

Properties of excited baryons in a deformed oscillator quark model

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Contents

Systematics in the masses of excited baryons

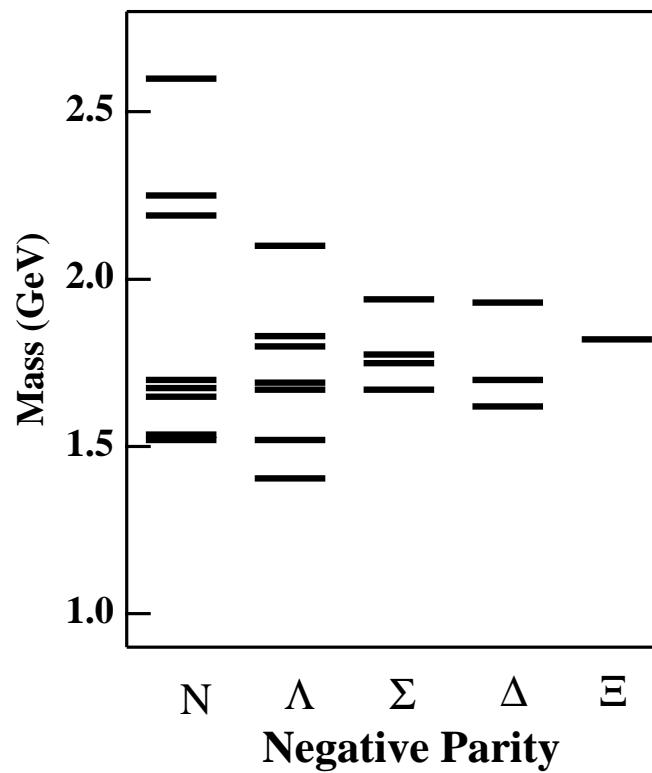
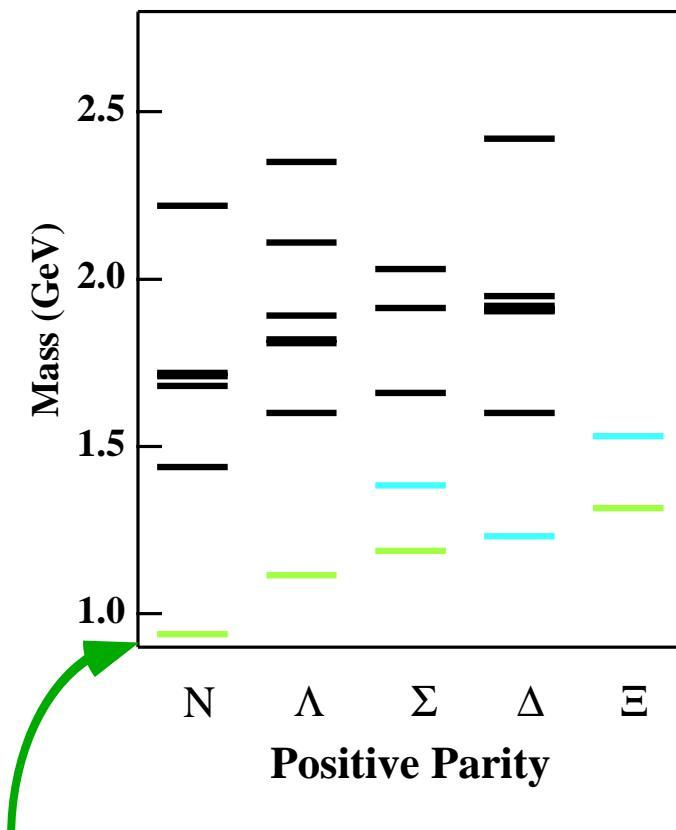
Deformed Oscillator Quark Model (DOQ)

Intra-band transitions

Summary

Systematics in the mass spectra

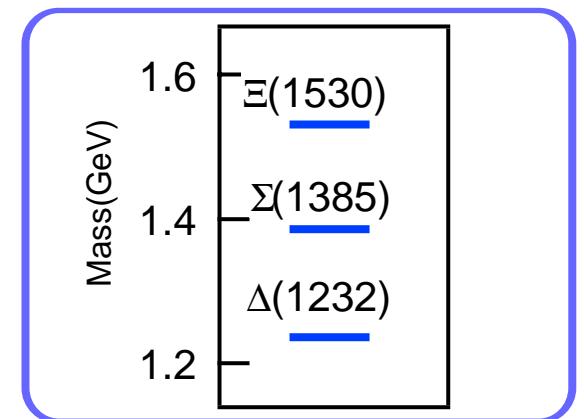
- Observed mass spectra [PDG (2000)]



● Gell-Mann-Okubo Mass Formula

$$\frac{m_{\Sigma} + 3m_{\Lambda}}{2} = m_{\Xi} + m_N$$

● Equal Mass Splittings of Baryon Decuplet



Ground states

$$m_N = m_0 + 3m_u$$

$$m_{\Sigma} = m_0 + 2m_u + m_s$$

$$m_{\Lambda} = m_0 + 2m_u + m_s$$

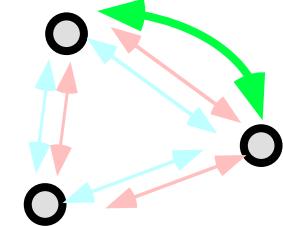
$$m_{\Xi} = m_0 + 2m_s + m_u$$

Constituent Quark Model

- Constituent Quark model

Baryon \rightarrow three quark system

quark \rightarrow $m \sim 300$ MeV ,
Spin, Flavor, and Color degrees of freedom



$$|(Baryon)\rangle = |(Orbit)\rangle |(Spin)\rangle |(Flavor)\rangle |(Color)\rangle$$

The wave functions must be **Totally anti-symmetric**;
change overall sign under every single interchange of any quark pair.

For the ground state ----

Orbit \rightarrow Symmetric

Color \rightarrow Anti-symmetirc

Spin \rightarrow 1/2 (2) Mixed symmetric
3/2 (4) Symmetric

Flavor \rightarrow (8) Mixed symmetric
(10) Symmetric
(1) Anti-symmetirc

Symmetry of the product of two wave functions

S * S --- S

A * A --- S

S * A --- A

S * MS --- MS

A * MS --- MS

MS * MS --- S, A, MS

28 and **410** are allowed

Ground states are assigned to ...

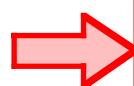
28 --- N, Λ, Σ(1189), Ξ(1317)

410 -- Δ, Σ(1385), Ξ(1530), Ω

Naming Scheme for Baryons

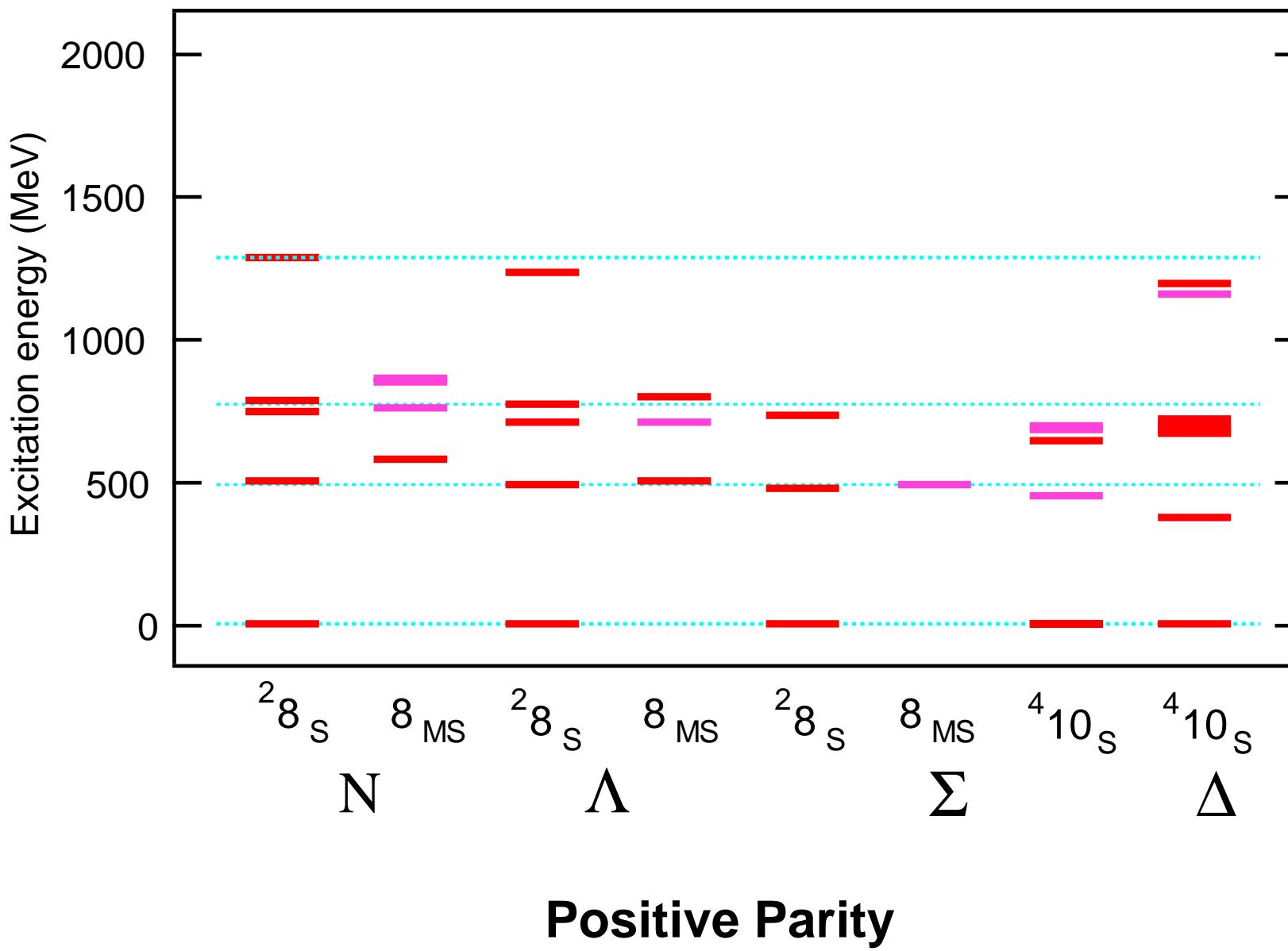
| Isospin | Strangeness | | | |
|-----------|-------------|-----------|----------|----------|
| | $S = 0$ | $S = -1$ | $S = -2$ | $S = -3$ |
| $I = 0$ | - | Λ | - | Ω |
| $I = 1/2$ | N | - | Ξ | - |
| $I = 1$ | - | Σ | - | - |
| $I = 3/2$ | Δ | - | - | - |

They are the ground states of each Spin-Flavor Multiplet

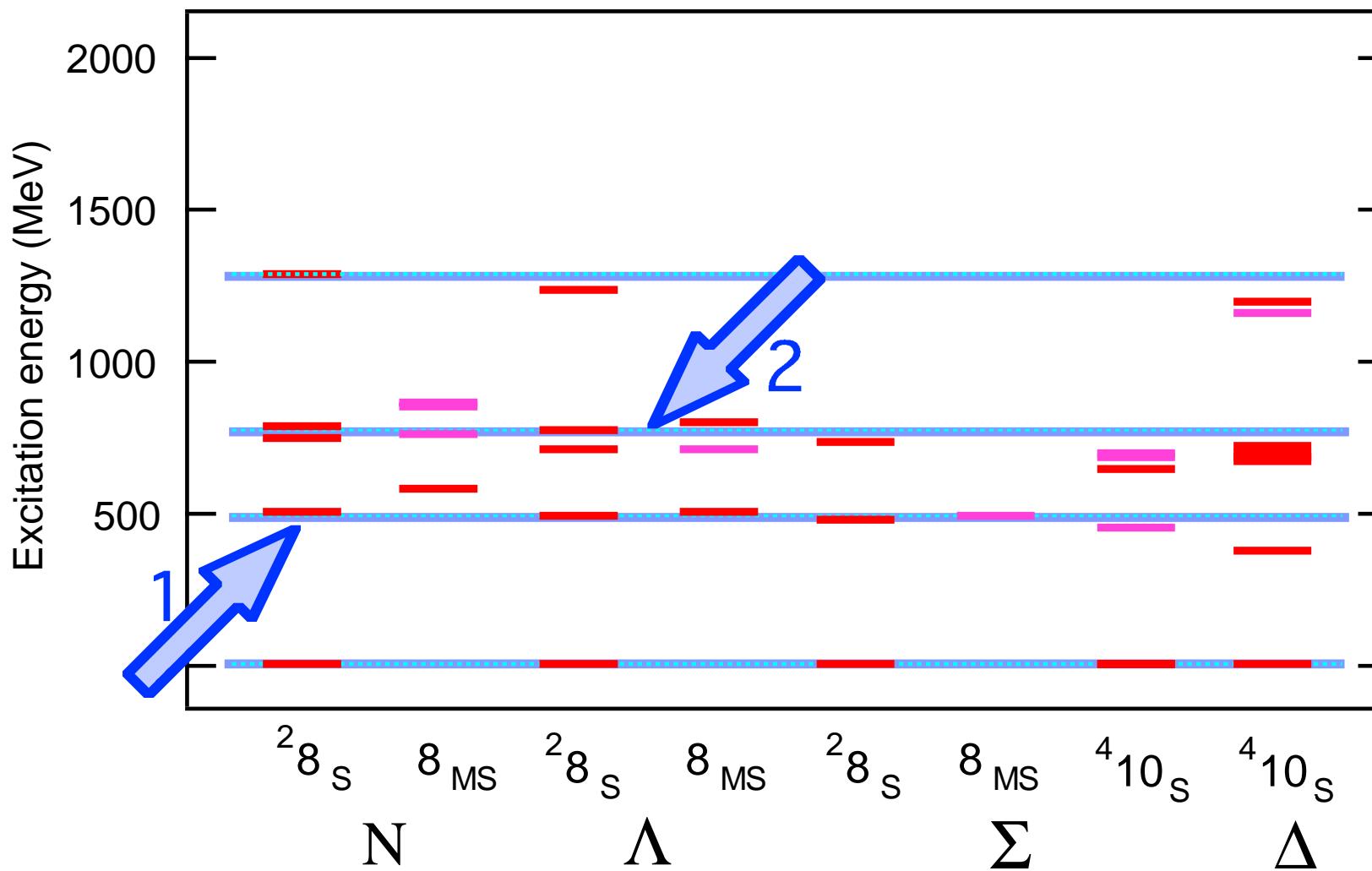


**Measure the Masses of Excited
Baryons from these Ground States**

- After classification ...



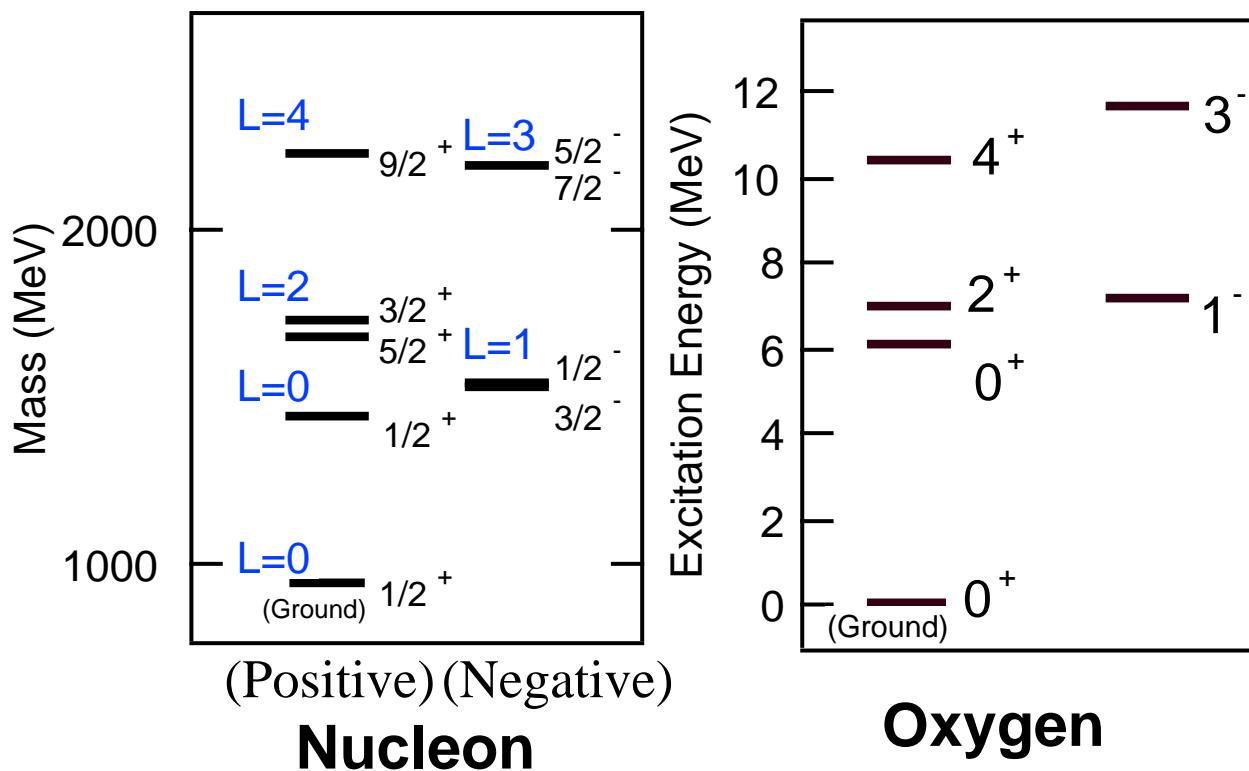
- After classification ...



Remarkable Systematics!

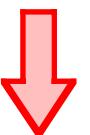
1. First Excited States --- 500 MeV
2. Second Excited States --- 700 MeV, Degeneracy of Spin States

- Nucleon spectra -- Excitation energy --



- ★ Feature of the spectra
= Rotational Band Structure
- Level Spacing (Positive parity) < Negative parity ($L(L+1)$ structure)
 - Degeneracy of Spin state $J=L+S$
 - The Roper state ($1/2^+$) is Identified with the Band Head

Analogy to Energy Spectrum of Deformed Nuclei



Are Excited Baryons Spatially Deformed ?

Deformed Oscillator Quark Model (DOQ)

- Hamiltonian

$$H_{DOQ} = \sum_{i=1}^3 \left(\frac{\mathbf{P}_i^2}{2m} + \frac{1}{2}m(\omega_x^2 x_i^2 + \omega_y^2 y_i^2 + \omega_z^2 z_i^2) \right) - H_{c.m.}$$

Murthy et al., PRD **30** ('84) 152
Hosaka et al., MPLA **13** ('98) 1699

$$E_{int} = (N_x + 1)\omega_x + (N_y + 1)\omega_y + (N_z + 1)\omega_z$$

$$N_x = n_{\rho_x} + n_{\lambda_x}$$

- Intrinsic state

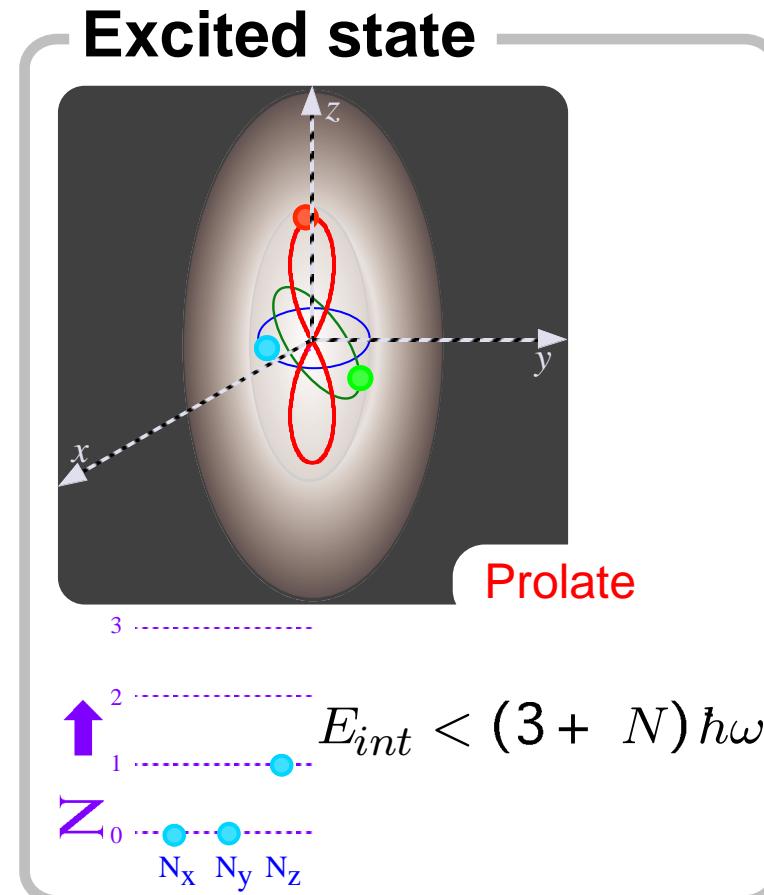
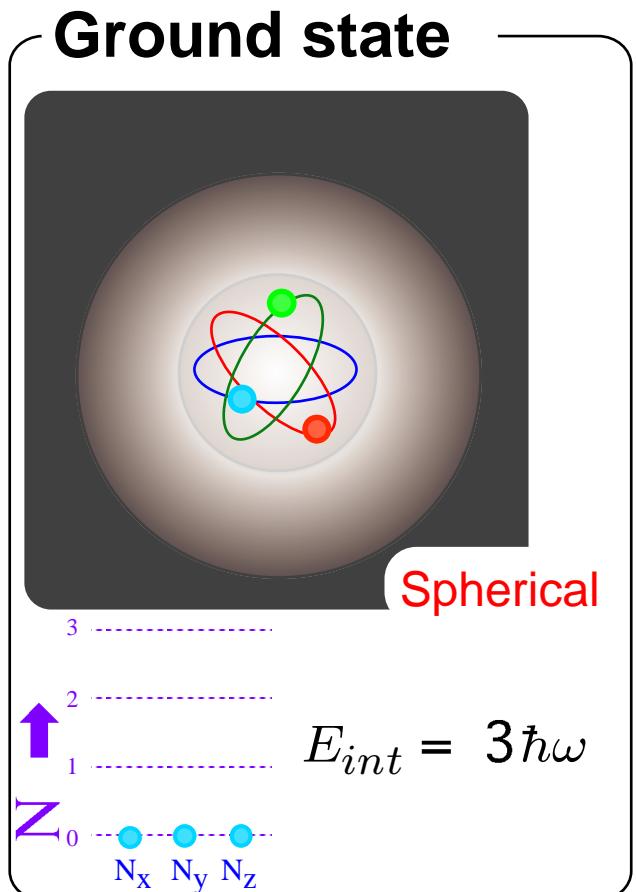
- Axial symmetry (for $N = N_x + N_y + N_z = 0, 1, 2$) $\rightarrow N_x = N_y$

- Energy Minimization (Volume Conservation Condition)

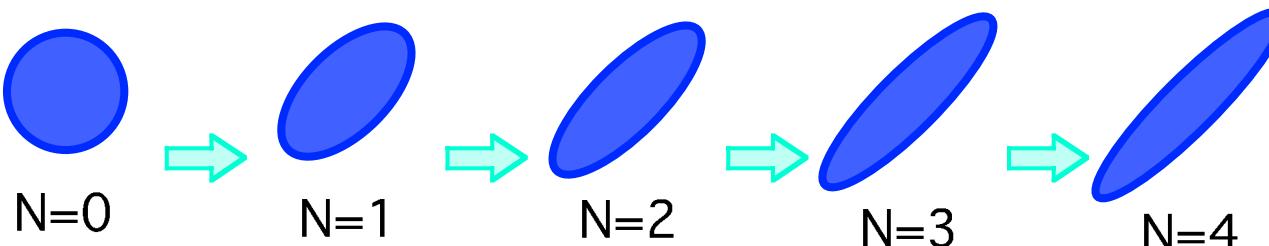
$$\frac{\partial E}{\partial \omega_x} = 0 \quad \frac{\partial E}{\partial \omega_z} = 0 \quad \longleftrightarrow \quad \omega_x \omega_y \omega_z = \omega^3 = \text{const.}$$

- Energy Minimization (Volume Conservation Condition)

$$\frac{\partial E}{\partial \omega_x} = 0 \quad \frac{\partial E}{\partial \omega_z} = 0 \quad \leftarrow \quad \omega_x \omega_y \omega_z = \omega^3 = \text{const.}$$



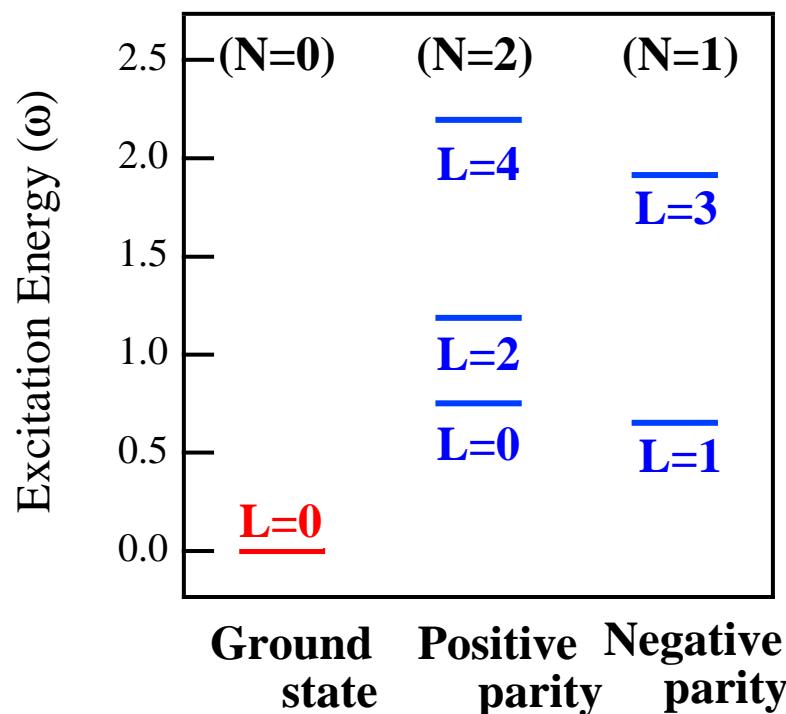
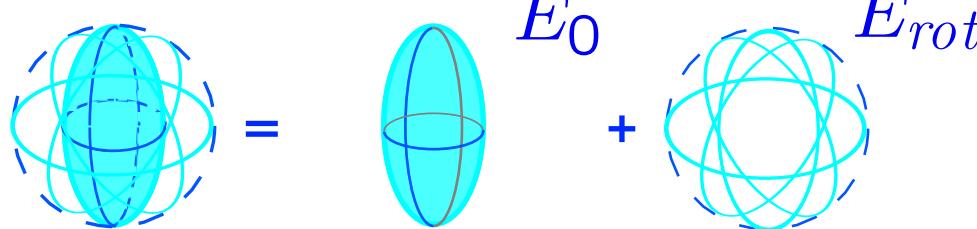
Prolate states have lower energy than oblate states.



$$d \equiv \frac{\omega_x}{\omega_z} \rightarrow N + 1$$

Total Energy

$$E_{total}^{N,L} = E_{int}^N - \frac{\langle L^2 \rangle}{2\mathcal{I}} + \frac{L(L+1)}{2\mathcal{I}}$$



\mathcal{I} : Moment of Inertia
(Cranking Model)

$\langle L^2 \rangle$: Expectation Value of L^2

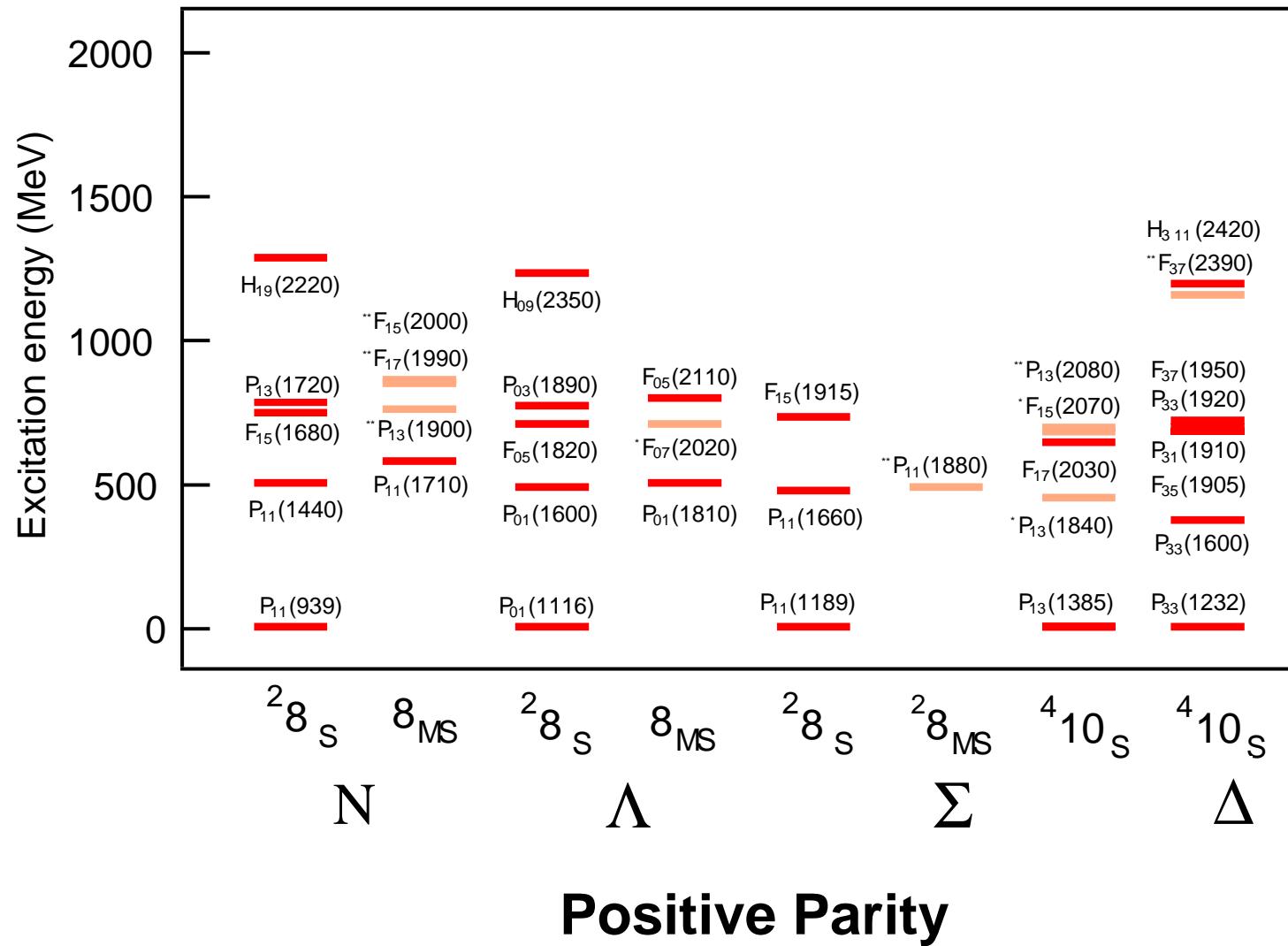
Moment of Inertia and the expectation value of L^2 can be obtained as a function of N , ω .

Energy of each state is written in units of ω

ω is our only one parameter!

- Comparison ---Data v.s. DOQ---

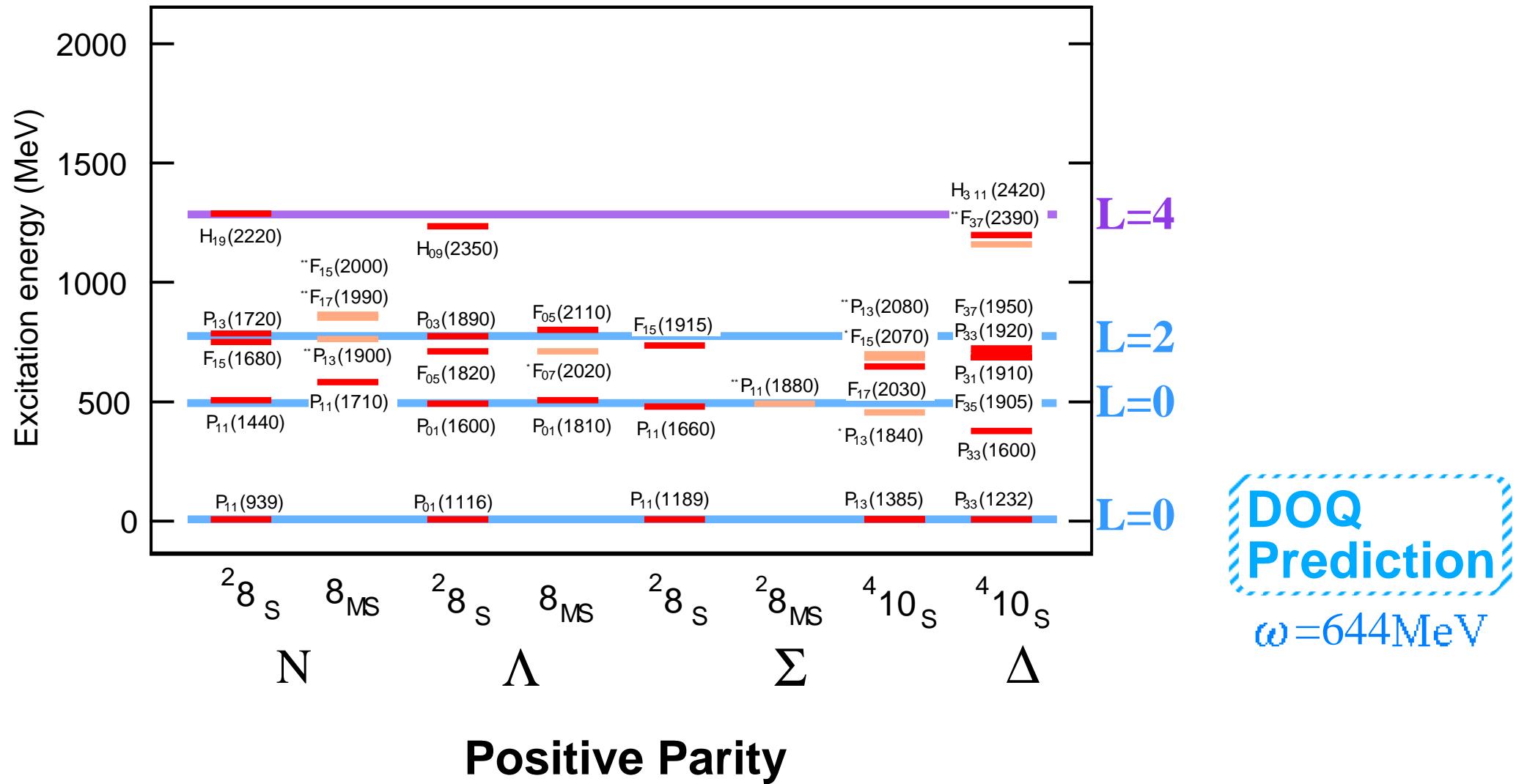
★All existing data with 3 or 4 stars — and some with 1 or 2 stars — are shown.

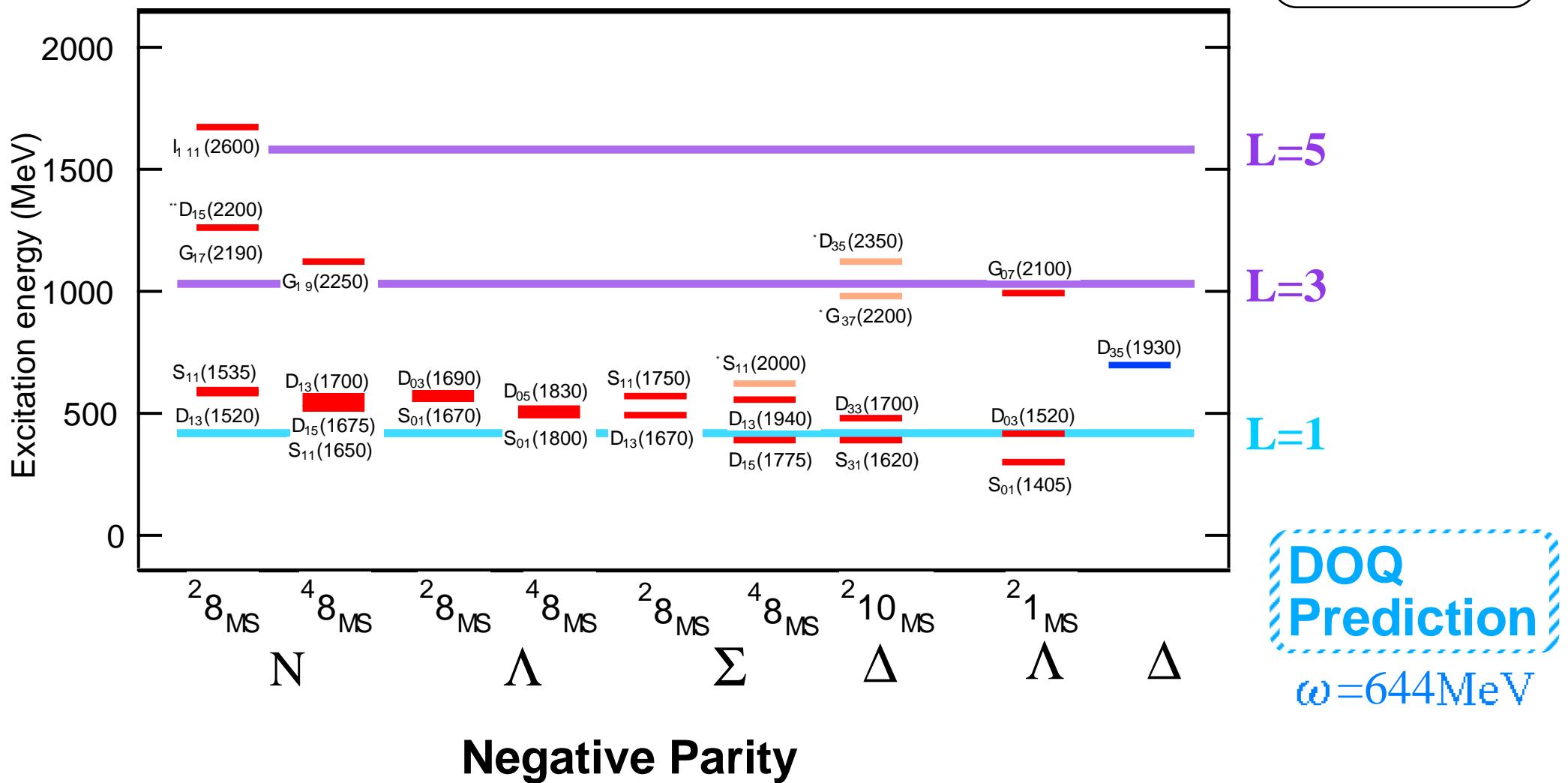


- Comparison ---Data v.s. DOQ---

★All existing data with 3 or 4 stars — and some with 1 or 2 stars — are shown.

Takayama et al.,
PTP 101('99)





cf. Regge trajectory (M^2 vs. J)

Total Spin J or Angular momentum L

Mass or Excitation Energy

Summary

- We have studied systematically all existing data of flavor SU(3) baryons in the scheme of the DOQ.

● Numbers of States

Almost all states fall into the DOQ systematics.

| | (***, *****) | (*, **) |
|---------------|--------------|---------|
| known baryons | 5 0 | 3 1 |
| DOQ | 4 9 | 1 3 |

Single Parameter!
 $\omega = 644 \text{ MeV}$

There is a very simple, flavor independent systematics in the masses of flavor SU(3) baryons.

The structure of the spectrum can be described as rotational bands of deformed excited state.

For further study ...

Is there any effect of deformation in the transition properties ?

- Inter-band transitions

- Intra-band transitions

Positive Parity
band

Negative Parity
band

Intra-band
transitions

Inter-band
transitions

Spherical Ground State

It is difficult to observe Intra-band transitions with a photon

-- Transitions with a pion

Intra-Band transitions in the DOQ Model

We calculate the decay width of $N^{*'} \rightarrow N^* \pi$
with DOQ wave functions and
Non-relativistic $\pi q\bar{q}$ interaction Hamiltonian

Decay Width:

$$W_{f,i} = 2\pi\delta(E_{P_f} + \omega_k - E_{P_i}) |\mathcal{M}|^2$$

Transition Amplitude: $\mathcal{M} = \langle P_f; J_f, M_f | H | P_i; J_i, M_i \rangle$

Momentum, spin and its z component: $P_{i(f)}; J_{i(f)}, M_{i(f)}$

Energy and momentum of emitted pion: \vec{k}, ω_k

Integrate (Average) over...

Spin z-component of Initial and final states
Momentum and charge of emitted pion

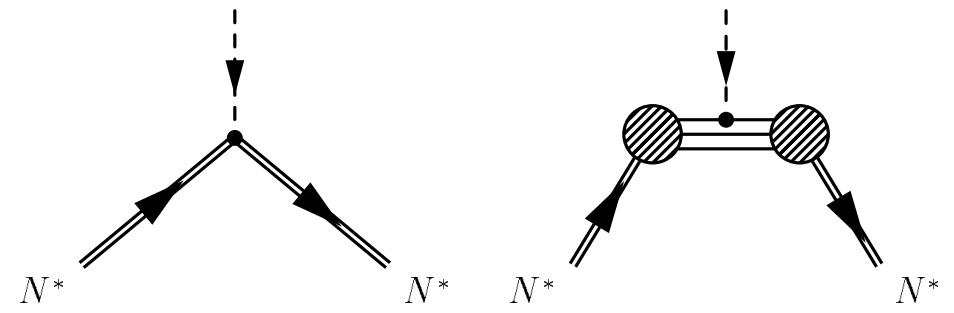
$$W = 6\pi \int \frac{d^3 k}{(2\pi)^3} \delta(E_{P_f} + \omega_k - E_{P_i}) \frac{1}{2J_i + 1} \sum_{M_i, M_j} |\mathcal{M}|^2$$

Non-relativistic πqq interaction Hamiltonian

$$\mathcal{M} = \langle P_f; J_f, M_f | H | P_i; J_i, M_i \rangle$$

$$\rightarrow 3 \int \Psi^*(P_f; J_f, M_f) H_{\pi qq} \Psi(P_i; J_i, M_i)$$

DOQ wave function



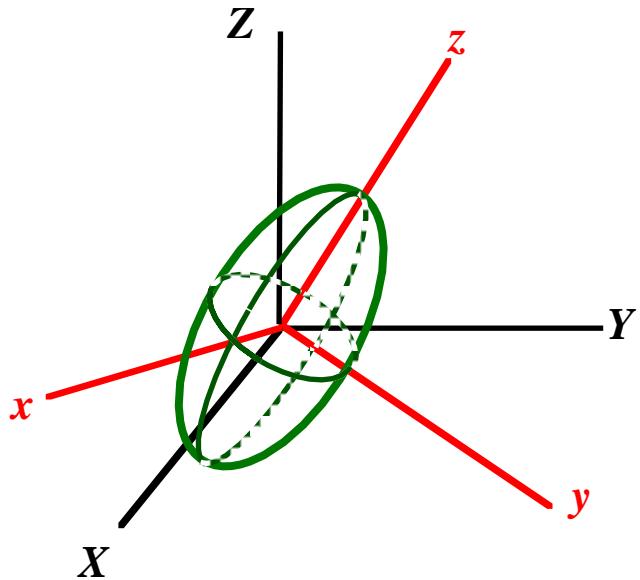
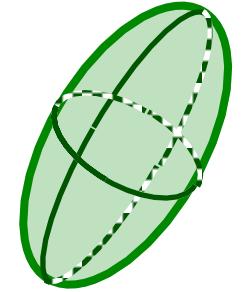
Permutation symmetry of quarks in
the baryons

$$H_{\pi qq} = -\frac{g}{2m} \frac{1}{\sqrt{2\omega_k}} \left(\vec{\sigma} \cdot \vec{\nabla} \exp(i\vec{k} \cdot \vec{x}) \right)$$

DOQ wave function --- Angular momentum projection ---

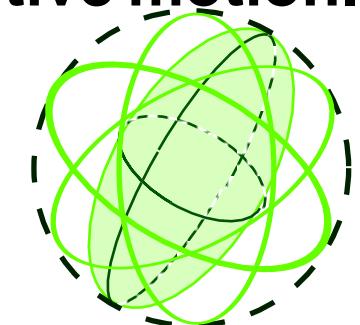
1) Intrinsic wave functions obtained by the variation are deformed.

$$\langle \{ \vec{x} \} | N^{(int)} \rangle \sim$$



2) Introduce the Euler angle $\Theta = (\alpha, \beta, \gamma)$ which connect the body-fixed frame and the Lab-frame.
Wave function can be written as a product of intrinsic wave function and collective motion.

$$D_{0L_z}^L(\vec{\Theta}) \langle \{ \vec{x} \} | N^{(int)} \rangle \sim$$



**Spin-Orbital wave function
with total spin J
and its z-component M**

$$\Psi(P; J, M) = \left[\sqrt{\frac{2L+1}{4\pi}} D_{0L_z}^L \langle \{ \vec{x}_i \} | N^{(int)} \rangle \otimes | \chi_\mu \rangle \right]_{JM}$$

Transition amplitudes

1) Separate the contribution from the deformed intrinsic state and the collective motion

$$\mathcal{M} = \int d\Theta \langle [L_f, 1/2]_{M_f}^{J_f} | 3\mathcal{O}_N^{int}(\Theta) | [L_i, 1/2]_{M_i}^{J_i} \rangle$$

2) Integration over the body-fixed coordinate
= Operator in the collective coordinate

$$\mathcal{O}_N^{int}(\Theta) \equiv \int d\{\vec{x}_i\} \langle N^{(int)} | \{\vec{x}_i\} \rangle \langle \{\vec{x}_i\} | \hat{R}^\dagger(\Theta) H_{\pi qq} \hat{R}(\Theta) | \{\vec{x}_i\} \rangle \langle \{\vec{x}_i\} | N^{(int)} \rangle$$

$$\mathcal{O}_N^{int}(\Theta, \vec{k}) = -\frac{(4\pi)^{3/2}g}{2m\sqrt{2\omega_k}} \sum_{l,m} i^l k Y_{lm}^*(\hat{k})$$

$$\times \left[\sqrt{\frac{l+1}{(2l+1)(2l+3)}} [Y_{l+1}\sigma]_m^l Q_N^{(l+1)}(k) + \sqrt{\frac{l}{(2l+1)(2l-1)}} [Y_{l-1}\sigma]_m^l Q_N^{(l-1)}(k) \right]$$

l -moment: $Q_N^{(l)}(k) \equiv \langle N^{(int)} | j_l(kx) Y_{l0}(\hat{x}) | N^{(int)} \rangle$

The l -moment contain all information of the deformed intrinsic states.

3) Integrate over the collective coordinate

$$\begin{aligned} \mathcal{M} = & \frac{3(4\pi)g}{2m\sqrt{2\omega_k}} \sum_{l,m} i^l k Y_{lm}^*(\hat{k}) (-)^m \begin{pmatrix} J_f & l & J_i \\ -M_f & m & M_i \end{pmatrix} \frac{\hat{J}_f \hat{J}_i (-1)^{J_f+L_f+1/2}}{\hat{l}} \\ & \times \begin{pmatrix} J_i & J_f & l \\ 1/2 & -1/2 & 0 \end{pmatrix} \left[\frac{(L_i - J_i)(2J_i + 1) + (L_f - J_f)(2J_f + 1) + l + 1}{\widehat{l+1}} Q_N^{(l+1)}(\boldsymbol{k}) \right. \\ & \quad \left. + \frac{(L_i - J_i)(2J_i + 1) + (L_f - J_f)(2J_f + 1) - l}{\widehat{l-1}} Q_N^{(l-1)}(\boldsymbol{k}) \right] \end{aligned}$$

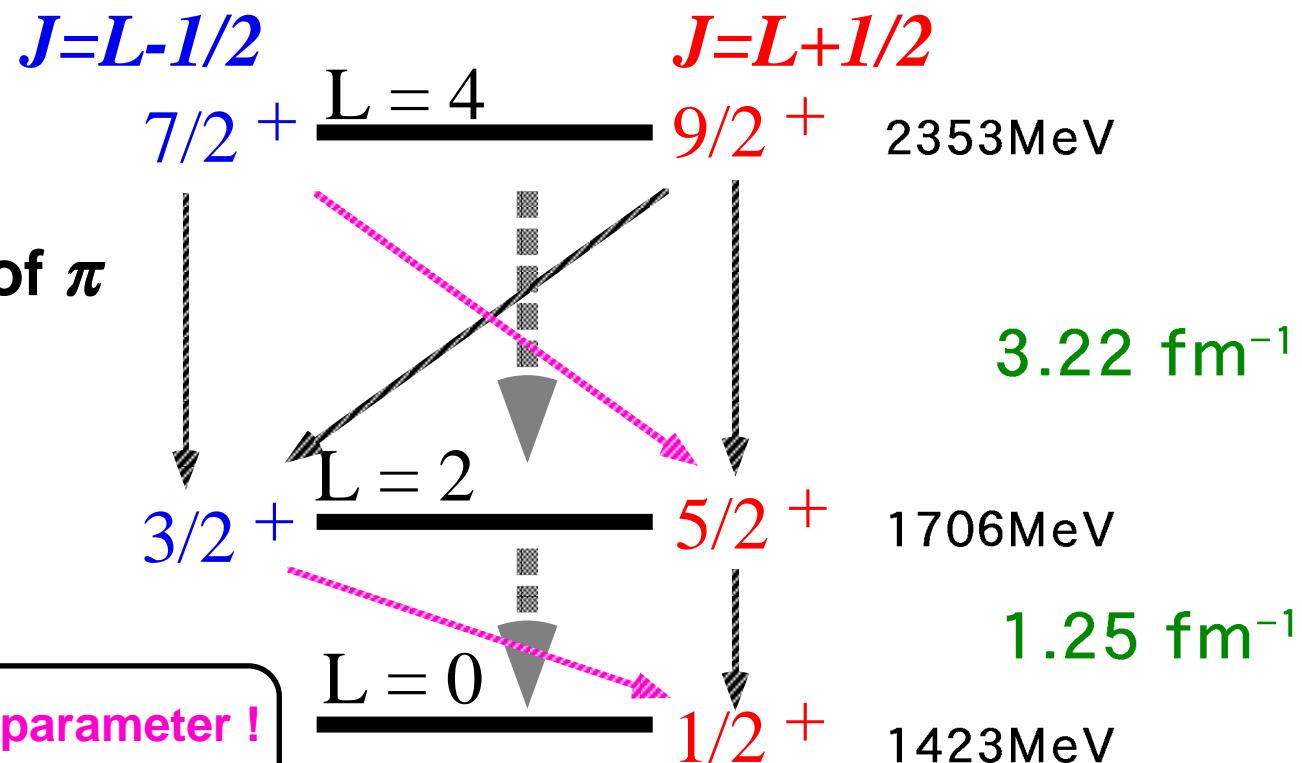
4) Substitute into the decay width (Integral over the momentum of π^-)

$$W = \frac{27g^2}{2m^2} \tilde{k} \sum_l \left| F_{l;L_i J_i L_f J_f}^{(N)}(\tilde{k}) \right|^2 \quad \boxed{\tilde{k} = \sqrt{(E_i - E_f)^2 - m_\pi^2}}$$

$$\begin{aligned} F_{l;L_i J_i L_f J_f}^{(N)}(\boldsymbol{k}) \equiv & i^l k \frac{\hat{J}_f (-1)^{J_f+L_f+1/2}}{(2l+1)} \begin{pmatrix} J_i & J_f & l \\ 1/2 & -1/2 & 0 \end{pmatrix} \quad (\text{Geometrical factor}) \\ & \times \left[\frac{(L_i - J_i)(2J_i + 1) + (L_f - J_f)(2J_f + 1) + l + 1}{\widehat{l+1}} Q_N^{(l+1)}(\boldsymbol{k}) \right. \\ & \quad \left. + \frac{(L_i - J_i)(2J_i + 1) + (L_f - J_f)(2J_f + 1) - l}{\widehat{l-1}} Q_N^{(l-1)}(\boldsymbol{k}) \right] \quad \times (\text{l-th moment}) \end{aligned}$$

Selection rule for N=2, Positive parity band

Angular momentum of π
 $l = 1, 3$
 (Parity conservation)



Parameters: No fit parameter !
 Pion mass : 138 MeV
 Quark mass: 300 MeV
 DOQ parameter: $\omega = 644 \text{ MeV}$
 Deformation parameter for N=2 $d=3$
 πqq coupling constant $g=7.14$

Momentum transfer from DOQ spectrum

Decay Width

| | | | | |
|-----------------------|-----------------------|-----------------------|-----------------------|-----------------------|
| $L = 4 \rightarrow 2$ | $9/2 \rightarrow 5/2$ | $9/2 \rightarrow 3/2$ | $7/2 \rightarrow 5/2$ | $7/2 \rightarrow 3/2$ |
| $\Gamma(\text{MeV})$ | 18.9 | 35.3 | 115 | 0.5 |
| $L = 2 \rightarrow 0$ | $5/2 \rightarrow 1/2$ | - | $3/2 \rightarrow 1/2$ | - |
| $\Gamma(\text{MeV})$ | 0.2 | - | 0.5 | - |

$L = 2 \rightarrow 0$ transitions are almost forbidden

$7/2 \rightarrow 5/2$ transition has large decay width
(a pion with $l = 1$ can contribute in this channel)

Long wave length limit

$$\tilde{k} \rightarrow 0$$

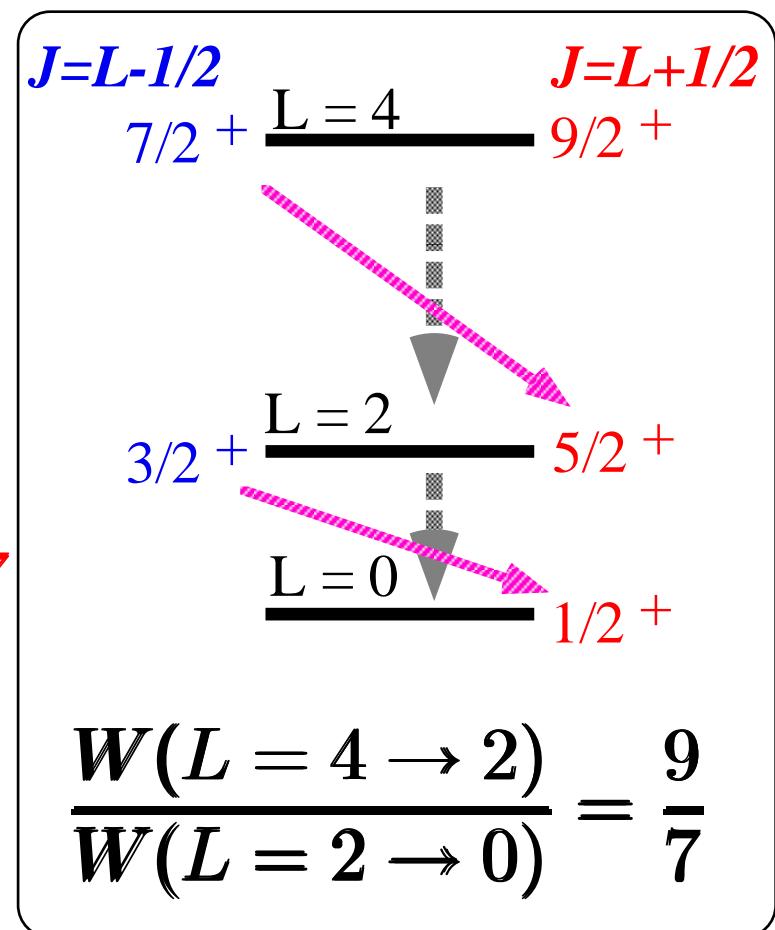
Compare $7/2^+ \rightarrow 5/2^+$, $3/2^+ \rightarrow 1/2^+$ transitions

Contribution of $l = 1$ pion
is dominant for both

$$W = \frac{27g^2}{2m^2} \tilde{k} \left(\tilde{k}^2 \frac{L-1}{2L-1} \frac{1}{5} (Q_N^{(2)}(\tilde{k}))^2 \right) \sim \tilde{k}^7$$

Transitions with low momentum transfer
are strongly suppressed.

Ratio of the width is
determined only by
geometrical factor with
the same momentum
transfer \tilde{k}



$$\frac{W(L=4 \rightarrow 2)}{W(L=2 \rightarrow 0)} = \frac{9}{7}$$

$$\frac{W(L+2 \rightarrow L)}{W(L \rightarrow L-2)} = \frac{(L-3)(2L-1)}{(2L-3)(L-1)}$$

(Alaga rule for Baryons)

l -th moment

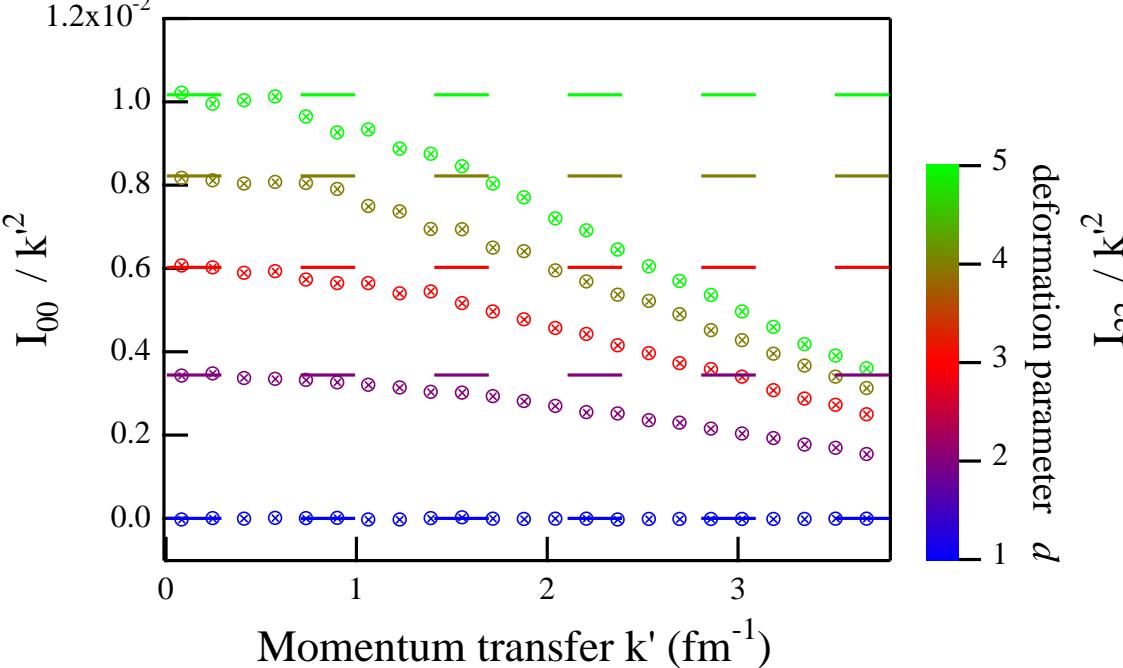
$$Q_{N=2,S}^{(l)}(k) \equiv \langle (2S)|j_l(kx)Y_{l0}(\hat{x})|(2S)\rangle$$

$$= \int d^3\rho d^3\lambda \Psi^{2S}(\vec{\rho}, \vec{\lambda}) j_l(k'\lambda) Y_{l0}(\hat{\lambda}) \Psi^{2S}(\vec{\rho}, \vec{\lambda})$$

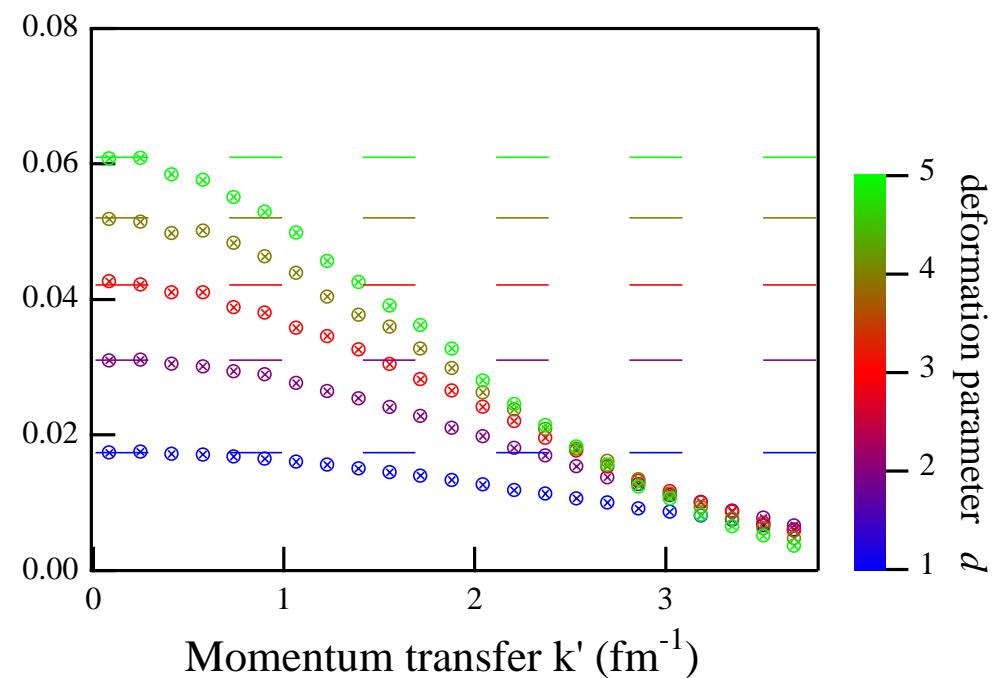
$$= \frac{1}{2} \int d^3\lambda \left[\left(\psi_{000}(\vec{\lambda}) j_l(k'\lambda) Y_{l0}(\hat{\lambda}) \psi_{000}(\vec{\lambda}) \right) + \left(\psi_{002}(\vec{\lambda}) j_l(k'\lambda) Y_{l0}(\hat{\lambda}) \psi_{002}(\vec{\lambda}) \right) \right]$$

$$\equiv \frac{1}{2} \left[I_{00}^{(l)} + I_{22}^{(l)} \right],$$

($l = 2$)



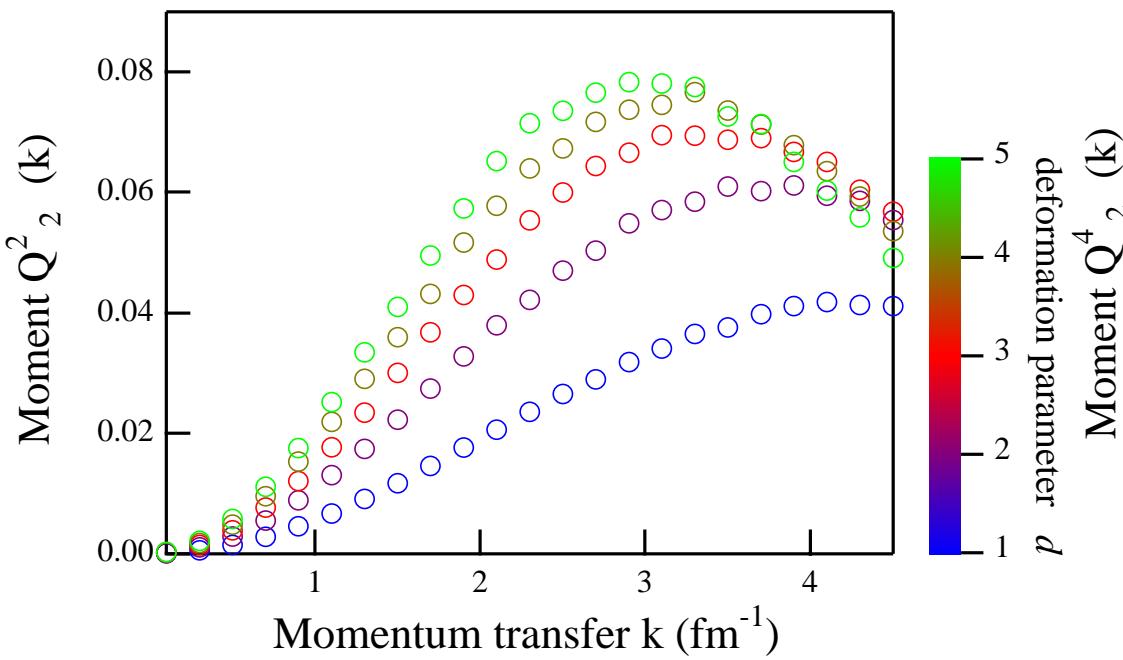
Symmetrized DOQ wave function
in the Jacobi coordinate



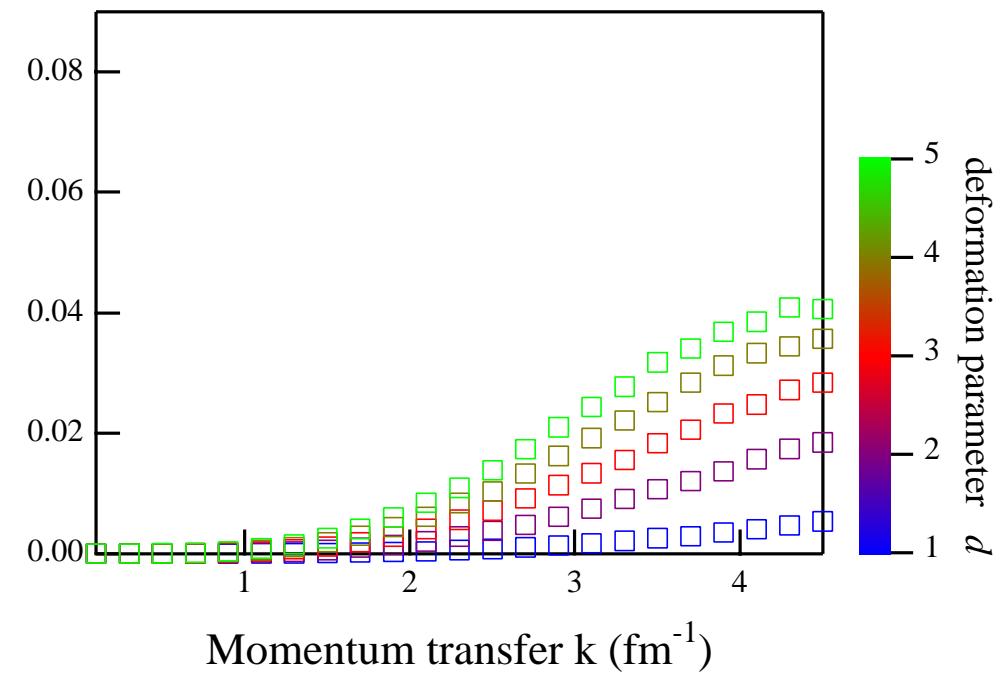
Dashed line -- long wave length limit

l -th moment

($l = 2$)



($l = 4$)

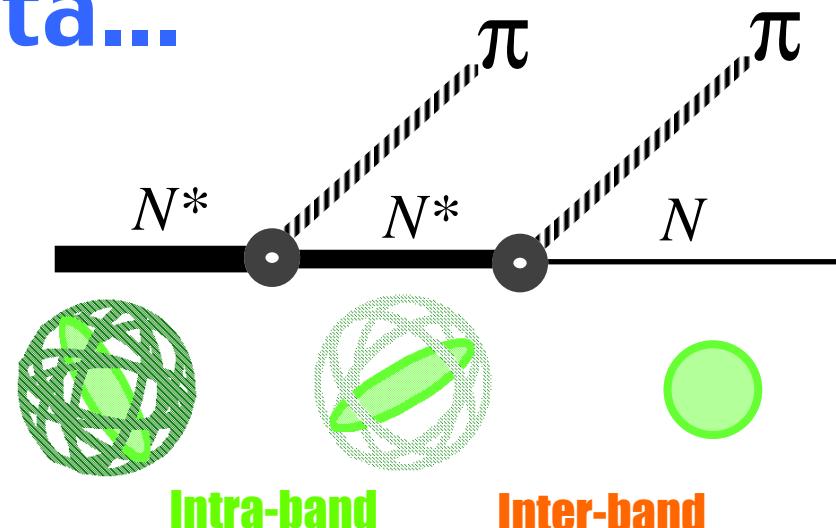


To compare with the data...

Experimentally, Intra-band transitions may be observed in the

$$N^* \rightarrow N\pi\pi$$

process



However, $L = 2 \rightarrow L = 0$ transitions, where some $N^* \rightarrow N\pi\pi$ processes are observed, are almost forbidden.

There is no analysis with the decay into $N\pi\pi$ for $L = 4$ baryons.
Moreover, $J = 7/2$ baryon is not identified yet.

(due to the large decay width from intra-band transitions?!)

$N(2220)$ DECAY MODES ($J = 2/9$)

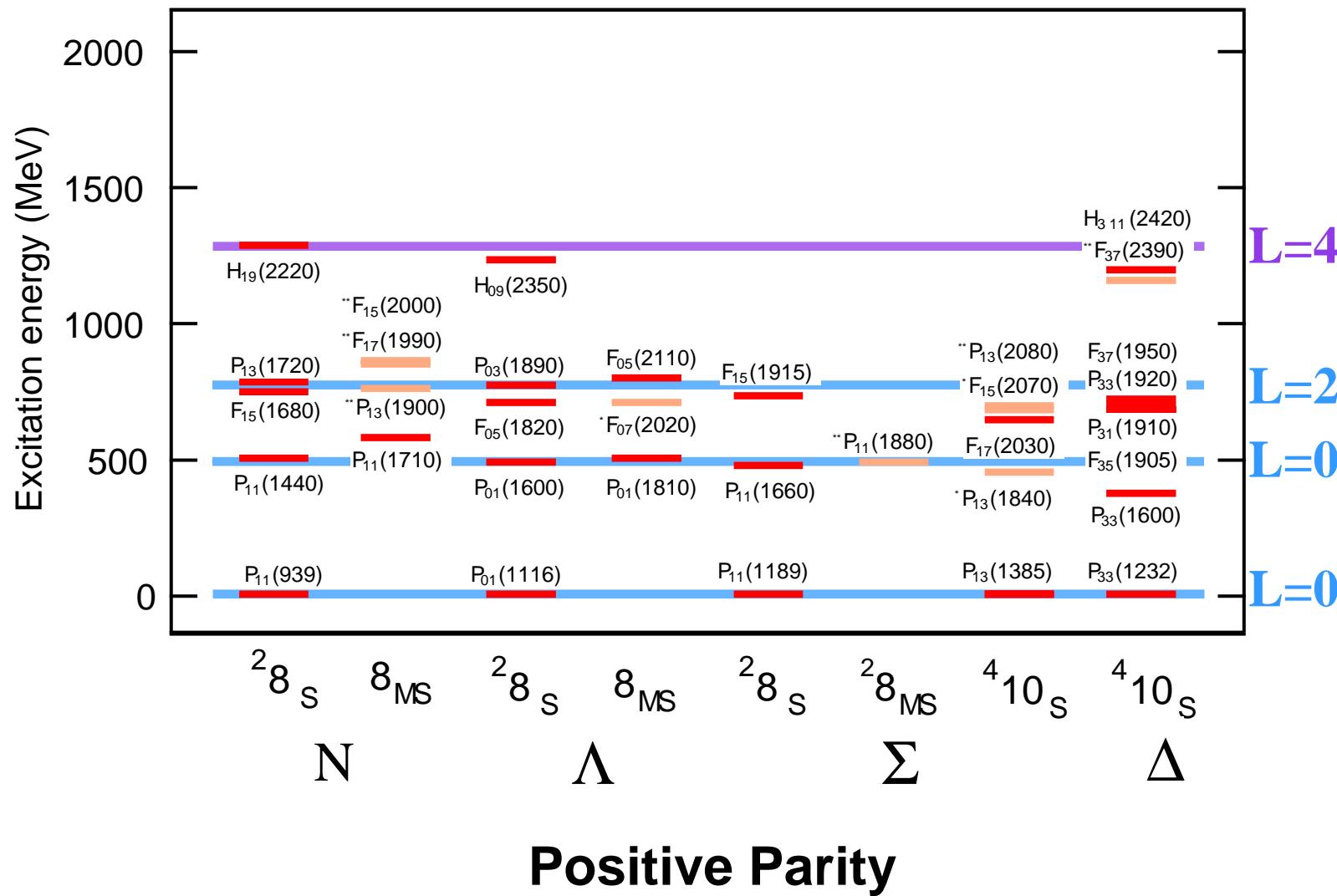
The following branching fractions are our estimates, not fits or averages.

| Mode | Fraction (Γ_i/Γ) |
|------------------------|--------------------------------|
| Γ_1 $N\pi$ | 10–20 % |
| Γ_2 $N\eta$ | |
| Γ_3 ΛK | |

- Comparison ---Data v.s. DOQ---

★All existing data with 3 or 4 stars — and some with 1 or 2 stars — are shown.

Takayama et al.,
PTP 101('99)



DOQ
Prediction
 $\omega = 644 \text{ MeV}$

Summary and Future Plan

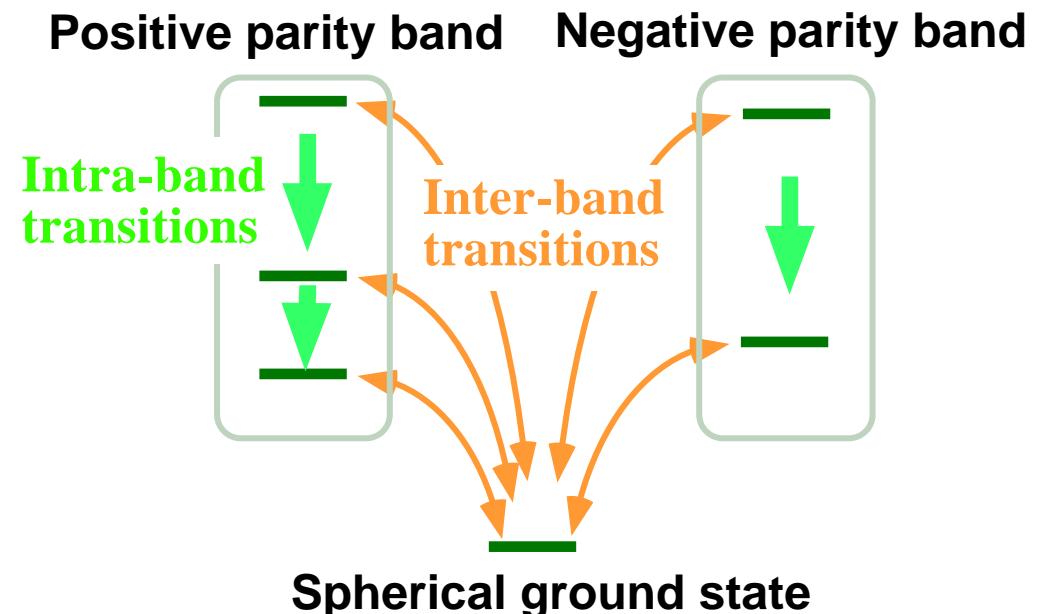
Mass spectrum of flavor SU(3) baryons

Deformation!

- Systematics for all flavor multiplets
- Rotational bands formed by deformed excited states

Is there any signs of the deformation in various transitions?

- Inter-band transitions
- Intra-band transitions



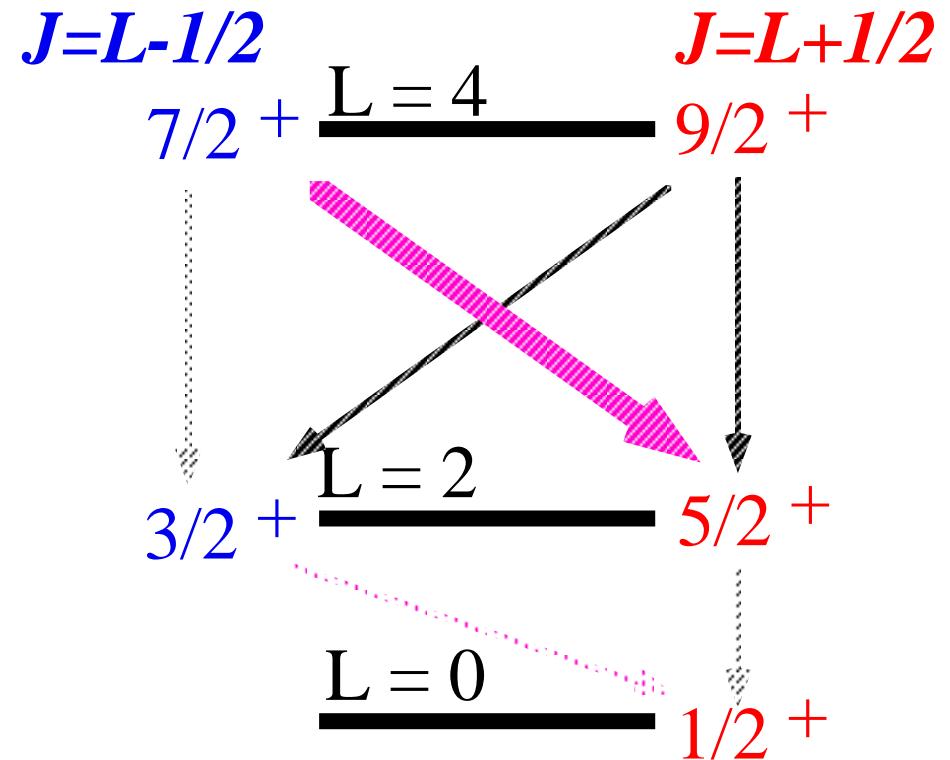
We have calculated the decay widths of intra-band transitions with pions in the DOQ model

Transitions from $J = 5/2, 3/2$
-- forbidden

Transitions from $J = 7/2$
-- Large decay width

Next, we will calculate

- Negative parity band
 - Inter-band transitions with pions
 - Decay into $\Delta(1232)\pi$
- = the decay width of $N^*\rightarrow N\pi\pi$ process
in the DOQ model



Angular momentum of π
 $l = 1, 3$
(Parity conservation)