

Phenomenology of spin 3/2 baryons with pentaquarks



Tetsuo Hyodo^a

and A. Hosaka^a

RCNP, Osaka^a

2005, Oct. 25th ₁

Contents

- ★ **Introduction**
- ★ **Pure antidecuplet case**
- ★ **$8-\overline{10}$ mixing case**
 - ★ **Mass spectra**
 - ★ **Decay widths**
- ★ **Two-meson coupling**
 - ★ **Coupling constants**
 - ★ **Meson induced Θ production**
- ★ **Summary**

Introduction : Flavor SU(3) symmetry

Existence of Θ^+ + Flavor SU(3) symmetry

→ **Existence of flavor partners of Θ^+**

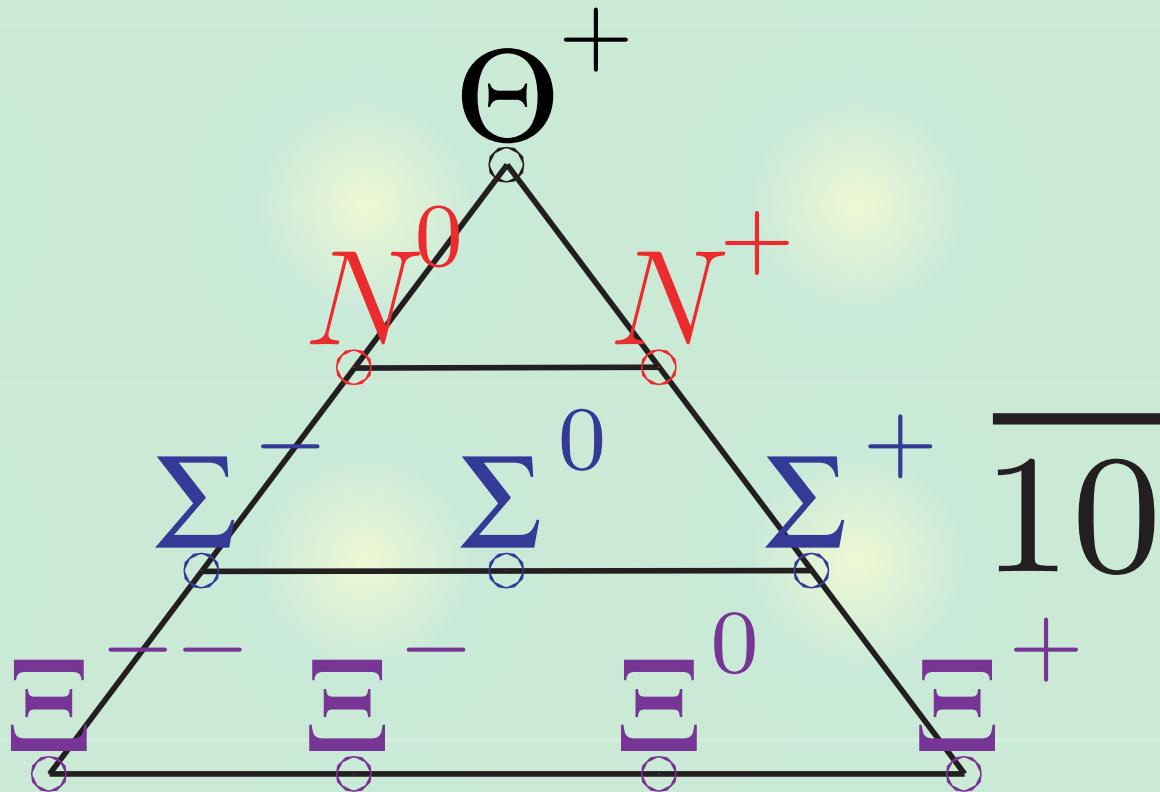
Assuming the flavor multiplet that Θ^+ belongs to, we examine its properties by symmetry relation, in connection with known baryon resonances.

→ **to determine the J^P of Θ^+**

Phenomenological but model independent analysis up to $O(m_s)$

Pure antidecuplet case

Simplest assignment for Θ^+



Test the masses and widths of partners via flavor SU(3) symmetry relations

Pure antidecuplet case

Mass and decay width [MeV]

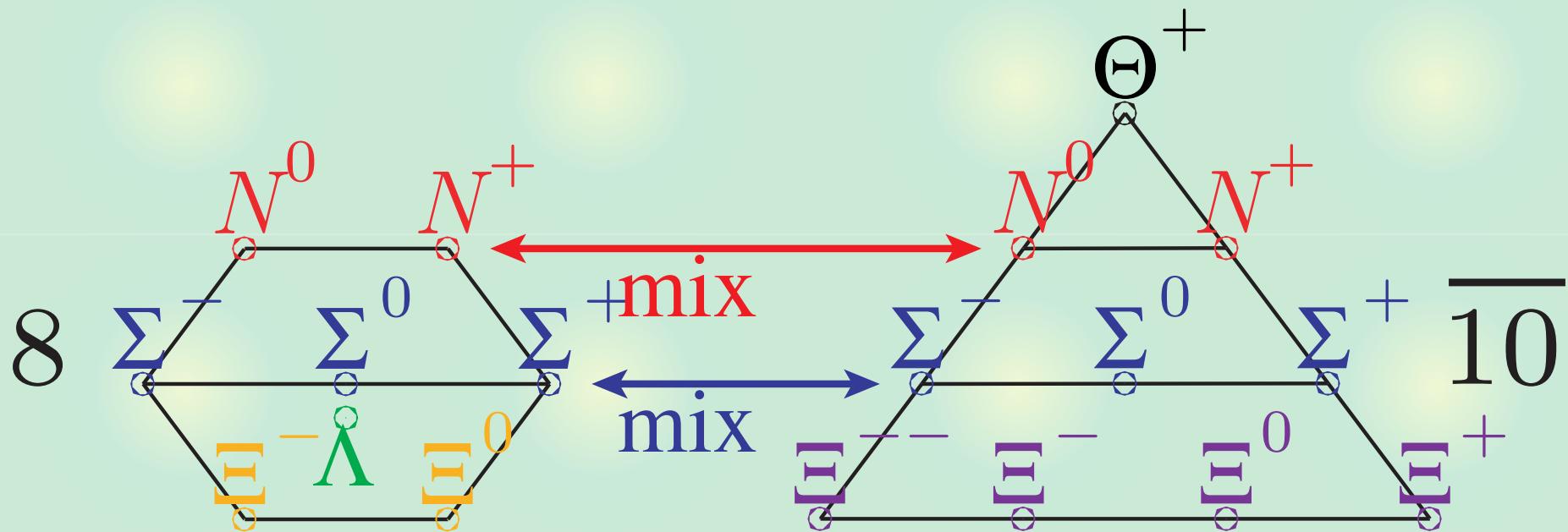
$$M(\overline{10}; Y) = M_{\overline{10}} - aY \quad g_{\Theta KN} = \sqrt{6}g_{N^*\pi N}$$

J^P	M_Θ	M_N	M_Σ	M_Ξ	Γ_Θ
$1/2^-$ exp.	1540 $\Theta(1540)$	1647 $N(1650)$	1753 $\Sigma(1750)$	1860 $\Xi(1860)$	156.1
$1/2^+$ exp.	1540 $\Theta(1540)$	1710 $N(1710)$	1880 $\Sigma(1880)$	2050 $\Xi(2030)$	7.2
$3/2^+$ exp.	1540 $\Theta(1540)$	1720 $N(1720)$	1900	2080	10.6
$3/2^-$ exp.	1540 $\Theta(1540)$	1700 $N(1700)$	1860	2020 $\Xi(2030)$	1.3

are not reproduced simultaneously.

Octet-antidecuplet mixing

Second simplest assignment for Θ^+



Mixing is induced by the $SU(3)$ breaking in mass term.

Octet-antidecuplet mixing

Mass formulae

$$M_\Theta = M_{\bar{\mathbf{10}}} - 2a$$

$$M_{\Xi_{\bar{\mathbf{10}}}} = M_{\bar{\mathbf{10}}} + a$$

$$M_\Lambda = M_{\mathbf{8}}$$

$$M_{\Xi_8} = M_{\mathbf{8}} + b + \frac{1}{2}c$$

$$M_{N_1} = \left(M_{\mathbf{8}} - b + \frac{1}{2}c \right) \cos^2 \theta_N + (M_{\bar{\mathbf{10}}} - a) \sin^2 \theta_N - \delta \sin 2\theta_N$$

$$M_{N_2} = \left(M_{\mathbf{8}} - b + \frac{1}{2}c \right) \sin^2 \theta_N + (M_{\bar{\mathbf{10}}} - a) \cos^2 \theta_N + \delta \sin 2\theta_N$$

$$M_{\Sigma_1} = (M_{\mathbf{8}} + 2c) \cos^2 \theta_\Sigma + M_{\bar{\mathbf{10}}} \sin^2 \theta_\Sigma - \delta \sin 2\theta_\Sigma$$

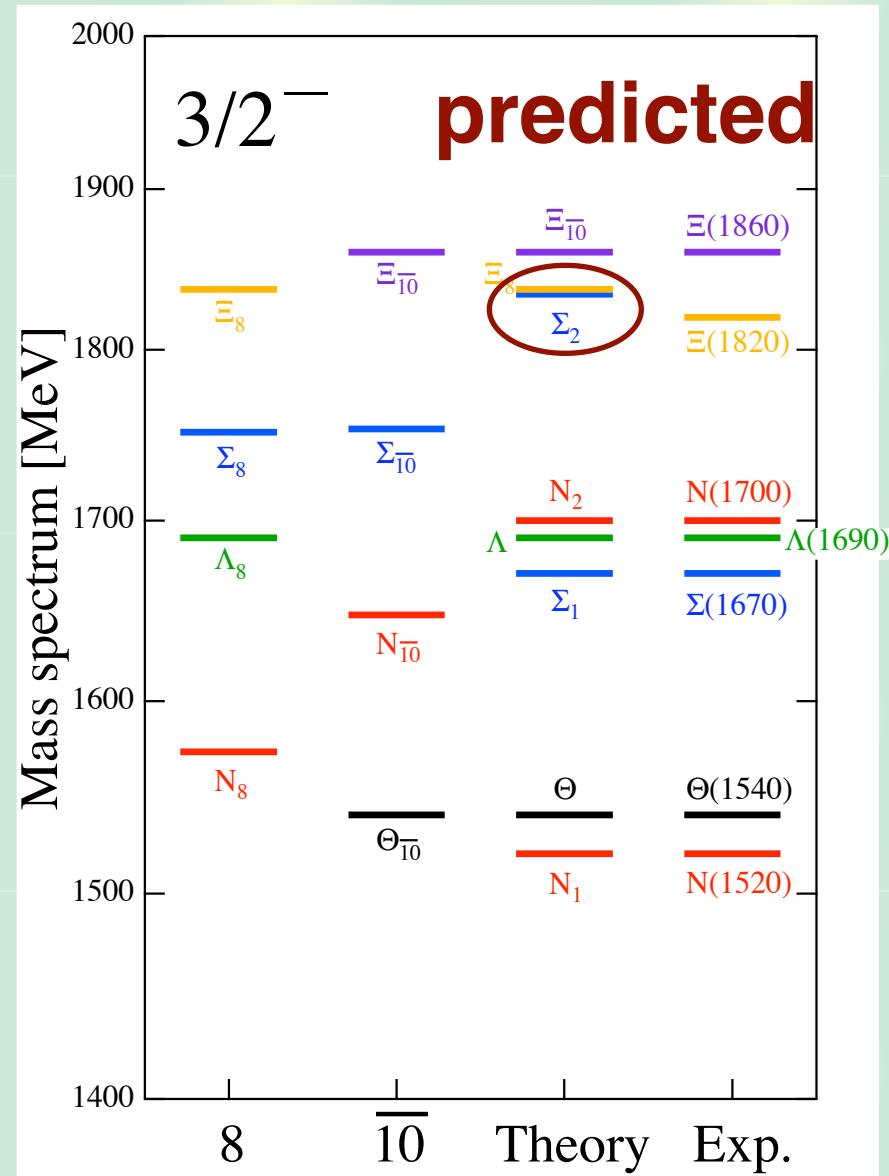
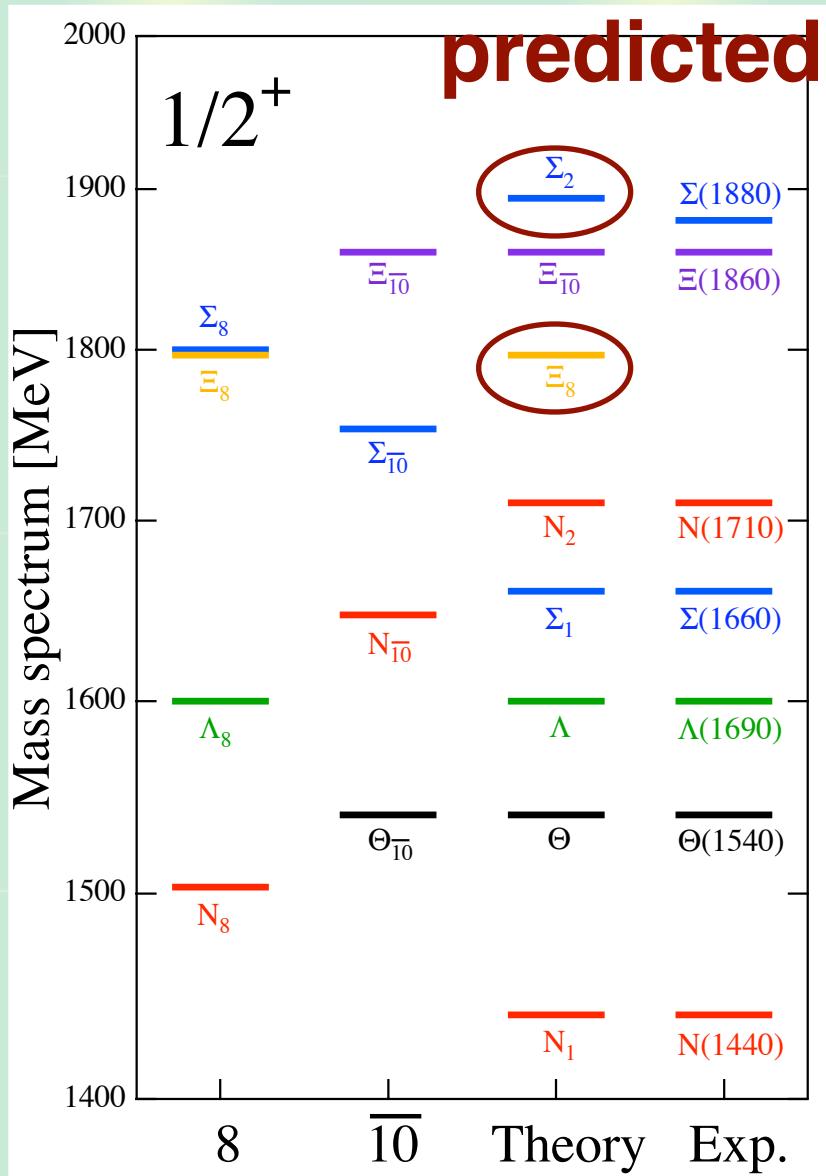
$$M_{\Sigma_2} = (M_{\mathbf{8}} + 2c) \sin^2 \theta_\Sigma + M_{\bar{\mathbf{10}}} \cos^2 \theta_\Sigma + \delta \sin 2\theta_\Sigma$$

8 masses v.s. 6 parameters

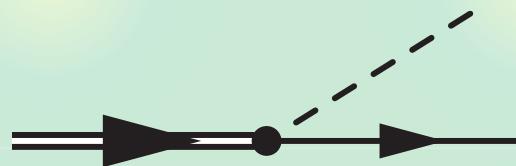
$J^P = 1/2^-$: too wide width

$J^P = 3/2^+$: states are not well established

Mass spectra



Decay width of Θ



Relation between coupling constants

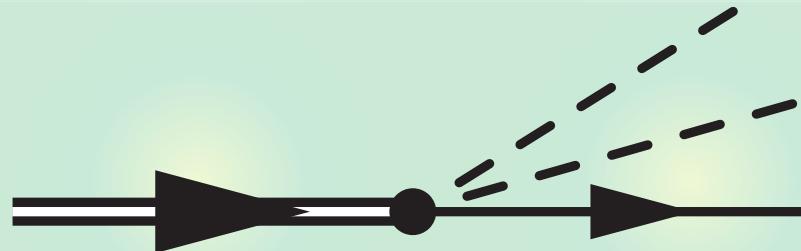
$$g_\Theta = \sqrt{6}(g_{N_2} \cos \theta_N - g_{N_1} \sin \theta_N)$$

C.G. Coeff. N* decay from masses

J^P	θ_N [deg]	Γ_Θ [MeV]
$1/2^+$	29	29.1
$3/2^-$	33	3.1

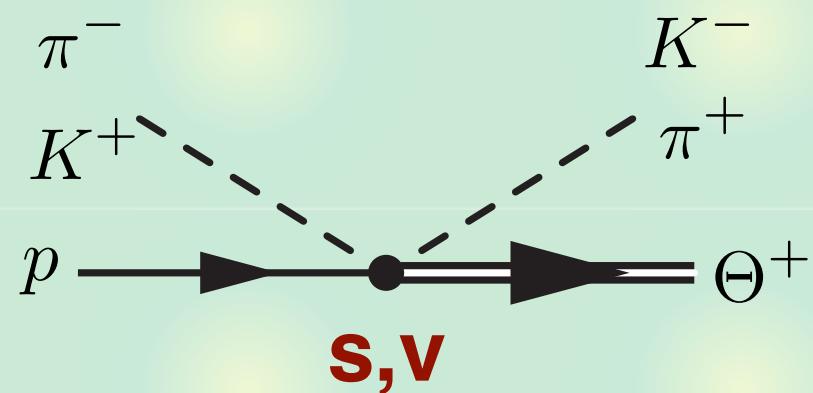
Two-meson coupling

Then, what about two-meson coupling?



SU(3) relation enable us to calculate

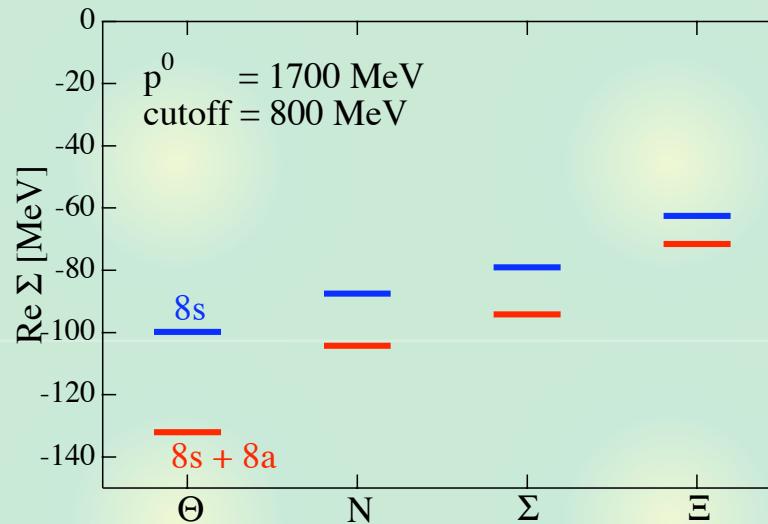
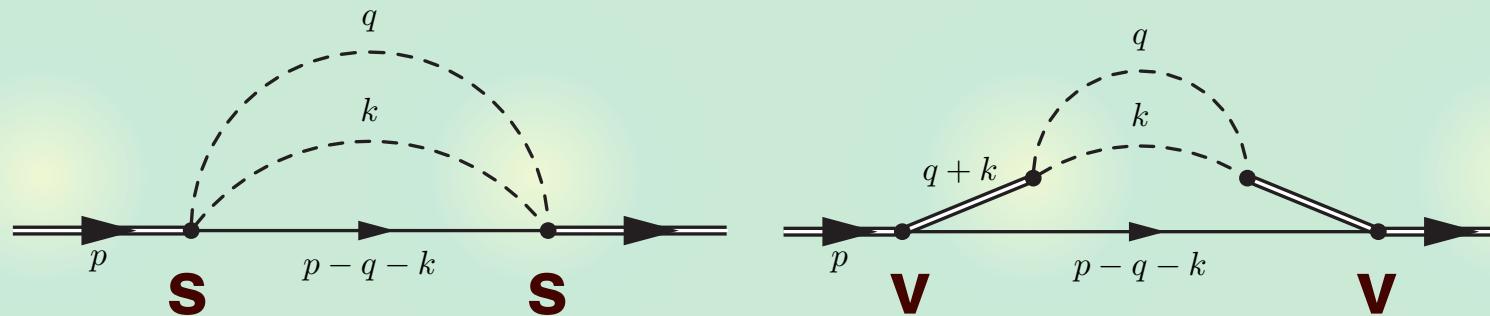
the cross section of



from the decay of nucleons into two pions.

Two-meson coupling

Contact interaction :

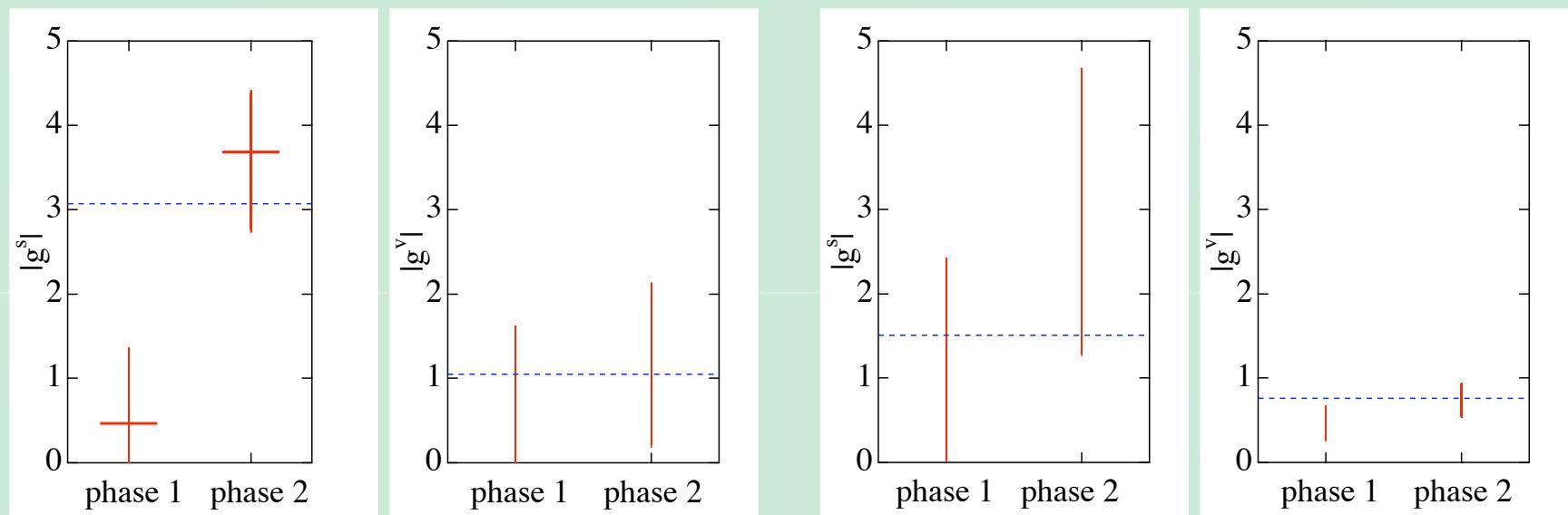


A. Hosaka, et al., Phys. Rev. C71 045205 (2005)
Talk by Llanes-Estrada

Two-meson coupling

Branching fraction [%]

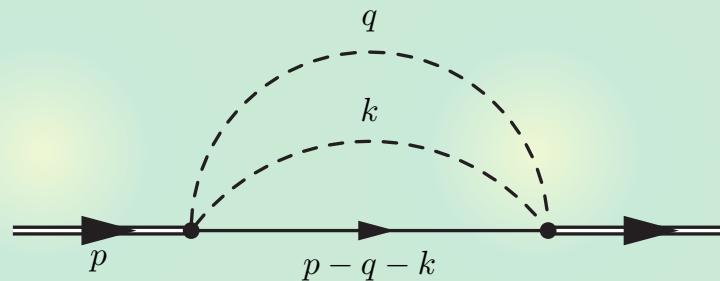
J^P	state	πN	$\pi\pi N(s)$	$\pi\pi N(v)$
$1/2^+$	N(1440)	65	7.5	<8
	N(1710)	15	25	15
$3/2^-$	N(1520)	55	25	20
	N(1700)	10	<85-95	<35



Still large uncertainty

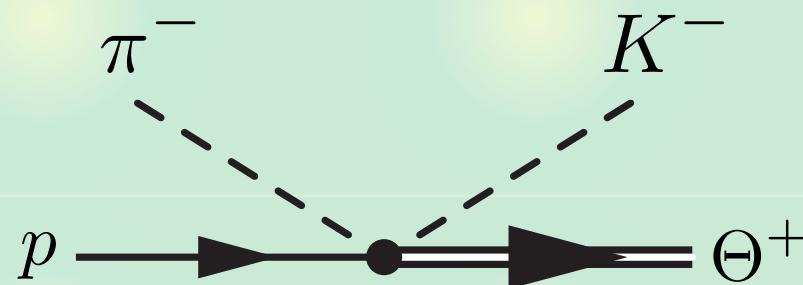
Constraints on the coupling

Self-energy : not too large, but not too small



$\sim 100 \text{ MeV}$

$\pi^- p \rightarrow K^- \Theta^+$ at KEK : upper limit is $\sim 4.1 \mu\text{b}$

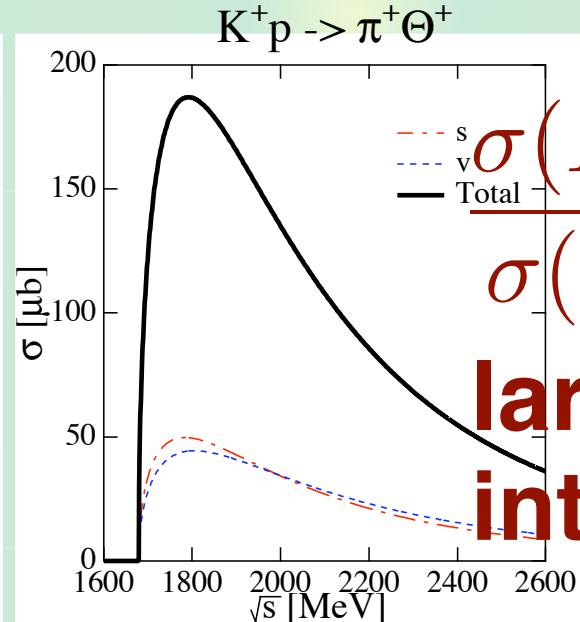
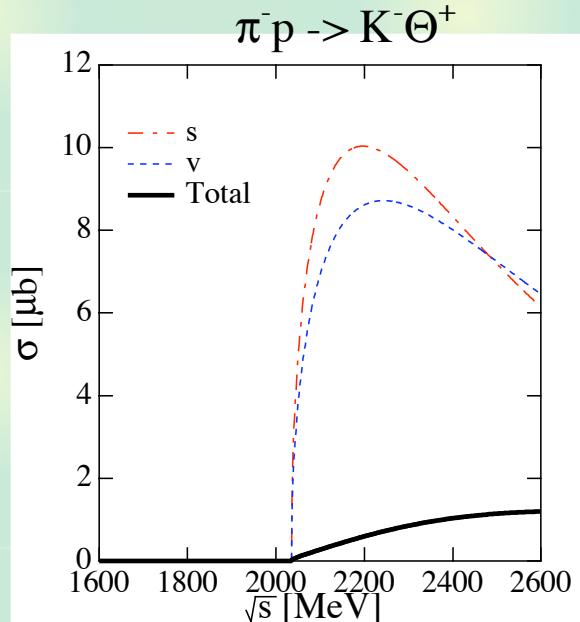


$< 4.1 \mu\text{b}$

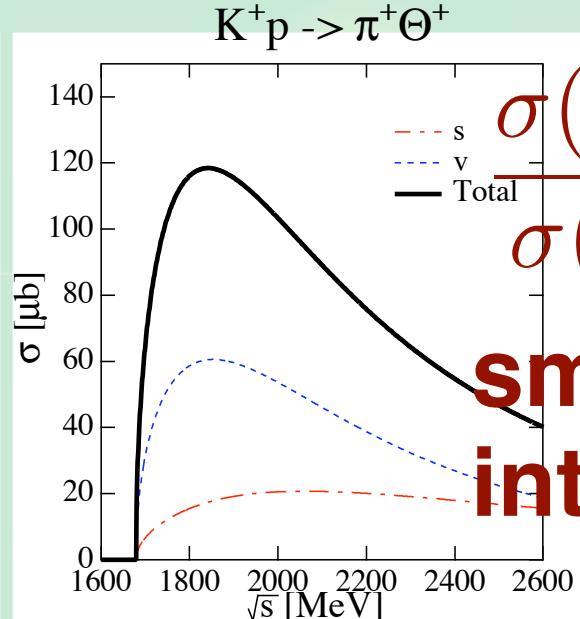
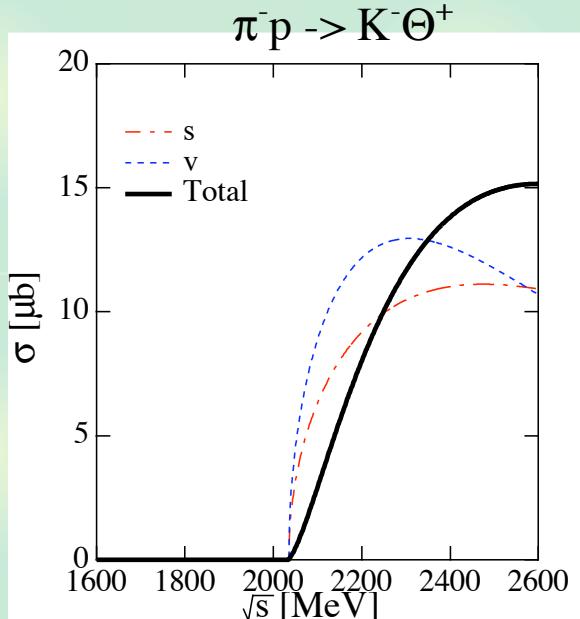
It is necessary to use the interference effect among two terms, s and v .

Θ production

$1/2^+$



$3/2^-$



$$\frac{\sigma(K^+)}{\sigma(\pi^-)} \sim 50$$

large interference

$$\frac{\sigma(K^+)}{\sigma(\pi^-)} \sim 3$$

small interference

Summary 1 : mixing scheme

We examine $8-\overline{10}$ mixing scheme for the exotic and non-exotic baryon resonances.

- Masses of $\Theta(1540)$ and $\Xi(1860)$ are well fitted in the $8-\overline{10}$ mixing scheme with $J^P = 1/2^+$ or $3/2^-$ baryons.
- The width of Θ is very narrow for the $J^P = 3/2^-$ case.
- For both cases, the mixing angle is close to the ideal angle.

Summary 2 : Two-meson coupling and Θ production

Based on the mixing scheme, we evaluate the two-meson coupling of Θ , and calculate the reaction process for Θ production



There is an interference effect between two amplitudes, which is prominent for $1/2^+$ case and rather moderate for $3/2^-$ case

J^P	g^s	g^v	σ_{π^-}	σ_{K^+}	$\text{Re}\Sigma_\Theta$
$1/2^+$	1.59	-0.27	4.1 μb	<1928 μb	-78 MeV
$3/2^-$	0.104	0.209	4.1 μb	<113 μb	-23 MeV