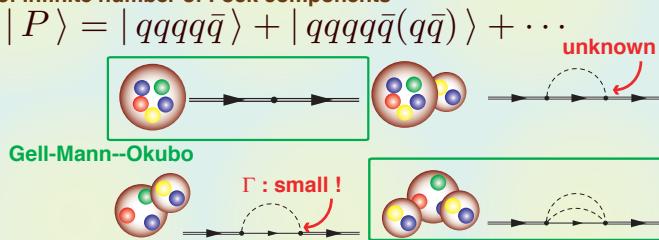


Flavor symmetry and the pentaquark $\Theta(1540)$

We study the properties of the exotic $\Theta(1540)$ baryon from the viewpoint of flavor SU(3) symmetry. The role of two-meson cloud in the baryon antidecuplet are investigated based on the effective interactions, assuming $J^P = 1/2^+$. Then we study the possibilities of having $J^P = 1/2^\pm, 3/2^\pm$ in connection with the known baryon resonances.

Motivation : Two-meson cloud

In general, an exotic state can be represented as linear combination of infinite number of Fock components



Possibilities of Θ as $K\pi N$ bound state

P. Bicudo, et al. Phys. Rev. C 69, 011503 (2004)
T. Kishimoto, et al. hep-ex/0312003
F. J. Llanes-Estrada, et al. Phys. Rev. C 69, 055203 (2004)

Self-energy results

Interaction Lagrangians

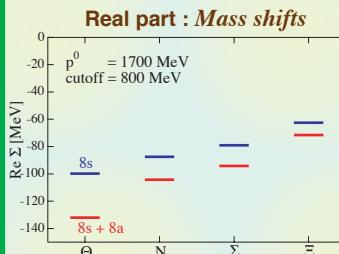
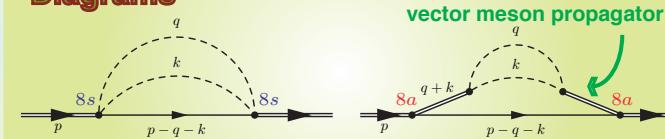
Lagrangian should be SU(3) symmetric.

$$\mathcal{L}^{8s} = \frac{g^{8s}}{2f} \bar{P}_{ijk} \epsilon^{lmk} \phi_l^a \phi_a^i B_m^j + \text{h.c.}$$

$$\mathcal{L}^{8a} = i \frac{g^{8a}}{4f^2} \bar{P}_{ijk} \epsilon^{lmk} \gamma^\mu (\partial_\mu \phi_l^a \phi_a^i - \phi_l^a \partial_\mu \phi_a^i) B_m^j + \text{h.c.}$$

Coupling constants are determined from the experimental data of the decay width of $N(1710) \rightarrow N\pi\pi$. Since the Lagrangians are SU(3) symmetric, self-energies of all the members of antidecuplet can be calculated.

Diagrams



Imaginary part : Decay widths

Decay widths (MeV)	$\Gamma^{(8s)}$	$\Gamma^{(8a)}$	Γ_{BMM}^{8s}
$N(1710) \rightarrow N\pi\pi$ (inputs)	25	15	40
$N(1710) \rightarrow N\eta\pi$	0.58	—	—
$\Sigma(1770) \rightarrow N\bar{K}\pi$	4.7	6.0	24
$\Sigma(1770) \rightarrow \Sigma\pi\pi$	10	0.62	—
$\Sigma(1770) \rightarrow \Lambda\pi\pi$	—	2.9	—
$\Xi(1860) \rightarrow \Sigma\bar{K}\pi$	0.57	0.46	2.1
$\Xi(1860) \rightarrow \Xi\pi\pi$	—	1.1	—

(in MeV)

Mass difference between Θ and $\Xi \sim 60 \text{ MeV} \sim 20\% \text{ of } 320 = 1860-1540$

Conclusions 1

We study the two-meson vertical cloud effect to the self-energy of baryon antidecuplet.

- Two-meson cloud effects are always **attractive**, and contribute to the antidecuplet mass splitting of the order of **20%**.
- Decay widths to the three body channel are relatively **small**.

A. Hosaka, T. H., F. J. Llanes-Estrada, E. Oset, J. R. Pelaez, and M. J. Vicente Vacas, Phys. Rev. C 71, 045205 (2005)

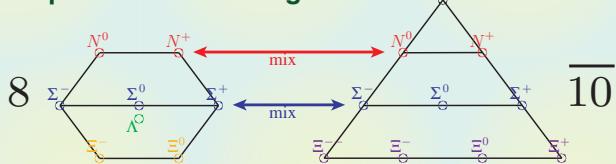
Tetsuo Hyodo^a

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What about other JP?

Properties of SU(3) multiplet in general framework

Representation mixing



We assign several baryon resonances in this framework together with $\Theta(1540)$ and $\Xi(1860)$.

Mass and width

Mass spectrum

$$M_\Theta = M_8^{1W} \quad M_{\Xi^{10}} = \dots$$

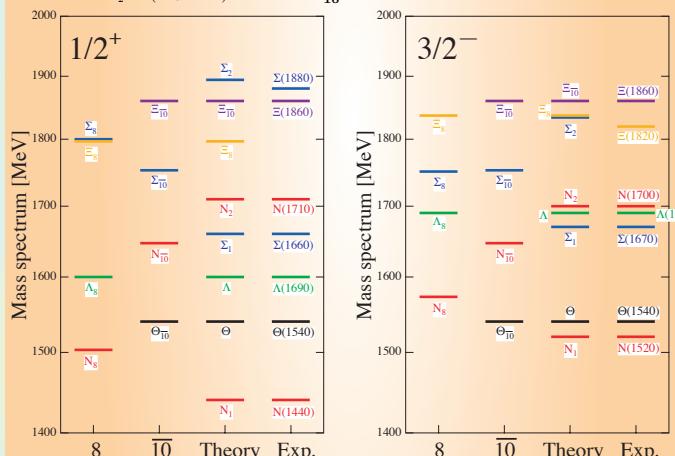
$$M_\Lambda = M_8 \quad M_{\Xi_8} = M_8 + b + \frac{1}{2}c$$

$$M_{N_1} = \left(M_8 - b + \frac{1}{2}c \right) \cos^2 \theta_N + (M_{\Xi^{10}} - a) \sin^2 \theta_N - \delta \sin 2\theta_N$$

$$M_{N_2} = \left(M_8 - b + \frac{1}{2}c \right) \sin^2 \theta_N + (M_{\Xi^{10}} - a) \cos^2 \theta_N + \delta \sin 2\theta_N$$

$$M_{\Sigma_1} = (M_8 + 2c) \cos^2 \theta_\Sigma + M_{\Xi^{10}} \sin^2 \theta_\Sigma - \delta \sin 2\theta_\Sigma$$

$$M_{\Sigma_2} = (M_8 + 2c) \sin^2 \theta_\Sigma + M_{\Xi^{10}} \cos^2 \theta_\Sigma + \delta \sin 2\theta_\Sigma$$



Decay width of Θ Exp. $N^* \rightarrow N\pi$ mass spectrum

$$g_\Theta = \sqrt{6} (g_{N_2} \cos \theta_N - g_{N_1} \sin \theta_N)$$

J^P	θ_N	Phase 1	Phase 2
$1/2^+$	29° (Mass)	29.1	103.3
$3/2^-$	33° (Mass)	3.1	20.0 (in MeV)

Conclusions 2

We examine 8-10 mixing scheme for the $\Theta(1540)$ and known resonances of $J^P = 1/2^\pm, 3/2^\pm$.

- The experimental mass spectrum of $1/2^+$ and $3/2^-$ are well reproduced.
- By observing the decay widths of $N^* \rightarrow N\pi$ with the mixing angle obtained from the masses, $J^P = 3/2^-$ assignment seems plausible.

T. H. and A. Hosaka, Phys. Rev. D 71, 054017 (2005)