## Parametrization of F<sub>2</sub> structure functions for inclusive γ\*p and γ\*A reactions in low Q<sup>2</sup> region

#### Hiroyuki Kamano (KEK)

#### **Collaborator: Shunzo Kumano (KEK)**

Workshop on "neutrino-nucleus interaction in a few GeV region" KEK Tokai Campus, November 18-19, 2017

#### Outline

**1.** Introduction: Neutrino collaboration at the J-PARC Branch of KEK Theory Center

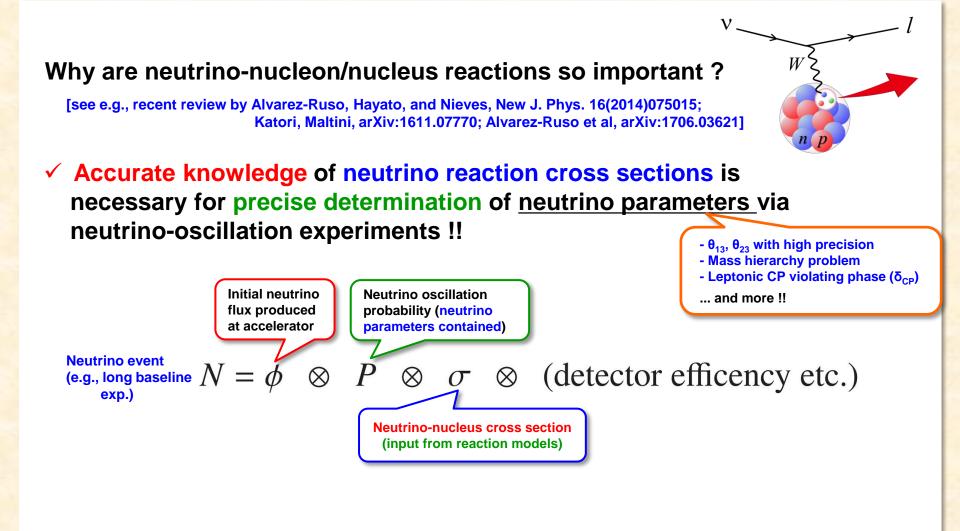
#### 2. Approach to describing F<sub>2</sub> structure functions in high-W & low-Q<sup>2</sup> region ("Regge" region)

#### **3.** Summary and future works

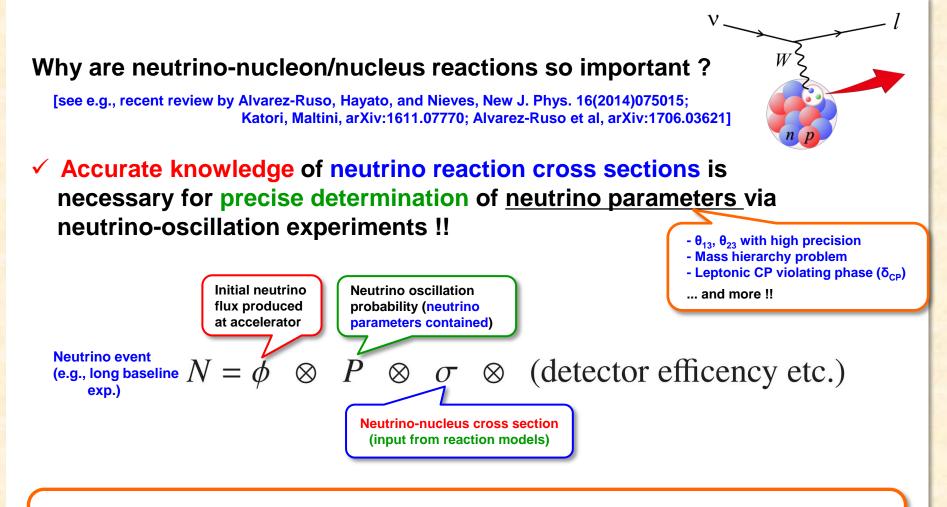
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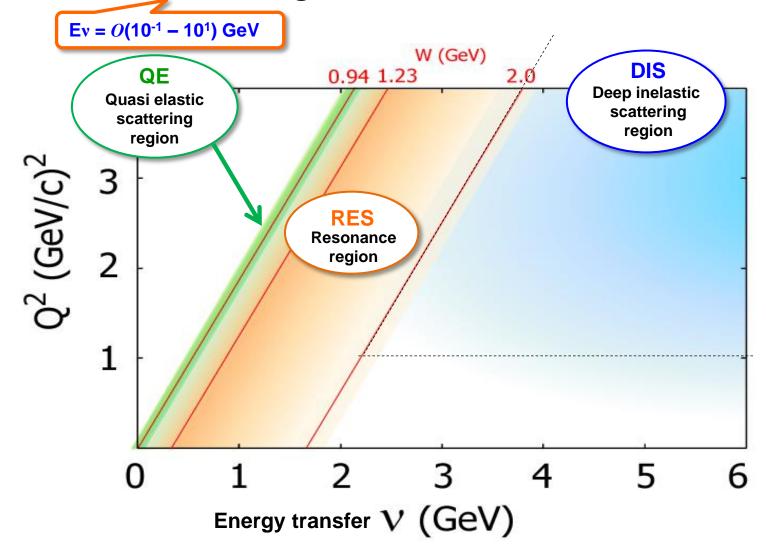
#### **Background and motivation**

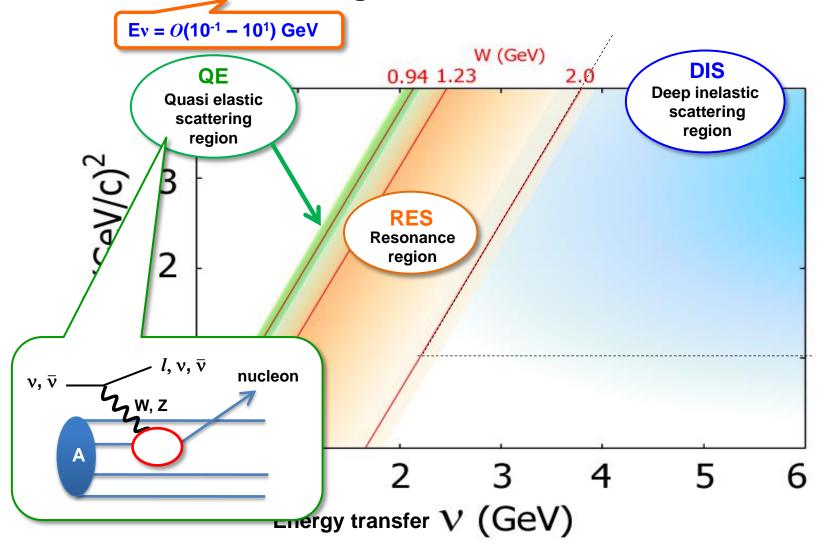


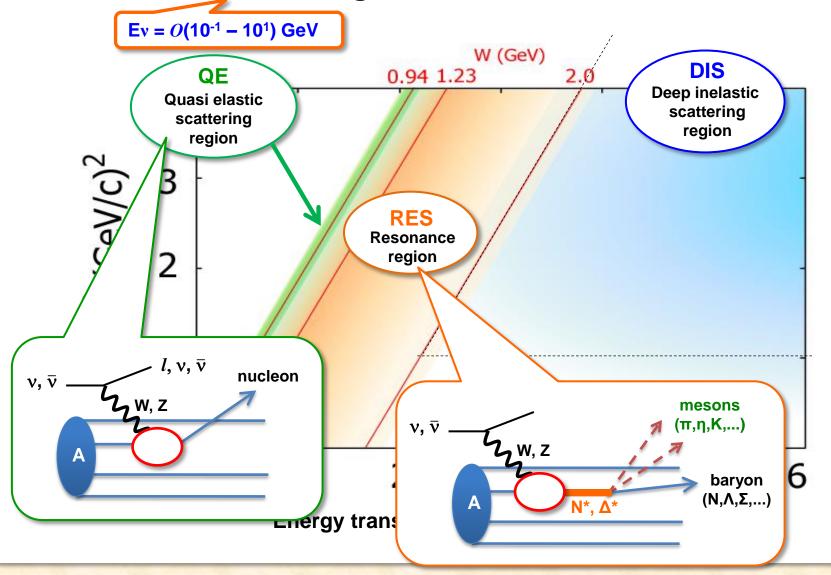
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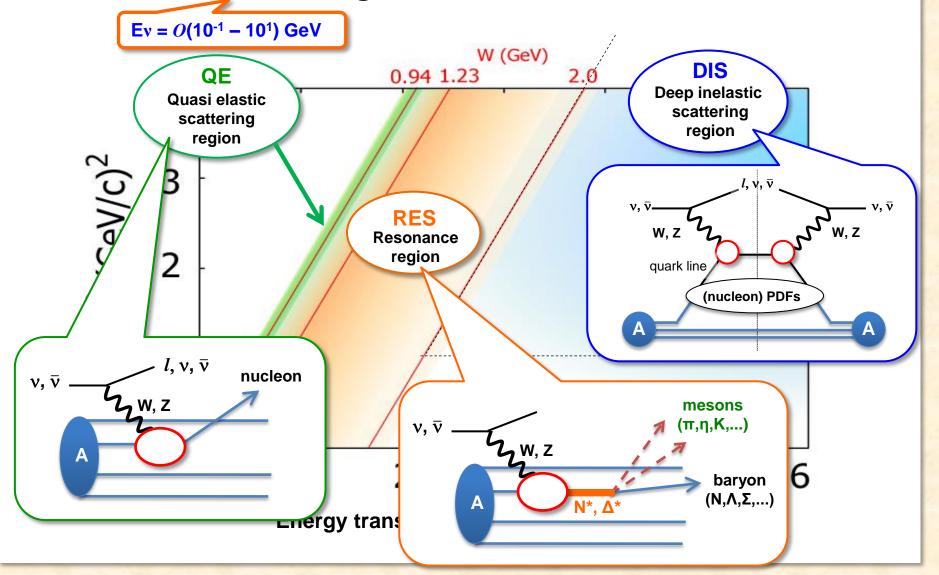


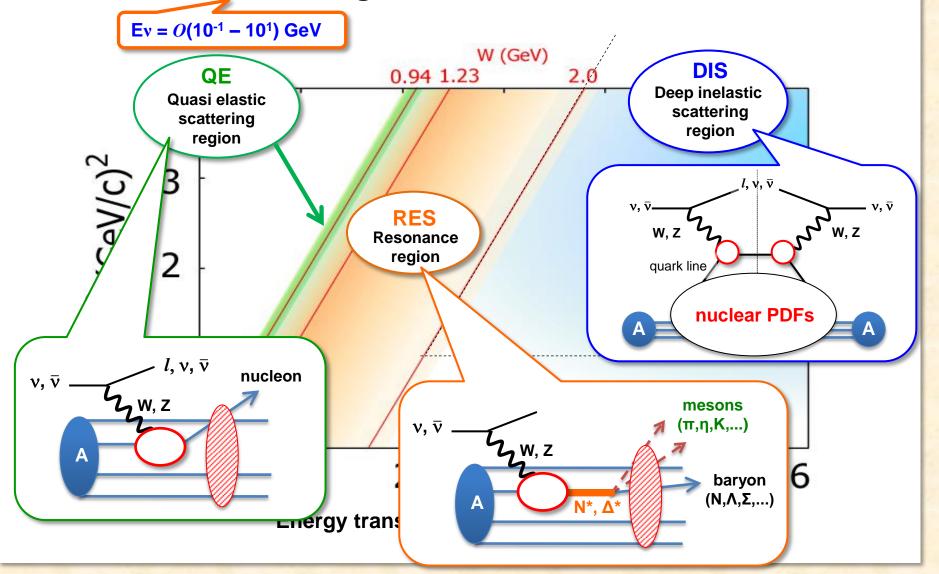
Need a reliable neutrino reaction model that describes various neutrino-nucleus reactions at the level of a few percent accuracy !!

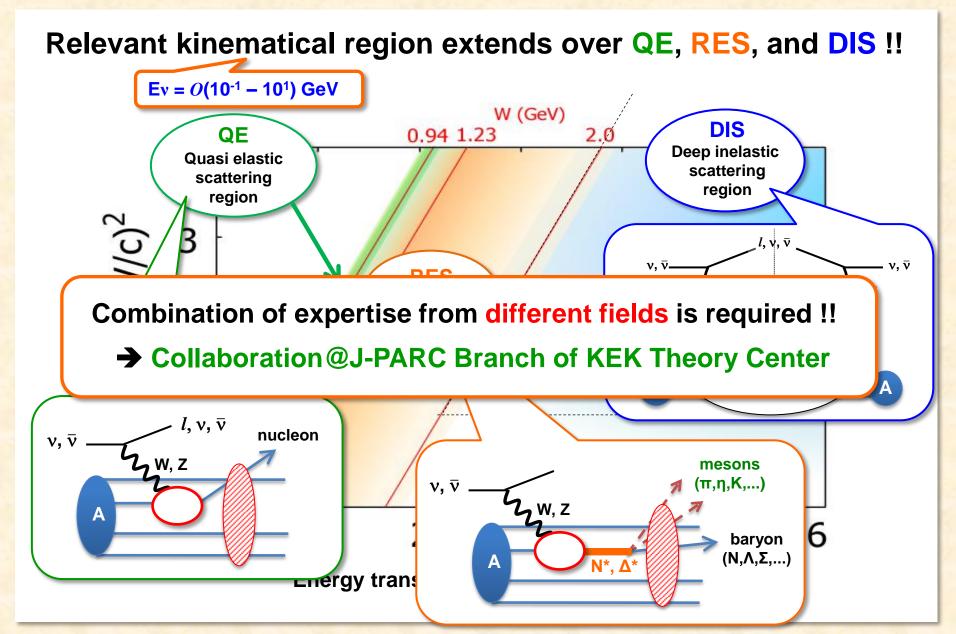








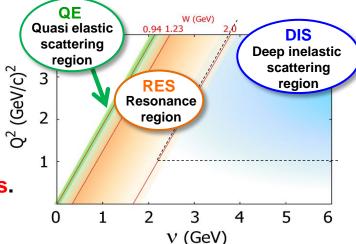




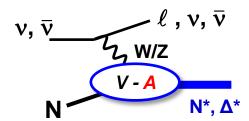
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GOAL: Construct a 'unified' model that comprehensively describes neutrino-nucleon/nucleus reactions over QE, RES, and DIS !!

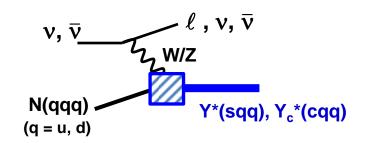
- Deepen our knowledge on complicated neutrino interactions with nucleus
- Help improve neutrino event generators.
- Investigate the internal structure of nucleon, baryon resonances, and nuclei with weak probes. (transition form factors, nuclear PDFs, ...)



Axial transition form factors



 $\checkmark$  |ΔS|=1(u→s), |ΔC|=1(d→c) transition form factors



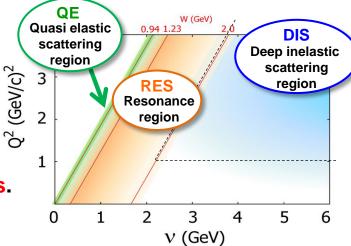
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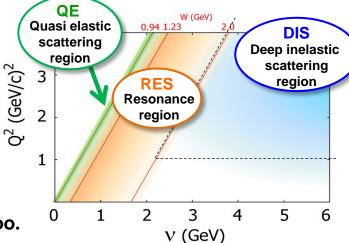
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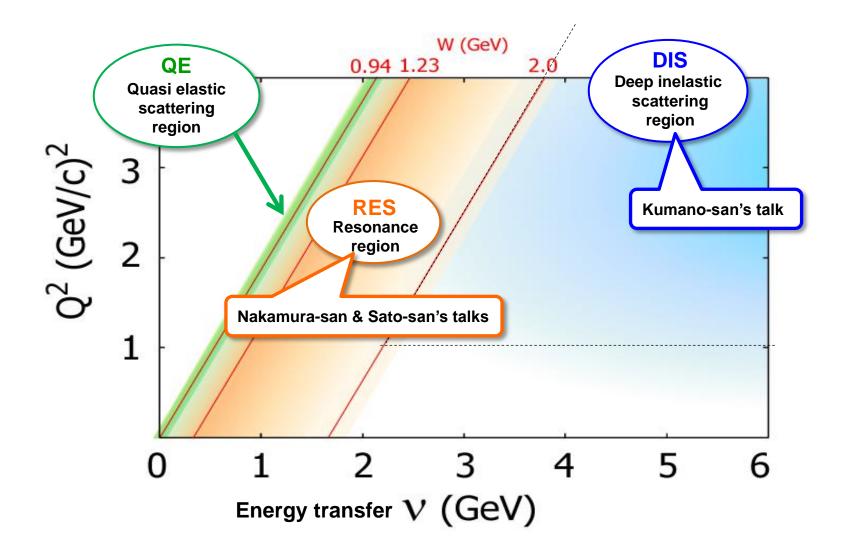
#### Strategy

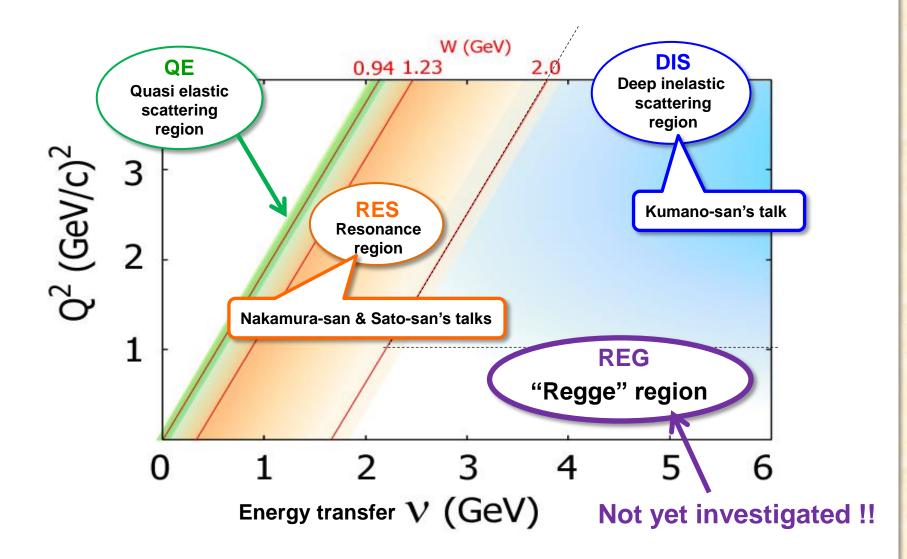
- 1. Construct 'baseline' models describing each of the kinematical regions individually.
  - constructed with appropriate effective d.o.f.s in each region
- 2. Connect the 'baseline' models
  - ➔ accomplished by matching observed quantities
  - → describes transition regions (QE⇔RES, RES⇔DIS), too.

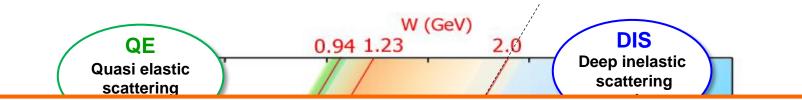
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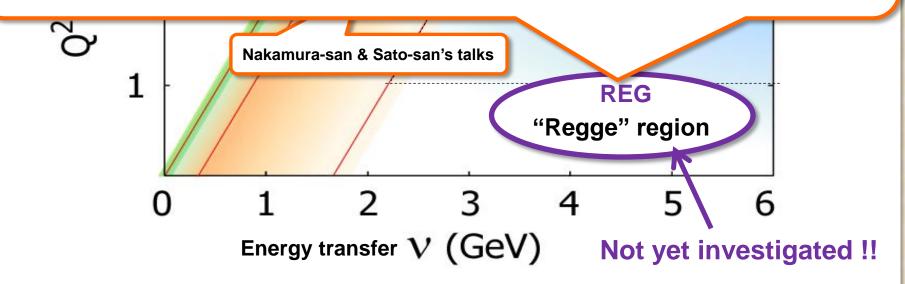




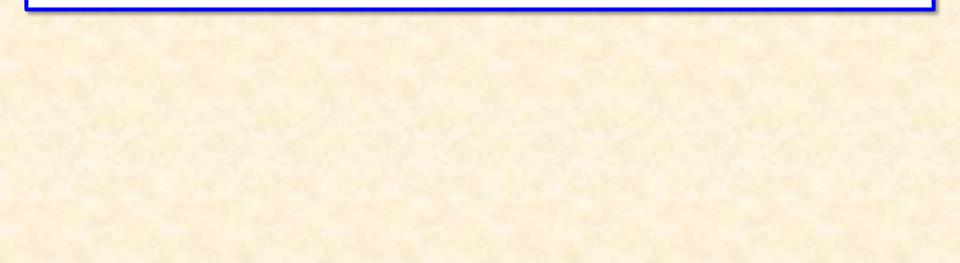




In this work, we focus on charged-lepton (e.m. current)  $F_2$  structure function before tackling on more complicated neutrino (CC and NC current) structure functions.



# 2. Approach to parametrizing F<sub>2</sub> structure functions in high-W & low-Q<sup>2</sup> region ("Regge" region)



- Purpose of the present work:
  - Develop a model/parametrization for e.m. F<sub>2</sub> structure functions in high-W and low-Q<sup>2</sup> region (= "Regge" region).
- ✓ (Practical) "requirement":
  - The model/parametrization should be as simple as possible for experimental use.

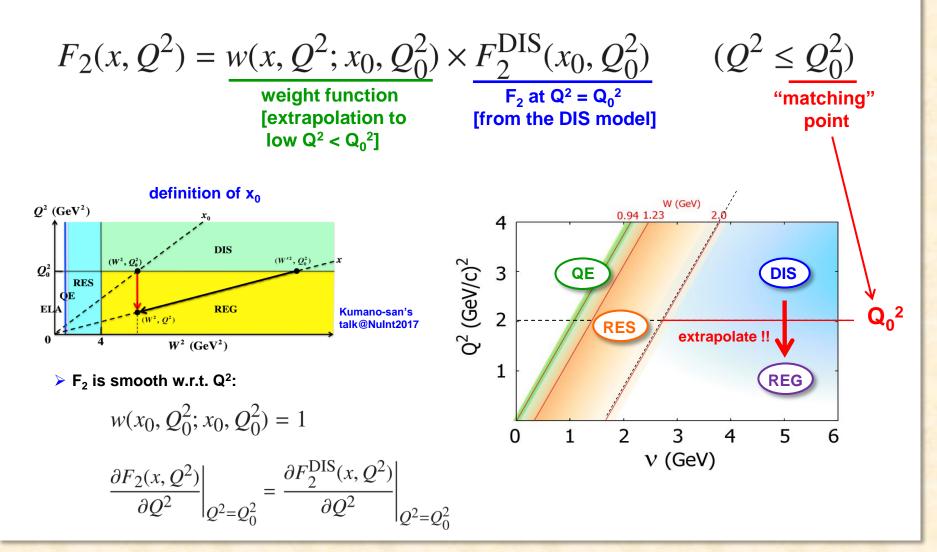
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     (> Treatment of multi-meson productions are very complicated.)
  - However, an accurate structure-function model is available in the DIS region (applicable at Q<sup>2</sup> > 1 GeV<sup>2</sup>; W > 2 GeV).
     [→ Hirai, Kumano, Nagai, PRC76(2007)065207]

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"Extrapolate" the DIS model to the low-Q<sup>2</sup> region !!

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$$F_{2}(x, Q^{2}) = \underbrace{w(x, Q^{2}; x_{0}, Q_{0}^{2})}_{\substack{\text{weight function}\\ [\text{extrapolation to}\\ \text{low } Q^{2} < Q_{0}^{2}]}} \times \underbrace{F_{2}^{\text{DIS}}(x_{0}, Q_{0}^{2})}_{\substack{\text{F}_{2} \text{ at } Q^{2} = Q_{0}^{2}\\ [\text{from the DIS model}]}} \qquad (Q^{2} \le Q_{0}^{2})$$

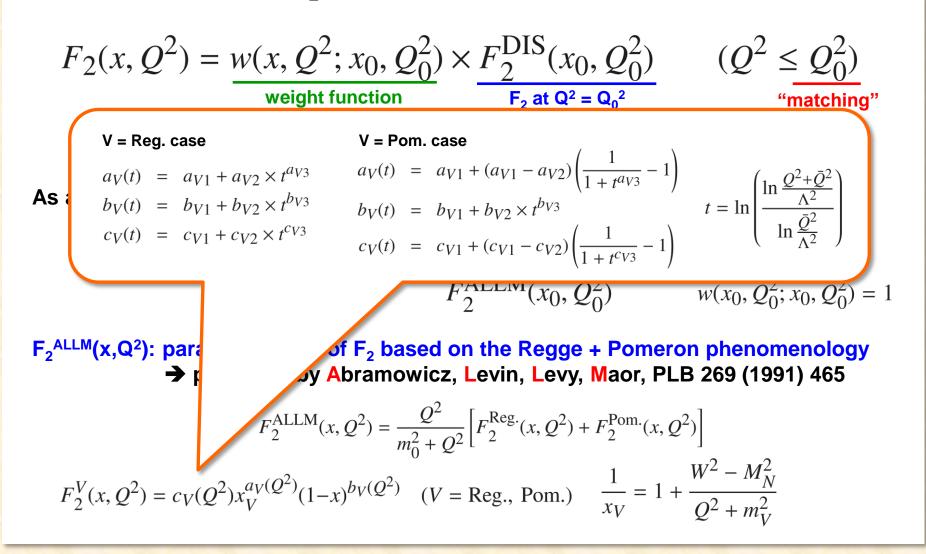
As a first attempt, we parametrize the weight function as:

$$w(x, Q^2; x_0, Q_0^2) \equiv \frac{F_2^{\text{ALLM}}(x, Q^2)}{F_2^{\text{ALLM}}(x_0, Q_0^2)} \Rightarrow \text{automatically satisfies}$$
$$w(x_0, Q_0^2; x_0, Q_0^2) = 1$$

F<sub>2</sub><sup>ALLM</sup>(x,Q<sup>2</sup>): parametrization of F<sub>2</sub> based on the Regge + Pomeron phenomenology → proposed by Abramowicz, Levin, Levy, Maor, PLB 269 (1991) 465

$$F_2^{\text{ALLM}}(x, Q^2) = \frac{Q^2}{m_0^2 + Q^2} \Big[ F_2^{\text{Reg.}}(x, Q^2) + F_2^{\text{Pom.}}(x, Q^2) \Big]$$
  
$$F_2^V(x, Q^2) = c_V(Q^2) x_V^{a_V(Q^2)}(1-x)^{b_V(Q^2)} \quad (V = \text{Reg., Pom.}) \quad \frac{1}{x_V} = 1 + \frac{W^2 - M_N^2}{Q^2 + m_V^2}$$

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- **1.** Determines parameters contained in  $w(x,Q^2; x_0,Q_0^2)$  by fitting to the existing data for the proton target:
  - > Electromagnetic  $F_2$  structure functions (for low  $Q^2 < Q_0^2$ )
  - > inclusive  $\gamma + p \rightarrow X$  total cross section (for  $Q^2 = 0$ )
  - > imposing "smoothness" of  $F_2$  function at  $Q^2 = Q_0^2$

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  - > imposing "smoothness" of  $F_2$  function at  $Q^2 = Q_0^2$
- 2. Estimates the nuclear  $F_2$  functions by replacing the DIS part:  $F_2^{\text{proton}}(x, Q^2) = w(x, Q^2; x_0, Q_0^2) \times F_2^{\text{proton, DIS}}(x_0, Q_0^2)$  $F_2^{\text{nucleus}}(x, Q^2) = w(x, Q^2; x_0, Q_0^2) \times F_2^{\text{nucleus, DIS}}(x_0, Q_0^2)$

#### $\checkmark$ Available F2 data for Q<sup>2</sup> < 2 GeV<sup>2</sup> (Q<sub>0</sub><sup>2</sup> = 2 GeV<sup>2</sup>) and W > 1.8 GeV taken from the following groups.

[Data are taken from http://hepdata.cedar.ac.uk/review/f2/ unless otherwise stated below.]

- **E665 (38 points)**
- > NMC (23 points)
- > H1 (31 points)
- H1ZEUS (101 points) [d09-158.nce+p.dat, d09-158nce-p.dat in https://www.desy.de/h1zeus/combined results/]
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- SLAC (117 points)
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(229 points) <u>http://pdg.lbl.gov/2016/hadronic-xsections/rpp2014-gammap\_total.dat</u>

Impose smoothness condition via  $\chi^2$ :

$$\chi^{2} = \left\{ 1 - \left( \left[ \frac{\partial F_{2}(x, Q^{2})}{\partial Q^{2}} \right]_{Q^{2} = Q_{0}^{2}} / \left[ \frac{\partial F_{2}^{\text{DIS}}(x, Q^{2})}{\partial Q^{2}} \right]_{Q^{2} = Q_{0}^{2}} \right) \right\}^{2} / ("\text{Error"})^{2}$$
Computed by
Computed with

taking difference DGLAP eq.

 $x = 1e-5, 0.025, 0.50, \dots, 0.45, 0.475, 0.5$  (21 points) 5% error is assigned for each point

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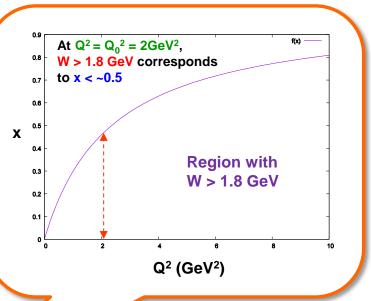
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Computed by taking difference

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**Total 796 points** 

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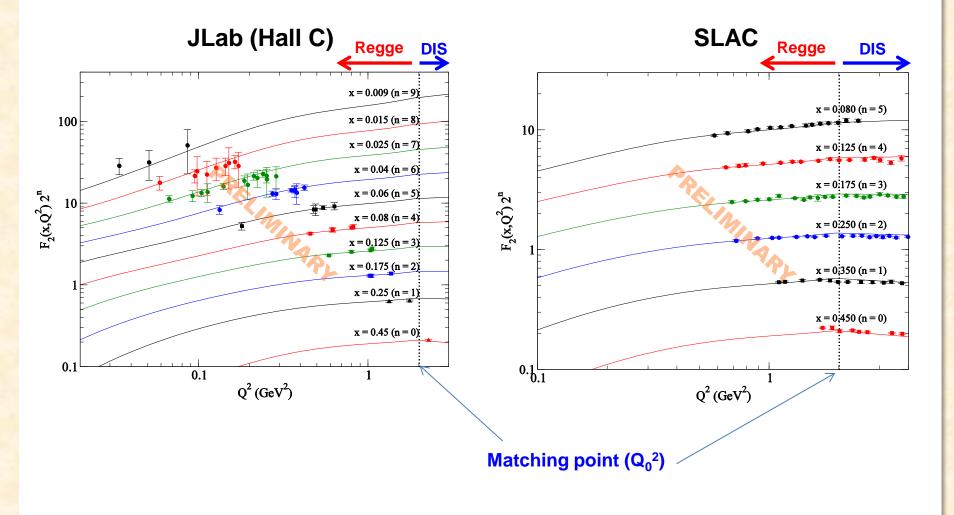
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Resulting  $\chi^2/d.o.f = 1.39$ 

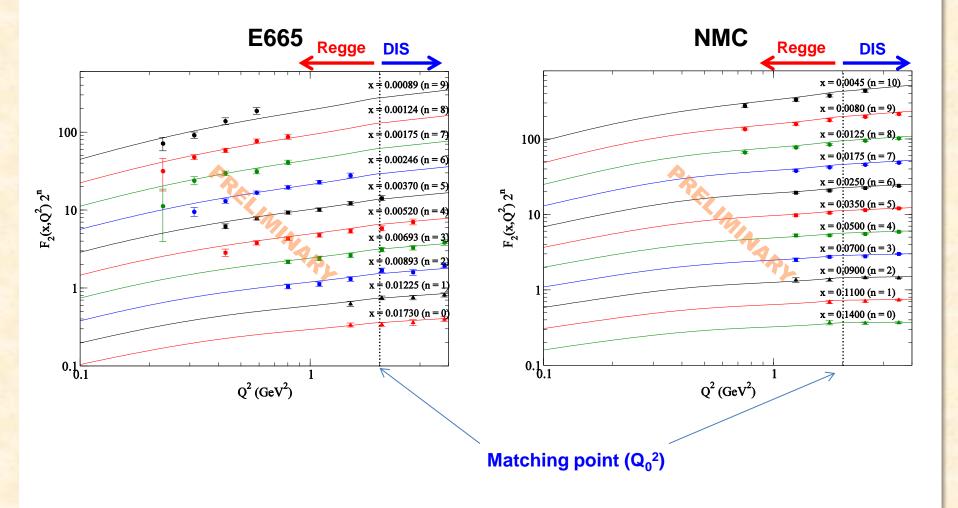
#### **Comparison with data at low Q<sup>2</sup>**

#### E.M. current F<sub>2</sub> structure function for the proton

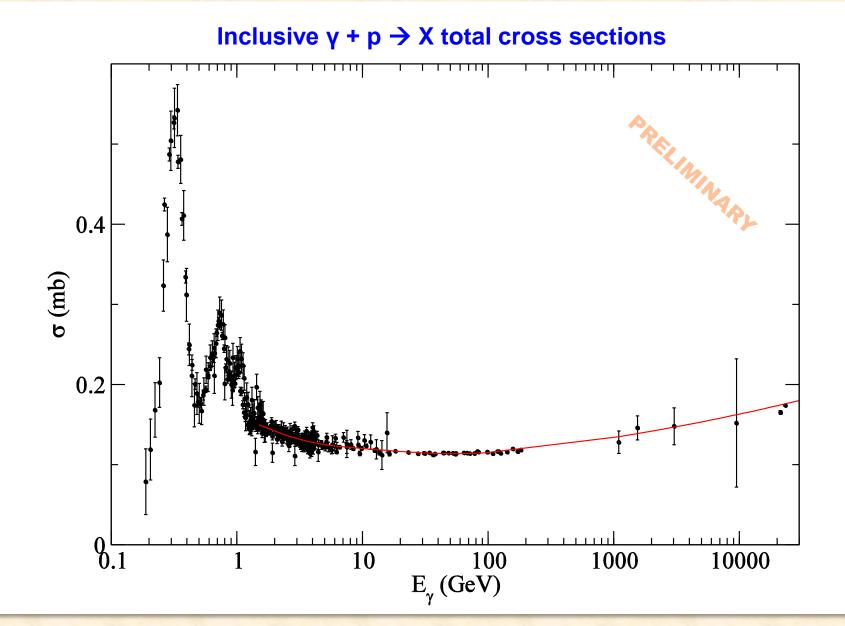


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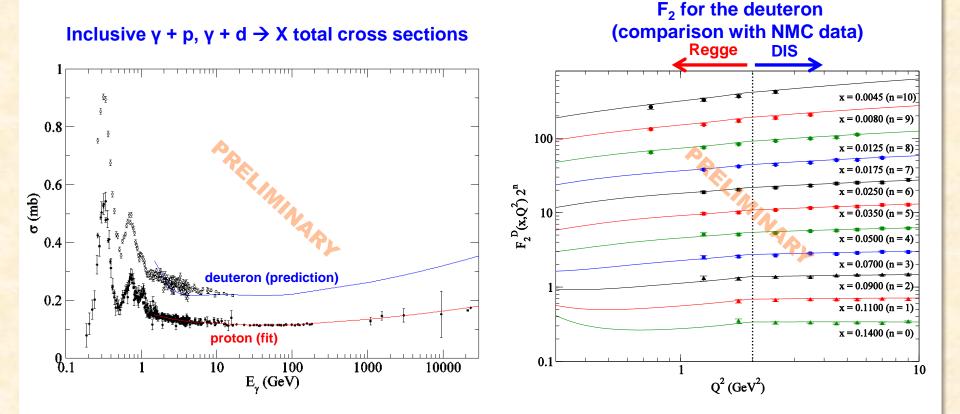
#### **Comparison with data at Q^2 = 0**



#### **Prediction for deuteron target**

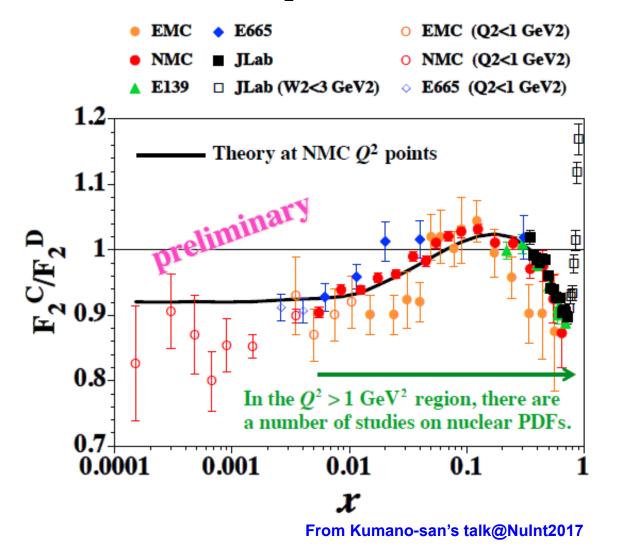
As a first attempt, try to estimate  $F_2^{deuteron}$  by replacing  $F_2^{proton, DIS}$  with  $F_2^{deuteron, DIS}$ :

$$F_2^{\text{deuteron}}(x, Q^2) = w(x, Q^2; x_0, Q_0^2) \times F_2^{\text{deuteron, DIS}}(x_0, Q_0^2)$$



#### **Prediction for nuclear target**

#### Ratio for Carbon/Deuteron F<sub>2</sub> structure functions



#### **Summary & future works**

#### Summary

 Investigated a possible, simple parametrization of the nucleon/nuclear structure functions in the Regge region (high-energy & low-Q<sup>2</sup> region).

→ important in analyzing the data from neutrino-oscillation experiments

 Accomplished by extrapolating a well-established model in the DIS (high energy & high Q<sup>2</sup>) region to the low Q<sup>2</sup> region:

$$F_2(x,Q^2) = w(x,Q^2;x_0,Q_0^2) \times F_2^{\text{DIS}}(x_0,Q_0^2) \qquad (Q^2 \le Q_0^2)$$

→ the DIS part is taken from Hirai, Kumano, Nagai PRC76(2007)065207
 → the weight function is parametrized by making use of the ALLM model [PLB269(1991)465]

#### Future works

Based on Reggeon, Pomeron phenomenology

- Better parameterization for the weight function w(x,Q<sup>2</sup>;x<sub>0</sub>,Q<sub>0</sub><sup>2</sup>) ??
- Further improvements of nuclear modification.
- Application to the neutrino-nucleus structure functions

## Back up

#### The ALLM parametrization of $F_2$ for $\gamma^{(*)}$ p inclusive reactions

Abramowicz, Levin, Levy, Maor, PLB269(1991)4656

$$F_2(x,Q^2) = \frac{Q^2}{Q^2 + m_0^2} \left( F_2^{\mathcal{P}}(x,Q^2) + F_2^{\mathcal{R}}(x,Q^2) \right)$$

 $\mathcal{P}$  = Pomeron-exchange contribution,  $\mathcal{R}$  = Regge-exchange contribution

$$F_2^{\mathcal{P}}(x,Q^2) = c_{\mathcal{P}}(t) x_{\mathcal{P}}^{a_{\mathcal{P}}(t)} (1-x)^{b_{\mathcal{P}}(t)},$$
  

$$F_2^{\mathcal{R}}(x,Q^2) = c_{\mathcal{R}}(t) x_{\mathcal{R}}^{a_{\mathcal{R}}(t)} (1-x)^{b_{\mathcal{R}}(t)}.$$

$$t = \ln\left(\frac{\ln\frac{Q^2 + Q_0^2}{\Lambda^2}}{\ln\frac{Q_0^2}{\Lambda^2}}\right)$$

Monotonically increasing in  $Q^2$ . At  $Q^2 = 0$ , t = 0.

$$\frac{1}{x_{\mathcal{P}}} = 1 + \frac{W^2 - M^2}{Q^2 + m_{\mathcal{P}}^2}, \\ \frac{1}{x_{\mathcal{R}}} = 1 + \frac{W^2 - M^2}{Q^2 + m_{\mathcal{R}}^2}.$$

If  $m_P = m_R = 0$ ,  $x_P$  and  $x_R$  reduces to x.  $x_P$  and  $x_R$  are nonzero even at  $Q^2 = 0$ . (cf. x = 0 at  $Q^2 = 0$ )

#### **The ALLM parametrization of F**<sub>2</sub> **for γ<sup>(\*)</sup>p inclusive reactions**

Abramowicz, Levin, Levy, Maor, PLB269(1991)4656

$$F_{2}^{\mathcal{P}}(x,Q^{2}) = c_{\mathcal{P}}(t)x_{\mathcal{P}}^{a_{\mathcal{P}}(t)}(1-x)^{b_{\mathcal{P}}(t)},$$

$$F_{2}^{\mathcal{R}}(x,Q^{2}) = c_{\mathcal{R}}(t)x_{\mathcal{R}}^{a_{\mathcal{R}}(t)}(1-x)^{b_{\mathcal{R}}(t)}.$$

$$a_{\mathcal{R}}(t) = a_{\mathcal{R}1} + a_{\mathcal{R}2} \times t^{a_{\mathcal{R}3}} \qquad a_{\mathcal{P}}(t) = a_{\mathcal{P}1} + (a_{\mathcal{P}1} - a_{\mathcal{P}2})\left(\frac{1}{1+t^{a_{\mathcal{P}3}}} - 1\right)$$

$$b_{\mathcal{R}}(t) = b_{\mathcal{R}1} + b_{\mathcal{R}2} \times t^{b_{\mathcal{R}3}} \qquad b_{\mathcal{P}}(t) = b_{\mathcal{P}1} + b_{\mathcal{P}2} \times t^{b_{\mathcal{P}3}}$$

$$c_{\mathcal{R}}(t) = c_{\mathcal{R}1} + c_{\mathcal{R}2} \times t^{c_{\mathcal{R}3}} \qquad c_{\mathcal{P}}(t) = c_{\mathcal{P}1} + (c_{\mathcal{P}1} - c_{\mathcal{P}2})\left(\frac{1}{1+t^{c_{\mathcal{P}3}}} - 1\right)$$

At t = 0 ( $Q^2 = 0$ ), only the first term survives.

 $a_R$ ,  $b_R$ ,  $c_R$ ,  $b_P \rightarrow$  Supposed to be monotonically increasing function of t. If  $a_{R3}$ ,  $b_{R3}$ ,  $c_{R3}$ ,  $b_{P3}$  positive (negative), then  $a_{R2}$ ,  $b_{R2}$ ,  $c_{R2}$ ,  $b_{P2}$ , must be positive (negative).

 $a_{P}, c_{P} \rightarrow$  Supposed to be monotonically decreasing function of t.  $a_{P1} > a_{P2}, a_{P3} > 0$  or  $a_{P1} < a_{P2}, a_{P3} < 0$ ;  $c_{P1} > c_{P2}, c_{P3} > 0$  or  $c_{P1} < c_{P2}, c_{P3} < 0$ 

 $c_R, c_P \rightarrow$  Must be positive for all t.  $c_{R1}, c_{P1}$  must be positive.