

# *Determining the $\square^+$ quantum numbers through the $K^+p \rightarrow \square^+ KN$ reaction*



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## Contents

### 1. Introduction

1. Status of theoretical studies
2. Photo-production process

### 2. The $K^+ p \rightarrow \square^+ K^+ n (K^0 p)$ reaction

1. Motivation & advantage
2. Model for the reaction
3. Numerical results
4. Conclusions

### 3. Future (current) works

1. Full calculation
2.  $K^*$  production

## Introduction

$\square^+$  : 5-quark (4 quark + 1 anti-quark)

LEPS, T. Nakano *et al.*, Phys. Rev. Lett. 91 (2003) 012002

**Quantum numbers are not yet determined**

### Theory prediction

D. Diakonov *et al.* (chiral quark soliton) :  $1/2^+$ , I=0

Naive quark model :  $1/2^\square$

S. Capstick *et al.* (isotensor formulation) :  $1/2^\square, 3/2^\square, 5/2^\square$ , I=2

A. Hosaka (chiral potential) :  $1/2^+$  (strong  $\square$ )

R. L. Jaffe *et al.* (qq-qq- $\bar{q}$  :  $\overline{10} + 8$ ) :  $1/2^+$ , I=0

J. Sugiyama *et al.* (QCD sum rule) :  $1/2^\square$ , I=0

F. Csikor *et al.* (Lattice QCD) :  $1/2^+ \rightarrow 1/2^\square$

S. Sasaki (Lattice QCD) :  $1/2^\square$

## Photo-production process

**Assuming the quantum numbers (spin, parity), we can calculate a reaction**



W. Liu *et al.* nucl-th/0308034

S. I. Nam *et al.* hep-ph/0308313

W. Liu *et al.* nucl-th/0309023

Y. Oh *et al.* hep-ph/0310117

- **Model (mechanism) dependence**

Initial cm energy  $\sim 2$  GeV ( $p_{cm} \sim 750$  MeV)

not low energy  $\rightarrow$  linear or nonlinear?

$N^*$  resonances,  $K^*$  exchange,  $\bar{\chi}_1$  exchange, ...

- **Form factor dependence**

Monopole, dipole... , value of  $\alpha$ , ...

- **Unknown parameters**

$\bar{\chi}\chi$  coupling,  $K^* p\bar{q}$  coupling, ...

## Motivation and advantage

We propose



- Low energy model is sufficient ( $p_{cm} \sim 350$  MeV)
- take decay into account  $\rightarrow$  background estimation  
 $\rightarrow$  Width independent
- Hadronic process : clear mechanism

to extract a qualitative behavior which depends  
on the quantum numbers of  $\square^+$ .

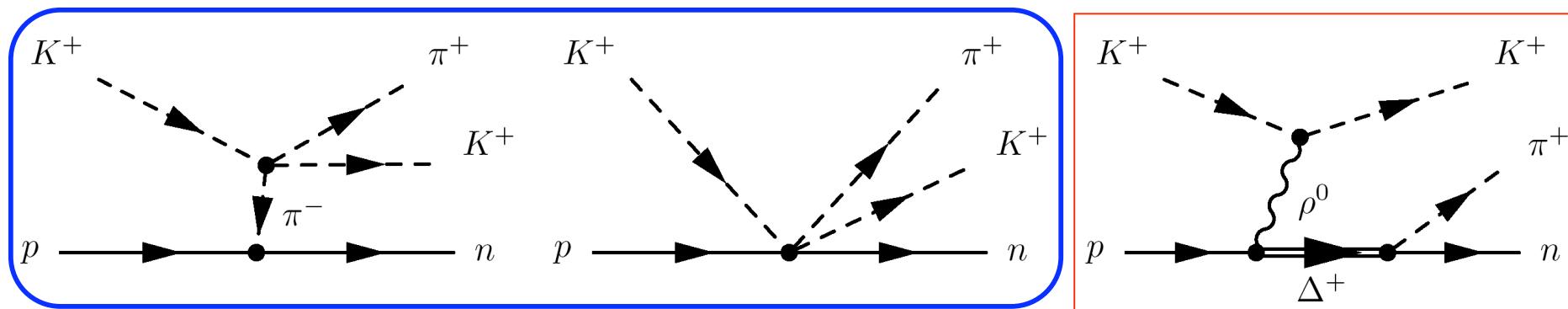


Determination of quantum numbers

## A model for $\bar{\psi}^+ p \rightarrow \bar{\psi}^+ K^+ n$

E. Oset and M. J. Vicente Vacas, PLB386, 39(1996)

Vertices are derived from the chiral Lagrangian



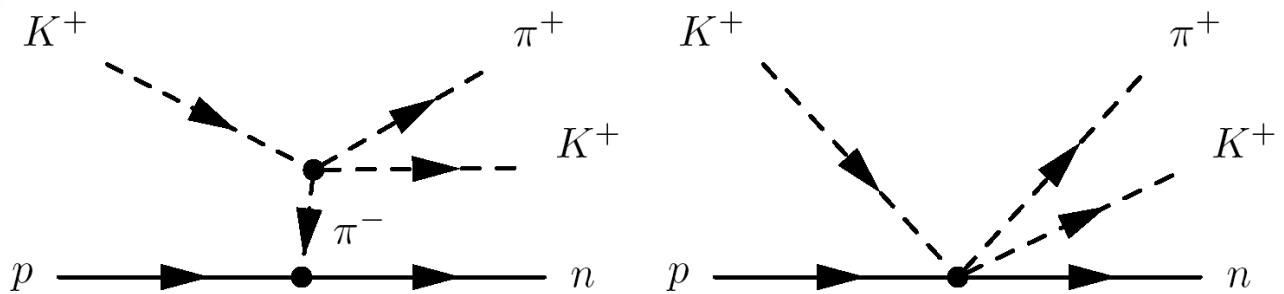
Dominant

Proportional to  $S \cdot p_{\pi^+}$   
vanishes

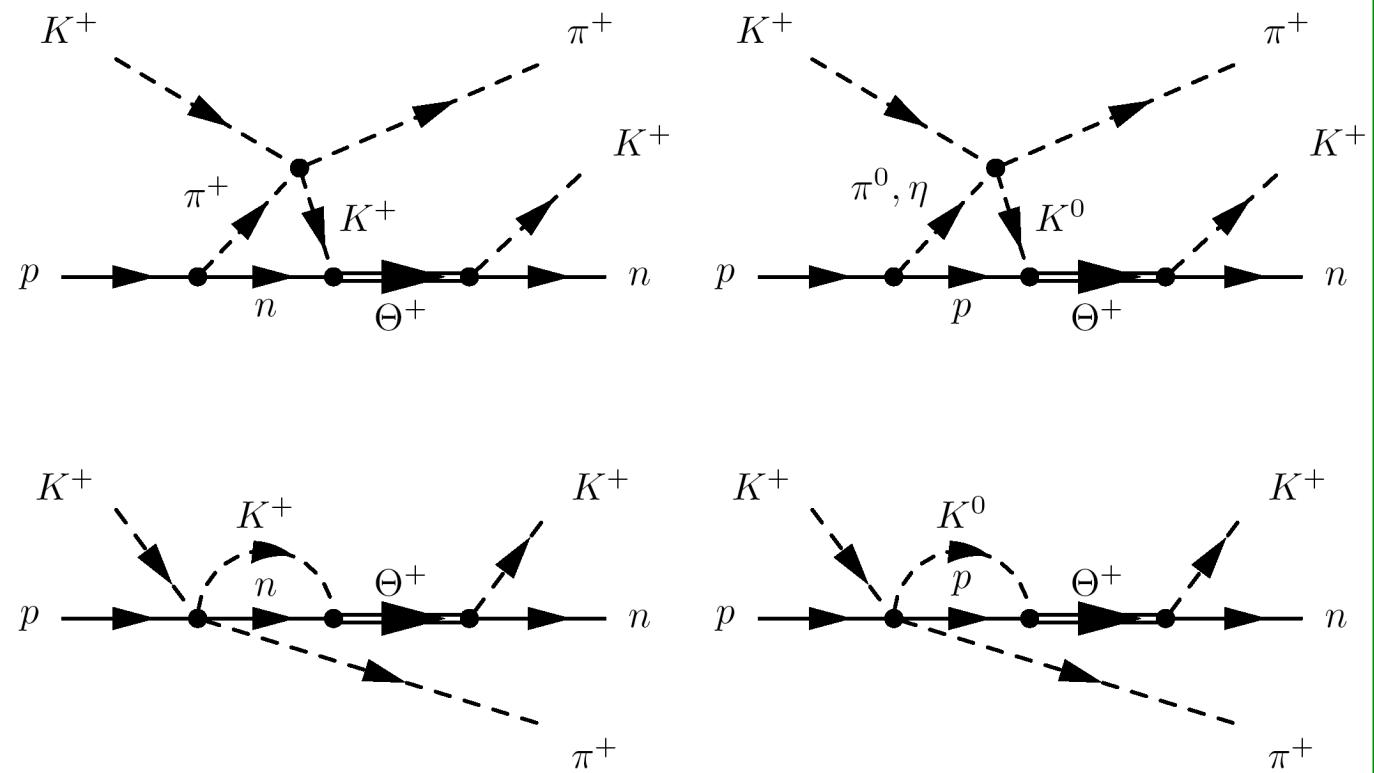
Assume final  $\bar{\psi}^+$  is almost at rest

# Diagrams

**Tree level  
(background)**



**One loop**



## Possibilities of spin & parity

**1/2 $\square$** (KN s-wave resonance)

**M<sub>R</sub> = 1540 MeV**

**1/2<sup>+</sup>, 3/2<sup>+</sup> (KN p-wave resonance)**

**$\square = 20 \text{ MeV}$**

$$t_{K^+n(K^0p) \rightarrow K^+n}^{(s)} = \frac{(\pm)g_{K^+n}^2}{M_I - M_R + i\Gamma/2} ,$$

$$t_{K^+n(K^0p) \rightarrow K^+n}^{(p,1/2)} = \frac{(\pm)\bar{g}_{K^+n}^2(\boldsymbol{\sigma} \cdot \mathbf{q}')(\boldsymbol{\sigma} \cdot \mathbf{q})}{M_I - M_R + i\Gamma/2} ,$$

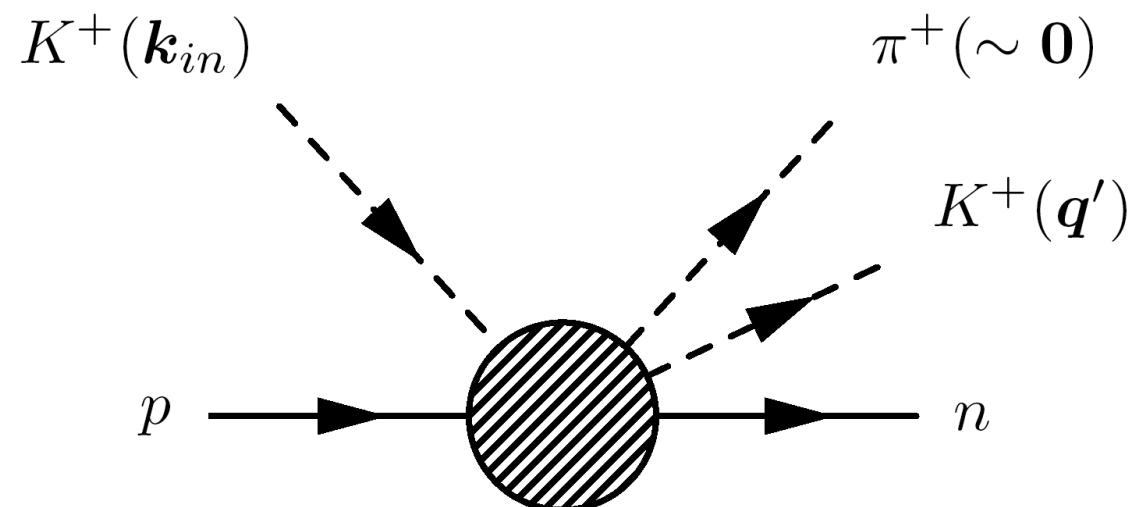
$$t_{K^+n(K^0p) \rightarrow K^+n}^{(p,3/2)} = \frac{(\pm)\tilde{g}_{K^+n}^2(\mathbf{S} \cdot \mathbf{q}')(\mathbf{S}^\dagger \cdot \mathbf{q})}{M_I - M_R + i\Gamma/2} ,$$

$$g_{K^+n}^2 = \frac{\pi M_R \Gamma}{M q} , \quad \bar{g}_{K^+n}^2 = \frac{\pi M_R \Gamma}{M q^3} , \quad \tilde{g}_{K^+n}^2 = \frac{3\pi M_R \Gamma}{M q^3}$$

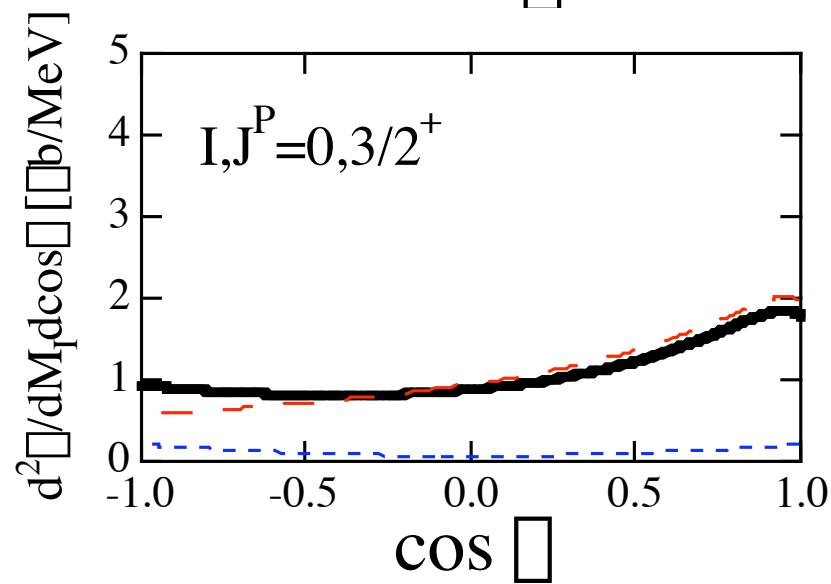
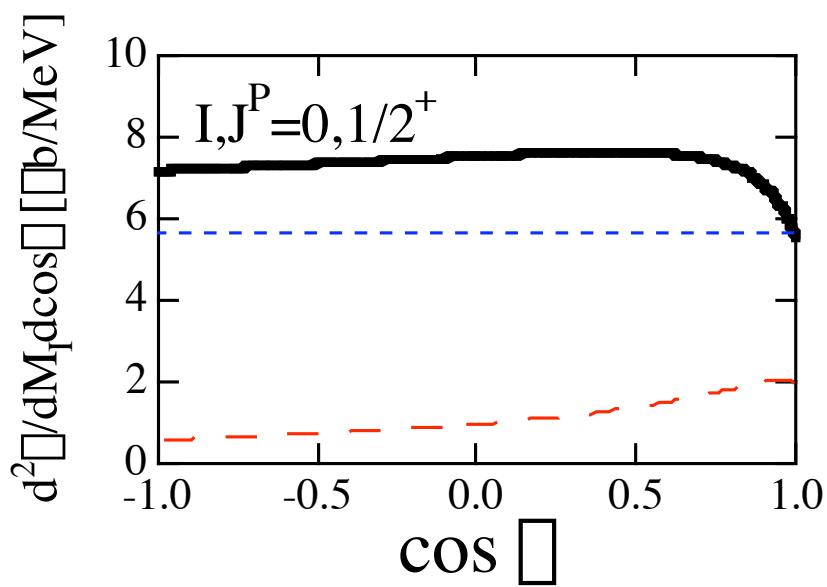
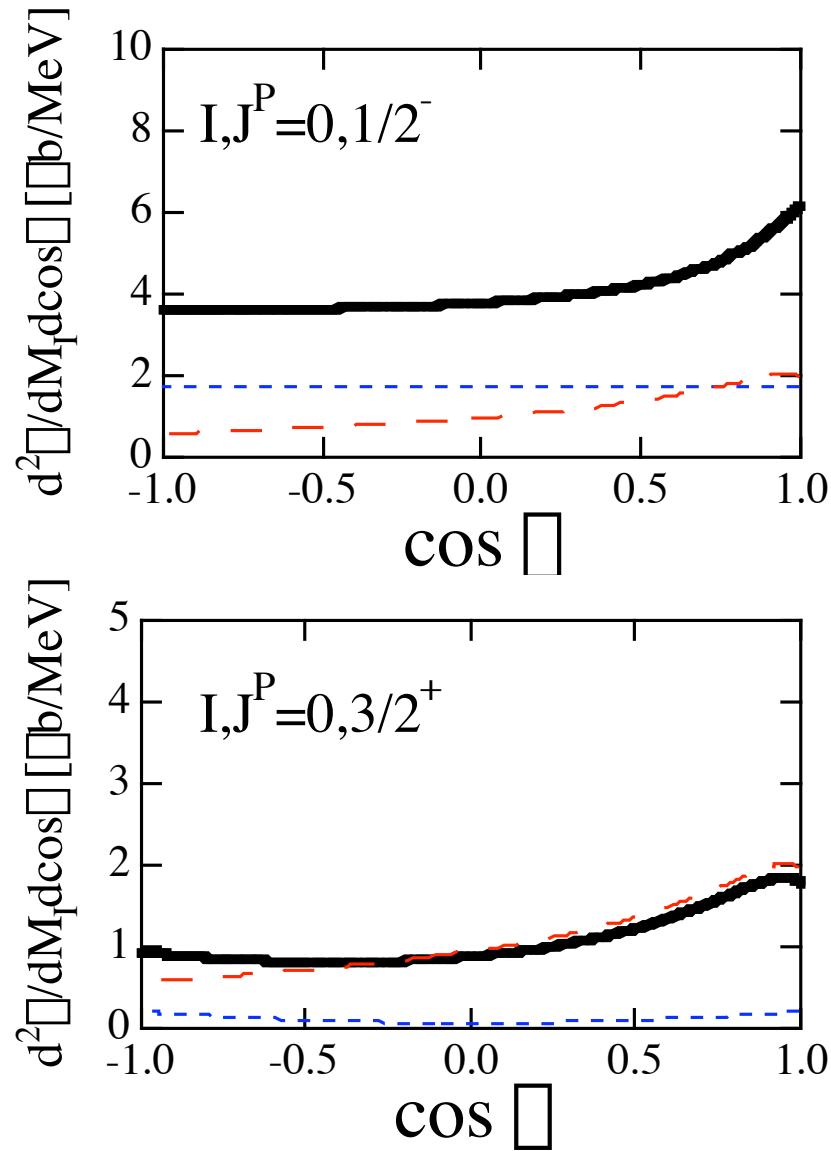
## Resonance term

**Amplitude of resonance term for  $K^+p \rightarrow \pi^+K^+n$  :**

$$\begin{aligned}
 -i\tilde{t}_i^{(s)} &= \frac{g_{K^+n}^2}{M_I - M_R + i\Gamma/2} \left\{ G(M_I)(a_i + c_i) - \frac{1}{3}\bar{G}(M_I)b_i \right\} \boldsymbol{\sigma} \cdot \mathbf{k}_{in} S_I(i) \\
 -i\tilde{t}_i^{(p,1/2)} &= \frac{\bar{g}_{K^+n}^2}{M_I - M_R + i\Gamma/2} \bar{G}(M_I) \left\{ \frac{1}{3}b_i \mathbf{k}_{in}^2 - a_i + d_i \right\} \boldsymbol{\sigma} \cdot \mathbf{q}' S_I(i) \\
 -i\tilde{t}_i^{(p,3/2)} &= \frac{\tilde{g}_{K^+n}^2}{M_I - M_R + i\Gamma/2} \bar{G}(M_I) \frac{1}{3}b_i \left\{ (\mathbf{k}_{in} \cdot \mathbf{q}')(\boldsymbol{\sigma} \cdot \mathbf{k}_{in}) - \frac{1}{3}\mathbf{k}_{in}^2 \boldsymbol{\sigma} \cdot \mathbf{q}' \right\} S_I(i)
 \end{aligned}$$



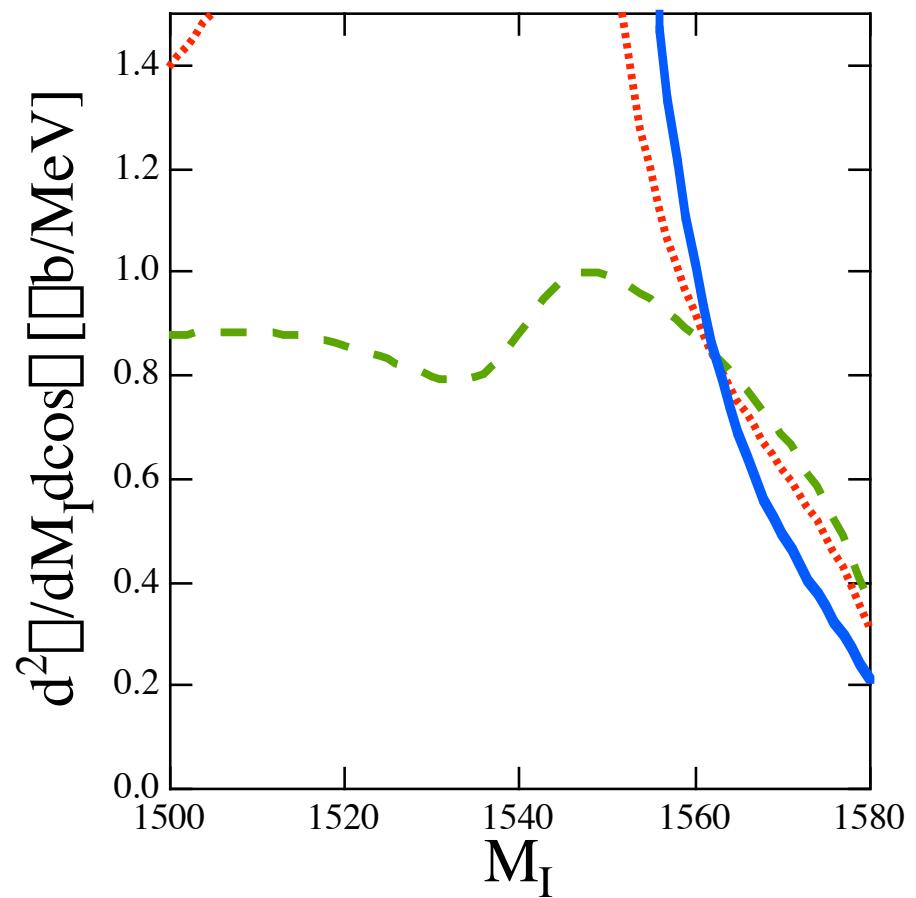
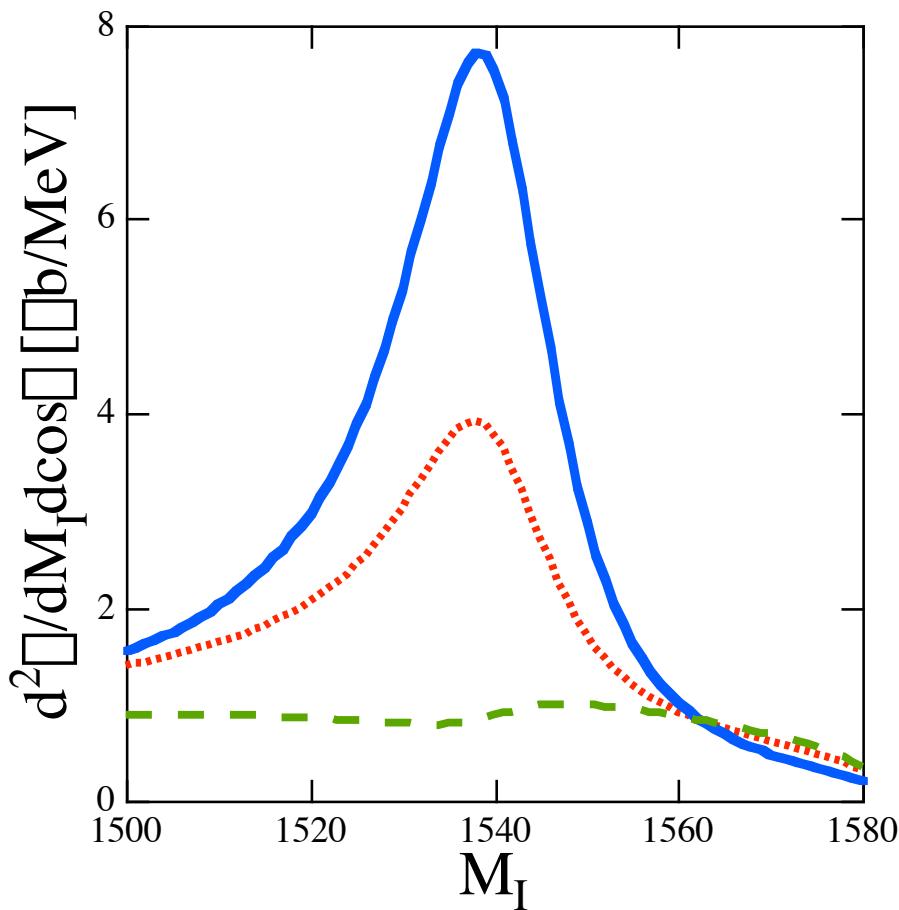
## Angular dependence

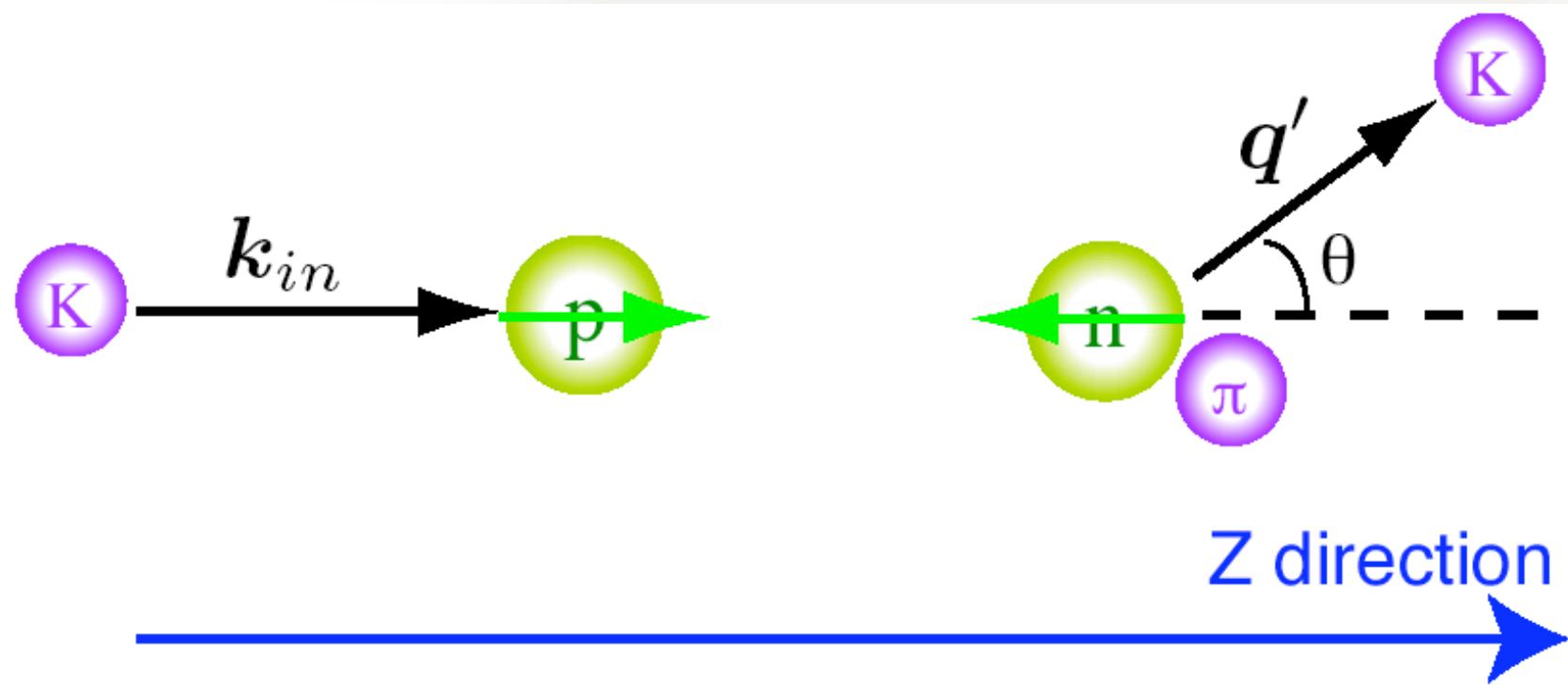


— total  
--- resonance  
- - - background

## Mass distributions

.....  $I, J^P = 0, 1/2^-$   
—  $I, J^P = 0, 1/2^+$        $k_{in}(\text{Lab}) = 850 \text{ MeV}/c$   
- -  $I, J^P = 0, 3/2^+$        $\theta = 90 \text{ deg}$



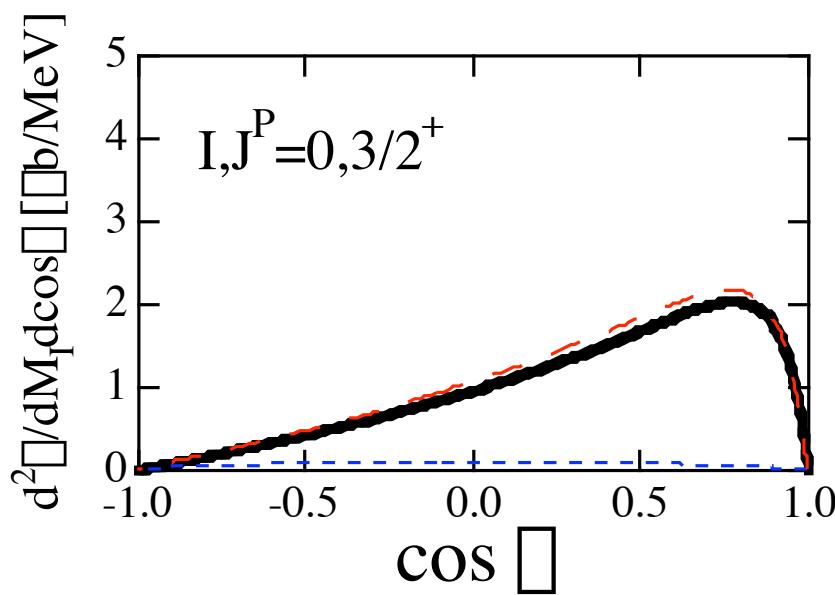
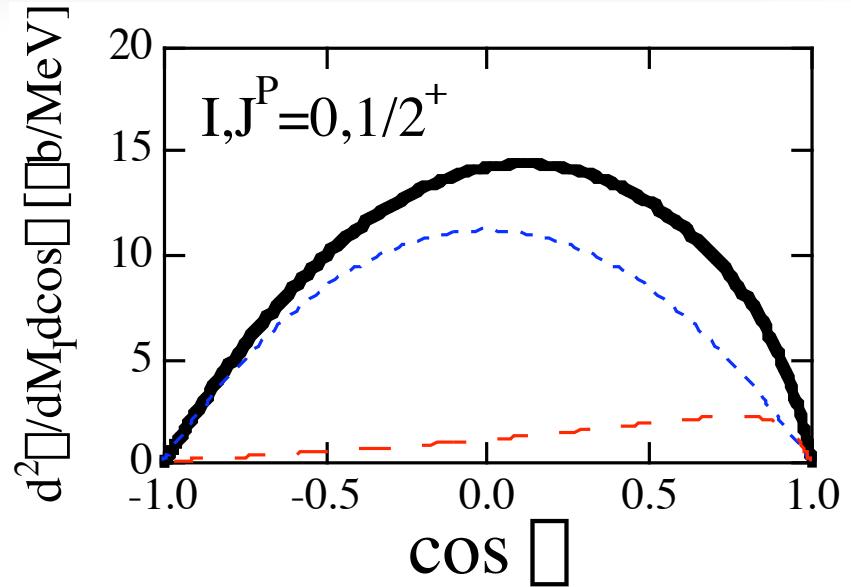
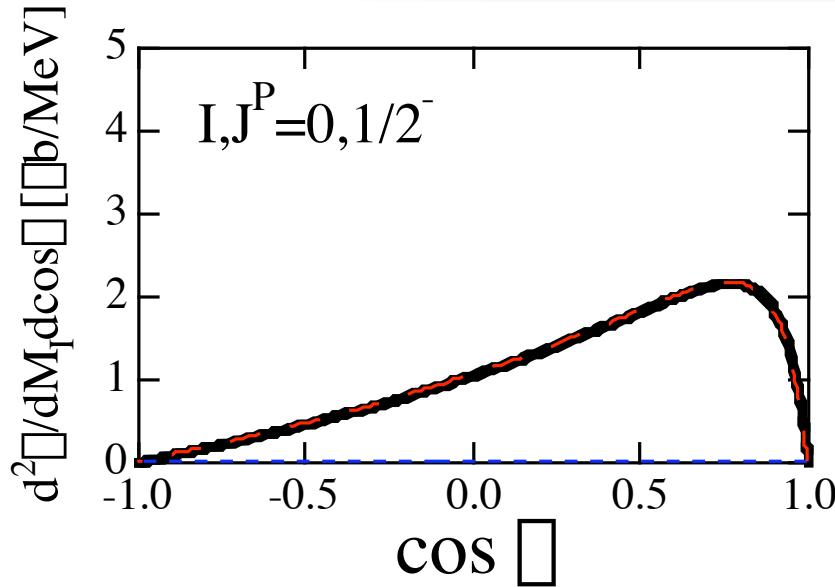


$$\langle -1/2 | \boldsymbol{\sigma} \cdot \mathbf{k}_{in} | 1/2 \rangle = 0$$

$$\langle -1/2 | \boldsymbol{\sigma} \cdot \mathbf{q}' | 1/2 \rangle \propto q' \sin \theta$$

# Same result is obtained for final pK<sup>0</sup>

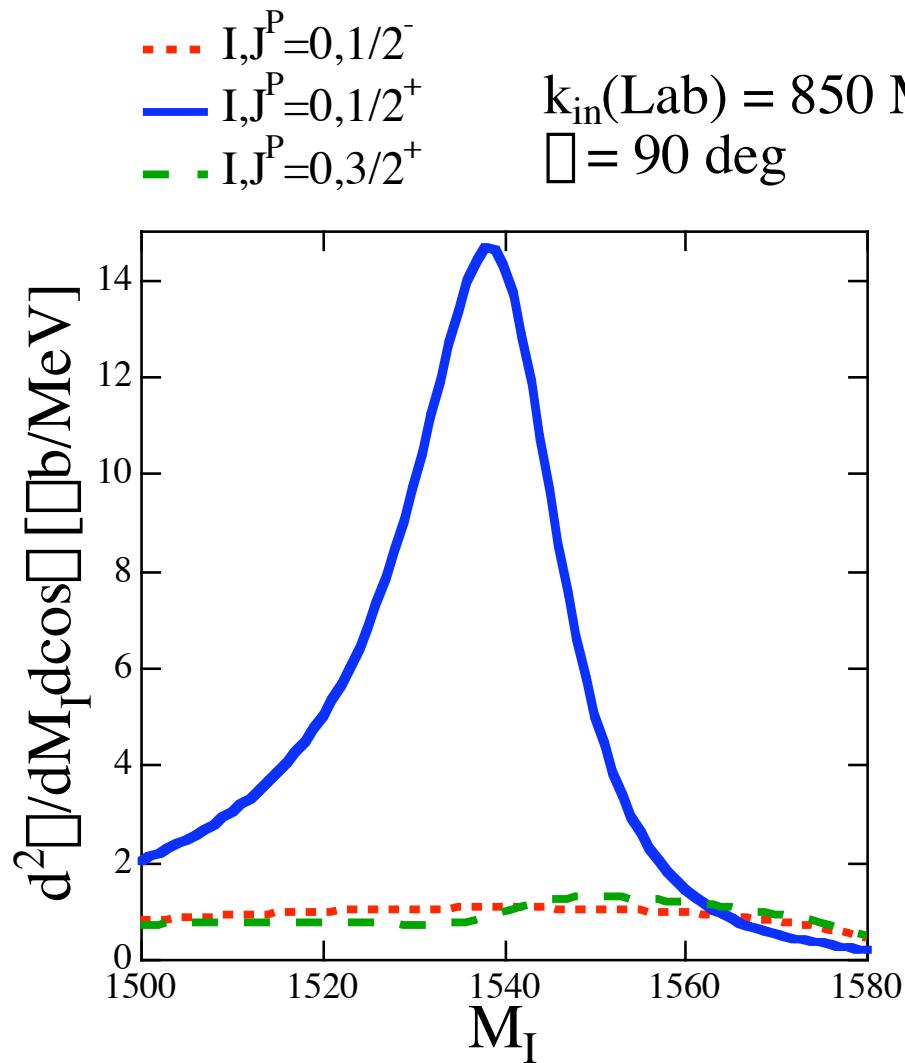
## Angular dependence : polarization test



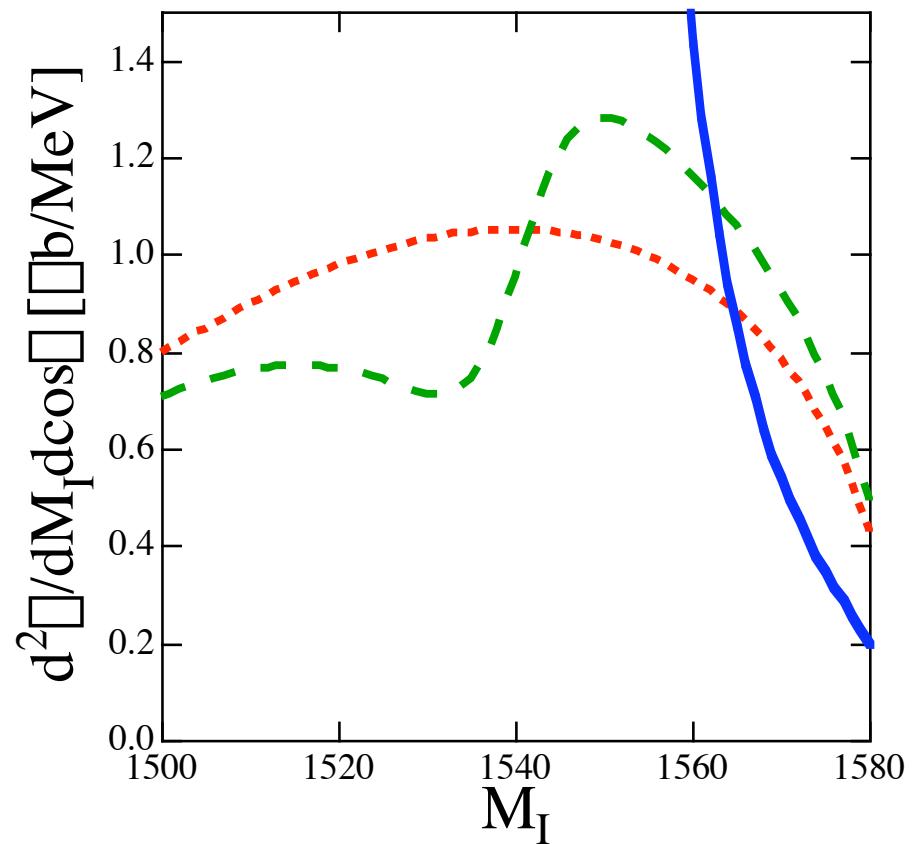
— total  
--- resonance  
- - - background

Polarization test

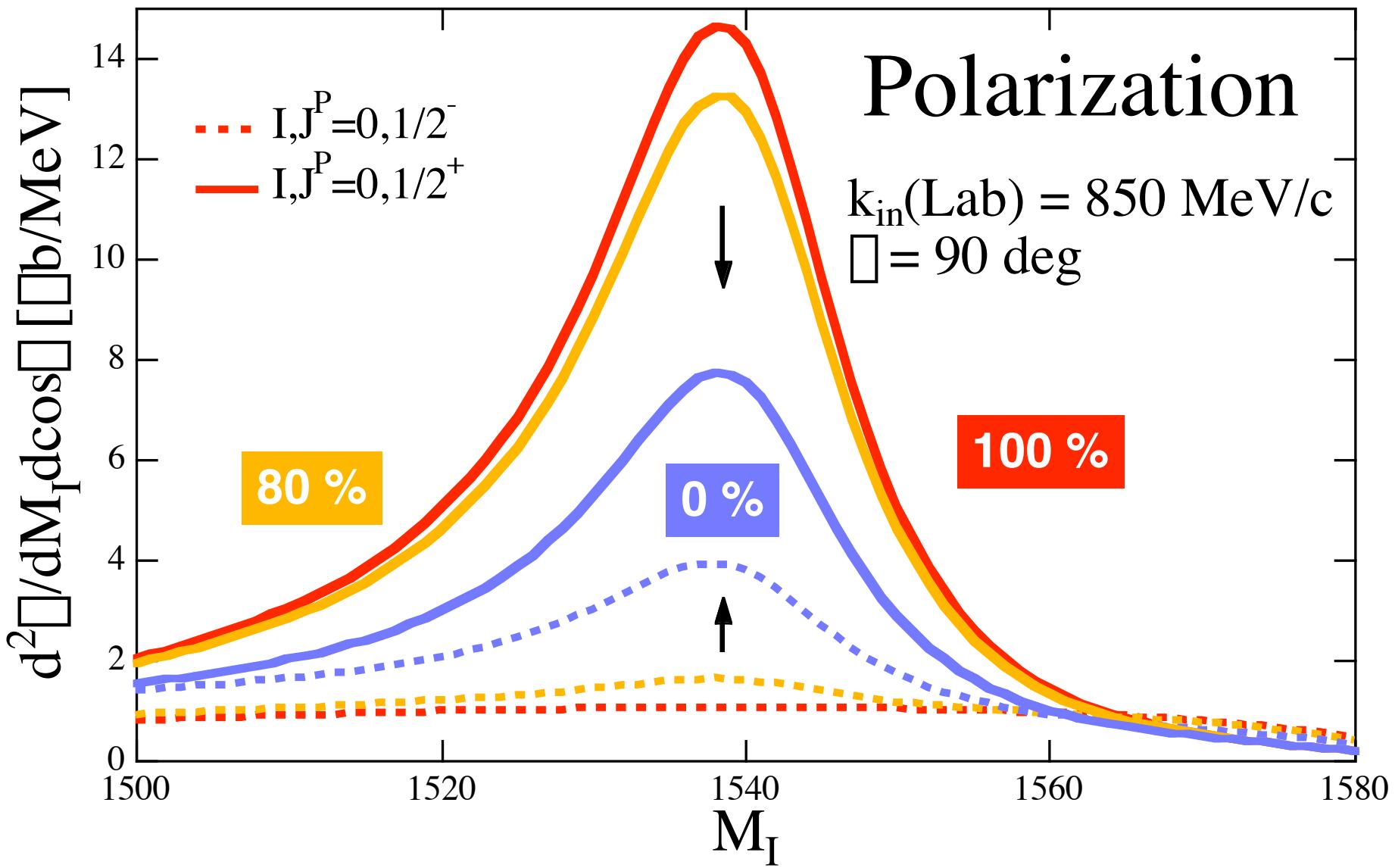
## Mass distributions : polarization test



## Polarization test



## Incomplete polarization



## Conclusion

We calculate the  $K^+ p \rightarrow \bar{\psi}^0 K N$  reaction using a chiral model, assuming the possible quantum numbers of  $\bar{\psi}^+$  baryon.

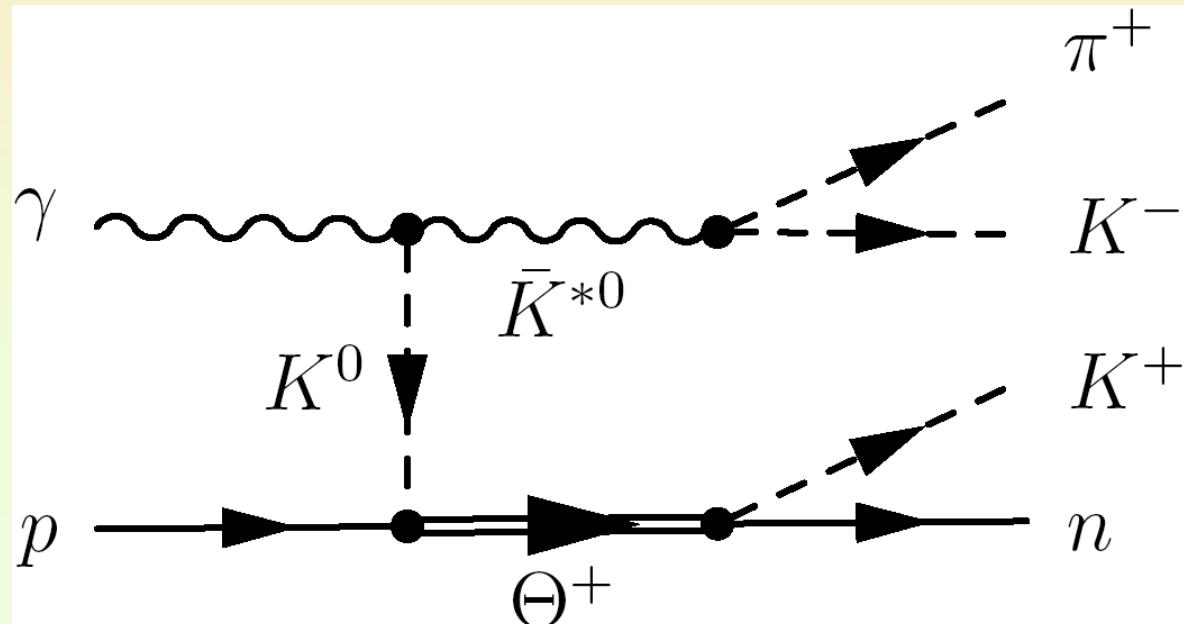
- If we find the resonance with polarization test, the quantum number of  $\bar{\psi}^+$  can be determined as  $|l|=0$ ,  $J^P=1/2^+$

T. Hyodo, A. Hosaka, and E. Oset, nucl-th/0307105

## Future work

- Full calculation of the present reaction without approximation of kinematics  
-> information from  $\Theta^+$  angular dependence
- photo-production of  $K^*$  and  $\bar{K}$

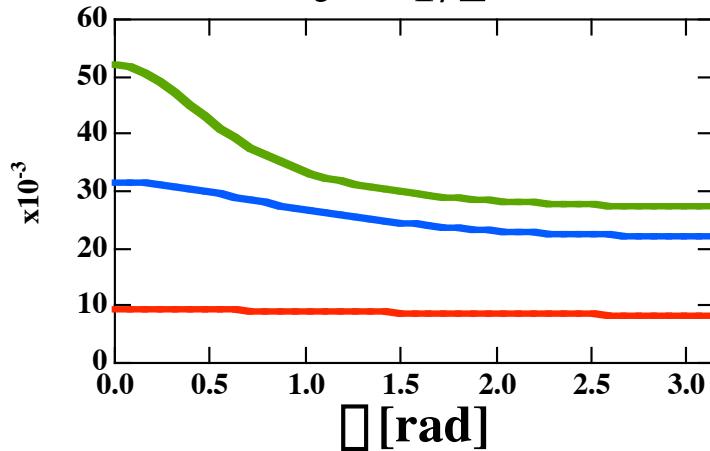
V. Kubarovskiy et al., hep-ex/0307088



# Photo-production of $K^*$ and $\bar{K}$

Angular dependence (upper) and  
integrated (lower) cross sections  
 $\bar{K}p \rightarrow K^* \bar{K}$   
K-exchange,  $\bar{K}$  1 GeV  
units : [ $\mu b$ ]

$$J^P = 1/2^+$$



## Preliminary

$$J^P = 1/2^-$$

