

# Five-body resonances of ${}^8\text{C}$ using the complex scaling method

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Recently, the new experiment on  ${}^8\text{C}$  has been reported [1]. The  ${}^8\text{C}$  nucleus is known as an unbound system beyond the proton drip-line and decays into many-body channels of  ${}^7\text{B}+p$ ,  ${}^6\text{Be}+2p$ ,  ${}^5\text{Li}+3p$  and  ${}^4\text{He}+4p$ . In this report, we present our recent study on the resonance spectroscopy of  ${}^8\text{C}$ . We employ the cluster-orbital shell model of the  ${}^4\text{He}+p+p+p+p$  five-body system, and describe the many-body resonances under the correct boundary conditions by using the complex scaling method. We adopt the Hamiltonian, the nuclear part of which reproduces the  ${}^4\text{He}-n$  scattering data and the  ${}^6\text{He}$  energy[2, 3]. The mirror nucleus of  ${}^8\text{C}$  is  ${}^8\text{He}$ , which is known as a neutron skin nucleus. It is interesting to examine the mirror symmetry between the proton-rich  ${}^8\text{C}$  and the neutron-rich  ${}^8\text{He}$ .

We show the level structures of  ${}^5\text{Li}$ ,  ${}^6\text{Be}$ ,  ${}^7\text{B}$  and  ${}^8\text{C}$  in Fig. 1. It is found that the present calculations agree with the observations and predict more energy levels. In Fig. 2, we compare the excitation energy spectra of proton-rich and neutron-rich sides. The good symmetry is confirmed between the corresponding nuclei. The differences of excitation energies for individual levels are less than 1 MeV.

We calculate the pair numbers of four valence protons in  ${}^8\text{C}$ , using the operator  $\sum_{\alpha<\beta} A_{J^\pi,S}^\dagger(\alpha\beta)A_{J^\pi,S}(\alpha\beta)$ . Here,  $\alpha$  and  $\beta$  represent the proton single-particle orbits and  $A_{J^\pi,S}^\dagger$  ( $A_{J^\pi,S}$ ) is the creation (annihilation) operator of a proton-pair with spin-parity  $J^\pi$  and the coupled intrinsic spin  $S$ . The total pair number is six for each state of  ${}^8\text{C}$ . This quantity helps us to understand the pair coupling behavior of four protons. Figure 3 shows the pair numbers for  ${}^8\text{C}$  ( $0_{1,2}^+$ ). In the  $0_1^+$  state, the  $2^+$  neutron pair is close to five and the  $0^+$  pair is almost unity. This is obtained from a main configuration of  $(p_{3/2})^4$  with the probability of 88% using CFP (1 and 5 for  $0^+$  and  $2^+$ , respectively). The  $0_2^+$  state has almost two  $0^+$  proton pairs in addition to the  $2^+$  pairs. This is consistent with the  $(p_{3/2})^2(p_{1/2})^2$  configuration with a probability of 93%; this configuration is decomposed into the pairs of  $0^+$ ,  $1^+$  and  $2^+$  with occupations of 2, 1.5, and 2.5, respectively. It is found that in the  $0_2^+$  state, the spin-singlet and the spin-triplet components are equally mixed in the  $0^+$  proton pair.

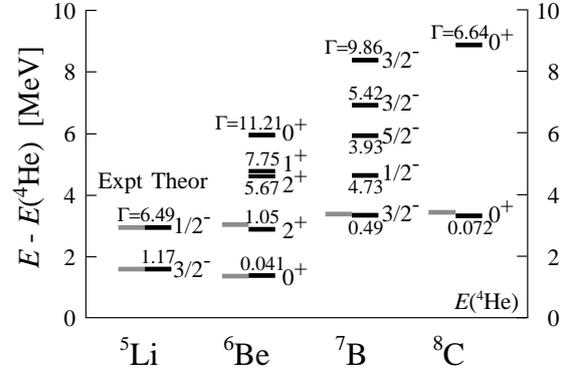


Figure 1: Energy levels of  ${}^5\text{Li}$ ,  ${}^6\text{Be}$ ,  ${}^7\text{B}$  and  ${}^8\text{C}$  measured from the energy of  ${}^4\text{He}$ . Small numbers are theoretical decay widths in units of MeV.

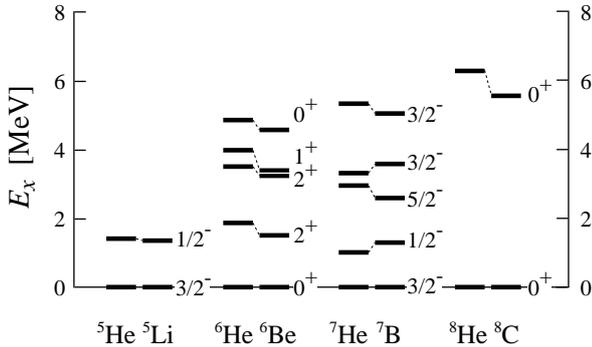


Figure 2: Excitation energy spectra of mirror nuclei.

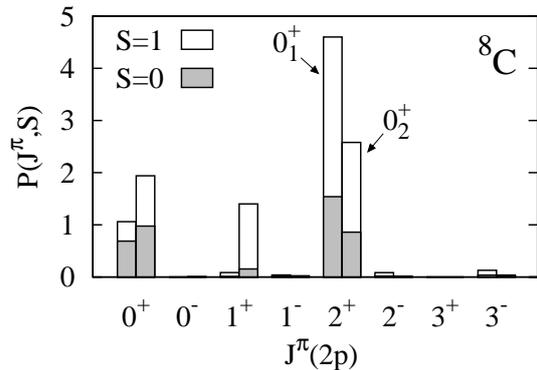


Figure 3: Pair numbers of the  $0_{1,2}^+$  states of  ${}^8\text{C}$ .

## References

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