon scatterings.

['10 Hisano, Ishiwata, Nagata]

The irreducible background from atmospheric neutrinos at about 10<sup>-48</sup>cm<sup>2</sup>. [arxiv:1003.5530]

### Direct Wino Production



✓ Main decay mode :  $\chi^{\pm} \rightarrow \chi^{0} + \pi^{\pm} : \tau_{wino} = O(10^{-10})$  sec.

Limits (disappearing track search):

<i>m<sub>wino</sub> &gt; 130GeV</i> ( <i>7TeV&amp;5fb</i> -1) using TRT	[arxiv:1210.2852]
<i>m<sub>wino</sub> &gt; <b>270GeV</b> (<b>8TeV&amp;20fb</b><sup>-1</sup>)</i> using SCT & TRT	[arXiv:1310.3675]

→ Prospects:  $M_{wino} \sim 500 \text{GeV} (14 \text{TeV}, 300 \text{fb}^{-1})$  $M_{wino} \sim 650 \text{GeV} (14 \text{TeV}, 3000 \text{fb}^{-1})$  $M_{wino} \sim 3 \text{TeV} (100 \text{TeV}, 3000 \text{fb}^{-1})$ 

[1407.7058, Cirelli, Sala, Taoso]

✓ stable stau search (GMSB with  $m_{3/2} > O(100) keV$ )



via gluino,squark production Mstau > 480 GeV (GMSB Nm = 3, Mgluino > 3.3TeV) via Drell-Yan production Mstau > 340 GeV (Run1) (Mgluino,squark >> TeV)

Gaugino Masses = Loop-suppressed (AMSB + Heavy Higgs Mediation)





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['04 Arkani-hamed,Dimopoulos,Giudice, Rommanio]

Physical CP-violation is suppressed by  $O(m_{3/2})$ .

#### For detailed analyses [`15 Hisano, Kobayashi,Kuramoto, Kuwahara]



EDM via one-loop slepton diagrams is suppressed by  $m_{slepton}^{-2}$ .

EDM are dominated by two-loop diagrams in which the light Higgs boson is circulating (suppressed by µH<sup>-1</sup>sin2β)

The current limits  $d_e/e < 8.7 \times 10^{-29}$  cm is reaching to  $\mu_H$  of O(10<sup>4</sup>) TeV ! [1310.7534 ACME : ThO]

# Wino Dark Matter Search (direct detections, $\chi N \rightarrow \chi N$ )



One-loop diagrams which contribute to the Wino-nucleon scatterings.

['10 Hisano, Ishiwata, Nagata]

Coupling to H and Z are highly suppressed for  $\mu_H = O(10-100)$  TeV at the tree-level.

Wino-Nucleon @ higher loop level  $\sigma_{p-N} = (10^{-47})cm^2$ ( much smaller than the current reach...) ['10 Hisano, Ishiwata, Nagata]

Darwin (multi-ton Argon/Xe detector) will reach down to 10<sup>-47</sup>cm<sup>2</sup> for WIMP mass below 300GeV.

The irreducible background from atmosphere neutrinos at about 10<sup>-48</sup>cm<sup>2</sup>. [arxiv:1003.5530]

### Indirect WIMP Detection (see more arXiv:1511.08787)

### Dwarf Spheroidal Galaxies !



http://astronomy.nmsu.edu/tharriso/ast110/class24.html

Object	$N_{\rm sample}$	RA(J2000) [hh:mm:ss]	DEC(J2000) [dd:mm:ss]	$M_V$	$D_{\odot}$ [kpc]	$b_*$ [pc]	q' (axial ratio)	$\operatorname{Ref.}^{a}$
Classical dwarfs								
Carina	776	06:41:36.7	-50:57:58	$-9.1 \pm 0.5$	$106 \pm 6$	$250 \pm 39$	$0.67 \pm 0.05$	1,6
Fornax	2523	02:39:59.3	-34:26:57	$-13.4\pm0.3$	$147 \pm 12$	$710 \pm 77$	$0.70 \pm 0.01$	1,6
Sculptor	1360	01:00:09.4	-33:42:33	$-11.1\pm0.5$	$86 \pm 6$	$283 \pm 45$	$0.68\pm0.03$	1,6
Sextans	445	10:13:03.0	-01:36:53	$-9.3\pm0.5$	$86 \pm 4$	$695 \pm 44$	$0.65\pm0.05$	1,6
Draco	468	17:20:12.4	+57:54:55	$-8.8\pm0.3$	$76 \pm 6$	$221\pm19$	$0.69\pm0.02$	1,7
Leo I	328	10:08:28.1	+12:18:23	$-12.0\pm0.3$	$254 \pm 15$	$251\pm27$	$0.79\pm0.03$	1,8
Leo II	200	11:13:28.8	+22:09:06	$-9.8\pm0.3$	$233 \pm 14$	$176 \pm 42$	$0.87\pm0.05$	1,9
Ultra faint dwarfs								
Segue 1	73	10:07:04.0	+16:04:55	$-1.5\pm0.8$	$32 \pm 6$	$29^{+8}_{-5}$	$0.53\pm0.10$	1,10
Segue 2	24	02:19:16.0	+20:10:31	$-2.5 \pm 0.3$	$35 \pm 2$	$35 \pm 3$	$0.85 \pm 0.13$	1,11
Boötes I	37	14:00:06.0	+14:30:00	$-6.3\pm0.2$	$66 \pm 2$	$242\pm21$	$0.61\pm0.06$	1,12
Hercules	18	16:31:02.0	+12:47:30	$-6.6\pm0.4$	$132 \pm 12$	$330^{+75}_{-52}$	$0.32 \pm 0.08$	1,13
Coma Berenices	59	12:26:59.0	+23:54:15	$-3.7\pm0.6$	$44 \pm 4$	$64 \pm 7$	$0.62\pm0.14$	1,14
Canes Venatici I	214	13:28:03.5	+33:33:21	$-7.9\pm0.5$	$224^{+22}_{-20}$	$554\pm 63$	$0.61 \pm 0.03$	1,14
Canes Venatici II	25	12:57:10.0	+34:19:15	$-4.8\pm0.6$	$151^{+15}_{-13}$	$132 \pm 16$	$0.48 \pm 0.11$	1,14
Leo IV	18	11:32:57.0	-00:32:00	$-5.1 \pm 0.6$	$158^{+15}_{-14}$	$152 \pm 17$	$0.51 \pm 0.11$	1,14
Leo V	5	11:31:09.6	+02:13:12	$-5.2 \pm 0.4$	$178 \pm 10$	$135 \pm 32$	$0.50\pm0.15$	1,15
Leo T	19	09:34:53.4	+17:03:05	$-7.1\pm0.3$	$417^{+20}_{-19}$	$170 \pm 15$	$\sim 1.00$	1,14
Ursa Major I	39	10:34:52.8	+51:55:12	$-5.6\pm0.6$	$106^{+9}_{-8}$	$308 \pm 32$	$0.20 \pm 0.04$	1,14
Ursa Major II	20	08:51:30.0	+63:07:48	$-3.8\pm0.6$	$32_{-4}^{+5}$	$127 \pm 21$	$0.37 \pm 0.05$	1,14
Reticulum II	25	03:35:42.1	-54:02:57	$-2.7 \pm 0.1$	$32 \pm 3$	$32^{+2}_{1}$	$0.41 \pm 0.03$	2,16
Draco II	9	15:52:47.6	+64:33:55	$-2.9 \pm 0.8$	$20 \pm 3$	$19^{+8}$	$0.76^{+0.27}$	3,17
Triangulum II	13	02:13:17.4	+36:10:42	$-1.8 \pm 0.5$	$30 \pm 2$	$34^{+9}$	$0.79^{+0.17}_{-0.01}$	4.18
Hydra II	13	12:21:42.1	-31:59:07	$-4.8 \pm 0.3$	$134 \pm 10$	-8 + 11	$0.99^{+0.01}$	5.19
Pisces II	7	22:58:31.0	+05:57:09	$-5.0 \pm 0.5$	$\sim 180$	$\sim 60$	$0.60 \pm 0.10$	1,19



['16 Hayashi, Ishikawa, Matsumoto, MI, Ishigaki, Sugai]

DM profile can be estimated from motions of stars.

We observe gamma ray flux from entire dwarf galaxies.

 $\rightarrow$  less sensitive to the structure of the core region!

Less active, and hence, less background gamma ray.

### Indirect WIMP Detection (see more arXiv:1511.08787)

#### **Constraints on continuous spectrum from dwarf Spheroidal Galaxy**

$$\frac{\mathrm{d}\Phi_{\gamma}}{\mathrm{d}E_{\gamma}}(E_{\gamma},\Delta\Omega) = \frac{1}{4\pi} \frac{\langle \sigma v \rangle}{2m_{\chi}^{2}} \frac{\mathrm{d}N_{\gamma}}{\mathrm{d}E_{\gamma}} \times \int_{\Delta\Omega} \int_{l.o.s} \rho_{\mathrm{DM}}^{2}(l,\Omega) dl d\Omega$$



#### WIMP cross section has been excluded for $m_{DM} < 100 \text{GeV}$ annihilating into bb!

### <u>The Higgs boson ~ 125GeV requires multi-TeV SUSY ?</u>

The simplest interpretation:  $m_H \sim 125 \text{ GeV}$  suggests that the sfermion (stop) masses are above O(10-1000)TeV ['91 Okada, Yamaguchi, Yanagida].



['12, MI, Matsumoto, Yanagida ( $\mu_H = O(M_{susy})$ )]

### The Higgs boson ~ 125GeV with a few TeV SUSY ?

A large stop A-term (trilinear coupling) can enhance the Higgs boson mass !



Having a large stop A-term (trilinear coupling) at the low energy scale is not very easy.

$$16\pi^{2} \frac{d}{dt} a_{t}$$

$$= a_{t} \Big[ 18y_{t}^{*}y_{t} + y_{b}^{*}y_{b} - \frac{16}{3}g_{3}^{2} - 3g_{2}^{2} - \frac{13}{15}g_{1}^{2} \Big] + 2a_{b}y_{b}^{*}y_{t}$$

$$+ y_{t} \Big[ \frac{32}{3}g_{3}^{2}M_{3} + 6g_{2}^{2}M_{2} + \frac{26}{15}g_{1}^{2}M_{1} \Big],$$

[e.g. GMSB realization '11 MI, Evans, Yanagida]
[see '15 Kitano, Murayama, Tobioka with KK-effect]
[Extra matter can also enhance A-term see '16 Moroi, Yanagida, Yokozaki]

 $\rightarrow$  As a whole,  $m_{squark} > TeV$  is required

### The Higgs boson ~ 125GeV with a few TeV SUSY ?

Extra Vector Matter (e.g. 10 representation ['92 Moroi, Okada])

$$10 = (Q_E, \bar{U}_E, \bar{E}_E)$$
,  $\bar{10} = (\bar{Q}_E, U_E, E_E)$ 

$$W = M_T \bar{Q}_E Q_E + M_T \bar{U}_E U_E + M_T \bar{E}_E E_E + \lambda_u H_u Q_E \bar{U}_E$$

Higgs quartic potential receives an additional radiative contributions !

$$\Delta V \simeq \frac{\Delta \lambda^2}{2} |H_u|^4 \qquad \Delta \lambda \simeq \frac{\lambda_u^2}{16\pi^2} \left( 6 \log \frac{M_{SUSY}^2}{M_T^2} - 5 \right)$$

→ 125GeV Higgs mass can be achieved as long as the extra matters have large SUSY soft masses in the TeV range.

If the soft masses of the MSSM and the extra matter have the same origin, the squarks are also in the *TeV* range though...

### <u>The Higgs boson ~ 125GeV with a few TeV SUSY ?</u>

Singlet Extension (NMSSM) [see '10 Maniatis for review]

We add a tree-level quartic term to the Higgs scalar potential :

$$W = \lambda N H_u H_d + \frac{1}{3} \kappa N^3 \quad \rightarrow \quad \Delta V \simeq \frac{\lambda^2}{4} \sin^2 2\beta \, (h^{\dagger} h)^2$$



Accordingly, the tree-level Higgs mass is lifted :

$$(m_{H_1}^{\text{NMSSM}})^2 < m_Z^2 \left( \cos^2(2\beta) + \frac{2|\lambda|^2 \sin^2(2\beta)}{g_1^2 + g_2^2} \right)$$

 $H_1$ 

b

Still, we need some radiative contributions to achieve *125GeV* from the squarks in a *TeV* range. ( $\lambda < 0.7$  from perturbativity)

[For the scenarios where the lightest Higgs is not the observed one see '13 Christensen, Han, Liu, Su].  $e^+$ 

 $e^{-}$ 

### <u>The Higgs boson ~ 125GeV with a few TeV SUSY ?</u>

Extended Gauge Symmetry ['06 Maloney, Pierce, Wacker]

If the gauge symmetry of the MSSM is extended, the Higgs obtains tree-level quartic *D-term* potential



We need a rather large *g* and at least *TeV* soft mass to the gauge symmetry breaking sector.

If the soft masses of the MSSM and the gauge breaking sector have the same origin, the squarks are again in the *TeV* range.