

Higher partial waves and resonance contributions in femtoscopy



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Introduction — Femtoscopy



Contribution from s-wave resonance

S. Watanabe, T. Hyodo, in preparation

- Origin of resonance peak



Contribution from higher partial waves

K. Murase, T. Hyodo, J. Subatomic Part. Cosmol. 3, 100017 (2025);

K. Murase, T. Hyodo, arXiv:2509.22844 [nucl-th]

- Correlation function with $l > 0$
- Regularized LL formula for $l > 0$

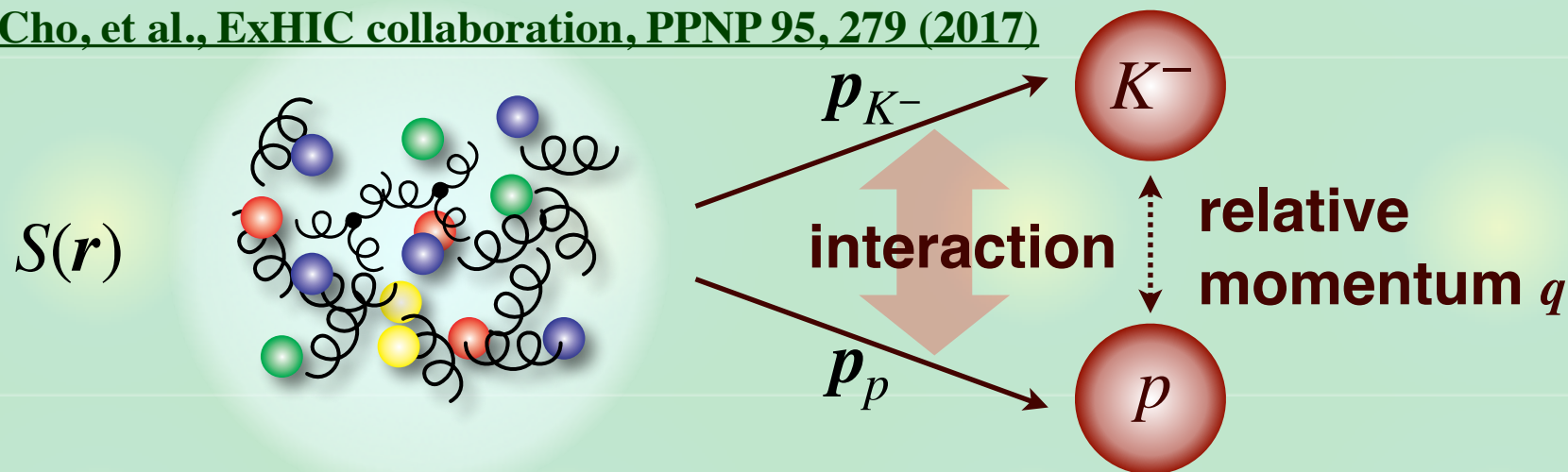


Summary

Correlation function and KP formula

High-energy collision: chaotic source $S(r)$ of hadron emission

S. Cho, et al., ExHIC collaboration, PPNP 95, 279 (2017)



- Definition

$$C(q) = \frac{N_{K^-p}(\mathbf{p}_{K^-}, \mathbf{p}_p)}{N_{K^-}(\mathbf{p}_{K^-})N_p(\mathbf{p}_p)} \quad (= 1 \text{ in the absence of FSI/QS})$$

- Theory: KP (Koonin-Pratt) formula

S.E. Koonin, PLB 70, 43 (1977); S. Pratt, PRD 33, 1314 (1986)

$$C(q) \simeq \int d^3r \, S(r) \, |\Psi_q^{(-)}(r)|^2$$

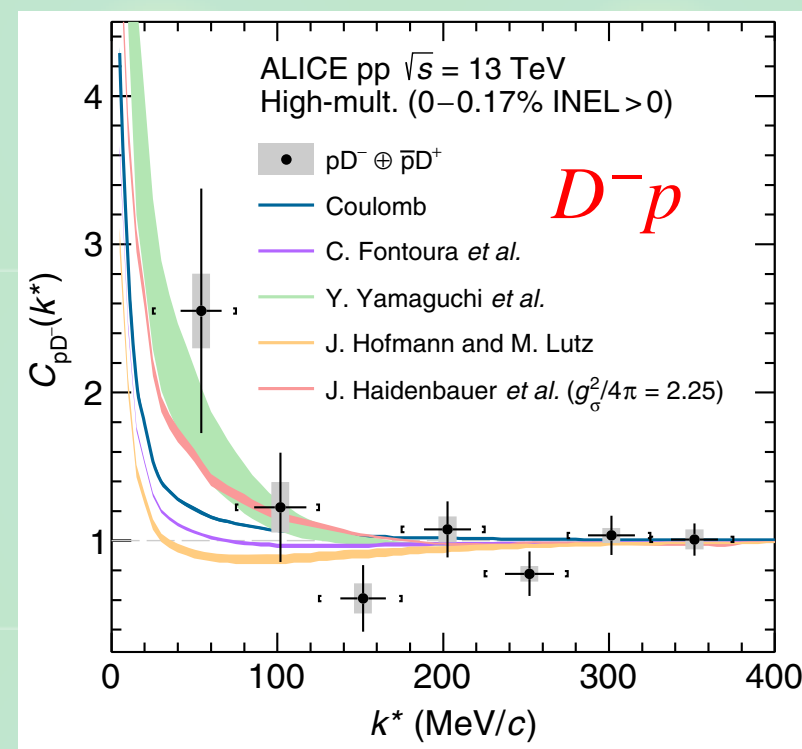
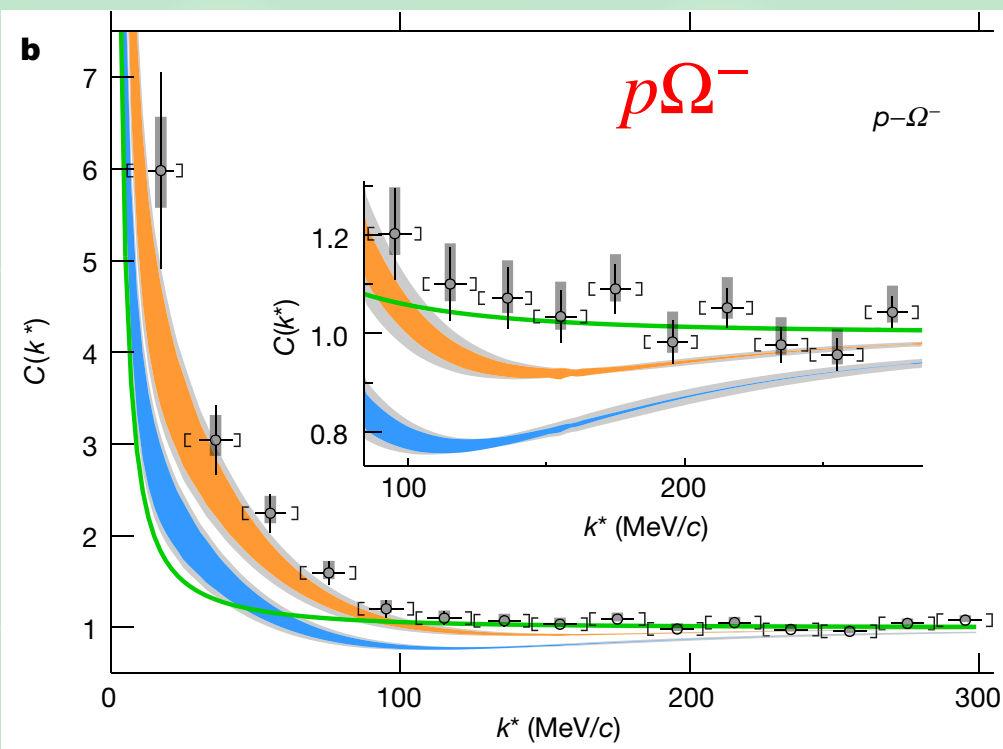
Source function $S(r) \longleftrightarrow$ wave function $\Psi_q^{(-)}(r)$ (interaction)

Experimental data with strangeness and charm

Correlation functions observed by ALICE@LHC

ALICE collaboration, Nature 588, 232 (2020);

ALICE collaboration, PRD 106, 052010 (2022)



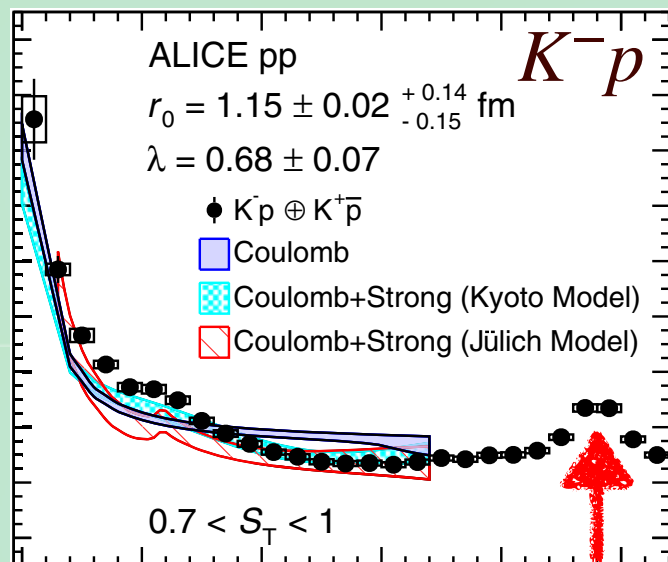
$\Omega^- \sim sss$: strangeness $S = -3$, $D^- \sim \bar{c}d$: charm $C = -1$

Almost impossible in scattering experiments

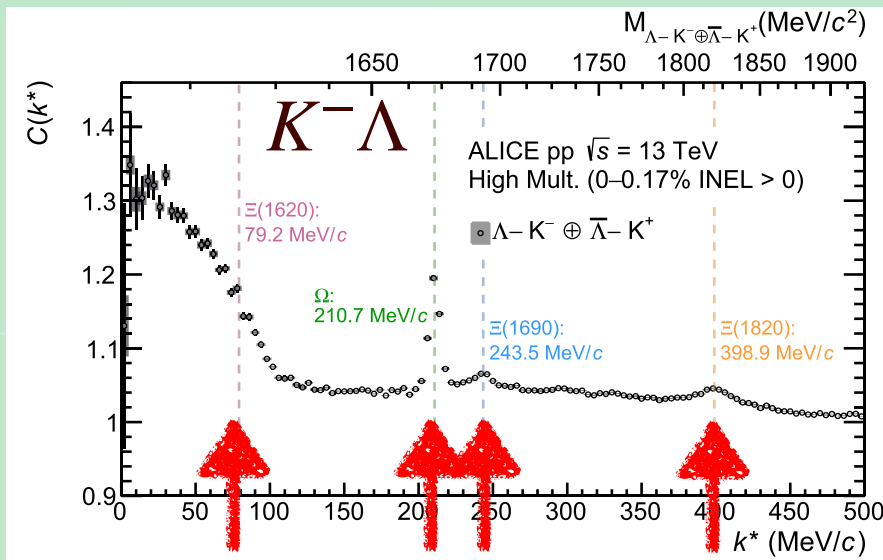
Resonance contributions

Resonance peaks in experimental data ($\ell = 0$ and $\ell \neq 0$)

ALICE collaboration, PRL 124, 092301 (2020); PLB845, 138145 (2023)



$\Lambda(1520)$: d wave



$\Xi(1620)$, $\Xi(1690)$: s wave,
 Ω : p wave (weak decay),
 $\Xi(1820)$: d wave

Questions:

- **Origin** of resonance peak?
- Contribution from **higher partial waves**?

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K. Murase, T. Hyodo, arXiv:2509.22844 [nucl-th]

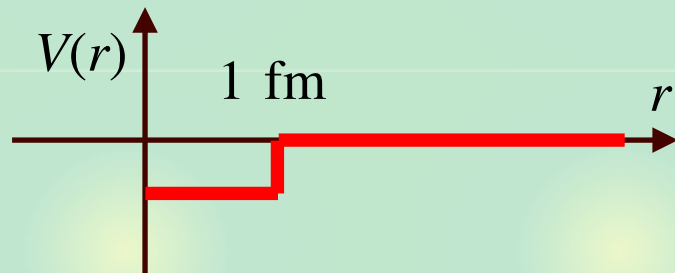
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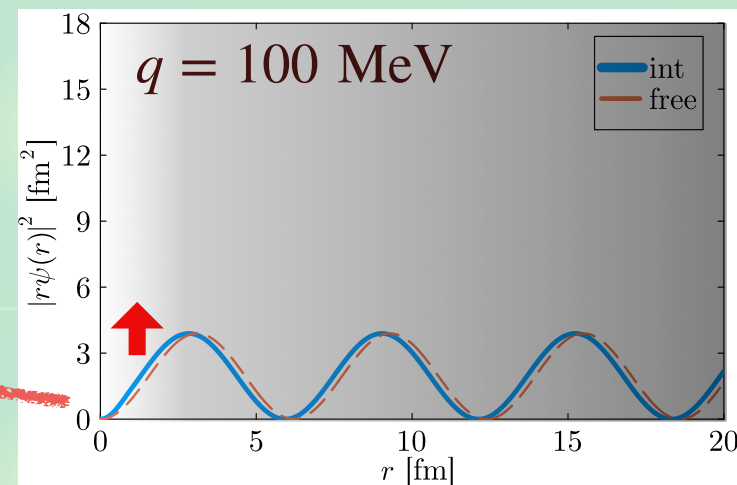
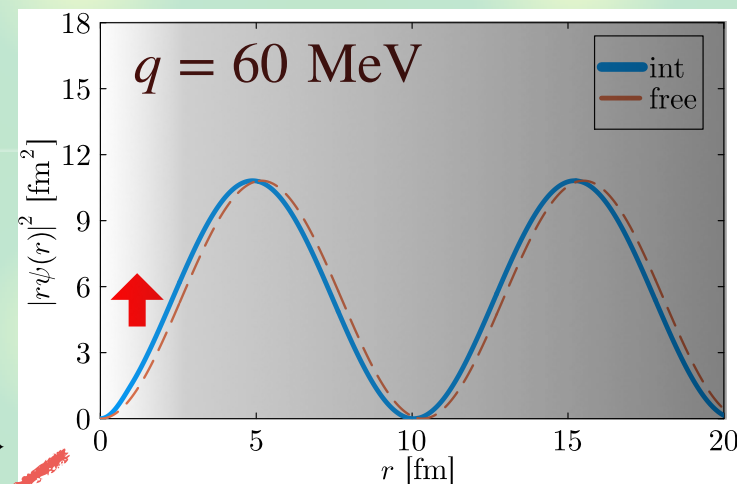
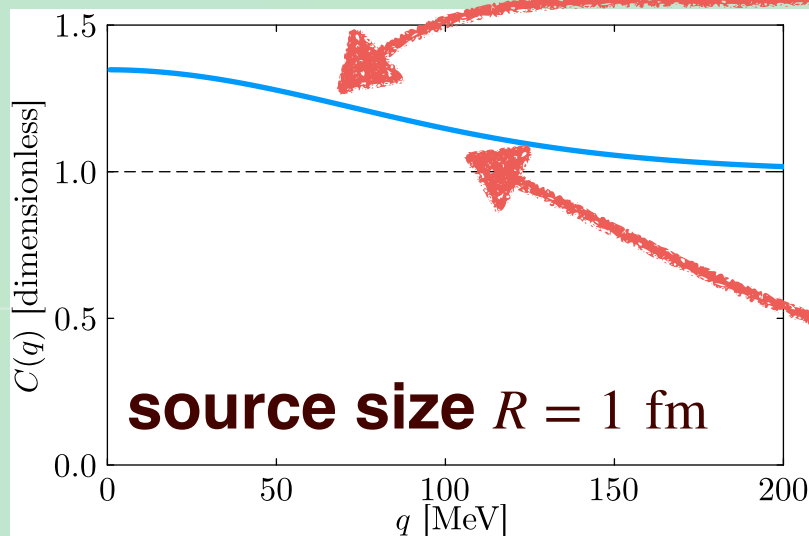
Summary

Wave function and correlation (attraction)

Attractive well (no bound state)



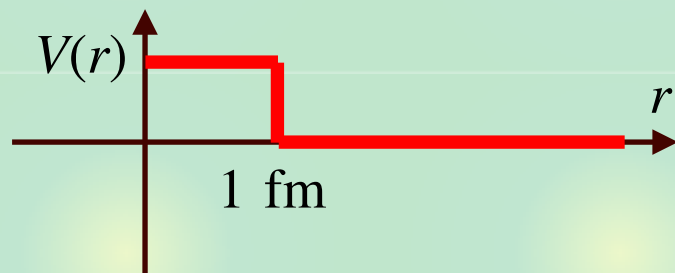
$$C(q) \simeq 1 + \int_0^\infty dr \frac{4\pi}{q^2} S(r) \{ ||rq\psi_q(r)||^2 - \sin^2(qr) \}$$



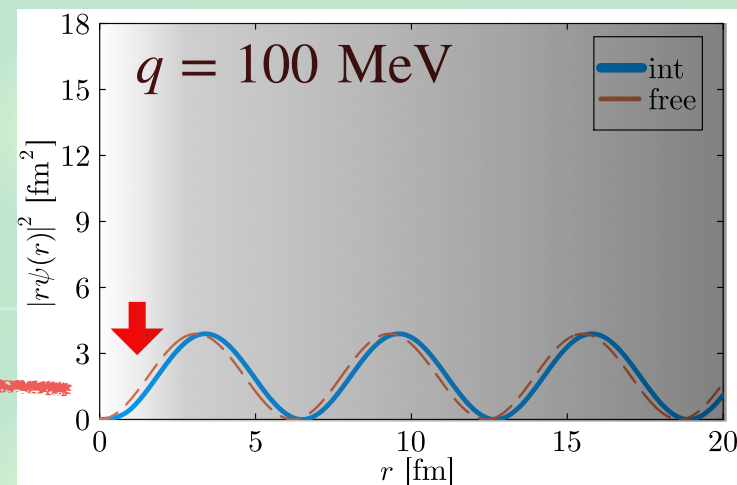
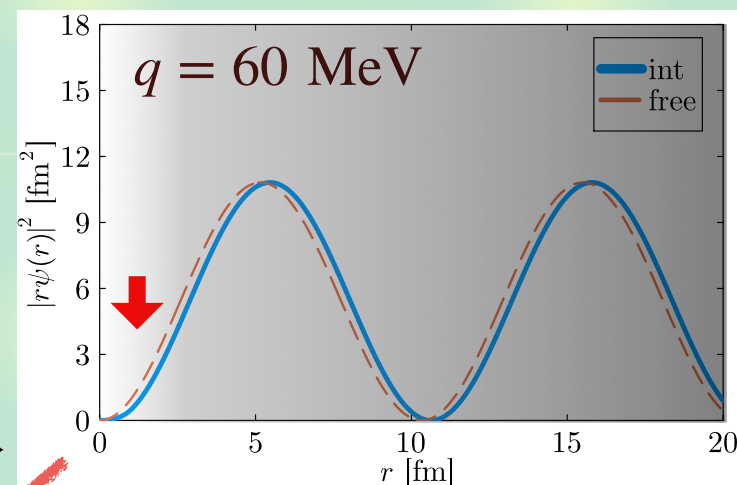
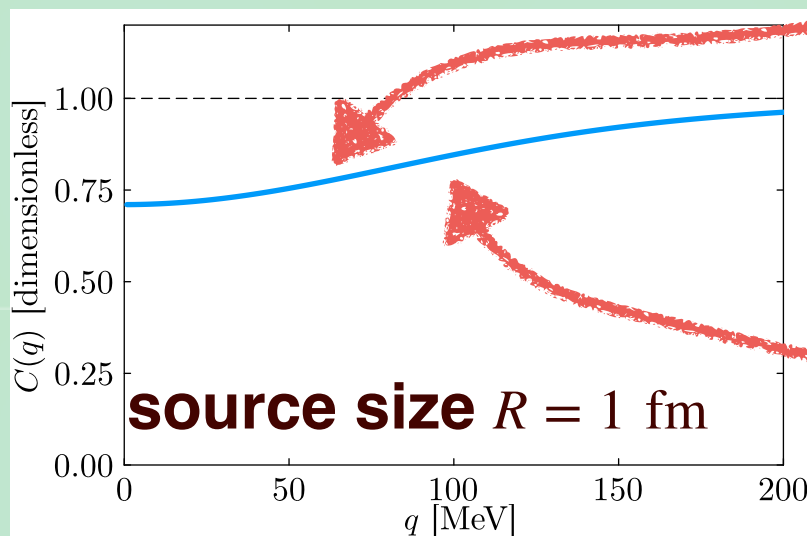
W.f. is pulled in, increased at $r \lesssim R \rightarrow C(q)$ enhancement

Wave function and correlation (repulsion)

Repulsive rectangular potential



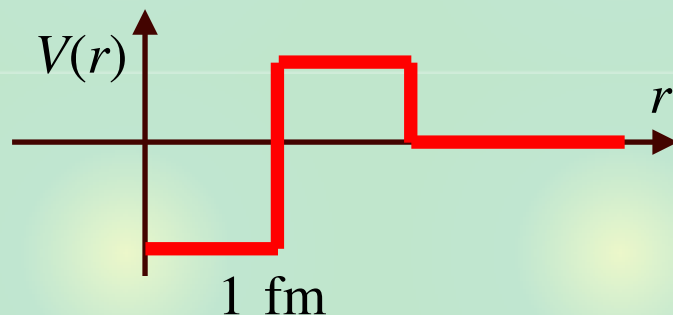
$$C(q) \simeq 1 + \int_0^\infty dr \frac{4\pi}{q^2} S(r) \{ ||rq\psi_q(r)||^2 - \sin^2(qr) \}$$



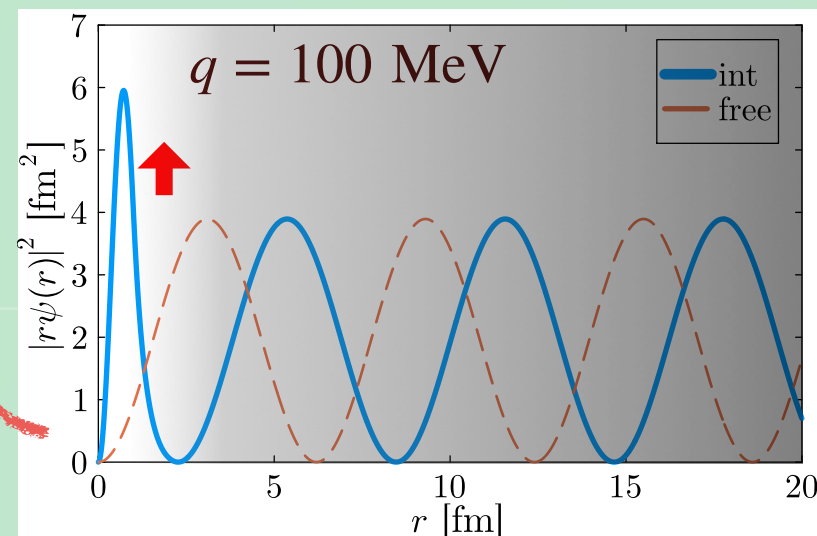
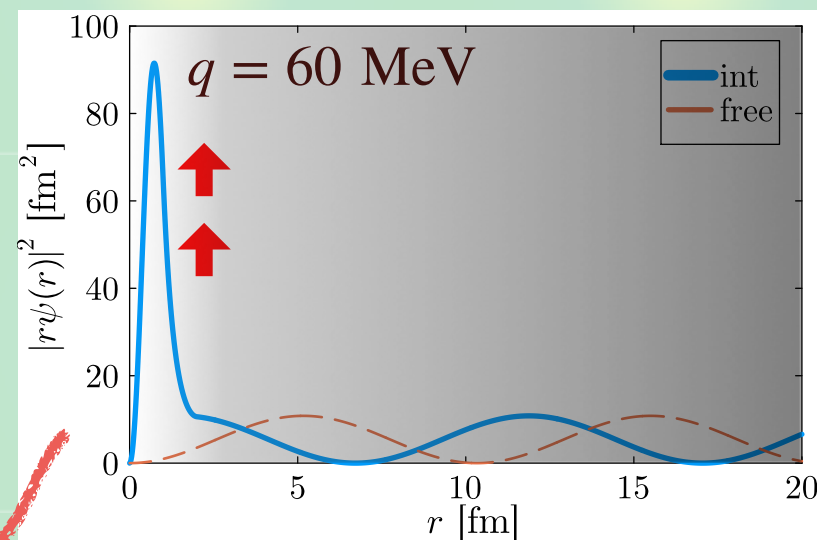
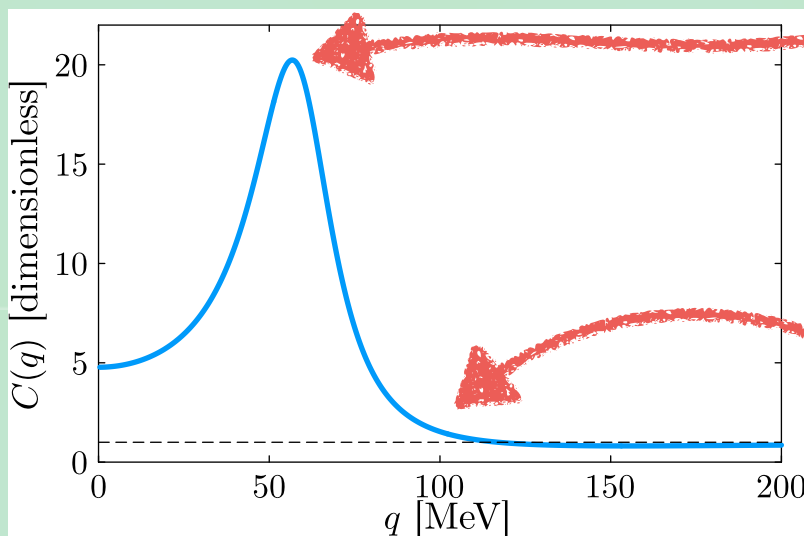
W.f. is pushed out, decreased at $r \lesssim R \rightarrow C(q)$ suppression

Wave function and correlation (resonance)

Well + barrier potential



- resonance @ $q^- = 59 - 14i$ MeV



W.f. is **localized** in $r \lesssim R$ at pole momentum \rightarrow peak in $C(q)$

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Summary

Higher partial waves

KP formula with $l > 0$ (spherical source)

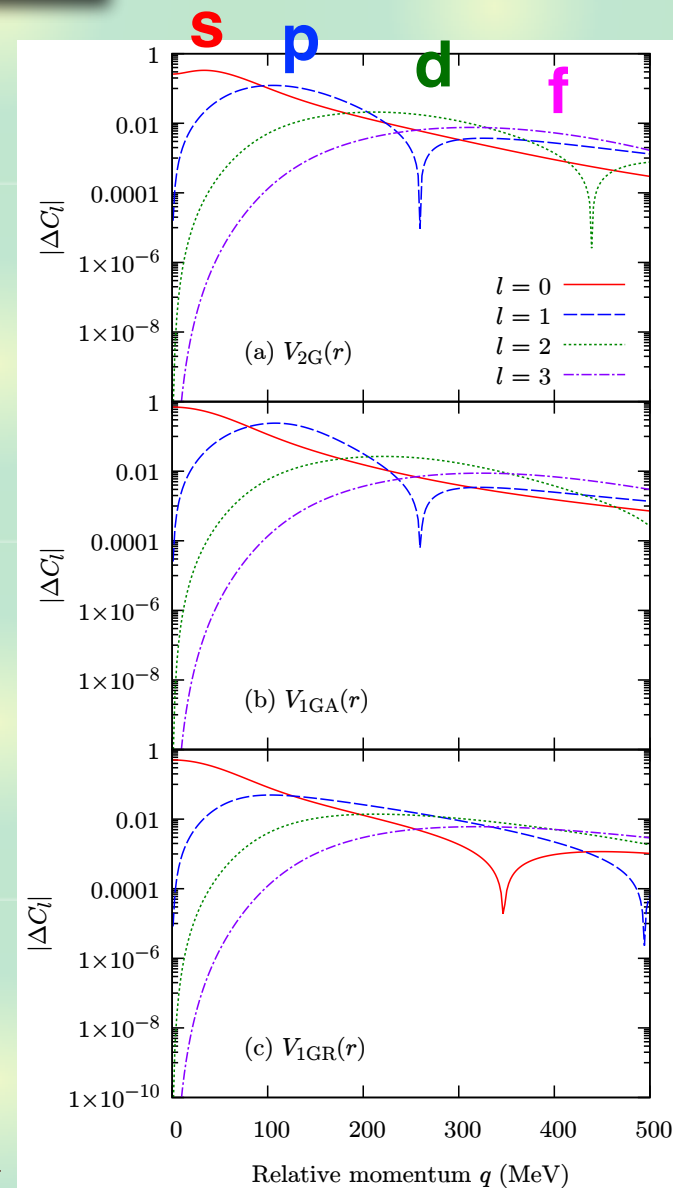
$$C(q) = 1 + \sum_{l=0}^{\infty} \Delta C_l(q)$$

$$\Delta C_l(q) = (2l+1) \times \int_0^{\infty} dr 4\pi r^2 S(r) [|R_l(r)|^2 - |j_l(qr)|^2]$$

- sum of partial wave contributions
- interacting w.f. $R_l(r)$ — free w.f. $j_l(qr)$

Gaussian potentials (range ~ 1.25 fm)

- $l > 0$ components at **larger** q
- l -th wave dominant at $q \sim l/r \sim 160l$ MeV



Resonance contribution

With $q \rightarrow \infty$, $C(q)$ approaches unity

- How can resonances be seen?

Resonances in higher partial waves

- p-wave resonance at

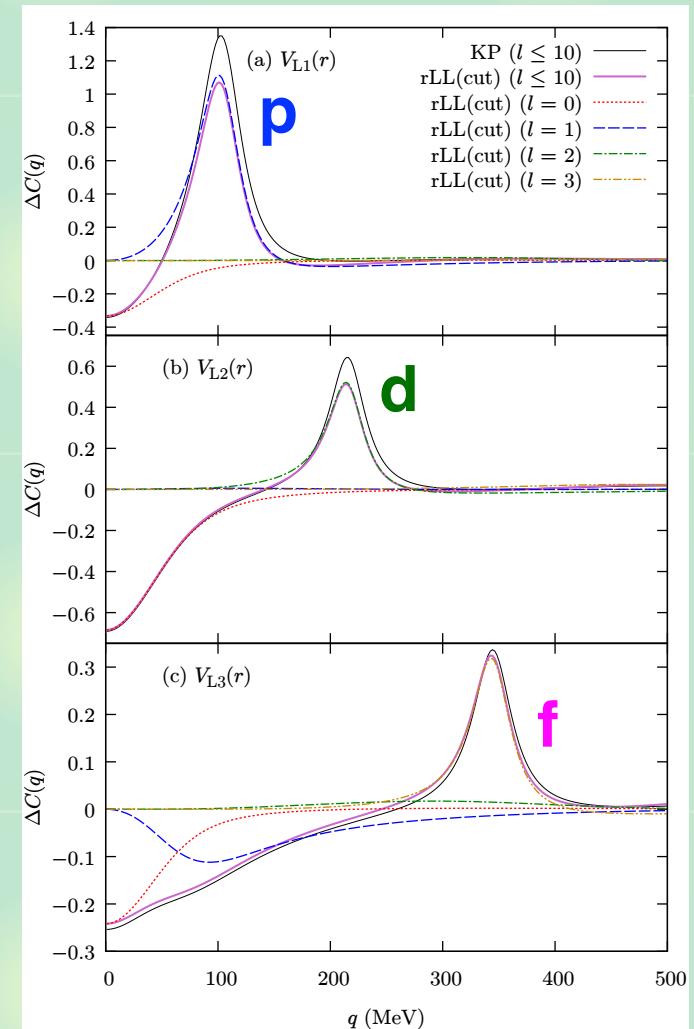
$$q^- \sim 105 - 23i \text{ MeV}$$

- d-wave resonance at

$$q^- \sim 216 - 20i \text{ MeV}$$

- f-wave resonance at

$$q^- \sim 345 - 21i \text{ MeV}$$



With **resonances**, $l > 0$ components is enhanced

LL formula for higher partial wave?

Lednický-Lyuboshits (LL) formula for s-wave

R. Lednický, V.L. Lyuboshits, Yad. Fiz. **35**, 1316 (1981);

K. Murase, T. Hyodo, J. Subatomic Part. Cosmol. **3**, 100017 (2025)

- Asymptotic w.f. for entire $r \leftarrow$ **on-shell observable**

$$\Delta C(q) = \int_0^\infty dr \frac{4\pi}{q^2} S(r) [\sin^2(qr - \delta(q)) - \sin^2(q)]$$

- corresponding to zero range limit (point-like interaction)

Naive generalization for $l > 0$ is **not possible**

$$\Delta C_l(q) = (2l + 1) \int_0^\infty dr 4\pi r^2 S(r) [|R_l^{\text{aym}}(r)|^2 - |j_l(qr)|^2]$$

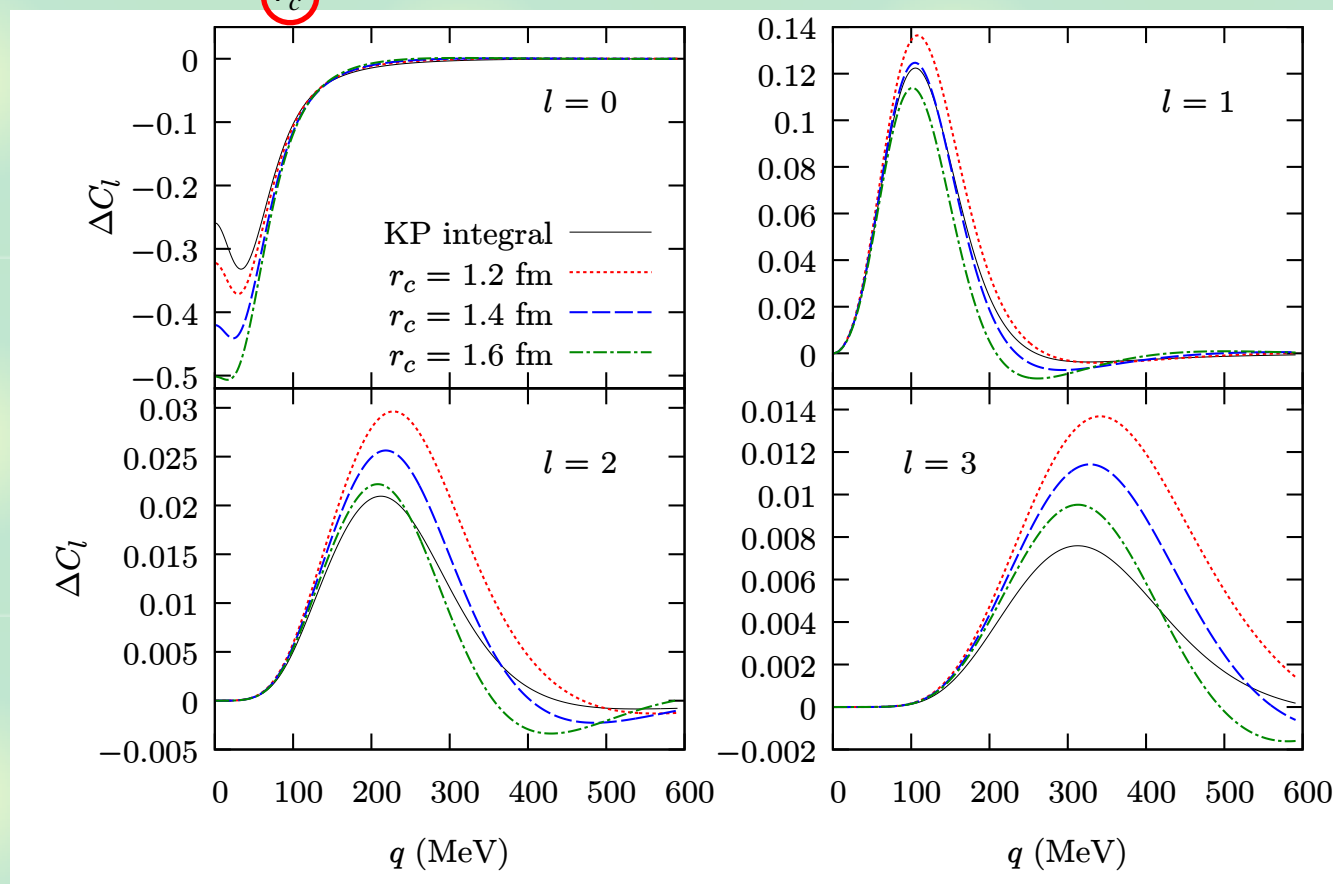
- $R_l^{\text{aym}}(r)$ is too singular for $l > 0$ ($n_l(qr) \sim r^{-l-1}$) at $r \rightarrow 0$

Some **regularization** for $r \rightarrow 0$ is needed

Regularized LL formula


Simplest choice: introducing **cutoff** r_c

$$\Delta C_l(q) = (2l + 1) \int_{r_c}^{\infty} dr 4\pi r^2 S(r) [|R_l^{\text{aym}}(r)|^2 - |j_l(qr)|^2]$$




Works for $l > 0$ with $r_c \sim$ interaction range (1.25 fm)

Summary




Resonance peaks in correlation functions




s-wave resonance peak \leftarrow **localization** of wave function at interacting region

S. Watanabe, T. Hyodo, in preparation



Higher partial wave ($l > 0$) contributions becomes important for larger q



Regularized LL formula with suitable cutoff r_c works for $l > 0$

K. Murase, T. Hyodo, J. Subatomic Part. Cosmol. 3, 100017 (2025);

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