

Spin-isospin responses by charge-exchange reactions and implications for astrophysics

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and Dark Matter in Nuclear Physics (NDM12)**

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Spin-Isospin Modes

1. **Importance of studying GT^+ in *fp*-shell nuclei**
2. **Experimental method**
3. **Case Study: ^{58}Ni**
4. **Measurements on several *fp*-shell nuclei**
5. **Measurements on 2β -decaying nuclei**
(Dieter Frekers' Talk)
6. **Conclusions and outlook**

Why are Gamow-Teller transitions in *fp*-shell nuclei important ?

- Role of *fp*-shell nuclei in supernova explosions: Core of supernova star is composed of *fp*-shell nuclei.
⇒ electron capture
- Neutrino absorption cross sections by *fp*-shell nuclei are essential in understanding of nuclear synthesis in Supernova explosions in cosmos.
- Difficulties in shell-model calculations for *fp*-shell nuclei.
- Importance to measure spin-isospin responses of *fp*-shell nuclei to gauge theoretical calculations.

Charge-exchange probes

(p,n) -type ($\Delta T_z = -1$)

- β^- -decay
- (p,n)
- $(^3\text{He},t)$
- heavy ion

(n,p) -type ($\Delta T_z = +1$)

- β^+ -decay
- (n,p)
- $(d,^2\text{He})$
- $(t,^3\text{He})$
- heavy ion; ($^7\text{Li}, ^7\text{Be}$)

- Energy per nucleon (>100 MeV/u)
- Spin-flip versus non-spin-flip
- Complexity of reaction mechanism
- Experimental considerations

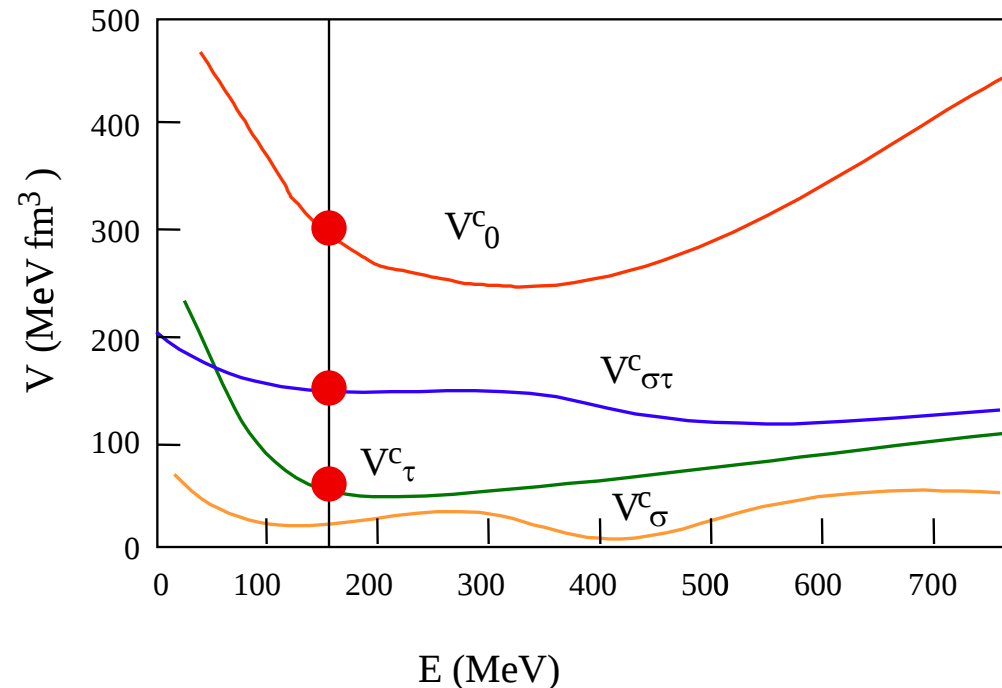
$(^3\text{He}, t)$ Reaction ≥ 100 MeV/u

- Energy dependence of effective interactions.

- At RCNP, Osaka

$E(^3\text{He}) \approx 150$ MeV/u

- V_0 part: Minimum.
- $V_{\sigma t}$ part: Relatively large.
- V_t part: Minimum.



The ($^3\text{He}, t$) reaction at 0 degree

- Cross sections at $E(^3\text{He})=450$ MeV, $q=0$ for ($^3\text{He}, t$) reactions

$$\frac{d\sigma}{d\Omega} = \frac{\mu_i \mu_f}{(\pi \hbar^2)^2} \left(\frac{k_f}{k_i} \right) (N_\tau^D | J_\tau |^2 B(F) + N_{\sigma\tau}^D | J_{\sigma\tau} |^2 B(GT))$$

T. N. Taddeucci *et al.*, Nucl. Phys. A469, 125 (1987)

I. Bergqvist *et al.*, Nucl. Phys. A469, 648 (1987)

- Neutrino absorption cross sections

$$\sigma = \frac{1}{\pi \hbar^4 c^3} \left[G_V^2 B(F) + G_A^2 B(GT) \right] \times F(Z, E_e) p_e E_e$$

$F(Z, E_e)$ is the relativistic Coulomb barrier factor

Importance of charge-exchange reactions at intermediate energies

Measuring GT strengths

$$\frac{d\sigma}{d\Omega}(q=0) = KN_D |J_{\sigma\tau}|^2 B(GT)$$

Diagram illustrating the components of the cross-section formula:

- $\frac{d\sigma}{d\Omega}$: kinematic factor
- K : distortion factor
- N_D : nucleon-nucleus interaction
- $|J_{\sigma\tau}|^2$: Gamow-Teller strength
- $B(GT)$: Gamow-Teller strength

**Calibration of $B(GT)$ to cross section for known transitions
(e.g. from β -decay)**

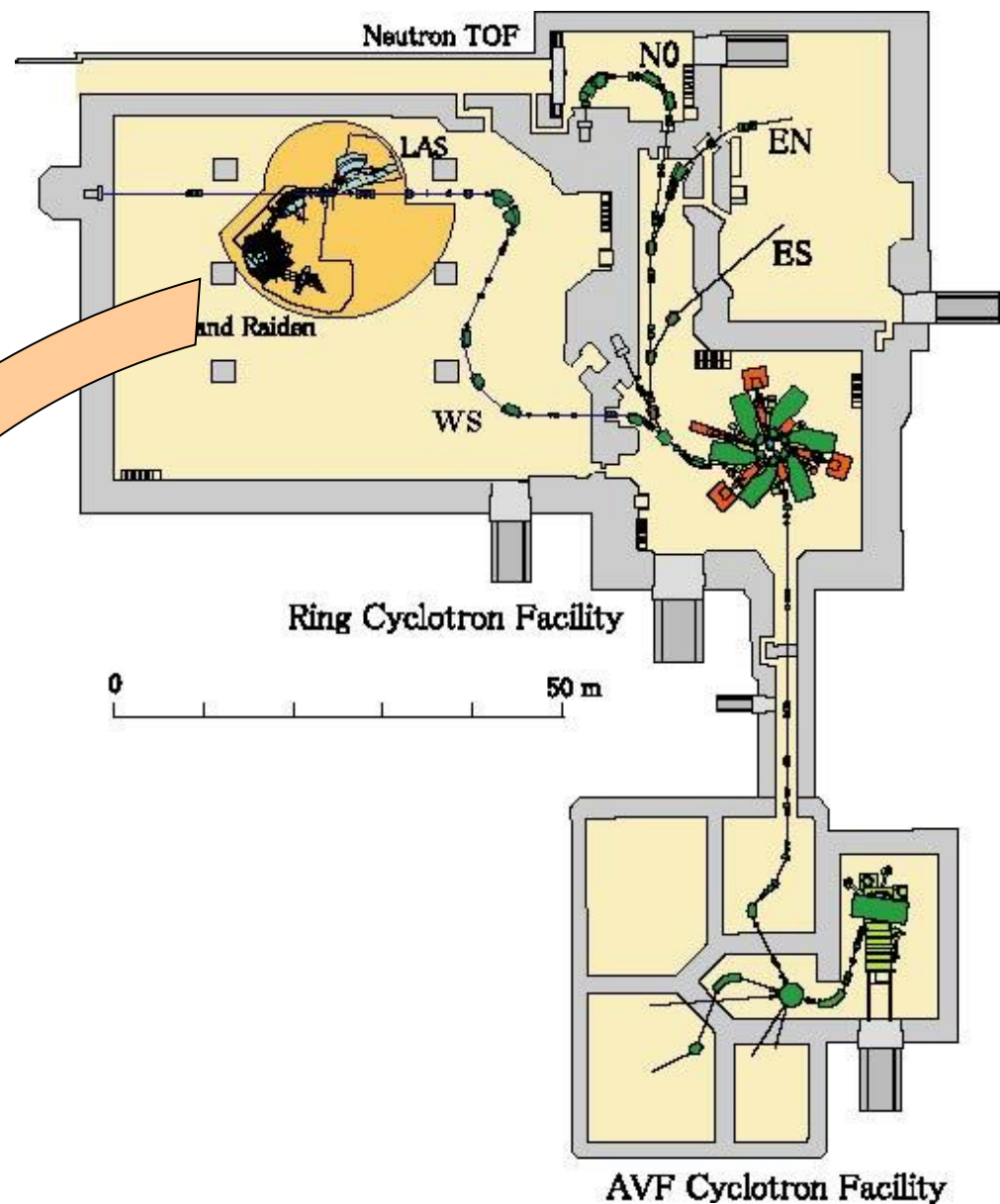
Experiments

- RCNP facility

 - K=400 MeV ring cyclotron

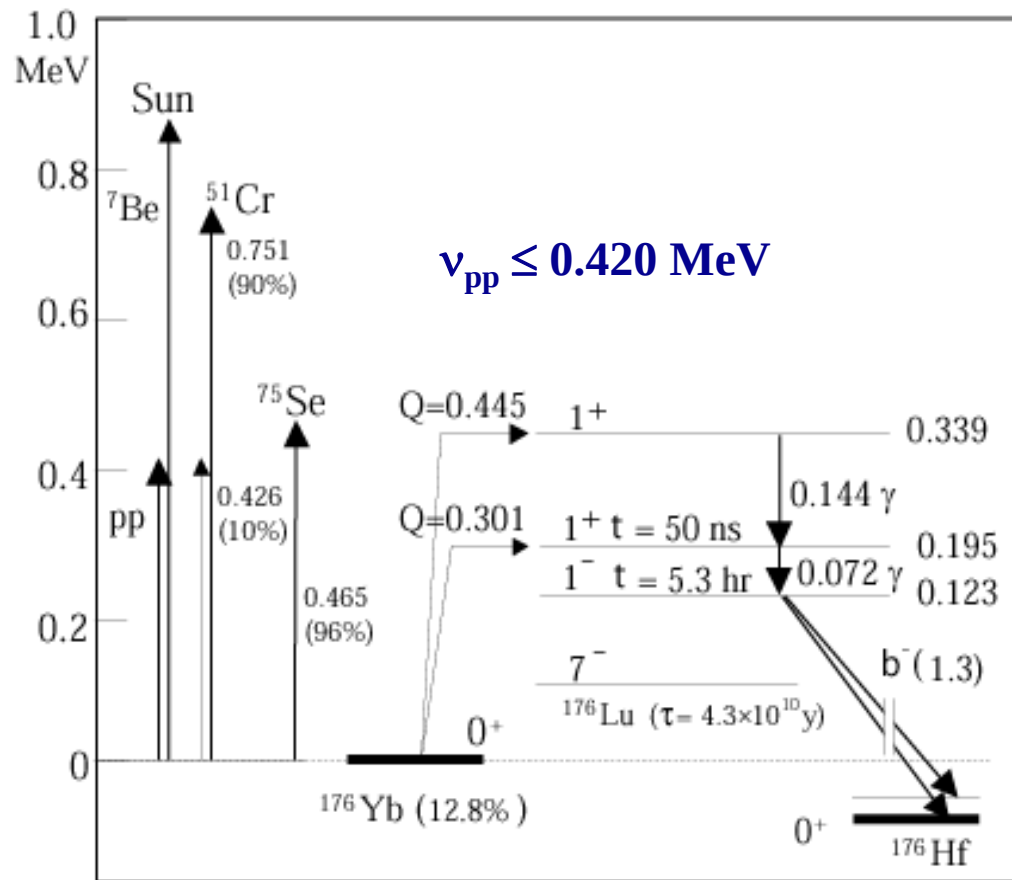
 - Grand Raiden spectrometer

- Beam: $^3\text{He}^{++}$, 450 MeV

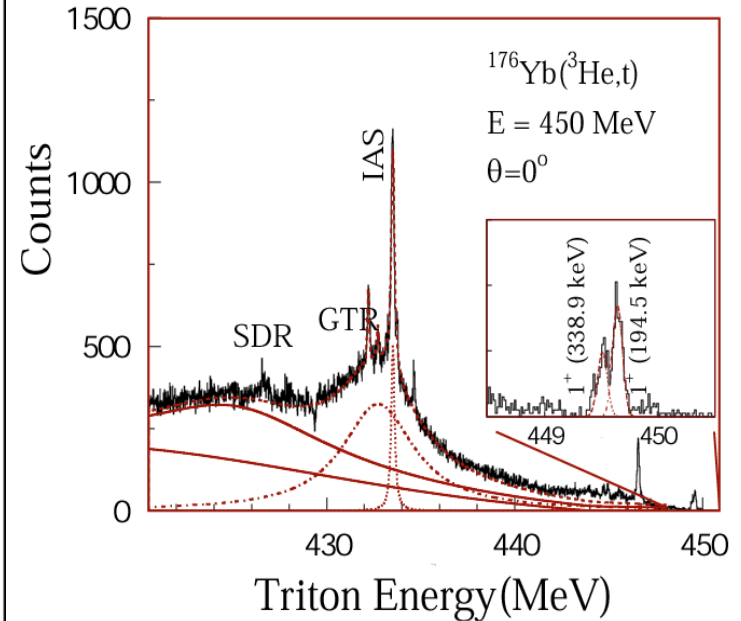


M. Fujiwara *et al.*, NIM A422 (1999) 484

Used $^{164}\text{Dy}(^3\text{He}, t)^{164}\text{Ho}$ (g.s., 1+) reaction for calibration: $\log ft$ 4.6 $\rightarrow$ $B(\text{GT}) = 0.293 \pm 0.006$



M. Fujiwara *et al.*, PRL 85 (2000) 4442



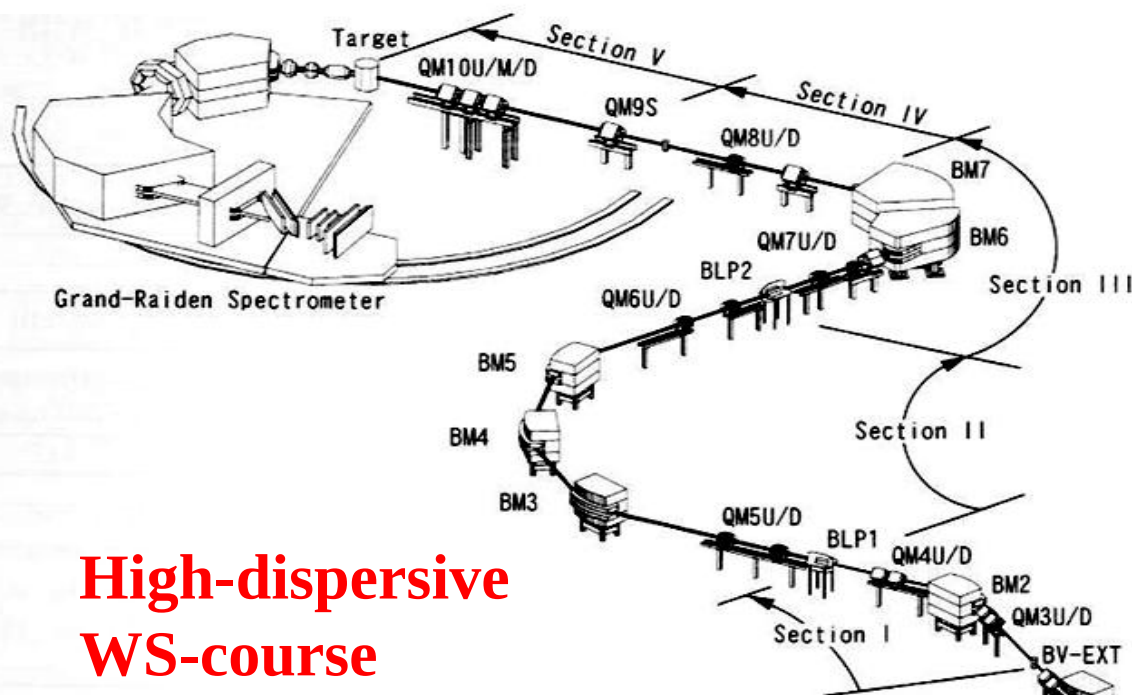
Resolution ≈ 100 to 130 keV

E_x (MeV)	0.195+0.339 (p, n)	0.195 ($^3\text{He}, t$)	0.339 ($^3\text{He}, t$)
$B(\text{GT})$	0.32 ± 0.04	0.20 ± 0.04	0.11 ± 0.02

Beam line WS-course

Grand-Raiden Spectrometer

M. Fujiwara *et al.*, NIM A422 (1999) 484



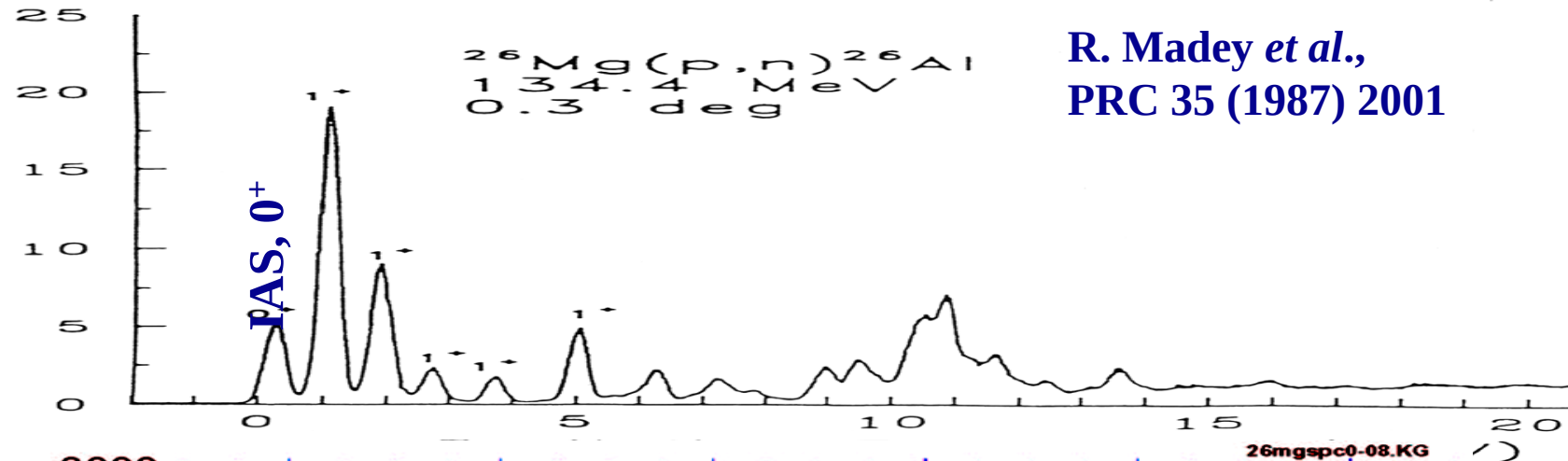
**High-dispersive
WS-course**

T. Wakasa *et al.*, NIM A482 (2002) 79

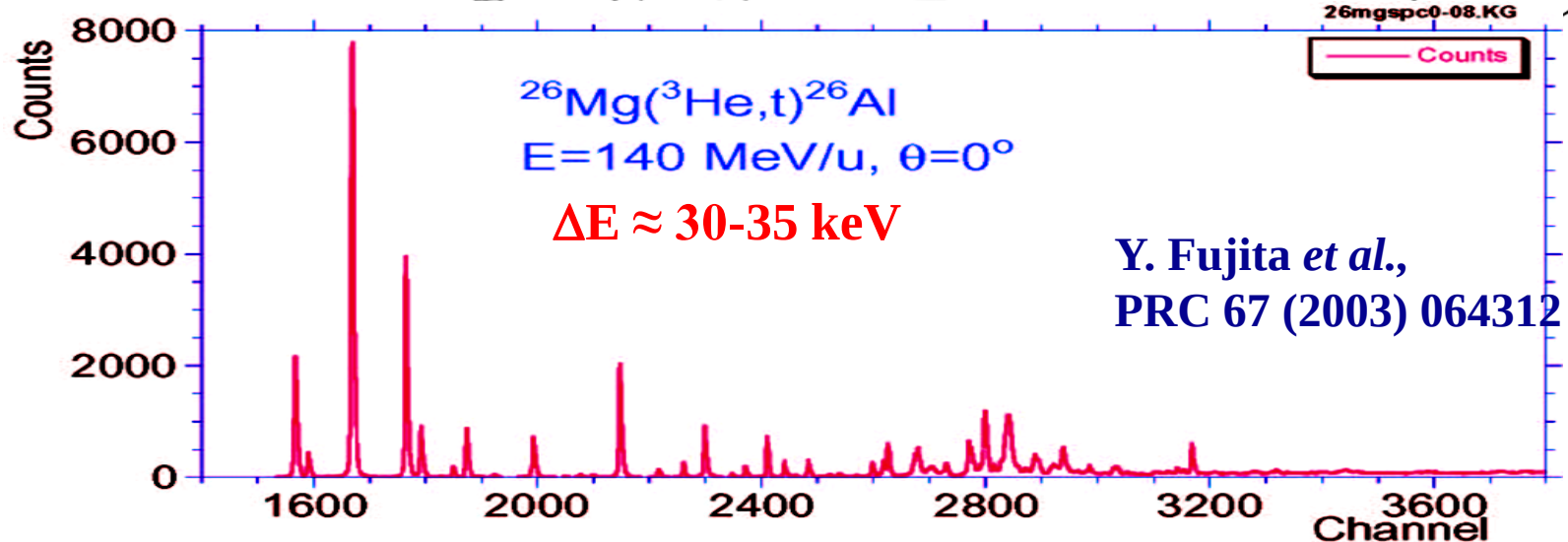
**RCNP Ring
Cyclotron**

Ring Cyclotron

$^{26}\text{Mg}(p,n)^{26}\text{Al}$ & $^{26}\text{Mg}(^3\text{He},t)^{26}\text{Al}$ spectra



R. Madey *et al.*,
PRC 35 (1987) 2001



Y. Fujita *et al.*,
PRC 67 (2003) 064312

Prominent states are GT states and the IAS !

$^{136}\text{Xe}(^3\text{He},t)^{136}\text{Cs}$

$E(^3\text{He}) = 420 \text{ MeV}$

$\Delta E = 42 \text{ keV}$

$B_{\text{exp}}(\text{GT}+) =$

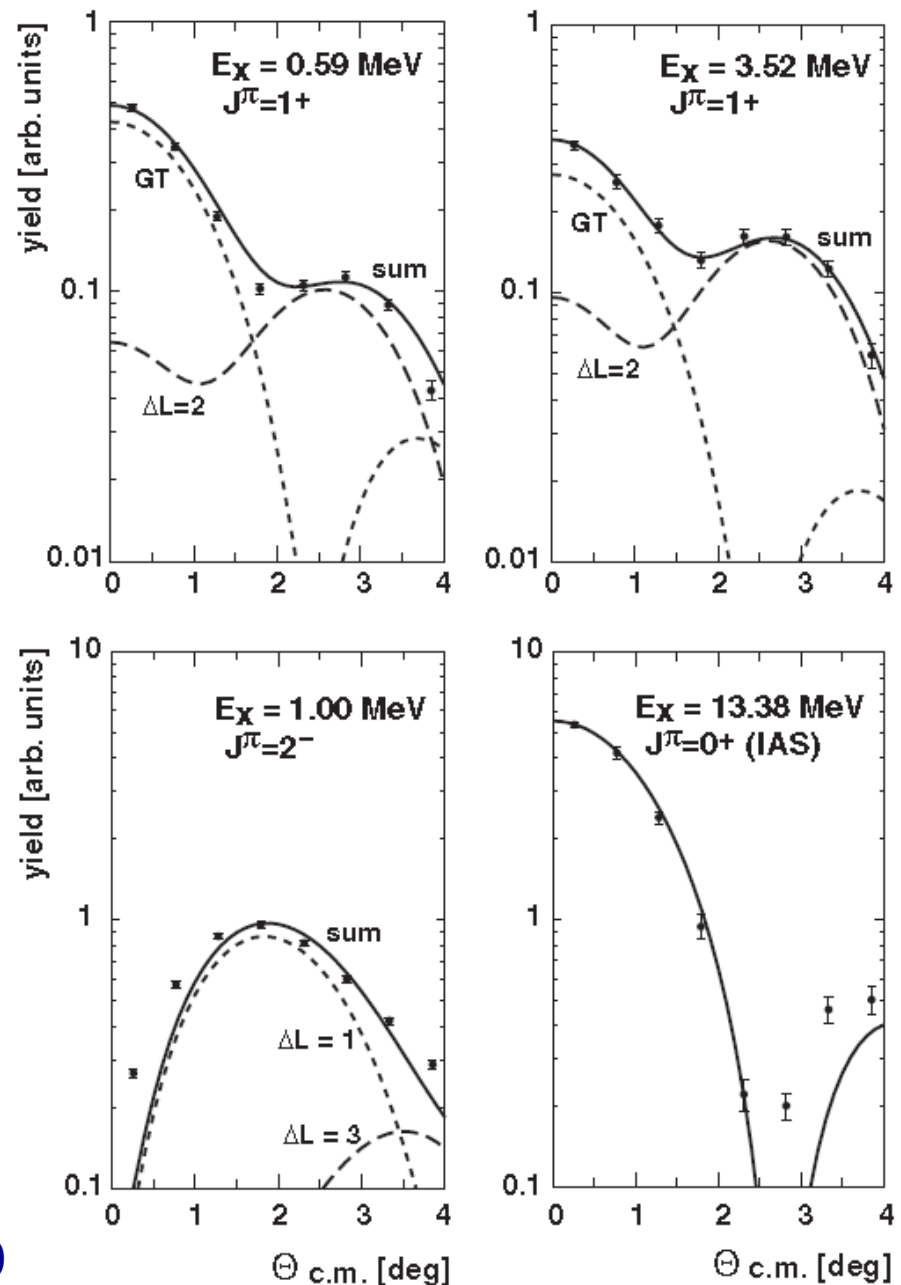
$$\frac{d\sigma(q=0)}{d\Omega} \cdot \left[\frac{d\hat{\sigma}(\text{GT})}{d\Omega} \right]^{-1}$$

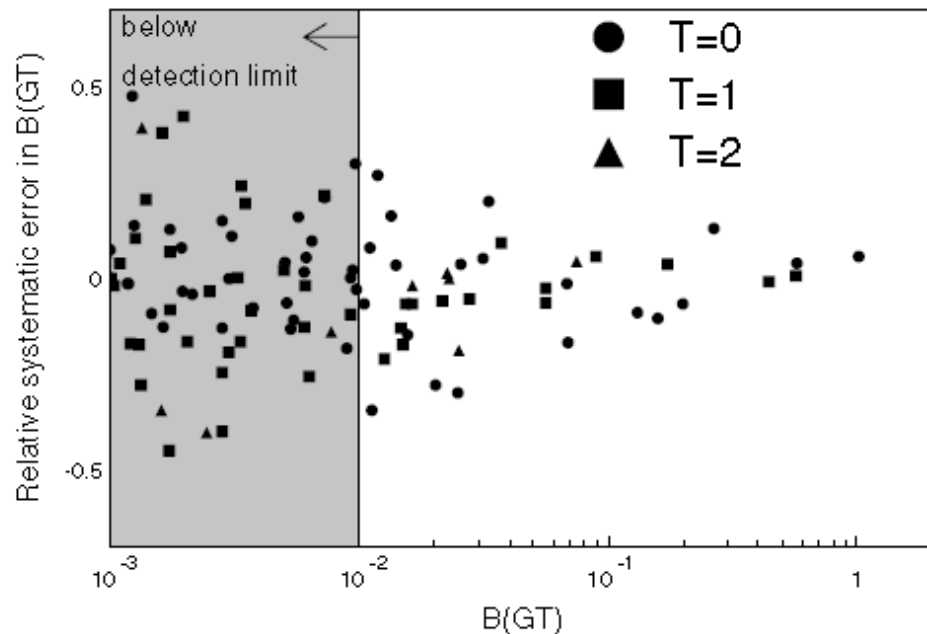
extrapolated
(DWBA)

unit cross section

$\Delta L = 2$ & $\Delta L = 0$ incoherent

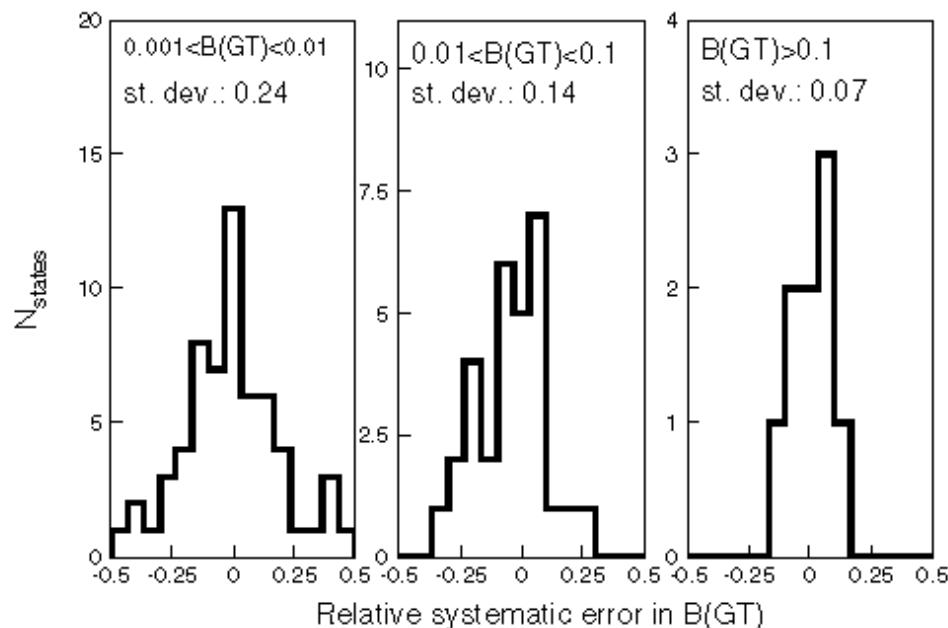
P. Puppe *et al.*, PRC 84 (2011) 051305(R)





Theoretical Study $^{26}\text{Mg}(^3\text{He}, t)^{26}\text{Al}$

Effects of $\Delta L = 2$, $\Delta S = 1$ contributions mediated via the T_τ interaction that interfere with $\Delta L = 0$, $\Delta S = 1$ contributions to Gamow-Teller transitions.



$$\text{Rel. syst. error} = \frac{B(GT)_{\text{DWBA}} - B(GT)_{\text{SM}}}{B(GT)_{\text{SM}}}$$

R.G.T. Zegers *et al.*, PRC74 (2006) 024309

Determination of GT^+ Strength and its Astrophysical Implications

In supernova explosions, **electron capture (EC)** on *fp*-shell nuclei plays a dominant role during the last few days of a heavy star with $M > 10 M_{\odot}$

Presupernova stage; deleptonization $\Rightarrow$ core collapse $\Rightarrow$ subsequent type IIa Supernova (SN) explosion

H.A. Bethe *et al.*, Nucl. Phys. A324 (1979) 487

Electron capture in *fp*-shell

- The rate for **EC** is governed by the **GT⁺ strength distribution at low excitation energy**; not accessible to β -decay.
- **Fuller, Fowler and Newman (FFN) (1982-1985)**; estimates of stellar rates in stellar environments using s.p. model.
- **Caurier *et al.*, Martínez-Pinedo & Langanke (1999), Otsuka *et al.* $\Rightarrow$ Large shell-model calculations $\Rightarrow$ marked deviations from FFN EC rate; generally smaller EC rates.**
- Experiments and theory relied on (n,p) data (**TRIUMF**) which have a **rather poor energy resolution**.

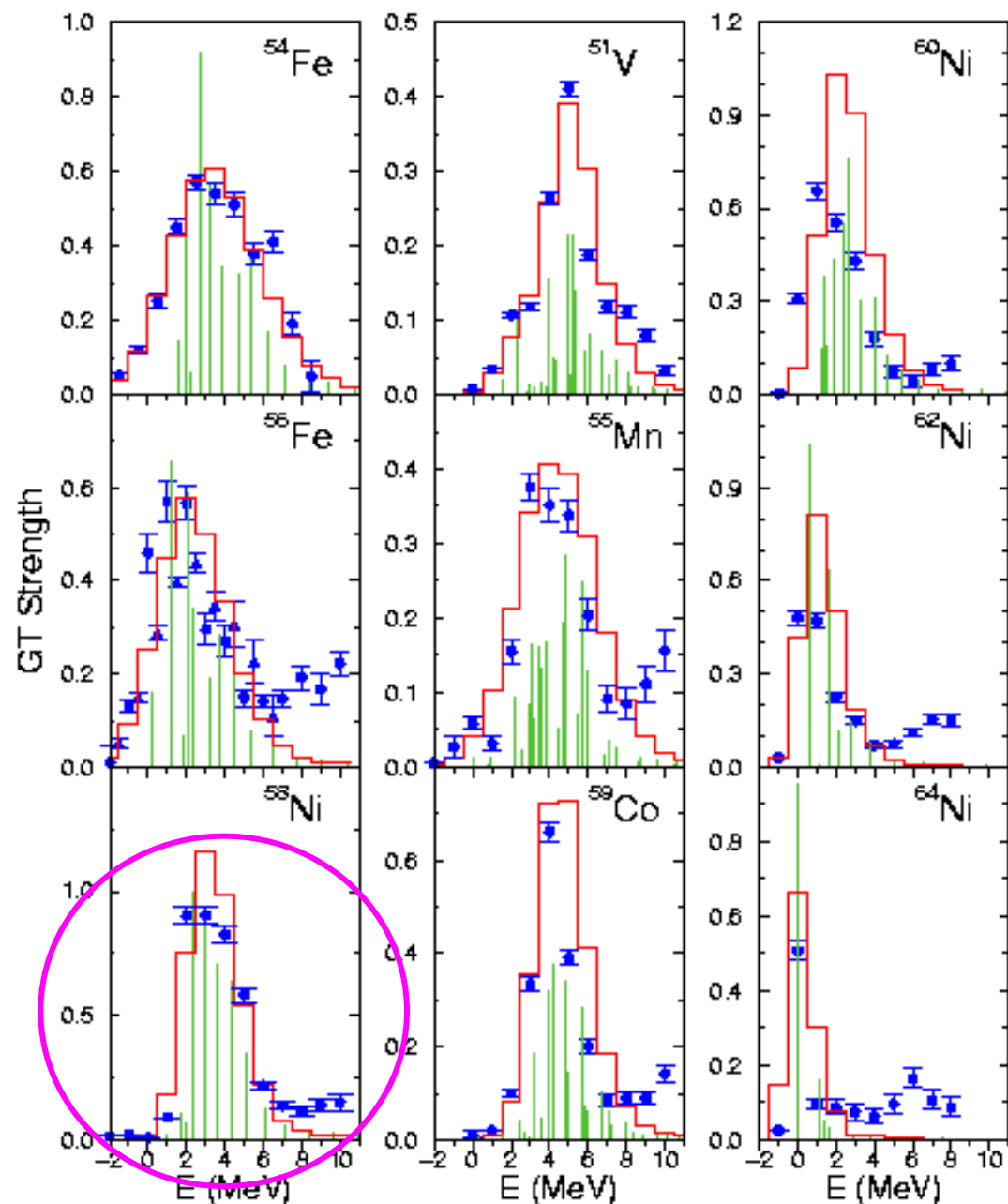
fp-shell nuclei: large scale shell model calculations

E. Caurier *et al.*

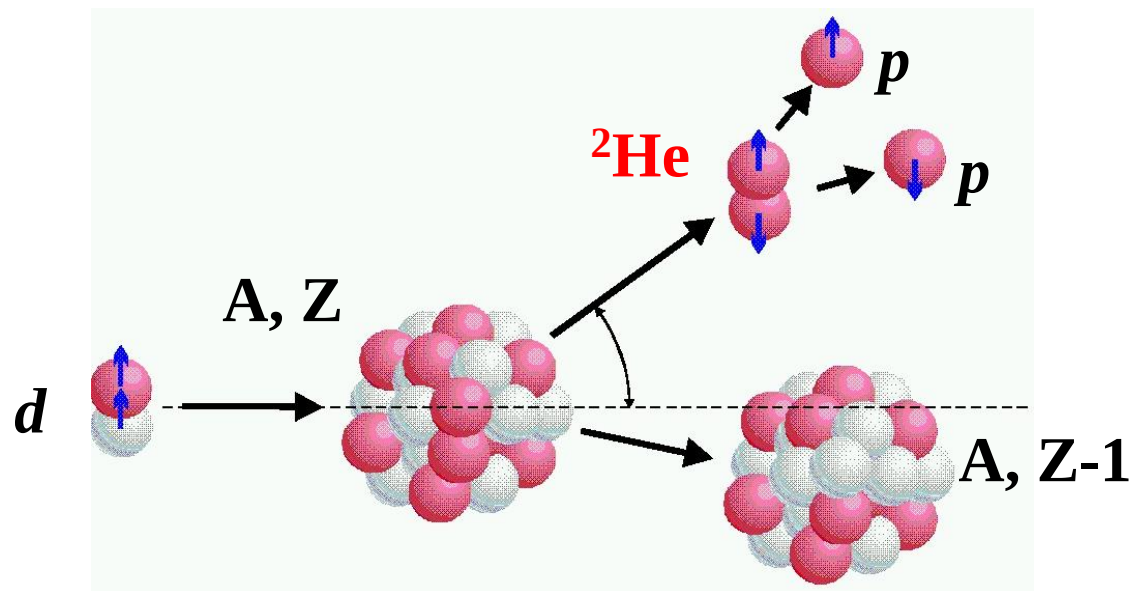
NPA 653 (1999) 439

- Stellar weak reaction rates with improved reliability
- Large scale shell model (SM) calculations
- Tuned to reproduce GT^+ strength measured in (n,p)
- (n,p) data from TRIUMF
- GT^+ strength from SM
- Folded with 1 MeV energy resolution

Case study: ^{58}Ni



Exclusive excitations $\Delta S=\Delta T=1$: ($d, {}^2\text{He}$)

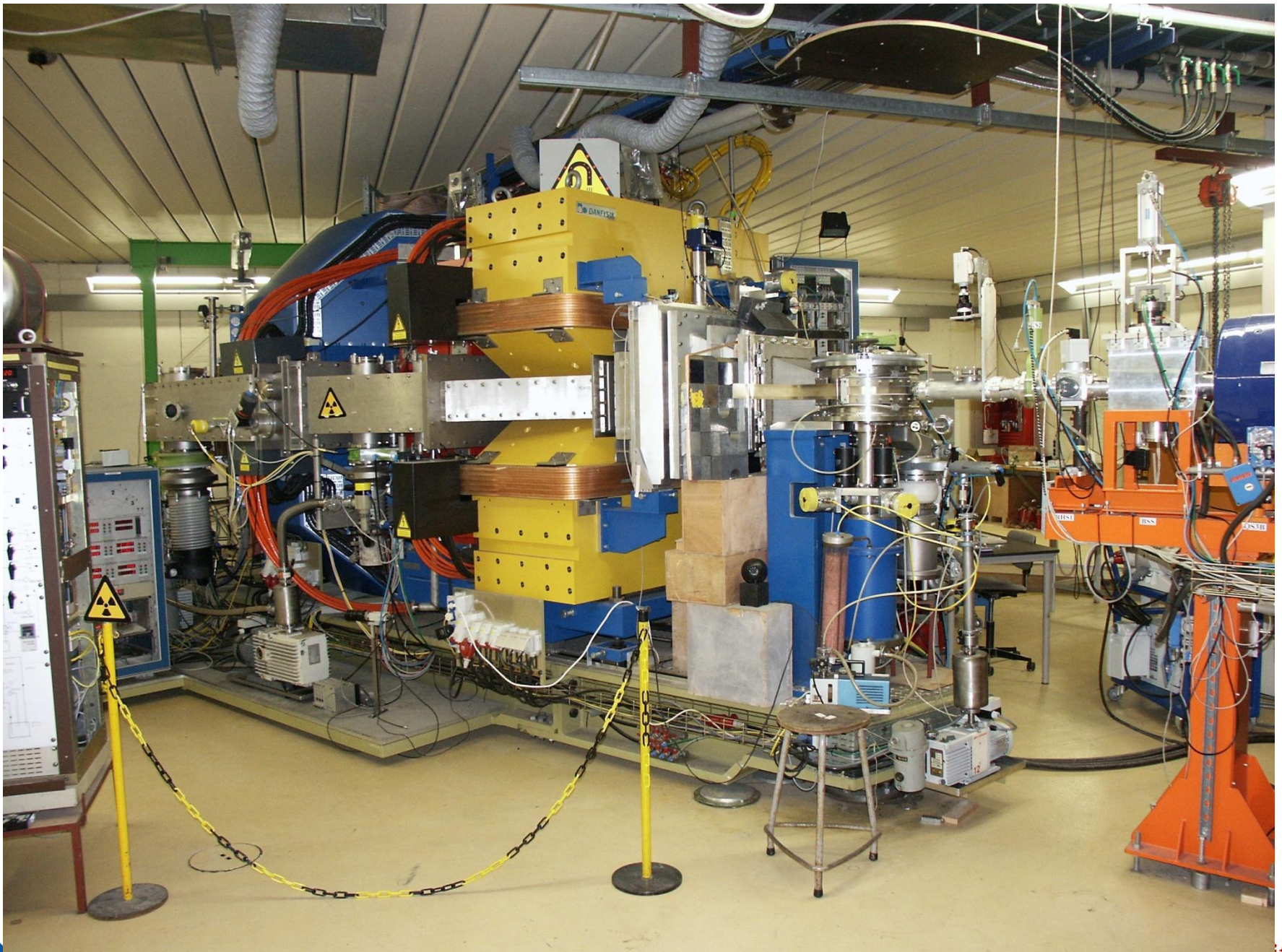


${}^3\text{S}_1$ deuteron $\Rightarrow$ ${}^1\text{S}_0$ di-proton (${}^2\text{He}$)

${}^1\text{S}_0$ dominates if (relative) 2-proton kinetic energy $\varepsilon < 1$ MeV

(n,p) -type probe with exclusive $\Delta S=1$ character (GT^+ transitions)

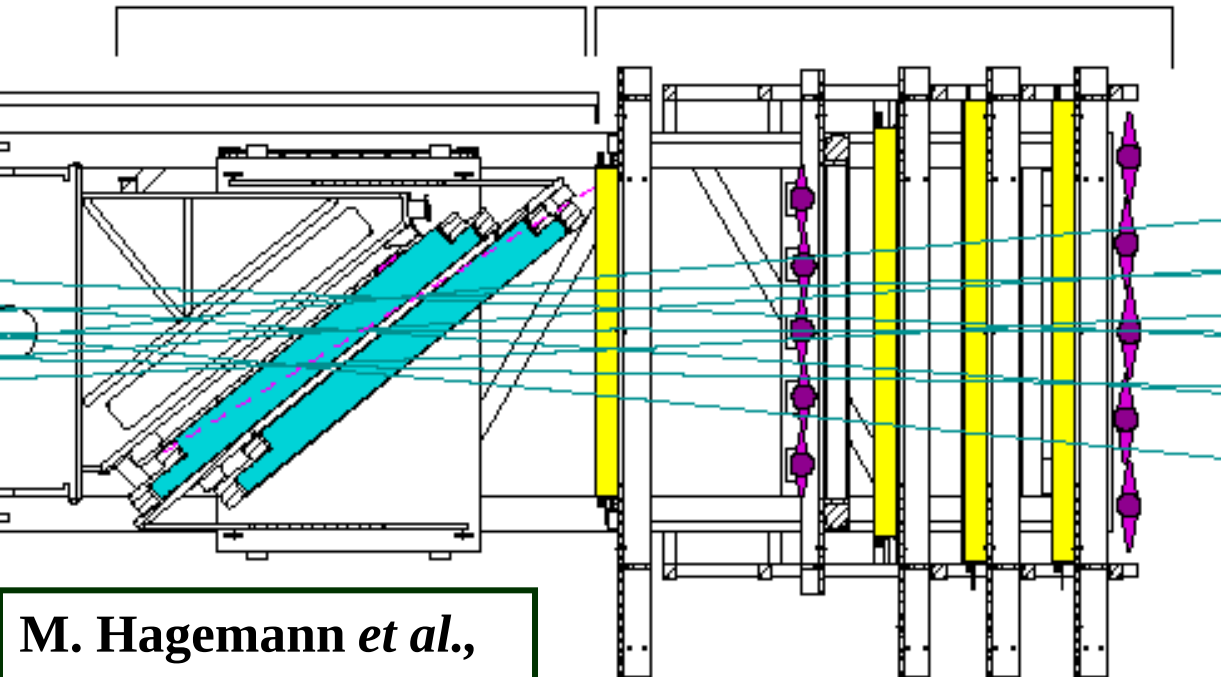
But near 0° : tremendous background from d -breakup



Setup: ESN detector

FPDS

FPP



Focal-Plane Detector:
(FPDS): 2 VDCs

Focal-Plane Polarimeter:
**(FPP): 4 MWPCs &
graphite analyzer**

features a.o.:

fast readout

VDC readout pipeline

TDC's

**VDC decoding using
imaging techniques**

DSP based online analysis

**M. Hagemann *et al.*,
NIM A437 (1999) 459
V.M. Hannen *et al.*,
NIM A500 (2003) 68**

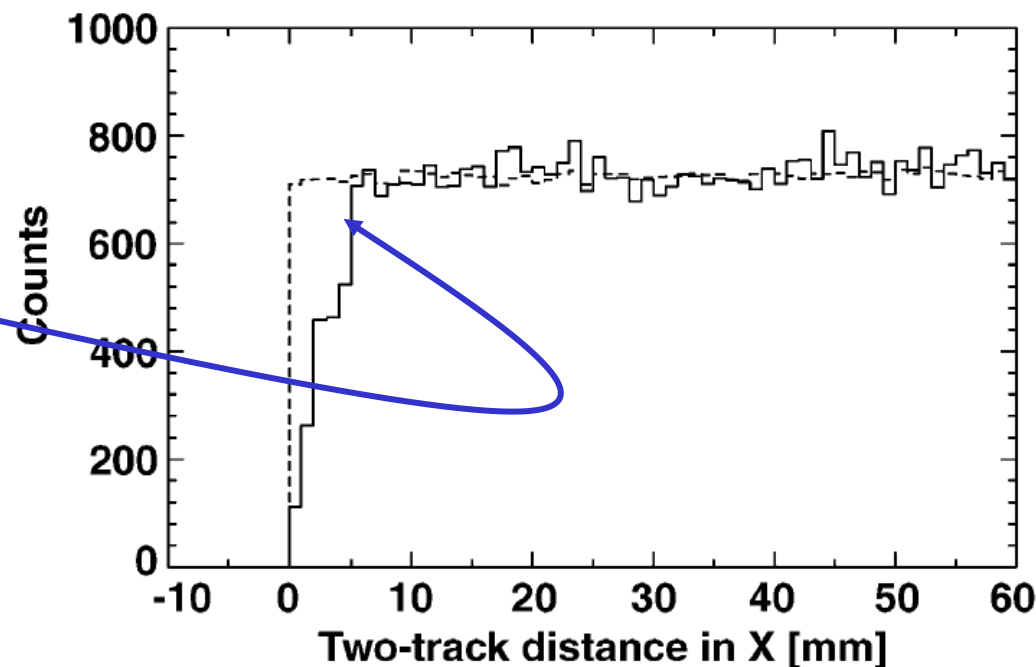
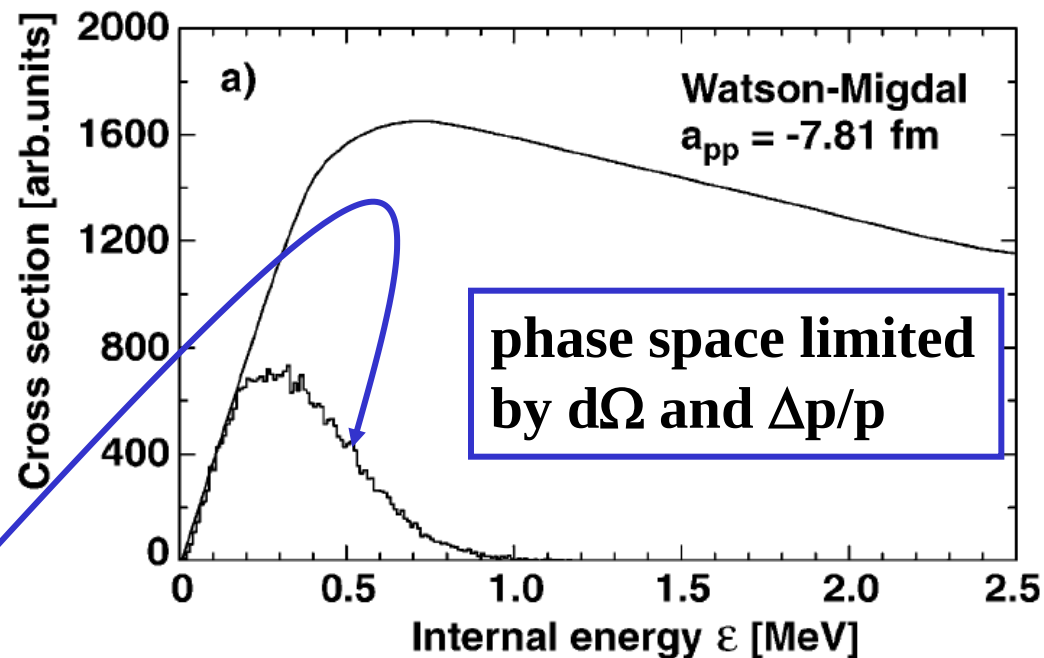
ESN

Bari, Darmstadt, Gent, Iserlohn, KVI, Milano, Münster, TRIUMF

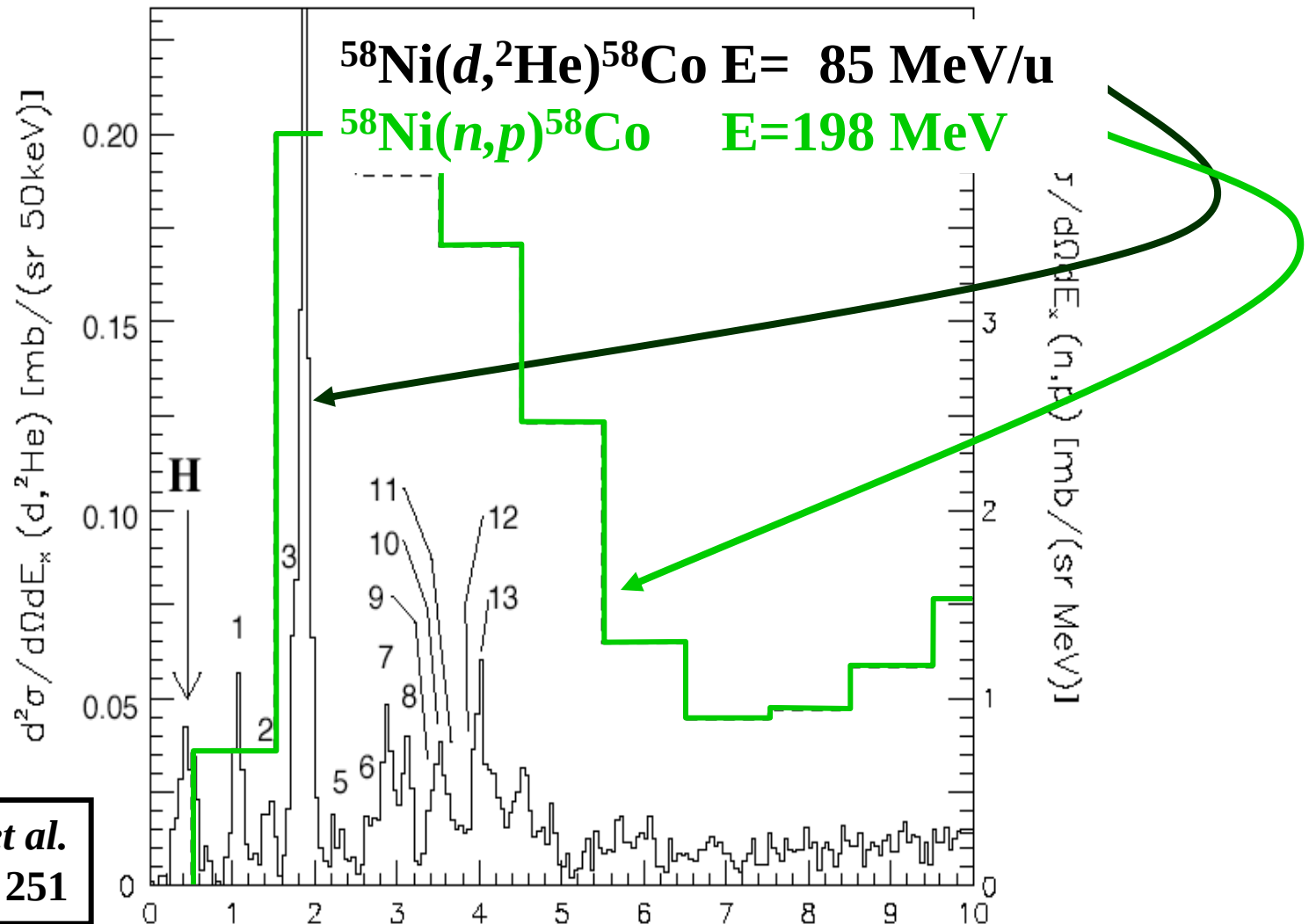
- Good double tracking
- **Use VDC information**
- Good phase-space coverage for small relative proton energies

S. Rakers *et al.*,
NIM A481 (2002) 253

measured

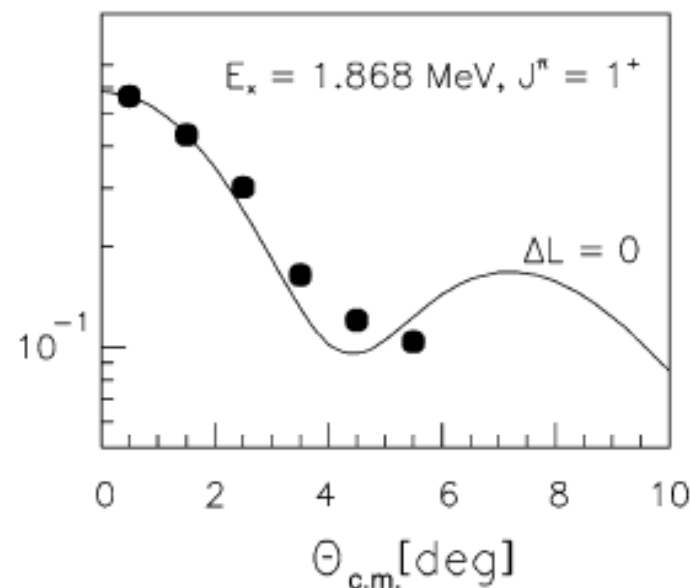
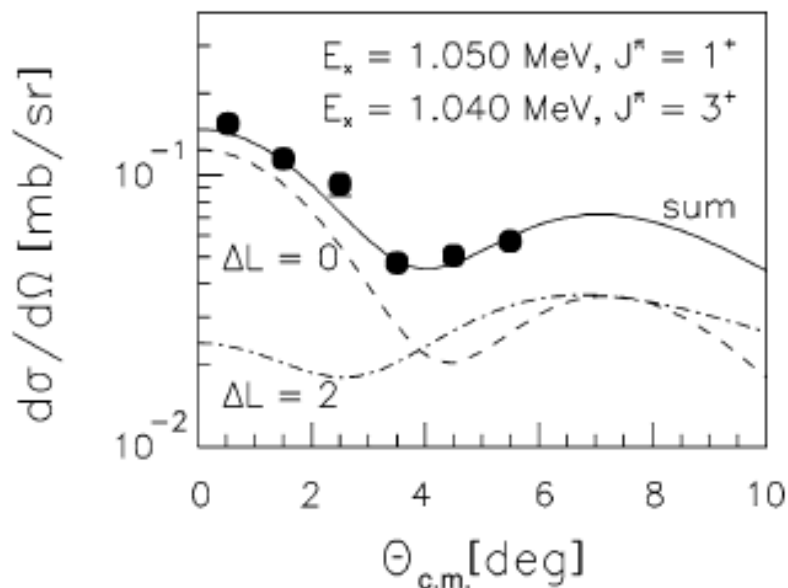


$(d, {}^2\text{He})$ as GT^+ probe in fp -shell nuclei



M. Hagemann *et al.*
 PLB 579 (2004) 251

$^{58}\text{Ni}(d, ^2\text{He})^{58}\text{Co}$ $E=85\text{A MeV}$



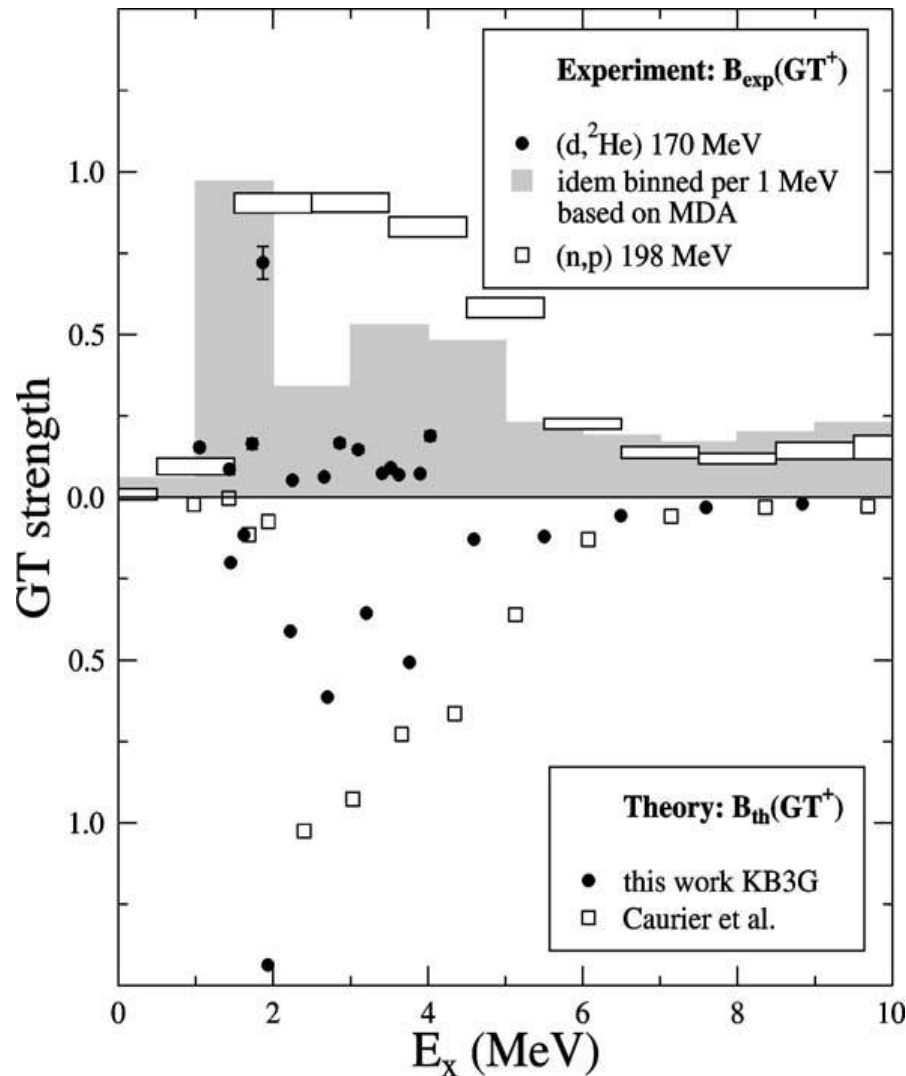
$$\mathbf{B}_{\text{exp}}(\text{GT}+) =$$

$$\frac{d\sigma(q=0)}{d\Omega} \cdot \left[\frac{d\hat{\sigma}(\text{GT})}{d\Omega} \right]^{-1}$$

GT Strength in ^{58}Co from $(d, ^2\text{He})$ reaction

E_x	$d\sigma/d\sigma(0.5^\circ)$	$\sigma(L=0)/\sigma(\tau\circ\tau)$	$B(\text{GT}+)$
[MeV]	[mb/sr]		
1.050	0.159 ± 0.009	0.88	0.15 ± 0.01
1.435	0.078 ± 0.006	1.00	0.09 ± 0.01
1.729	0.148 ± 0.014	1.00	0.16 ± 0.02
1.868	0.648 ± 0.020	1.00	0.72 ± 0.05
2.249	0.047 ± 0.004	1.00	0.05 ± 0.01
2.660	0.057 ± 0.005	0.96	0.06 ± 0.01
2.860	0.145 ± 0.009	0.99	0.17 ± 0.01
3.100	0.126 ± 0.008	0.99	0.15 ± 0.01
3.410	0.065 ± 0.007	0.96	0.07 ± 0.01
3.520	0.080 ± 0.009	0.95	0.09 ± 0.01
3.625	0.067 ± 0.007	0.87	0.07 ± 0.01
3.900	0.062 ± 0.006	0.97	0.07 ± 0.01
4.030	0.155 ± 0.010	1.00	0.19 ± 0.01
4.05-5.00	0.381 ± 0.061		0.49 ± 0.09

GT⁺ strength: comparison (n,p), (d,²He) & theory



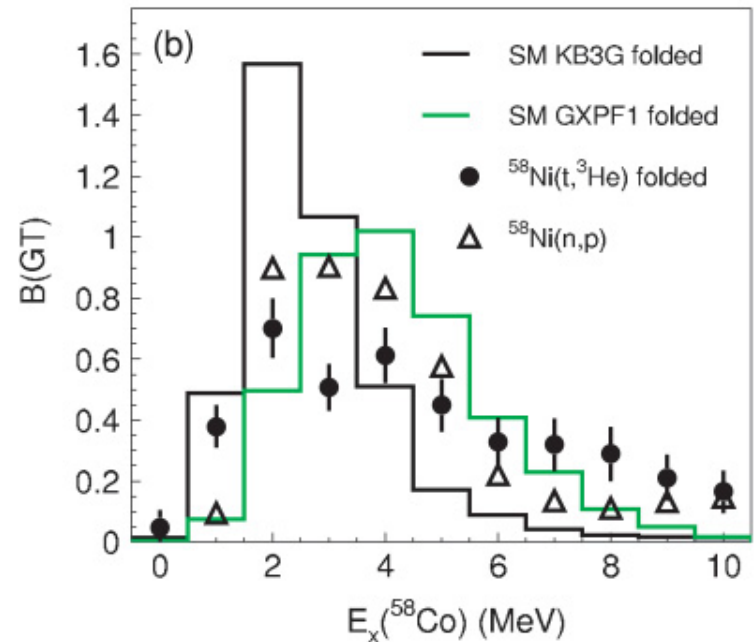
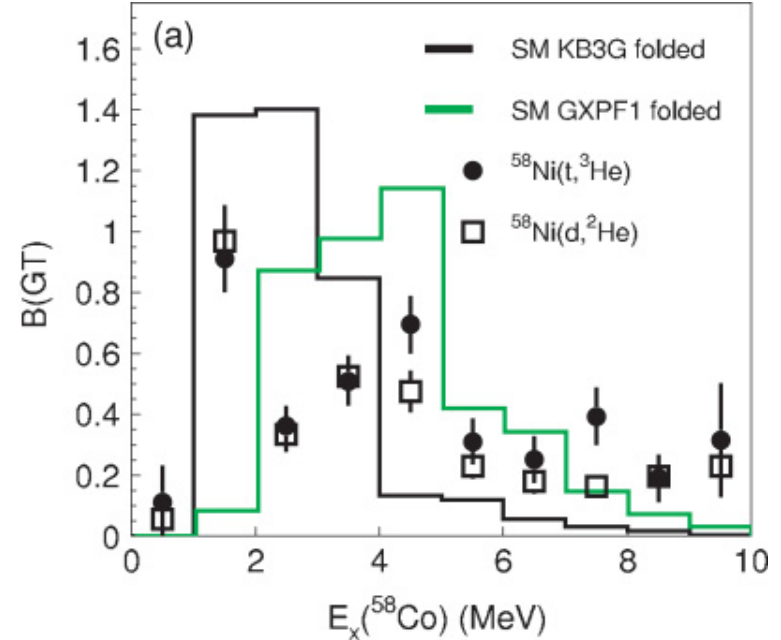
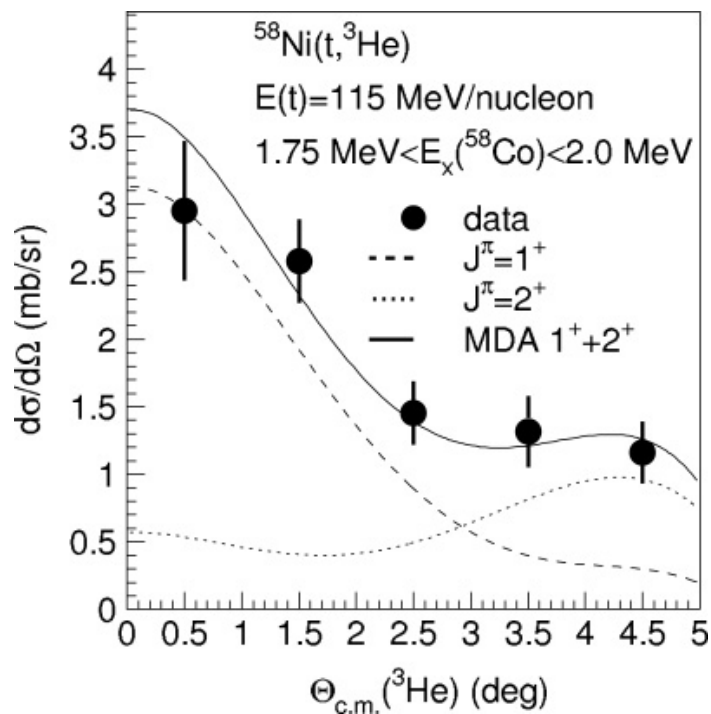
Up to 4 MeV excitation:

**13 GT transitions measured
(d,²He)**

Strength re-binned in 1 MeV bins

Significant differences

**Updated shell model calculations
by Martínez-Pinedo/Langanke
using KB3G interaction**



$^{58}\text{Ni}(t, ^3\text{He})^{56}\text{Co}$

$E_t = 115 \text{ MeV/u}$

Resolution = 250 keV

A.L. Cole *et al.*, PRC74 (2006) 034333

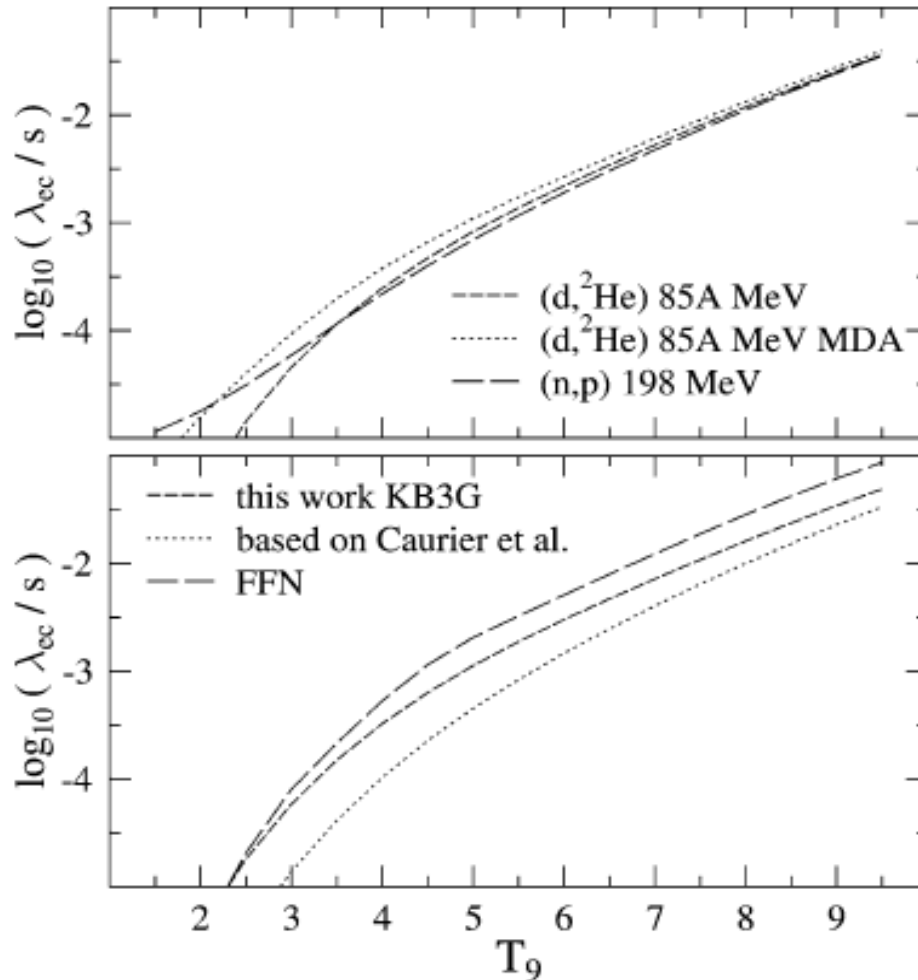
Electron capture rate

$$\lambda_{ec} \approx \sum_i B_i(GT) \int_{\omega_l}^{\infty} \omega p \left(Q_i + \omega \right)^2 F(Z, \omega) S_e(\omega, T) d\omega$$

With

- $B_i(GT)$ Gamow-Teller strength distribution
- ω and p energy and momentum of electrons
- $F(Z, \omega)$ is the relativistic Coulomb barrier factor
- $S_e(\omega, T)$ Fermi-Dirac distribution electron gas at temperature T

e^- -capture rates using experimental strengths (Martínez-Pinedo, Langanke)



Evolution of core of
25 $M_{\odot}$ star. Conditions
following silicon depletion.

$$T_9 = 4.05$$

$$\rho = 3.18 \times 10^7 \text{ g/cm}^3$$

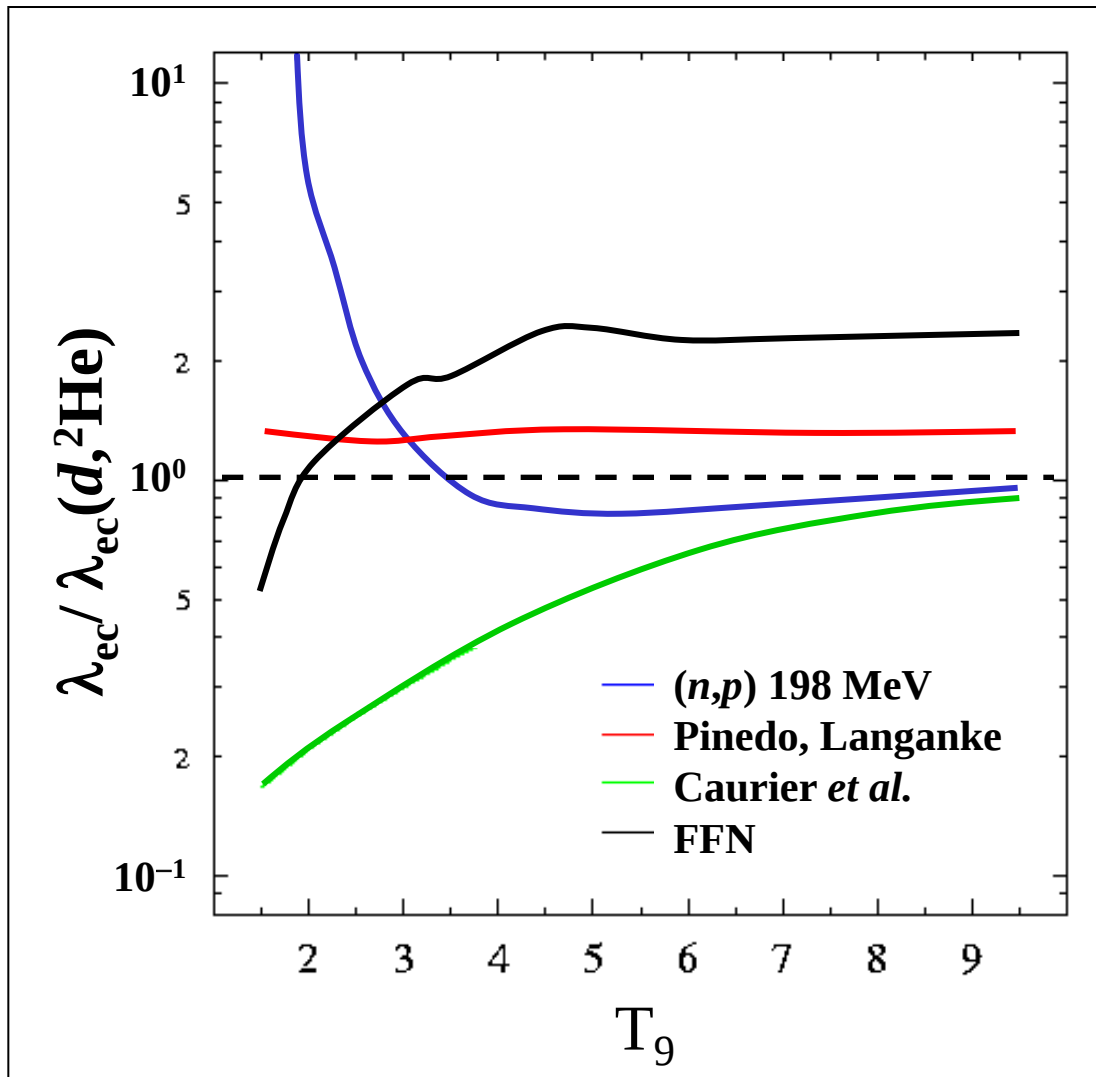
$$Y_e = 0.48$$

[Heger *et al.*, *Astrophys. J.* 560
(2001) 307]

Calculate EC rates as
function of T_9 for GT
transitions from ${}^{58}\text{Ni}_{\text{g.s.}}$

Strength deviations at low excitation $\Rightarrow$ rates deviation at low T

^{58}Ni : comparison of e -capture rates theory/experiment

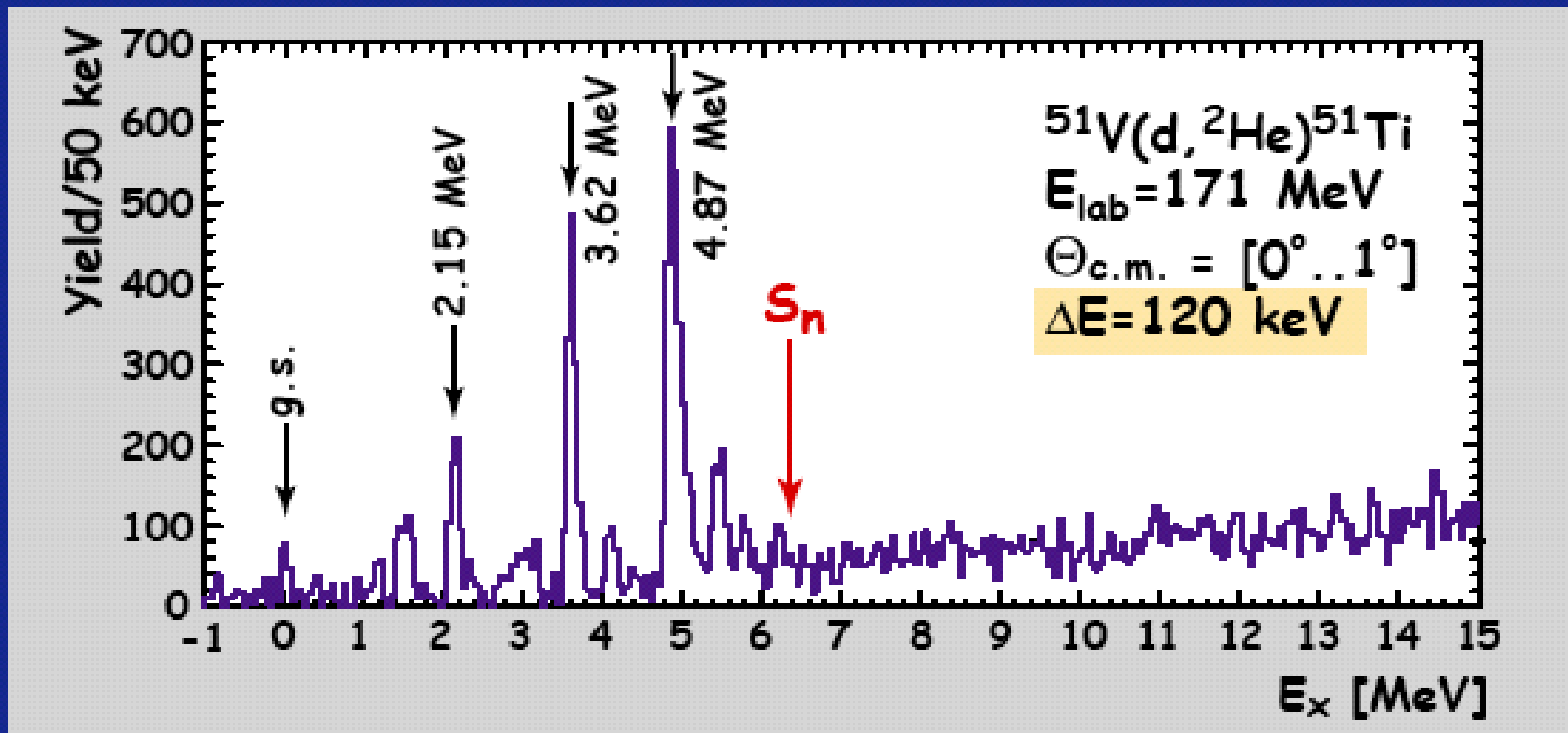


- Influence of GT strength distribution on calculated capture rate is dramatic, especially at low temperatures
- rates vary up to a factor 5-6
- FFN not too far off
- large scale shell-model calculations fail at low T
- calculations with improved residual interaction (KB3G) in reasonable agreement

$^{51}\text{V}(d, ^2\text{He})^{51}\text{Ti}$: $B(\text{GT}^+)$ for proton-odd *fp*-shell nucleus

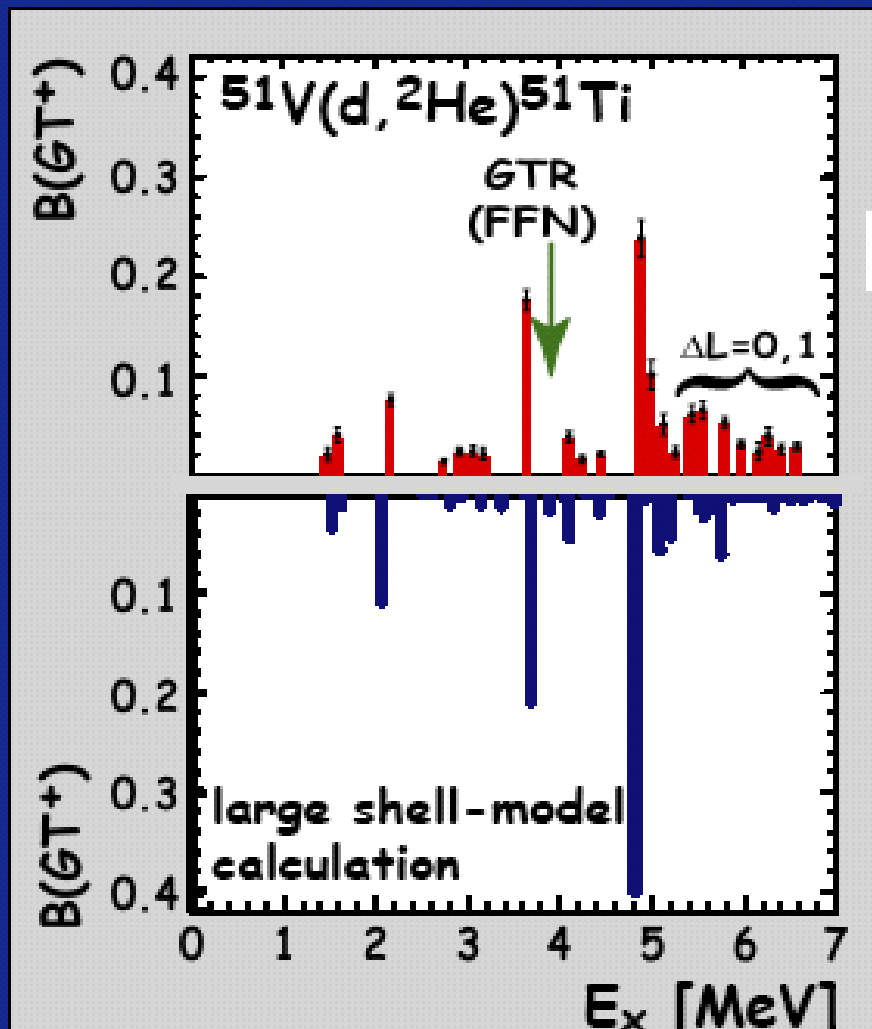
^{51}V g.s. ($J^\pi=7/2^-$, $T=5/2$) $\Rightarrow$ ^{51}Ti ($J^\pi=5/2^-, 7/2^-, 9/2^-, T=7/2$)

Independent single-particle model (FFN): $E_x(\text{GTR})=3.83$ MeV



C. Bäumer *et al.*, PRC **68**, 031303(R) (2003)

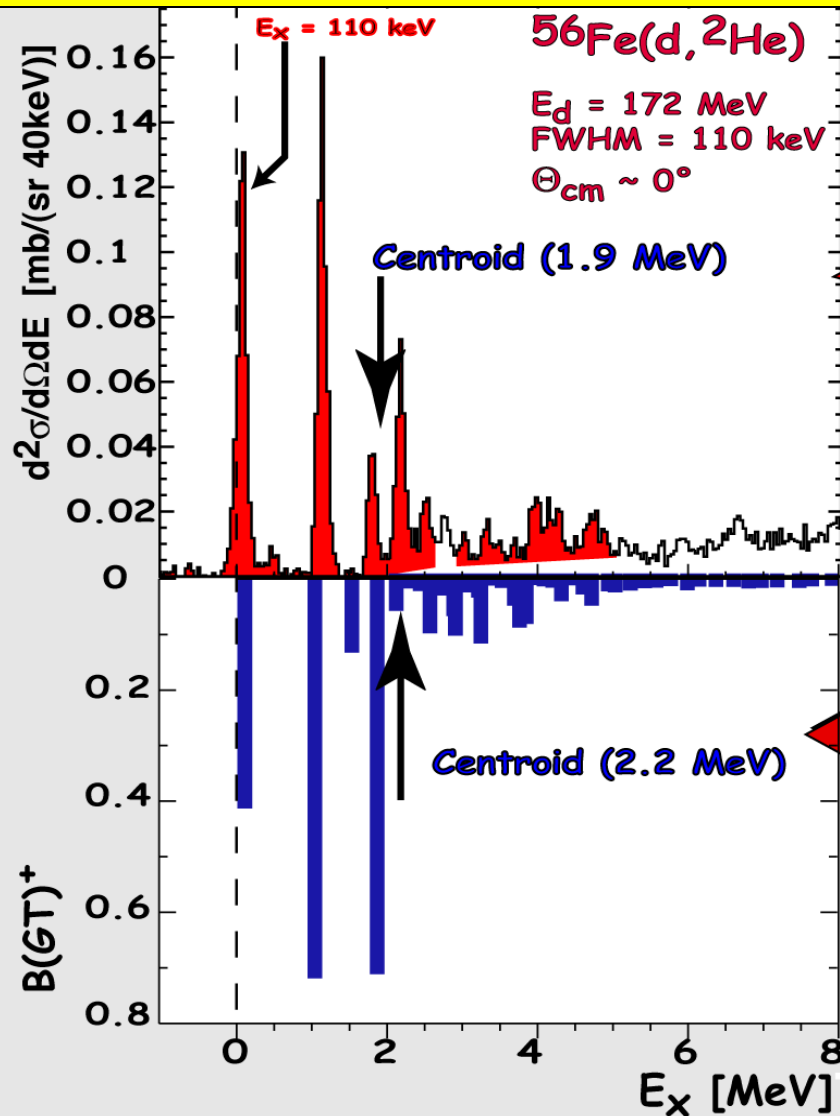
$^{51}\text{V}(d, ^2\text{He})^{51}\text{Ti}$: Comparison with shell-model calculations



← Experimental result

← Full *fp*-shell model calculations
quenching factor $(0.74)^2$
G. Martínez-Pinedo,
K. Langanke

$^{56}\text{Fe}(d,^2\text{He})$: Comparison with shell-model calculations



Experiment

Full *fp*-shell model
calculations (KB3G)
(G. Martínez-Pinedo)

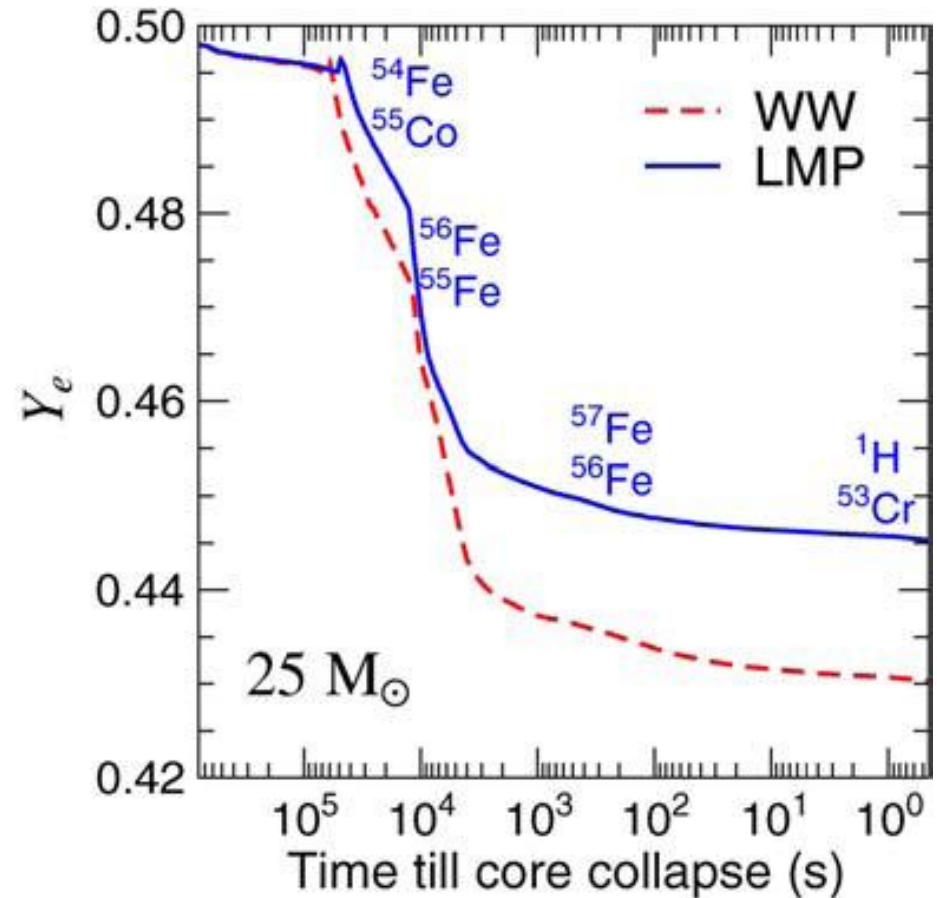
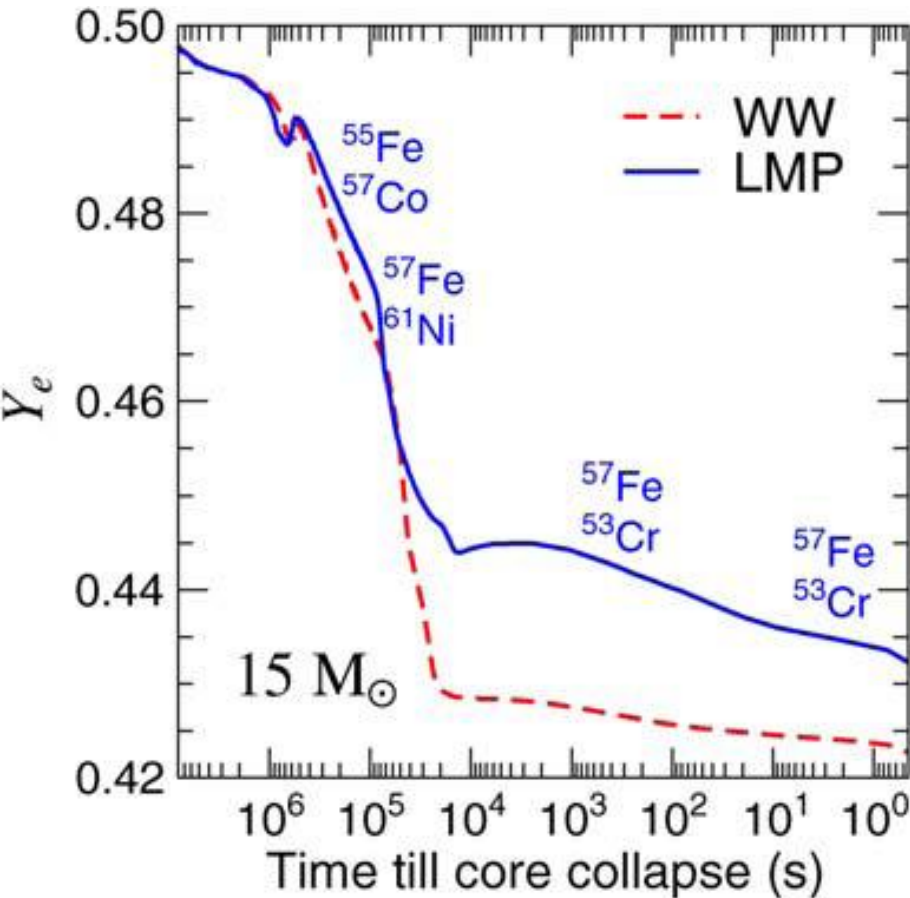
Comparison of centroids (MeV) of GT⁺ Strength distribution

	Nucleus	FFN	SM	Exp.
Even-Even	⁵⁶ Fe	3.8	2.2	1.9
	⁵⁸ Ni	3.8	3.6	3.4
Odd A-Odd <i>p</i>	⁵¹ V	3.8	4.7	4.1
Odd A-Odd <i>n</i>	⁵⁷ Fe	5.3	4.1	2.9
	⁶¹ Ni	3.5	4.6	4.2
	⁶⁷ Zn	4.4	--	3.4
Odd-Odd	⁵⁰ V	9.7	8.5	8.8

WW = Woosley-Weaver Model calculations (FFN rates)

LMP = Langanke-Martínez-Pinedo Large shell-model

calculations {G. Martínez-Pinedo *et al.*, NPA 777 (2006) 395}



Conclusions

- **Presupernova models depend sensitively on EC rates.**
- **GT⁺ transitions in *fp*-shell nuclei play a decisive role in determining EC rates and thus provide input into modeling of explosion dynamics of massive stars.**
- **Large shell-model calculations are needed especially as function of T. (Caurier *et al.*; Martínez-Pinedo & Langanke [KB3G]; Otsuka *et al.* [GXPF]) $\Rightarrow$ smaller EC rates for A=45-60 than FFN $\Rightarrow$ Larger Y_e (electron to baryon ratio) and smaller iron core mass (Heger *et al.*)**
- **New high resolution (*d*,²He) experiments provide essential tests for shell model calculations at 0 T.**

Outlook

Radioactive ion beams will be available at energies where it will be possible to study GT transitions (RIKEN, NSCL, FAIR, EURISOL)

- Determine $GT^\pm$ strength in unstable *sd* & *fp* shell nuclei
 - Electron capture rates (presupernova) and neutrino capture rates and inelastic scattering cross sections
- Use IV(S)GDR as tool to determine n-skin
 - Charge-exchange cross section proportional to n-skin

EuroSuperNova Collaboration

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M. Hunyadi³, **M. de Huu³**, **E. Jacobs⁴**, **B.C. Junk¹**, **A. Korff¹**,
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⁵ University Bari

⁶ Technische Universität Darmstadt

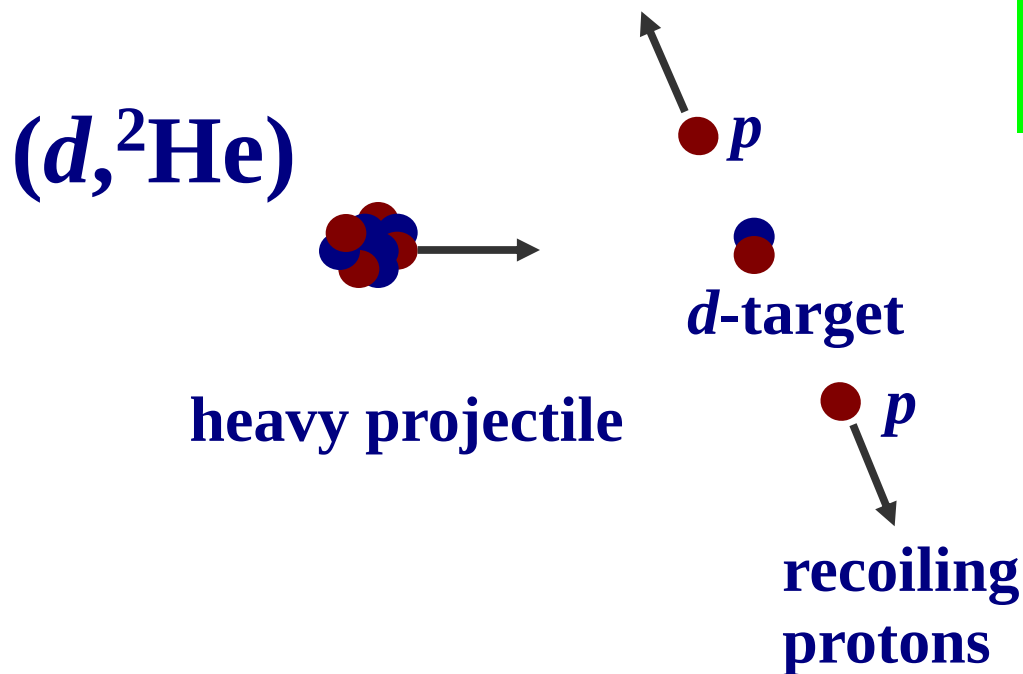
⁷ University Aarhus

⁸ Märkische Fachhochschule Iserlohn

Thank you for your attention

Nuclear structure studies with CE reactions in inverse kinematics

- Possible at FAIR and RIKEN
(intermediate beam energies are needed!)

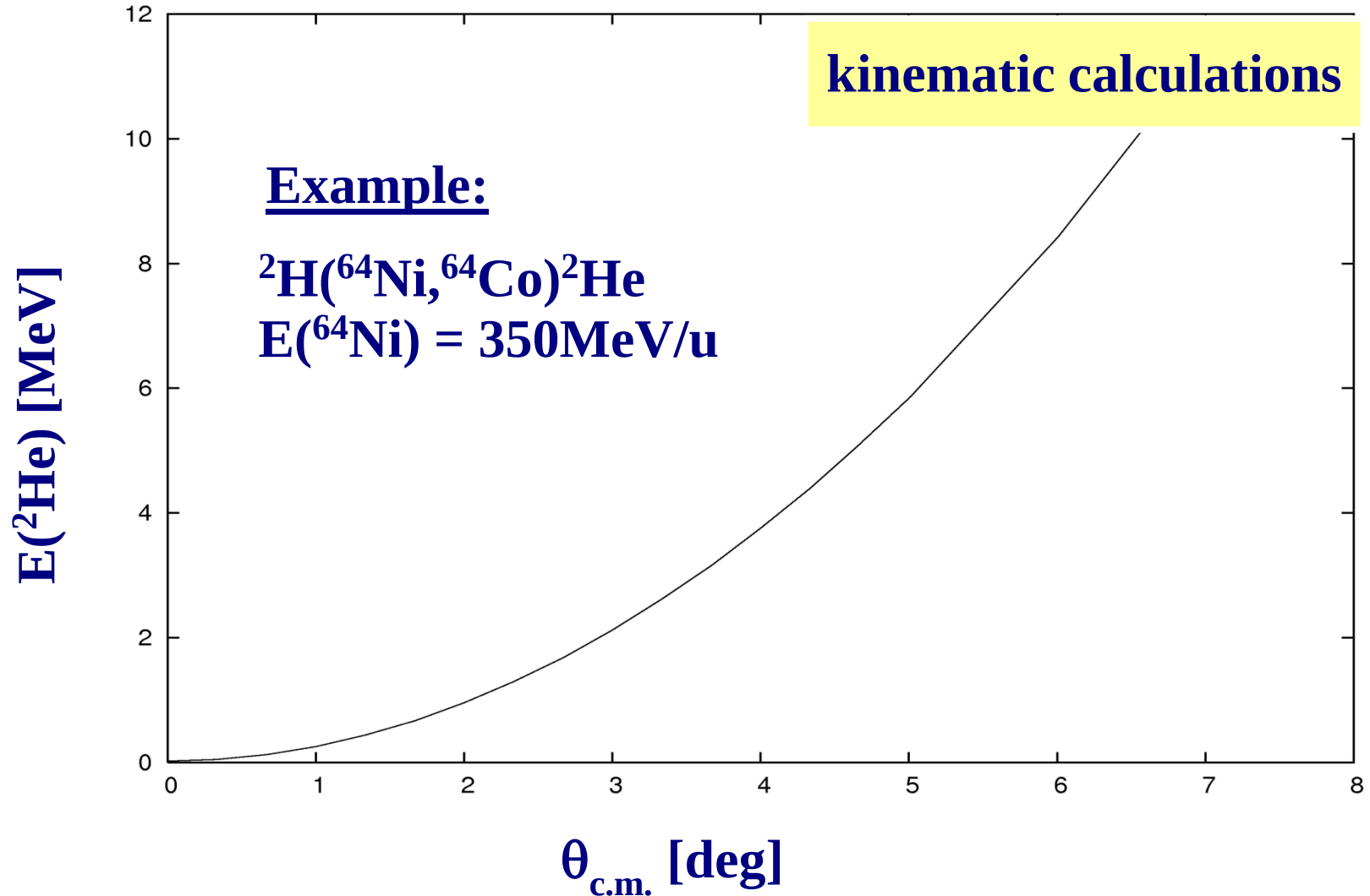


Approach (at FAIR):
measure the recoiling protons

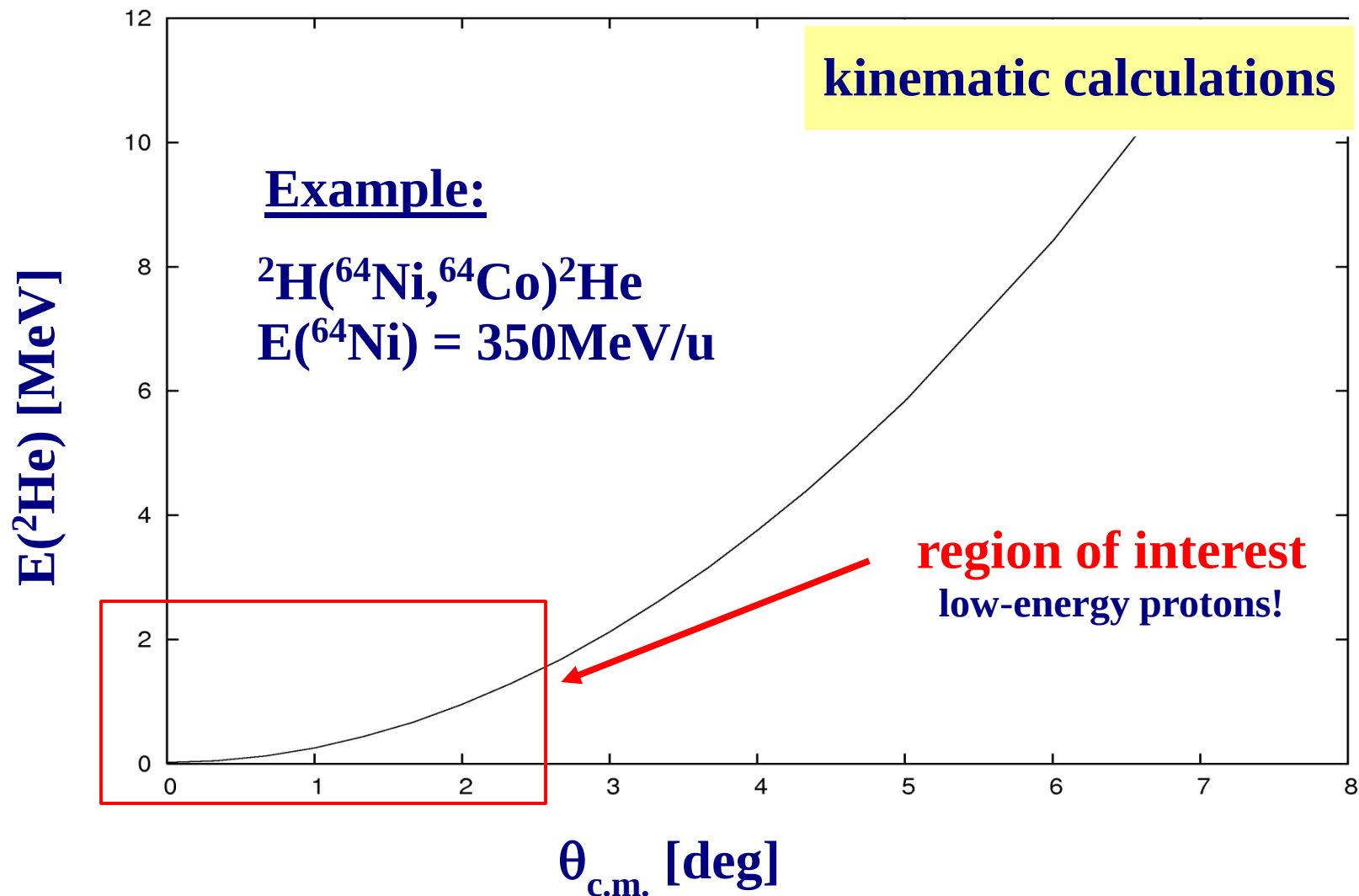
heavy ejectile

Inconvenience:
difficulty to detect the low-energy protons

How low?



How low?



Detection system @ FAIR

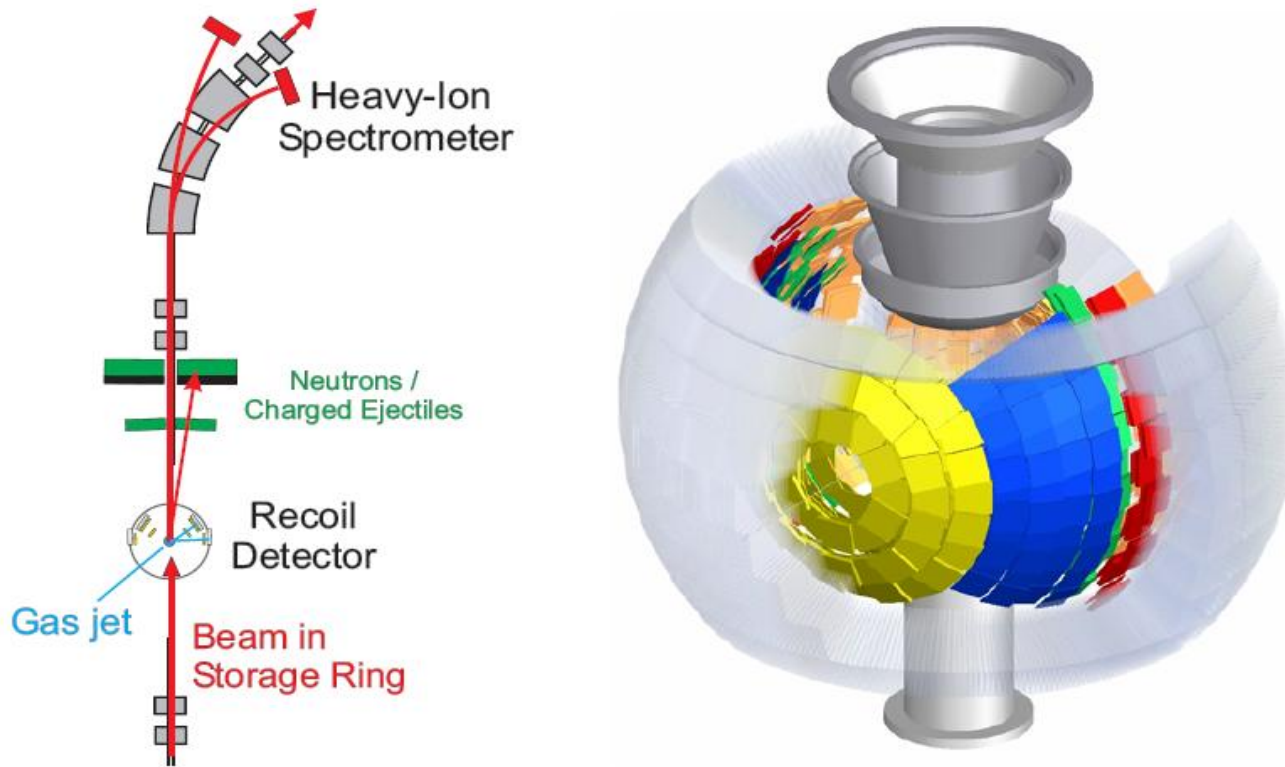


Figure 1: Schematic view of the EXL detection systems. Left: Set-up built into the NESR storage ring. Right: Target-recoil detector surrounding the gas-jet target.

- Use of EXL recoil detector is under evaluation
- Design & implementation of a dipole magnet for the momentum analysis of the protons