

Shell-model description for beta decays of pfg-shell nuclei

Workshop on
New Era of Nuclear Physics in the Cosmos
– the r-process nucleosynthesis
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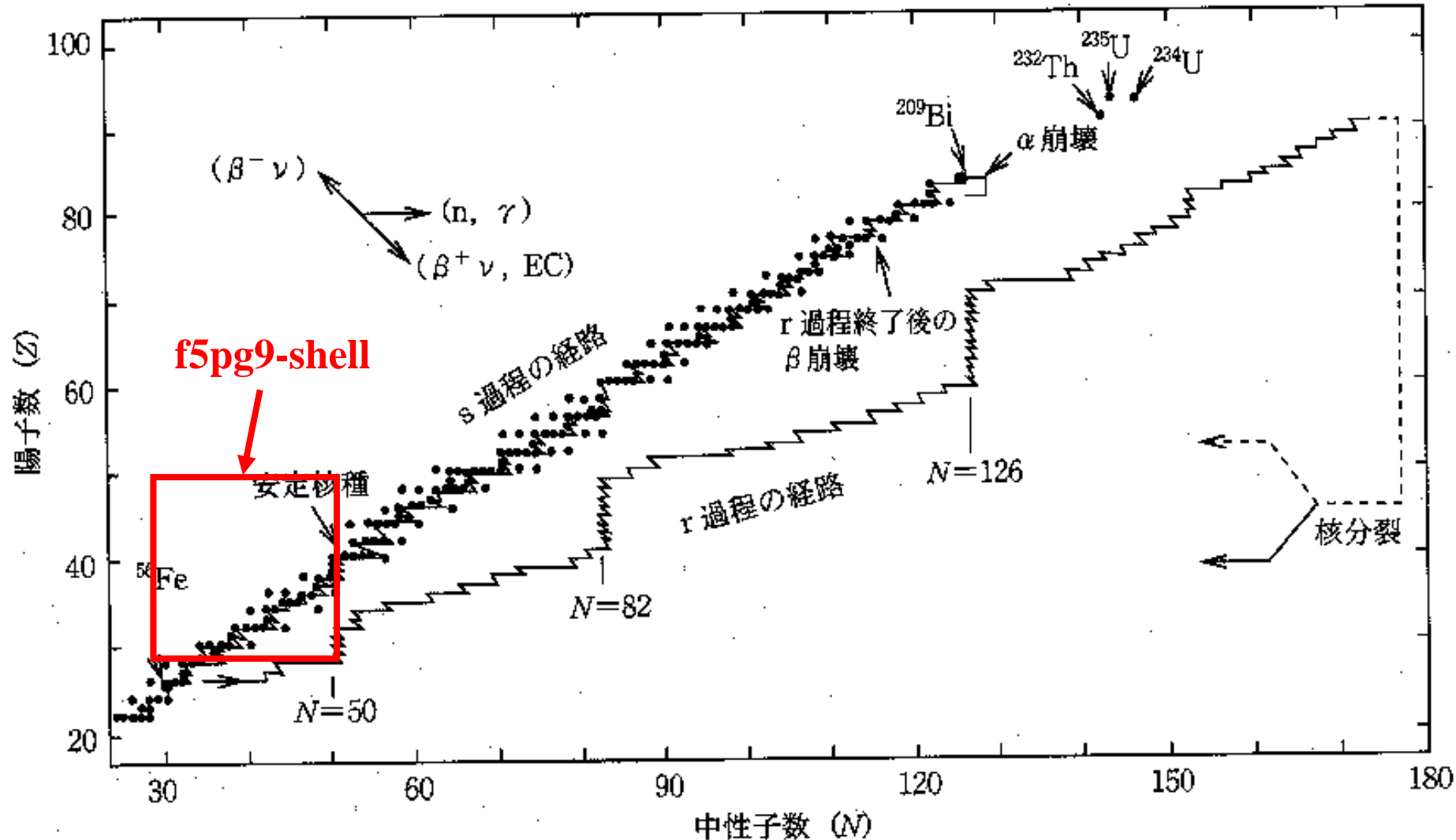
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M. Hjorth-Jensen (Oslo)

Introduction

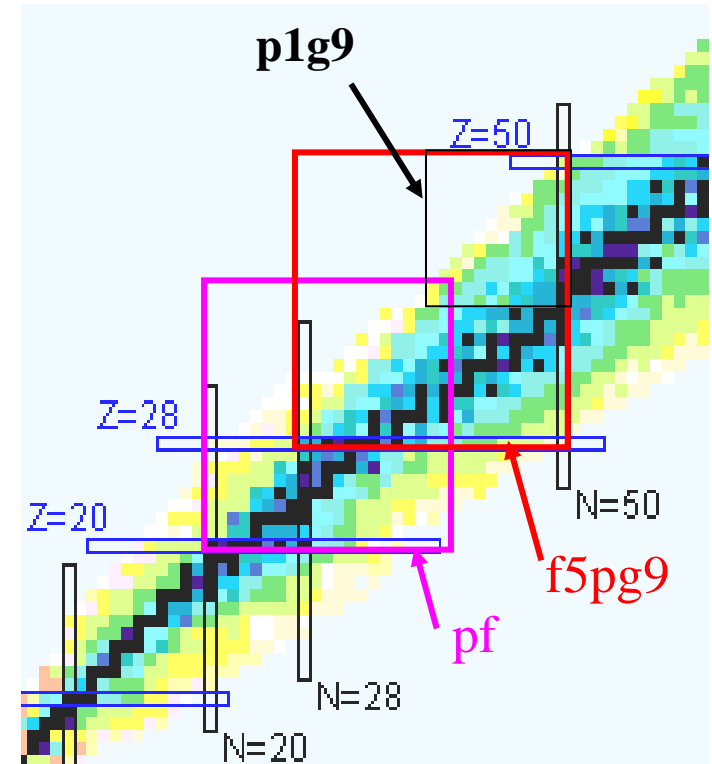
- Precise information of nuclear properties is needed for the analysis of the **r-process** nucleosynthesis which occur under extreme conditions which are not accessible by current experiments. Therefore, theoretical estimations are useful, and **predictions** with high accuracy are desired.
- Nuclear structure models based on the mean-field approximation such as **RPA** are widely used, which are applicable to any nuclei in the nuclear chart by taking sufficiently **large model space** and provide us reasonable description of gross properties. However, the results are not necessarily accurate because only **limited correlations** can be treated.
- The **shell model** can treat **any two-body correlations** and give accurate descriptions of nuclear structure, but its applicability is limited to relatively light nuclei or semi-magic nuclei because of too heavy numerical tasks. We have to take a **small model space** and introduce an **effective interaction**.
- **To what extent can shell model predict nuclear properties which are important for the study of the r-process?**

shell model for r-processes



f5pg9-shell

- f5pg9-shell
 - ^{56}Ni inert core
 - Valence orbits : $p_{3/2}$, $f_{5/2}$, $p_{1/2}$, $g_{9/2}$
 - No spurious center-of-mass motion
- Interests
 - Neutron-rich
 - Isomer
 - Shape-coexistence
 - Astrophysics
- Recent shell-model studies
 - S3V... J.Sinatkas, et al., J. Phys. G18, 1377; 1401 (1992)
 - Second order correction to the Sussex matrix elements
 - $N=50,49,48$ with severe truncation (weak coupling assumption)
 - Lisetskiy... A.F.Lisetskiy et al., PRC70, 044314 (2004)
 - Modify G-matrix interaction by least squares fit
 - $T=1$ part for proton and neutron (Ni isotopes and $N=50$ isotones)



Effective interaction

- JUN45 interaction

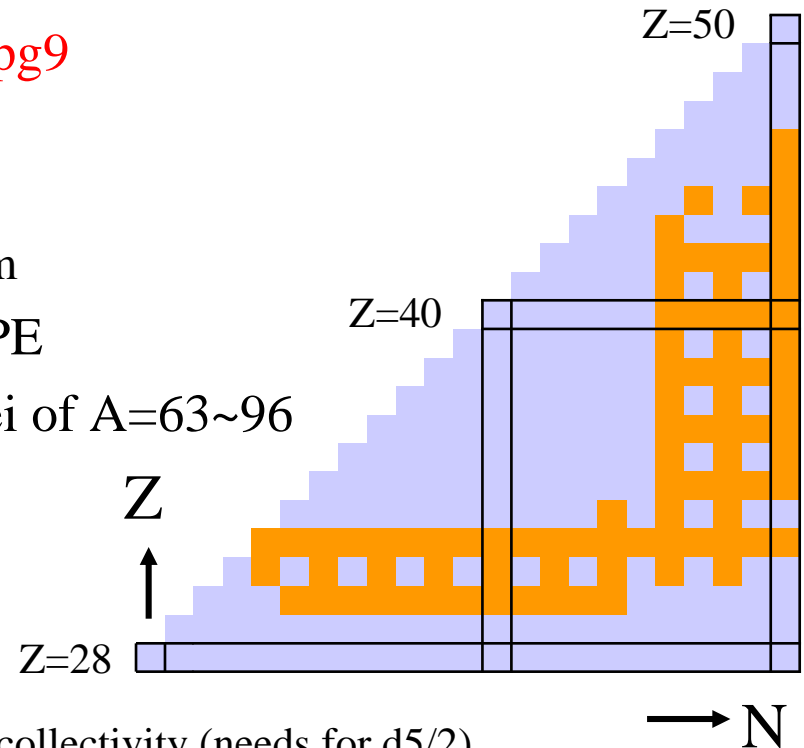
- Keep isospin symmetry
- Modify **microscopic interaction G-f5pg9**

M. Hjorth-Jensen, unpublished

Bonn-C potential

3rd order Q-box and folded diagram

- Vary 45 LC's of **133** TBME and **4** SPE
- Fit to 400 energy data out of 87 nuclei of $A=63\sim 96$
 - Include low-lying states of
 - even-Z nuclei
 - odd-A nuclei
 - Exclude
 - $N < 46$ for $Z > 33$... large quadrupole collectivity (needs for $d5/2$)
 - Ni, Cu isotopes ... large effects of $f7/2$ core-excitations
- Assume **$A^{-0.3}$** mass dependence
- Rms error of **185 keV**

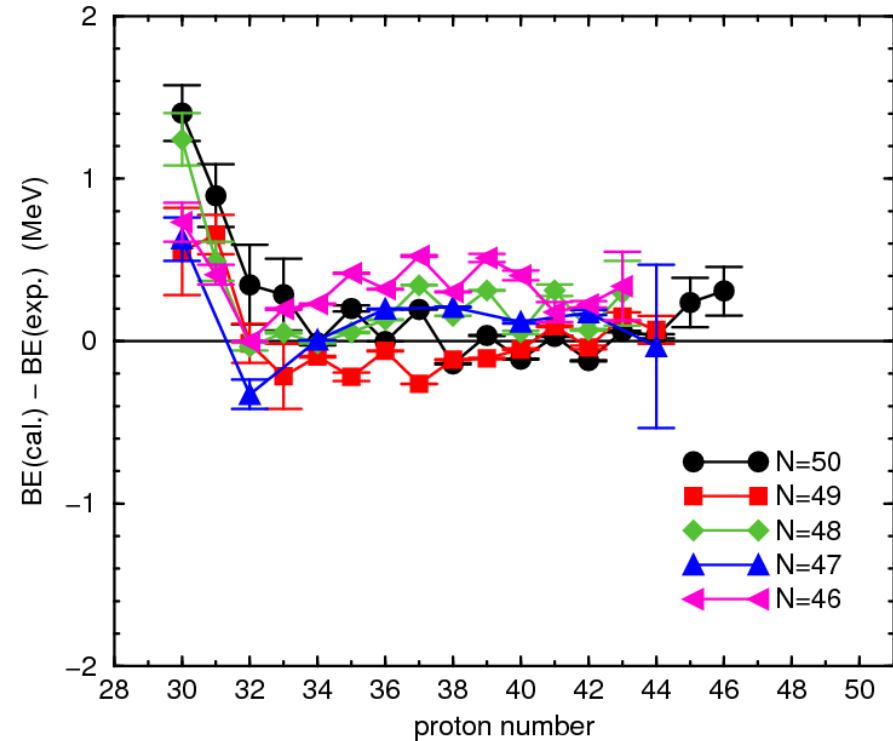
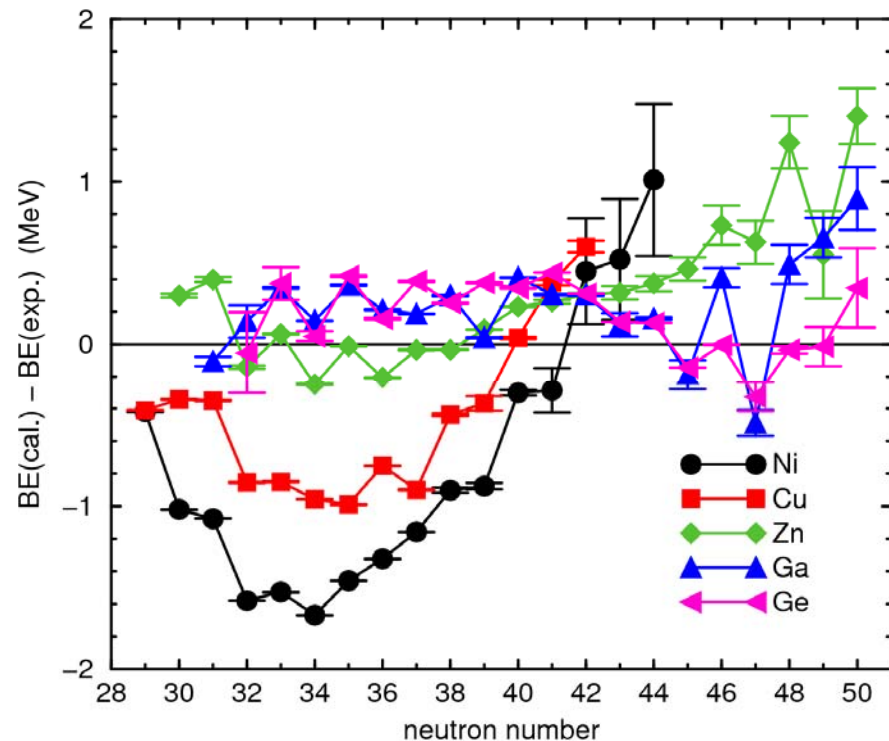


Binding energy

- Empirical Coulomb energy B.J.Cole, PRC59(1999)726

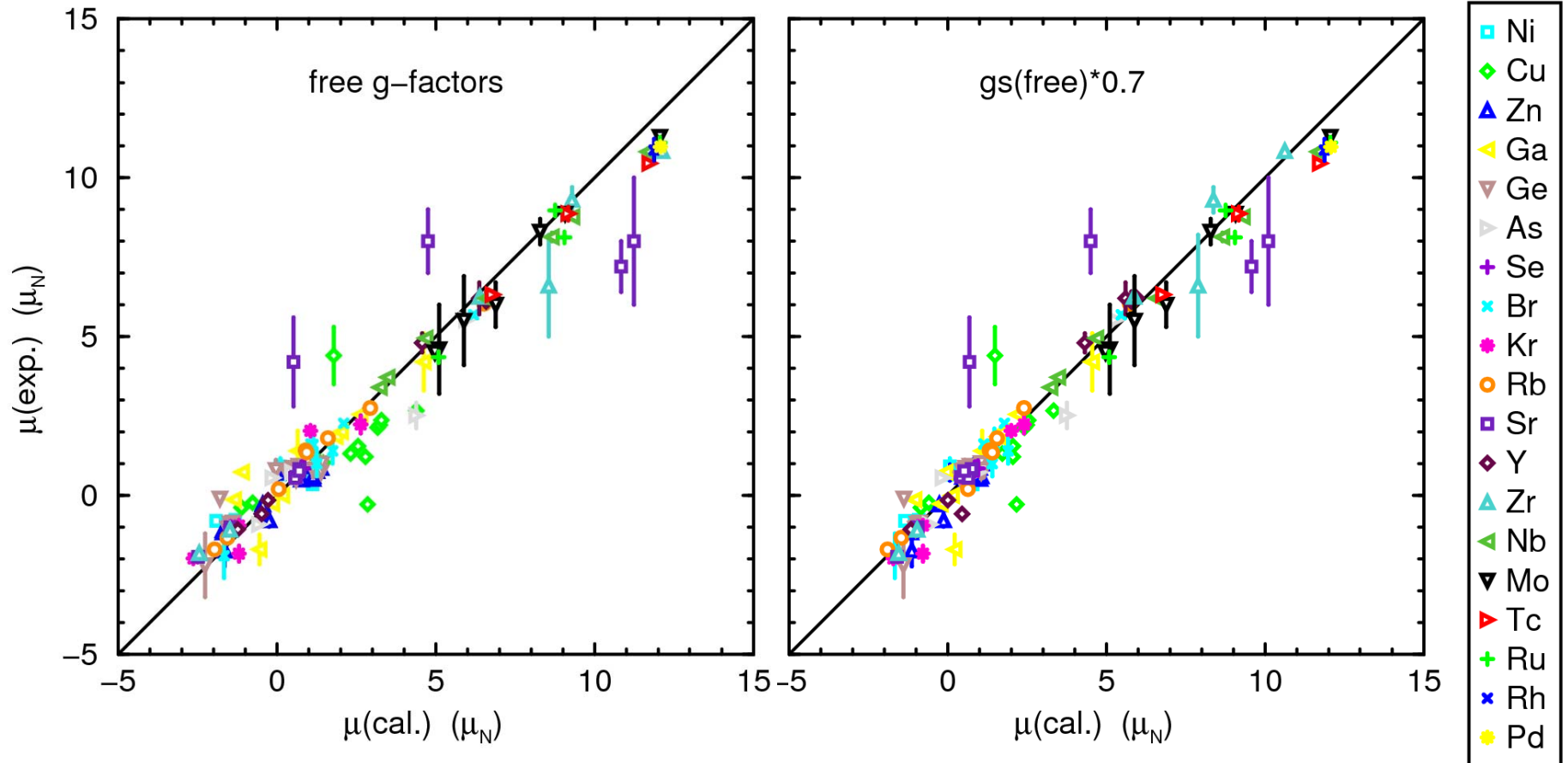
$$E_C(\pi, \nu) = \varepsilon_c \pi + V_c \frac{\pi(\pi-1)}{2} + b_c \left[\frac{1}{2} \pi \right] - \Delta_{\pi\nu} \pi \nu \quad \pi, \nu \dots \text{valence nucleon \#}$$

- f7/2 effects ?
 - Underbinding in Ni, Cu with N<40
 - Overbinding in neutron-rich Ni, Cu, Zn, Ga



Magnetic moments

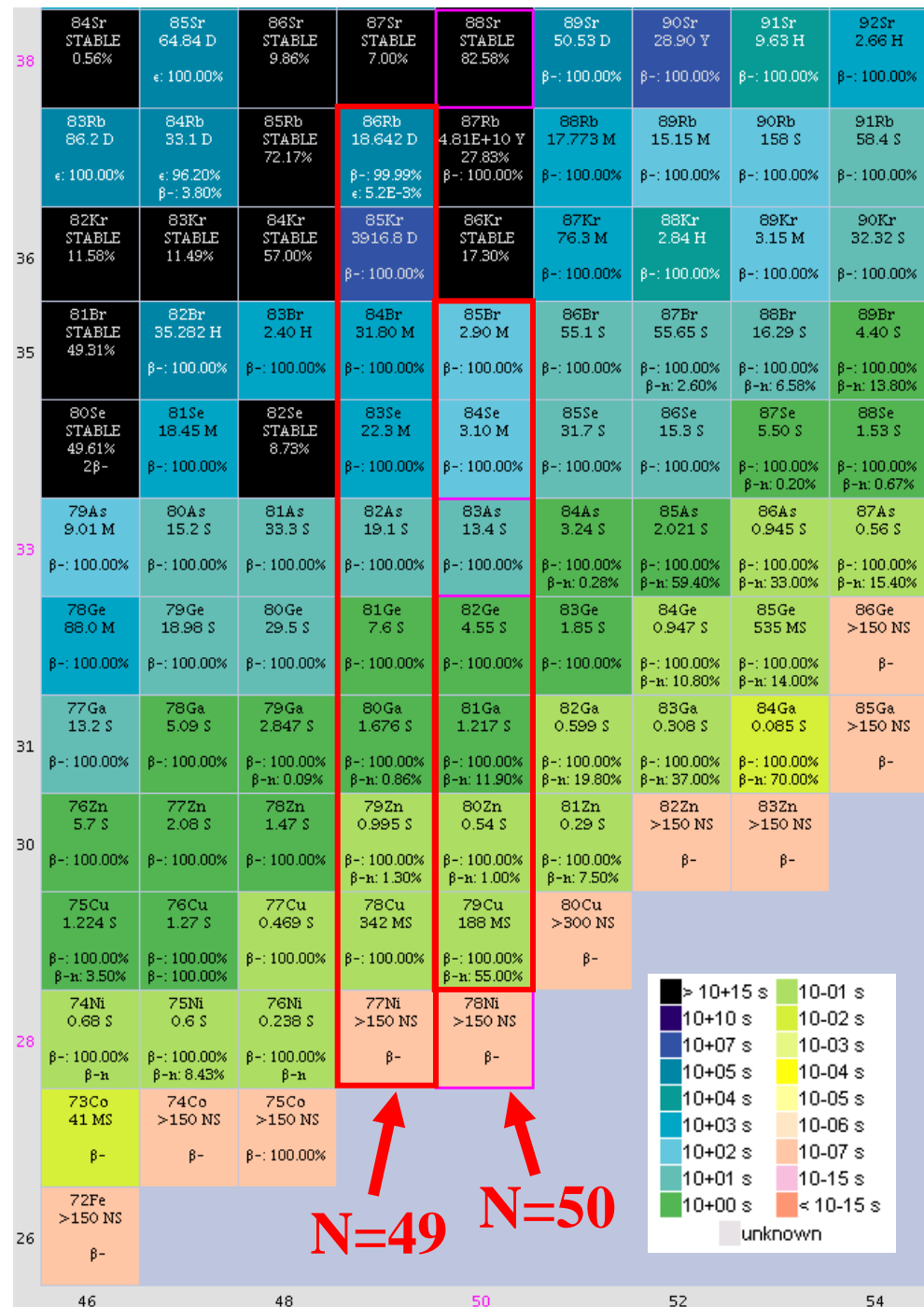
- Effective spin g-factors
 - Free-nucleon values are already good (also for sd- and pf-shell)
 - Slight improvement by effective $g_s(\text{eff})/g_s(\text{free}) \sim 0.7$
 - Note: $\sim 0.85 \dots$ GXPF1 for pf-shell



β^- decay

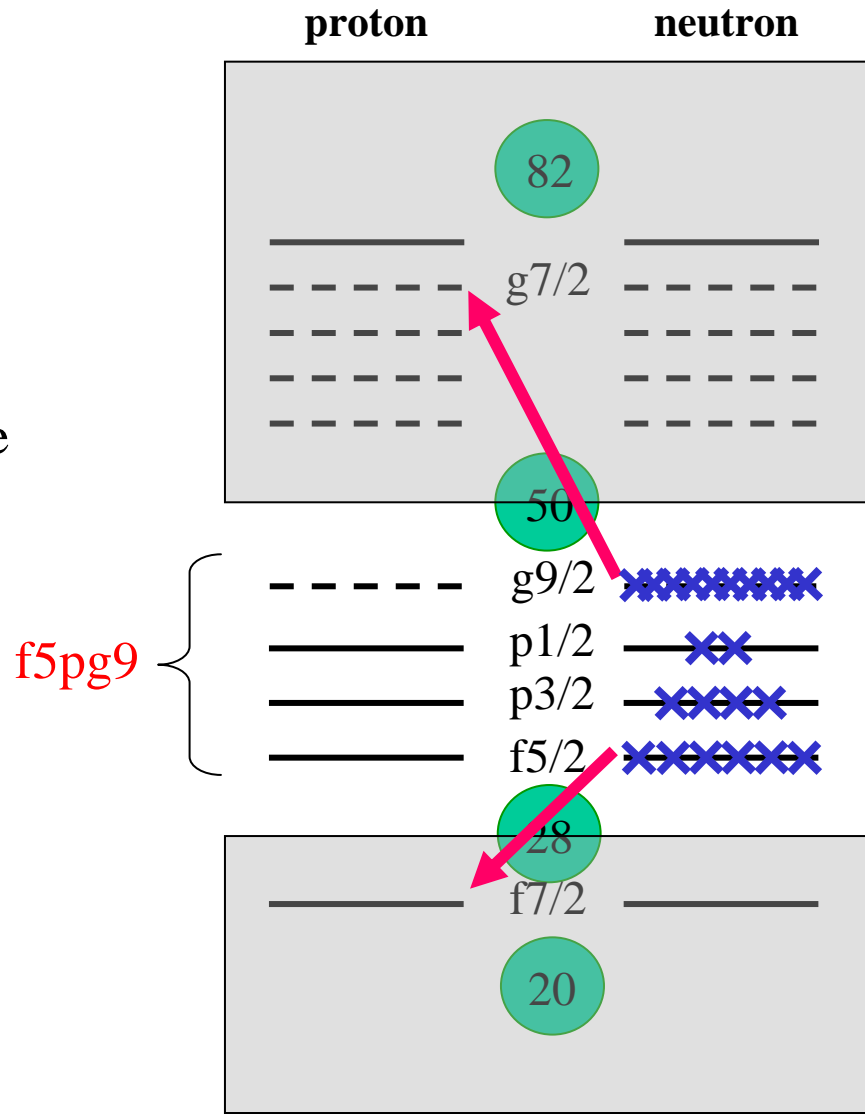
- Q_{β^-} and $T_{1/2}$
- $N=50$
 - Waiting point
 - $T=1$ TBME contributes
 - ^{78}Ni ... No valence particles/holes
- $N=49$
 - Low-lying isomer

NNDC Chart of Nuclides
<http://www.nndc.bnl.gov/chart/>



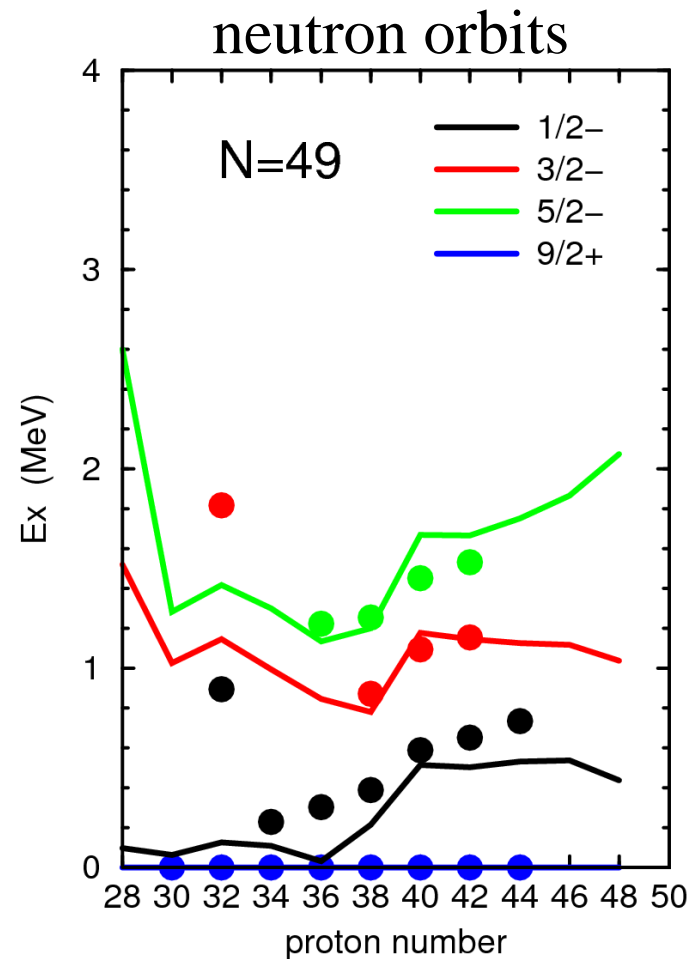
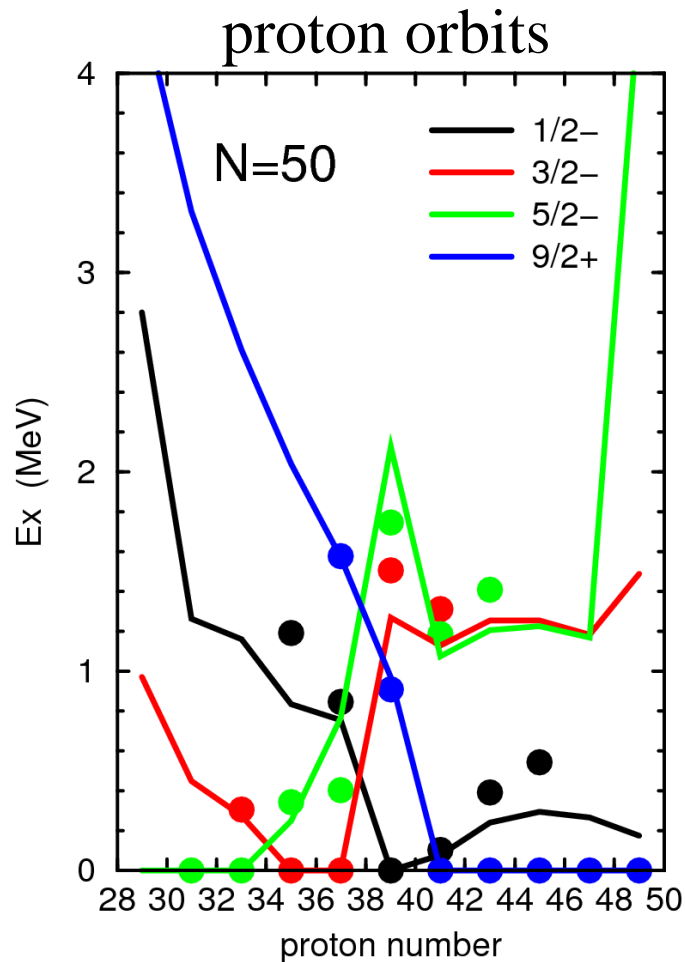
Neutron-rich nuclei

- f5pg9-shell
Spin-flip decay
 $\nu g_{9/2} \rightarrow \pi g_{7/2}$
 $\nu f_{5/2} \rightarrow \pi f_{7/2}$
are out of the model space
- Large Q-value
Higher excited states may be important
- Phase space factor
Low-lying states mainly contribute to $T_{1/2}$



Single-particle levels around N=50

- Low-lying states in odd-A nuclei



β -decay

- Q-value

- β^- $(A, Z) \rightarrow (A, Z+1) + e^- + \bar{\nu}_e$ $Q = ({}_Z^A M - {}_{Z+1}^A M)c^2$
- β^+ $(A, Z) \rightarrow (A, Z-1) + e^+ + \nu_e$ $Q = ({}_Z^A M - {}_{Z-1}^A M)c^2 - 2mc^2$
- EC $(A, Z) + e^- \rightarrow (A, Z-1) + \nu_e$ $Q = ({}_Z^A M - {}_{Z-1}^A M)c^2 - E_K$

- Allowed transition matrix elements

- Fermi

$$B(F) = \frac{\left| \left\langle f \left\| \Sigma_k \tau_{\pm}^k \right\| i \right\rangle \right|^2}{2J_i + 1}$$
 $\Delta J=0, \Delta\pi=\text{No}$
- Gamow-Teller

$$B(GT) = \frac{\left| \left\langle f \left\| \Sigma_k \sigma^k \tau_{\pm}^k \right\| i \right\rangle \right|^2}{2J_i + 1}$$
 $\Delta J=0, \pm 1, \Delta\pi=\text{No}$

- Lifetime

$$ft = \frac{6144.4 \pm 1.6 \text{ sec}}{B(F) + (g_A/g_V)^2 B(GT)}$$

- $g_A/g_V = -1.2720(18)$... ratio of axial-vector and vector c.c.
- f ... phase space factor

GT quenching

- **Quenching** of GT operator : $O_{\text{GT}} \rightarrow qO_{\text{GT}}$
 - Take into account the configurations outside the model space
 - p-shell ... $q = 0.82$ W.T.Chou et al, PRC47 (1993) 163
 - sd-shell ... $q = 0.77$ B.H.Wildenthal et al., PRC28 (1983) 1343
 - pf-shell ... $q = 0.74$ G.Martinez-Pinedo et al., PRC53 (1996) R2602
- f5pg9-shell
 - Violate Ikeda sum-rule : $S_- - S_+ = 3(N-Z)$
 - Low-lying transitions ... spin-flip contribution may be minor
- Fitting to β -decay data
 - Effective GT matrix element $M_{\text{eff}} = \sqrt{(2J_i + 1)B(GT)}$
 - Sum over final states relative to the sum-rule

$$T = \sqrt{\sum_i (M_{\text{eff}i})^2} / W \quad W = \begin{cases} |g_A/g_V| \sqrt{(2J_i + 1)3|N_i - Z_i|} & (\text{if } N_i \neq Z_i) \\ |g_A/g_V| \sqrt{(2J_f + 1)3|N_f - Z_f|} & (\text{if } N_i = Z_i) \end{cases}$$

Summed strength (f5pg9-shell)

- Exp. Data
 - setA ... all data
 - setB ... spin assigned
 - setC ... error in $E_x < 0.1\text{MeV}$

- Fit to setC

$$\Rightarrow q \sim 0.65$$

(with ^{64}Ge , ^{78}Ge)

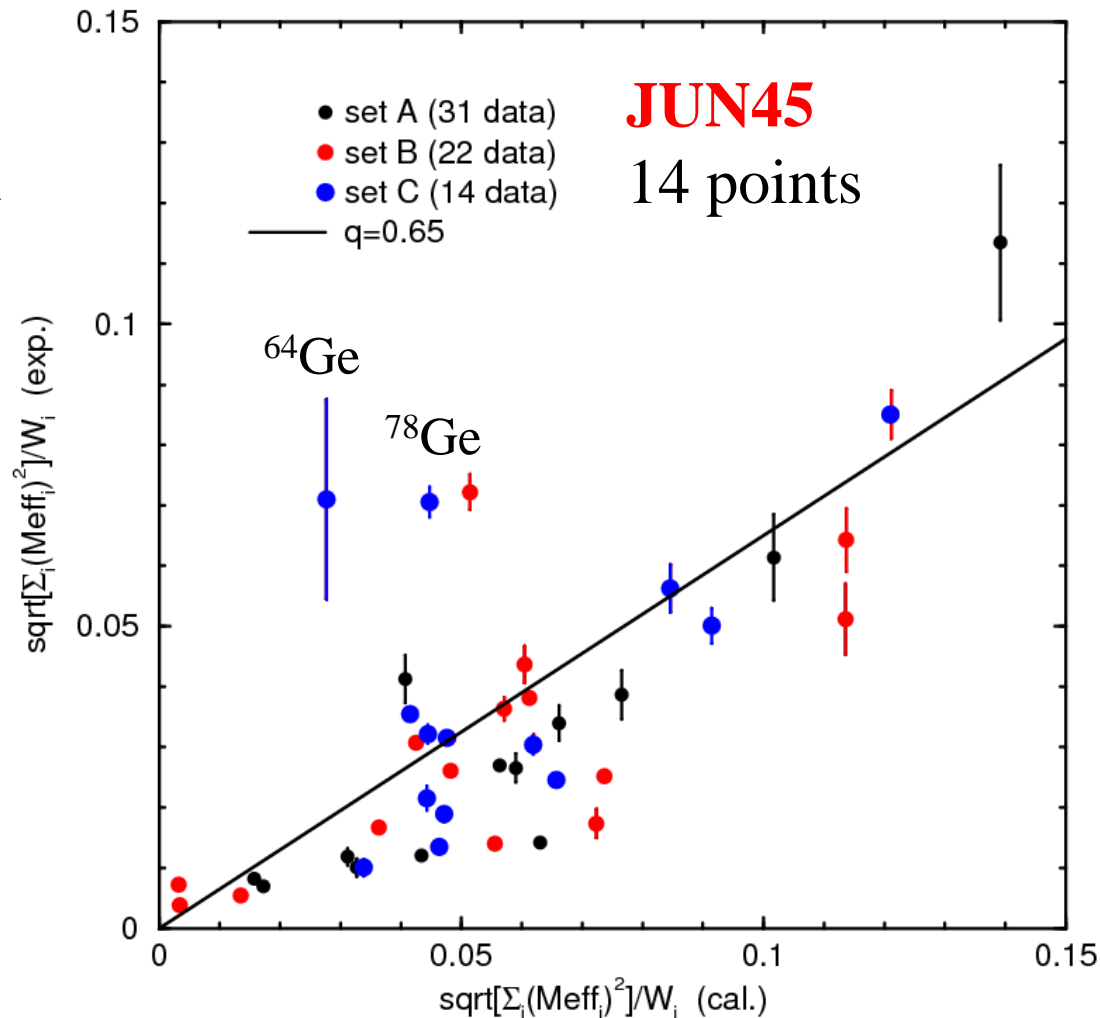
$$\Rightarrow q \sim 0.59$$

(without ^{64}Ge , ^{78}Ge)

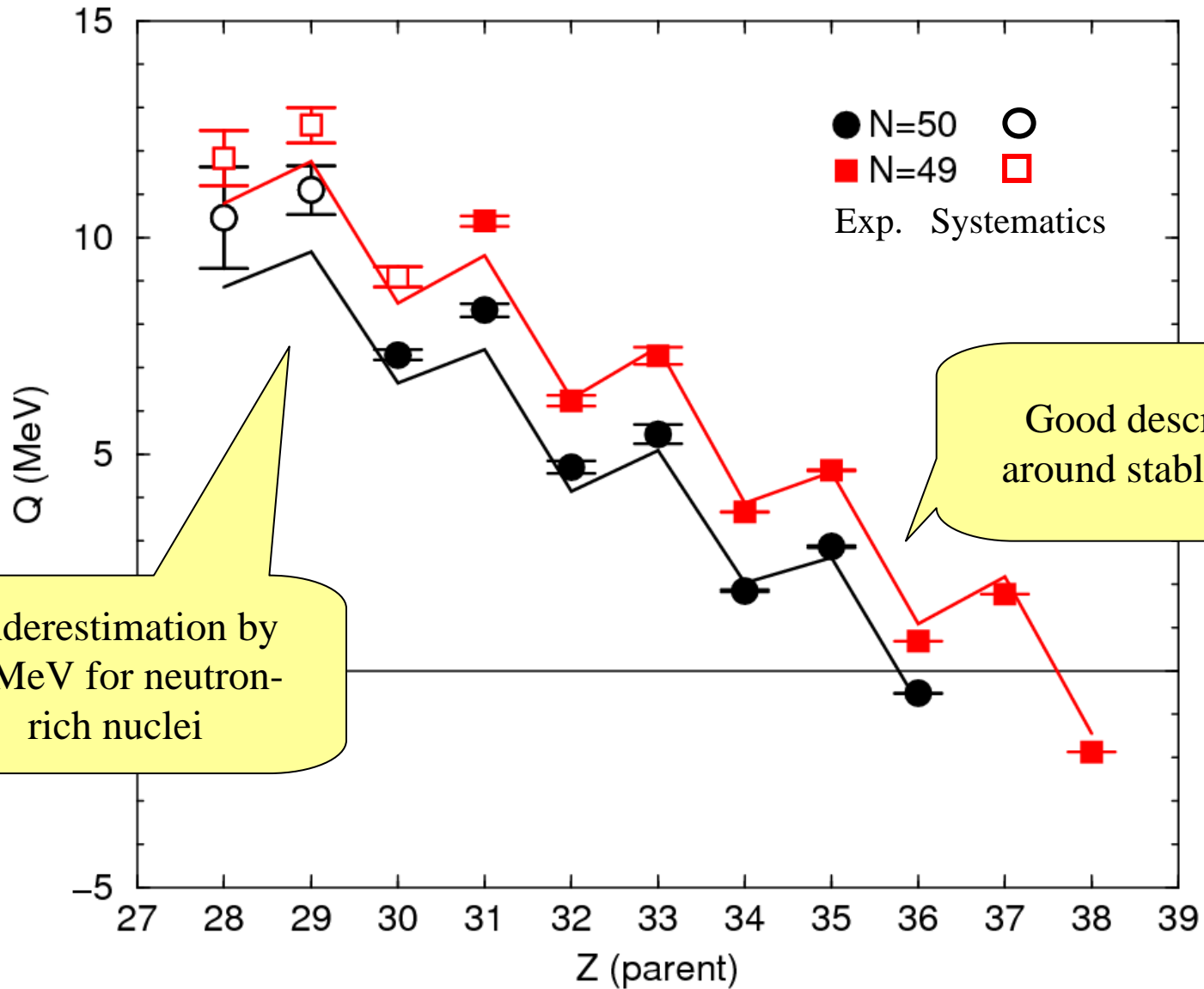
Adopt

$$q = 0.6$$

(large uncertainty)

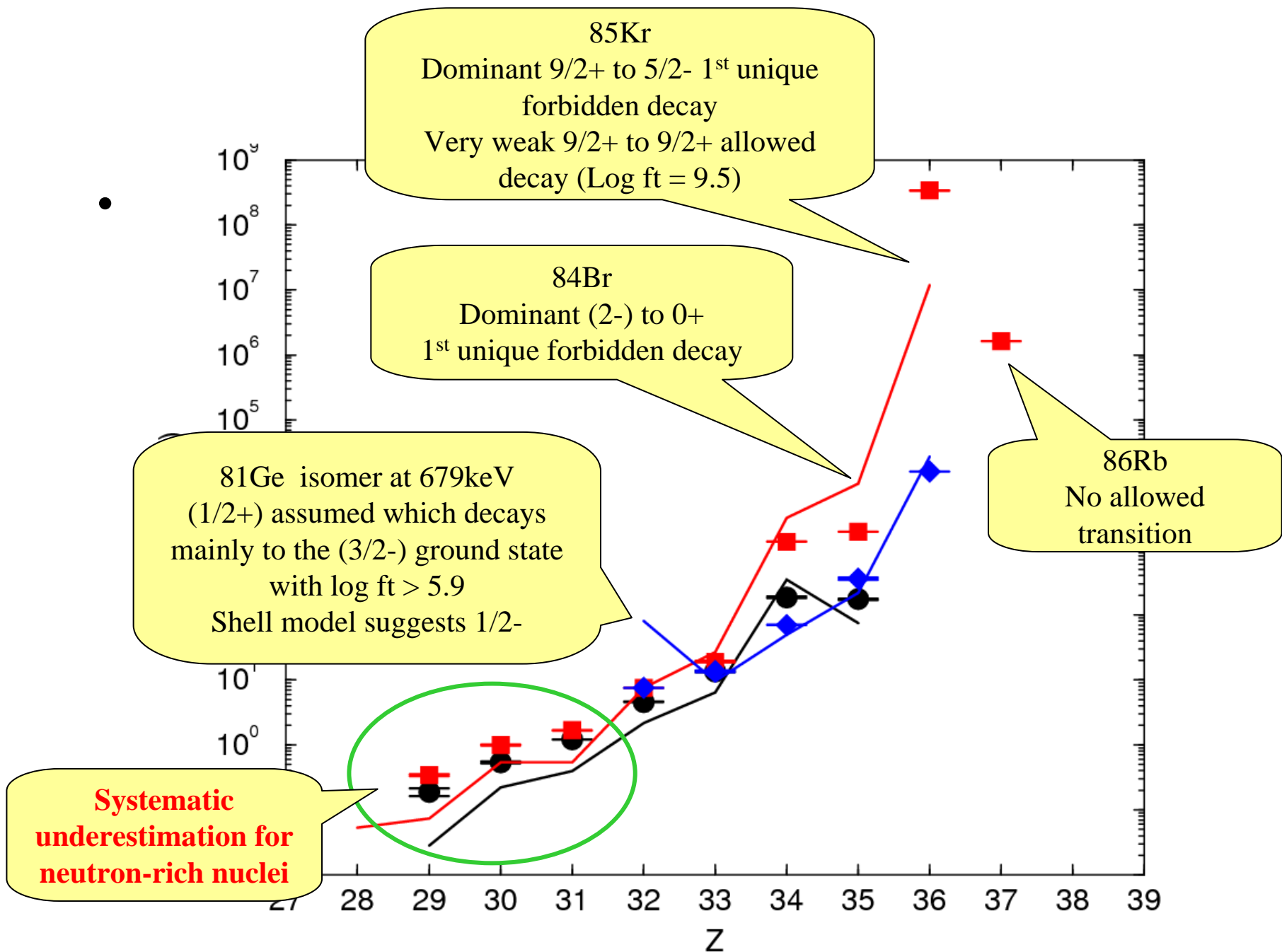


Q-value



Underestimation by
~1MeV for neutron-
rich nuclei

Good description
around stable nuclei



^{79}Cu

Shell model

5/2- ground state

$Q(\text{cal}) = 9.67\text{MeV}$

$T_{1/2}(\text{cal.}) = 0.028\text{ s}$

- No experimental level scheme

^{79}Cu β^- Decay (188 ms) [1991Kr15](#)

200206

Published: 2002 Nuclear Data Sheets.

^{79}Cu Parent: $E_x=0.0$; $T_{1/2}=188\text{ ms } 25$; $Q_{g.s.\rightarrow g.s.}=11742\text{ SY}$

History

Type	Author	Citation
Full evaluation	Balraj Singh	Nuclear Data Sheets 96, 1

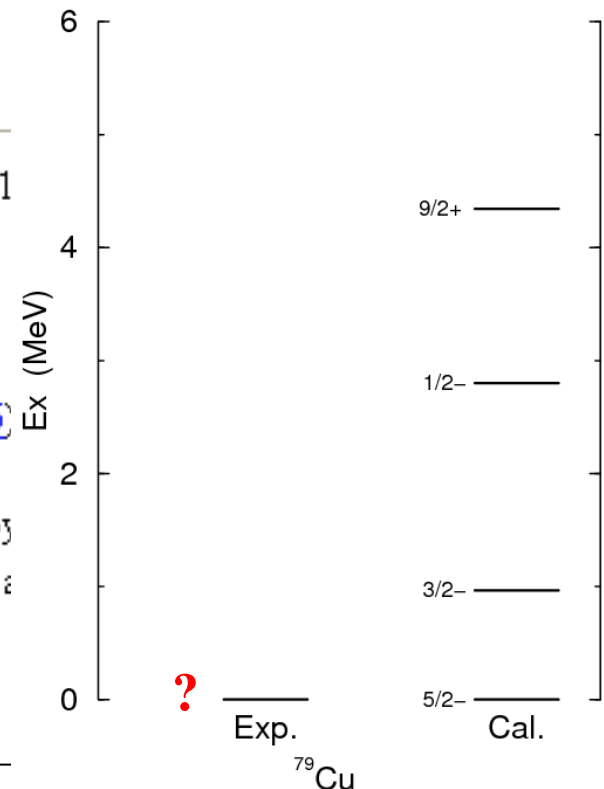
Full evaluation Balraj Singh Nuclear Data Sheets 96, 1

Measured $T_{1/2}$, % β^- n.

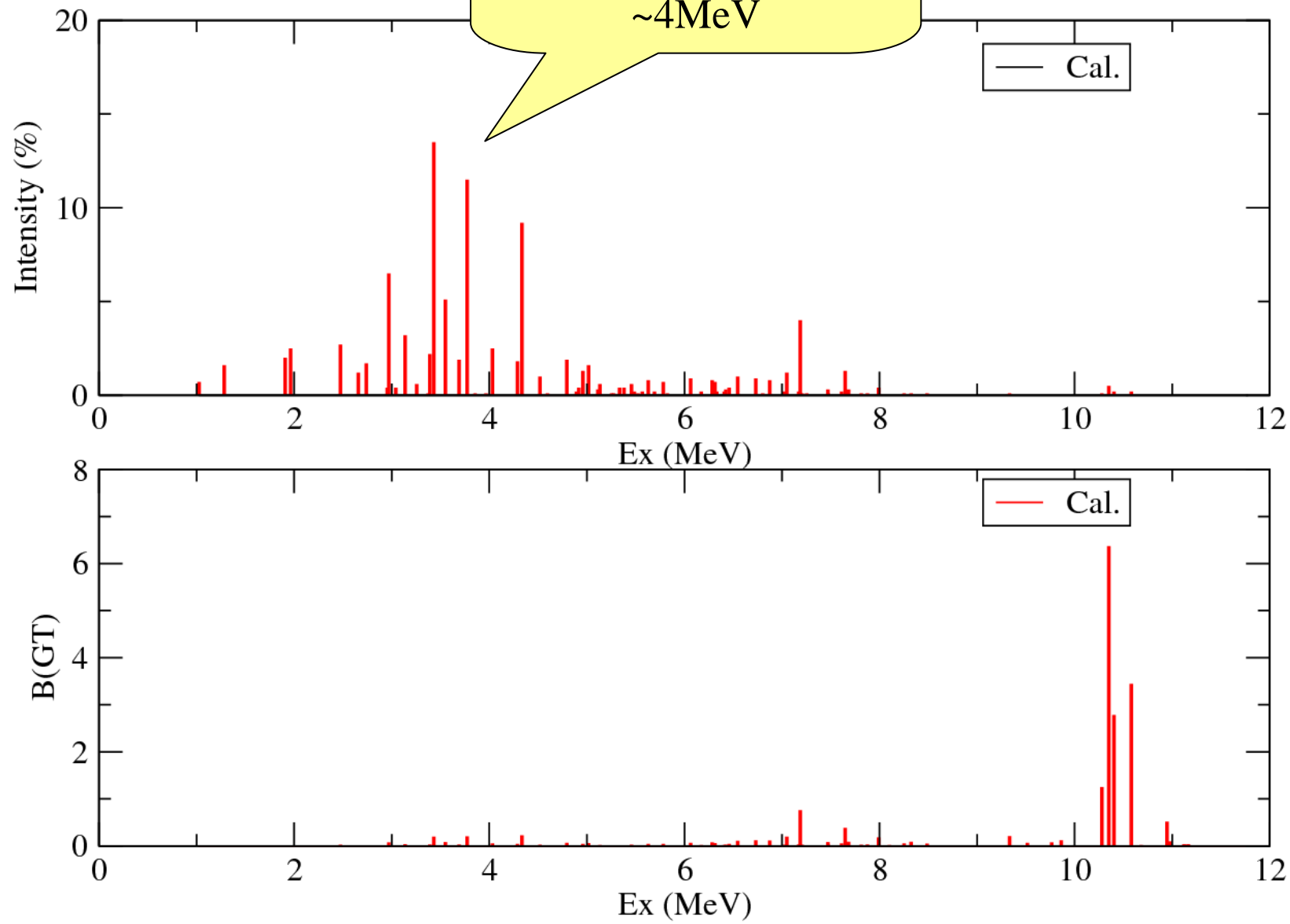
$Q(\beta^-)(^{79}\text{Cu})=11742\text{ 943 (syst, 1995Au04)}$; % β^- n=55 17 ([1991Kr15](#))

The isotope produced by $^{238}\text{U}(p,X)$ $E=600\text{ MeV}$ reaction followed by and mass separation techniques. [1995En07](#) use $^9\text{Be}(^{238}\text{U},F)$ reaction and magnetic methods to identify ^{79}Cu .

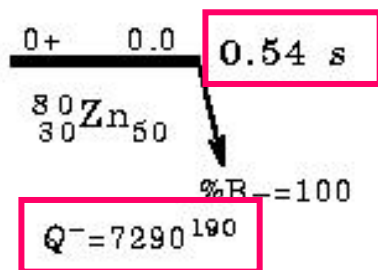
The details of this decay mode are not known.



^{79}Cu



^{80}Zn

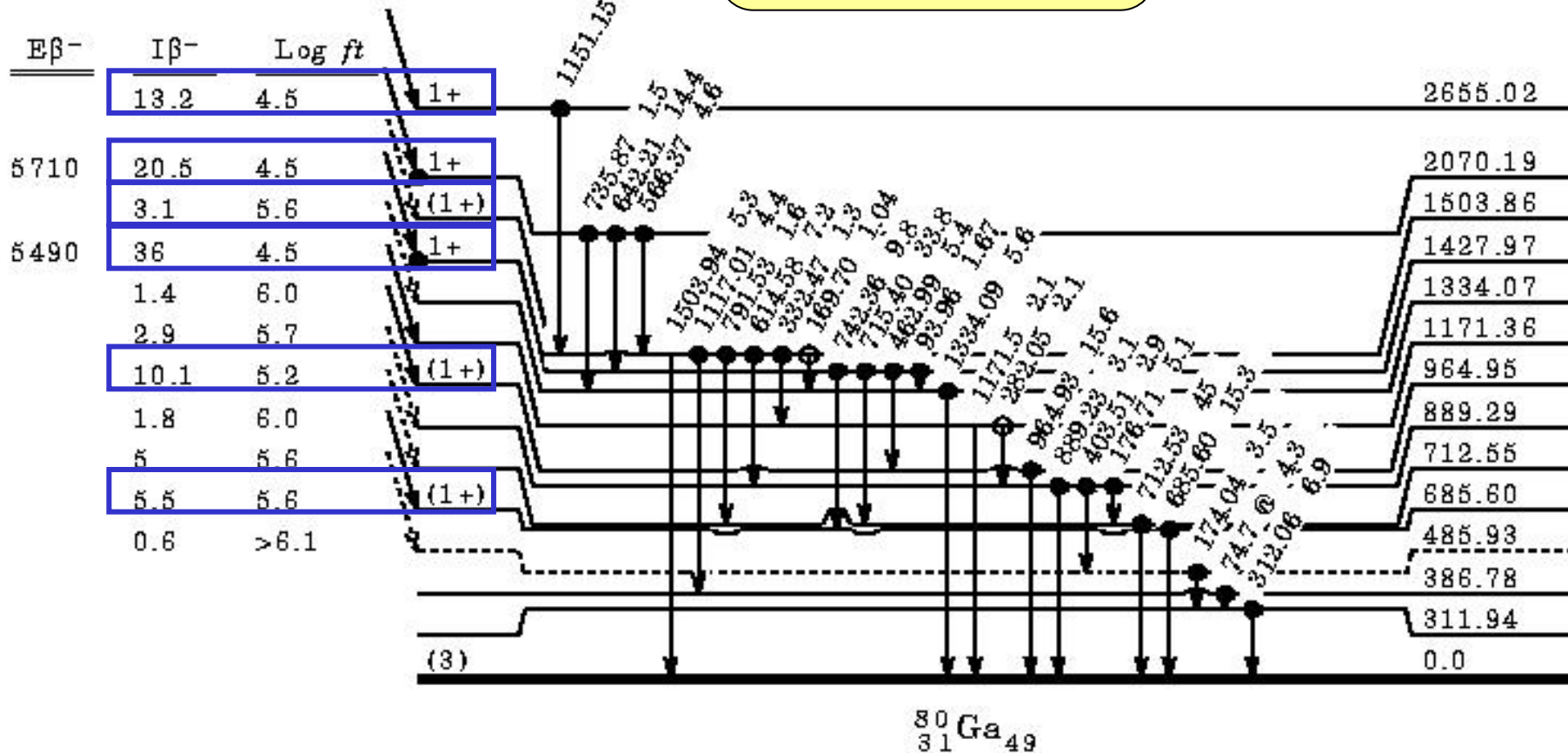


Intensities: I(γ +ce) per 100 parent decays suitably divided

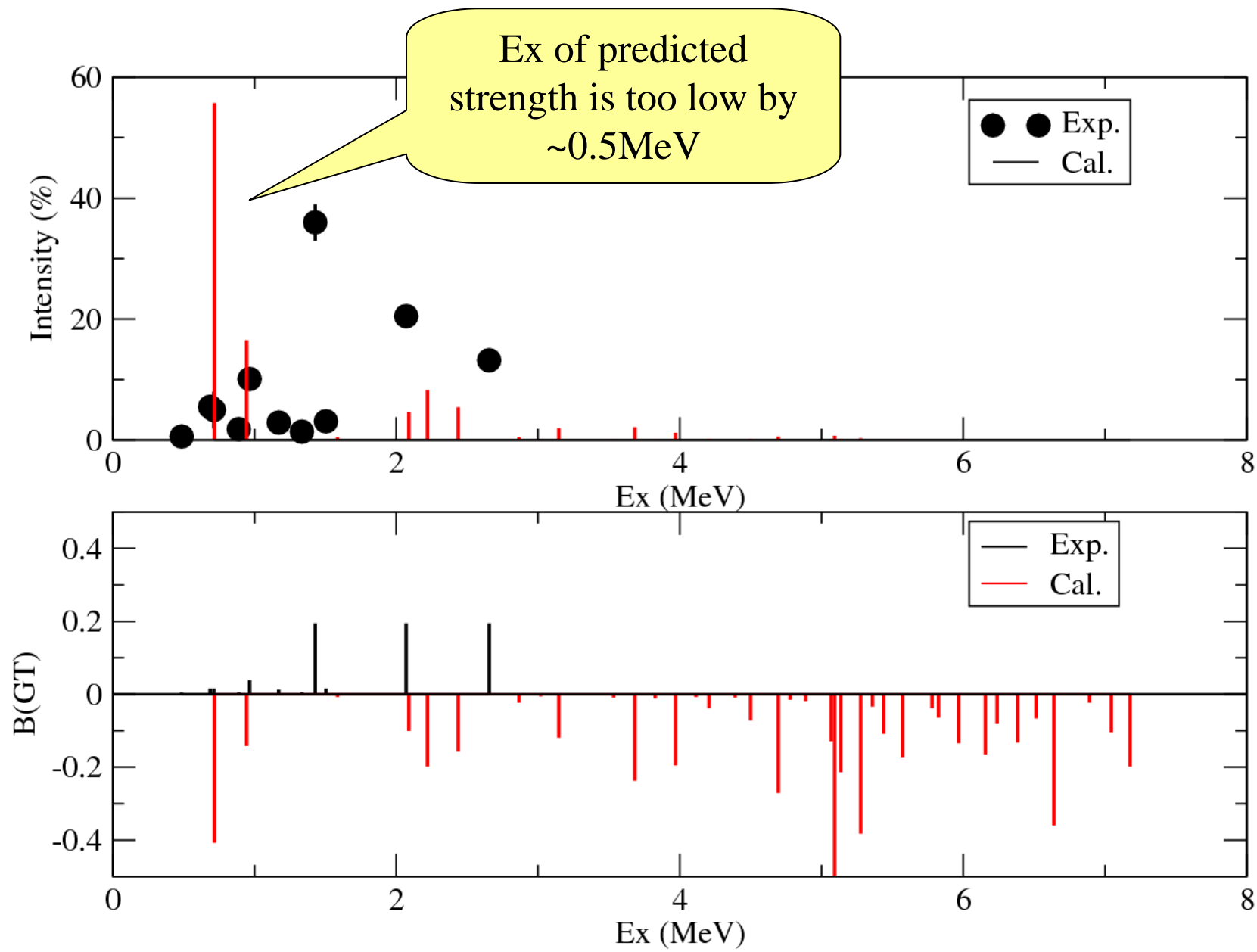
Shell model

$Q(\text{cal}) = 6.64\text{MeV}$

$T_{1/2}(\text{cal.}) = 0.22\text{ s}$



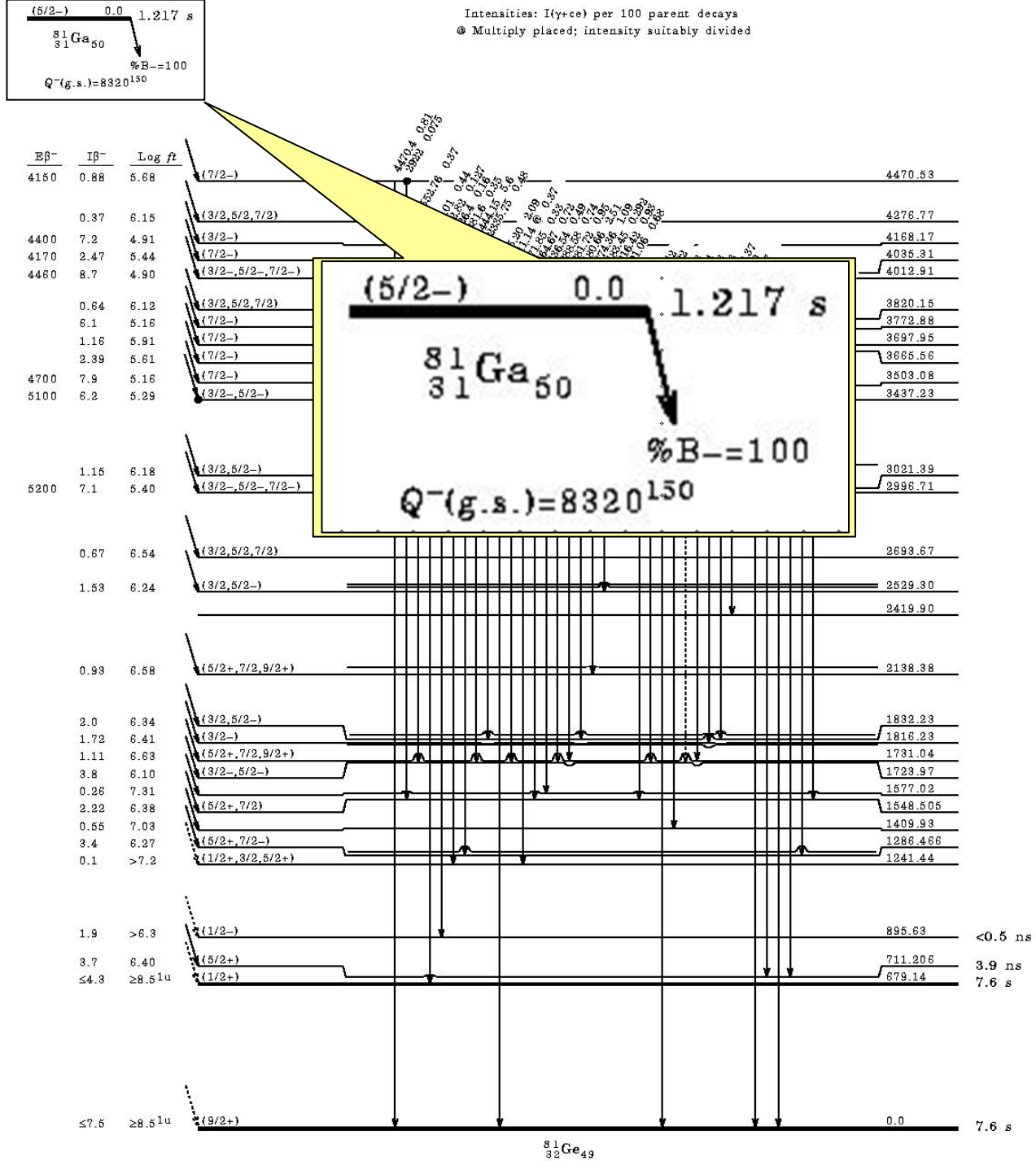
^{80}Zn



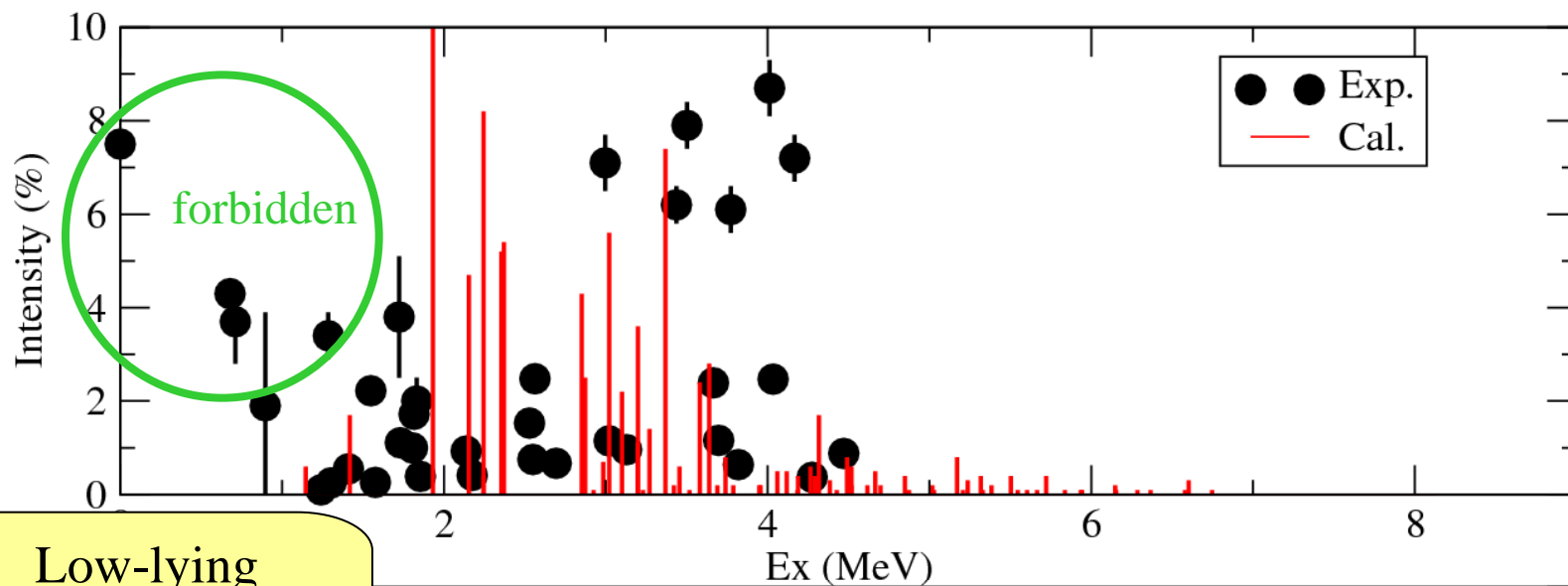
^{81}Ga

Shell model

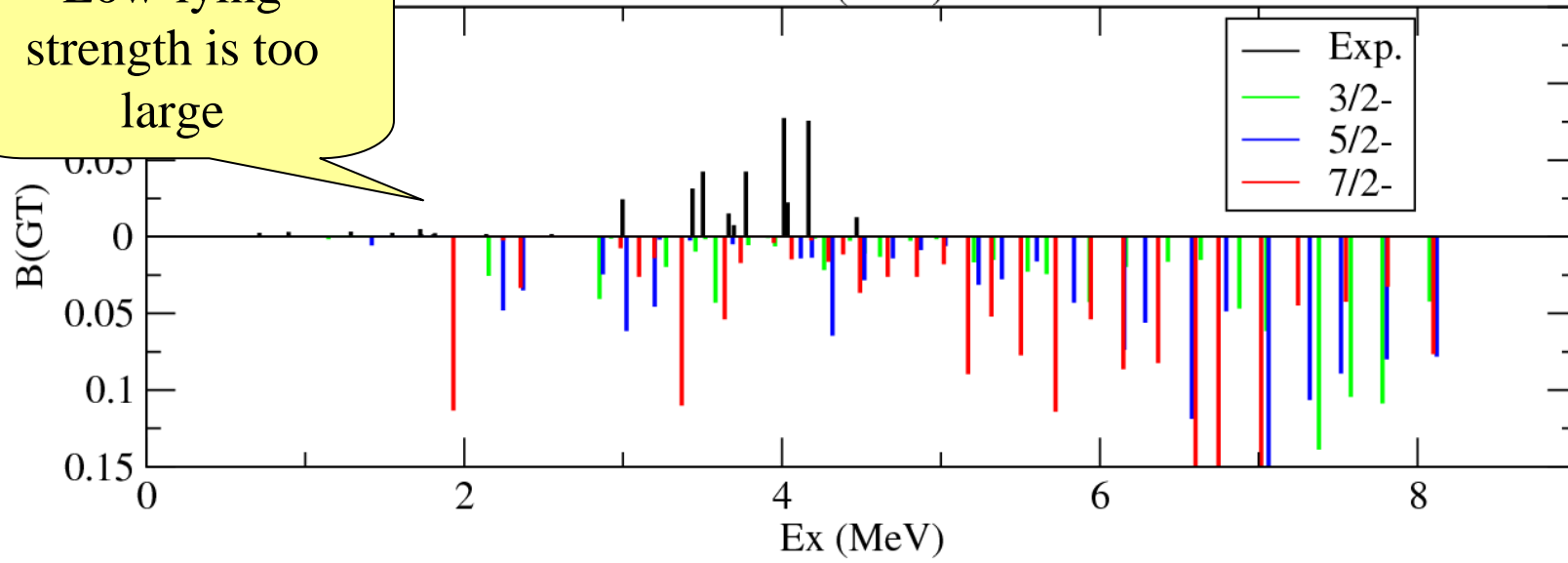
5/2- ground state
 $Q(\text{cal}) = 7.41\text{MeV}$
 $T_{1/2}(\text{cal.}) = 0.39\text{ s}$



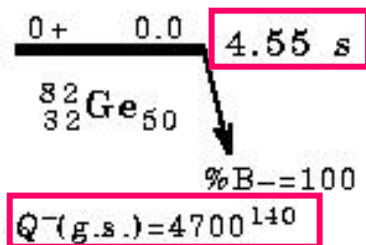
^{81}Ga



Low-lying strength is too large



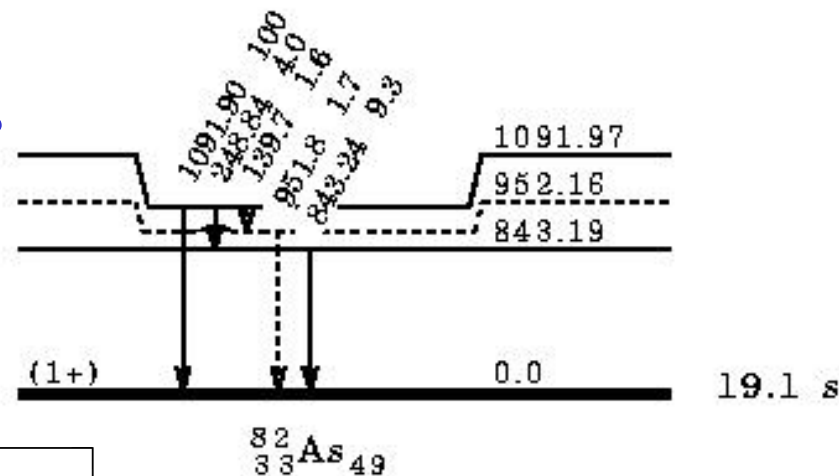
^{82}Ge



Intensities: relative I_y

1+ ?

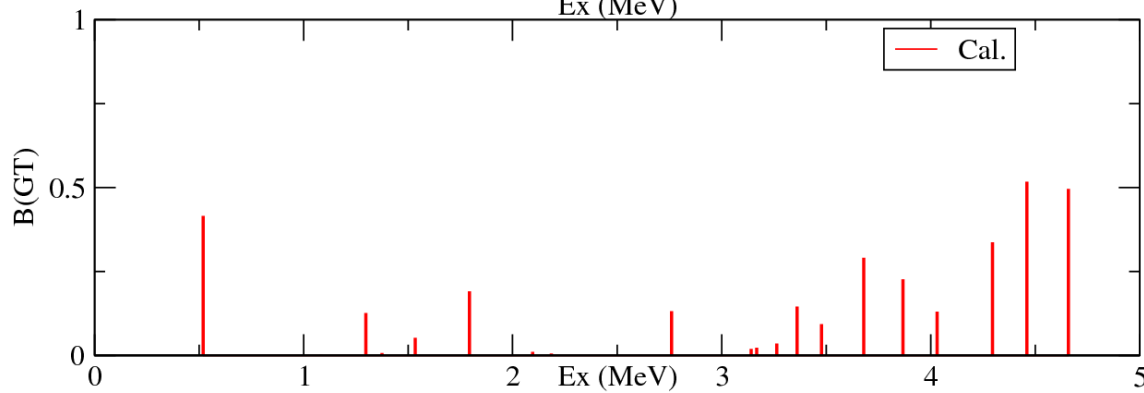
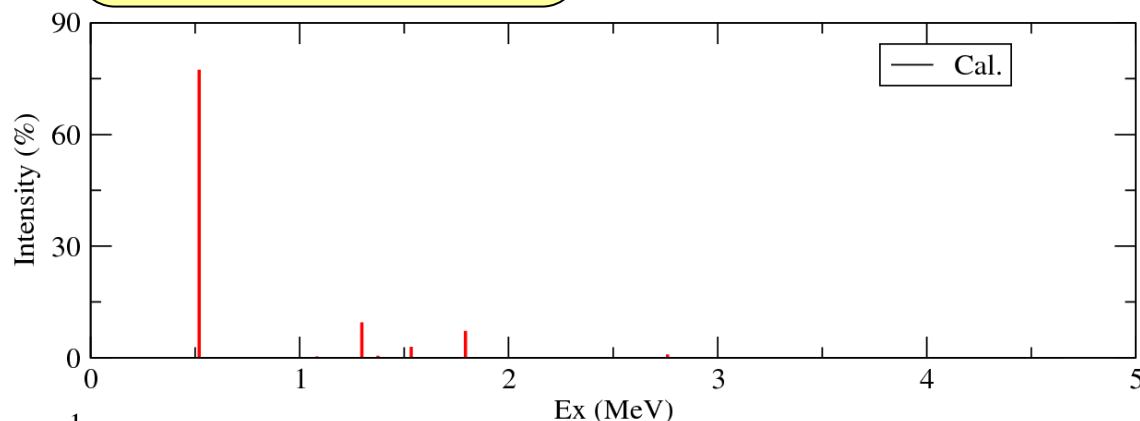
2- ?



Shell model

$Q(\text{cal}) = 4.14\text{MeV}$

$T_{1/2}(\text{cal.}) = 2.18\text{ s}$



Calculation predicts 2- or 5- ground state for ^{82}As .

Predicted 1+ state at 0.519 MeV almost explains the half life.

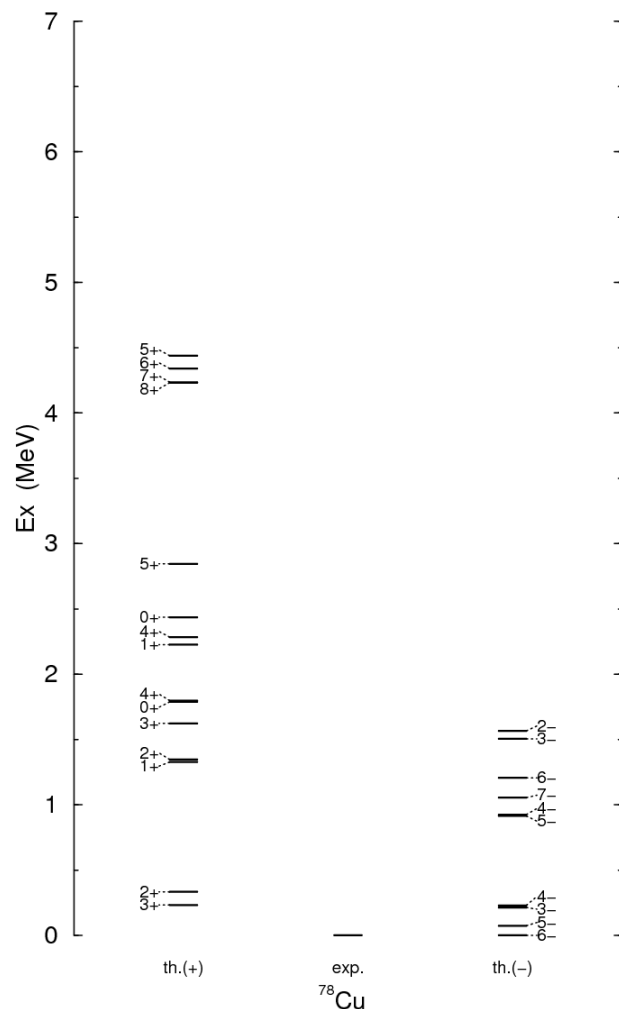
^{78}Cu

Shell model

6- ground state

$Q(\text{cal}) = 11.76\text{MeV}$

$T_{1/2}(\text{cal.}) = 0.074\text{ s}$



Adopted Levels

$Q_{\beta^-} = 13270\text{ SY}$ $S_n = 3540\text{ SY}$ $S_p = 14760\text{ SY}$ [1995Au04](#)

History

Type	Author	Cutoff Date	Comments
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Update J. Blachot	1-Jan-1997	CODING FROM NP A424, 1 (1997)	
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measured $T_{1/2}$ ([1987LuZX](#), [1997Au04](#)).

^{78}Cu levels

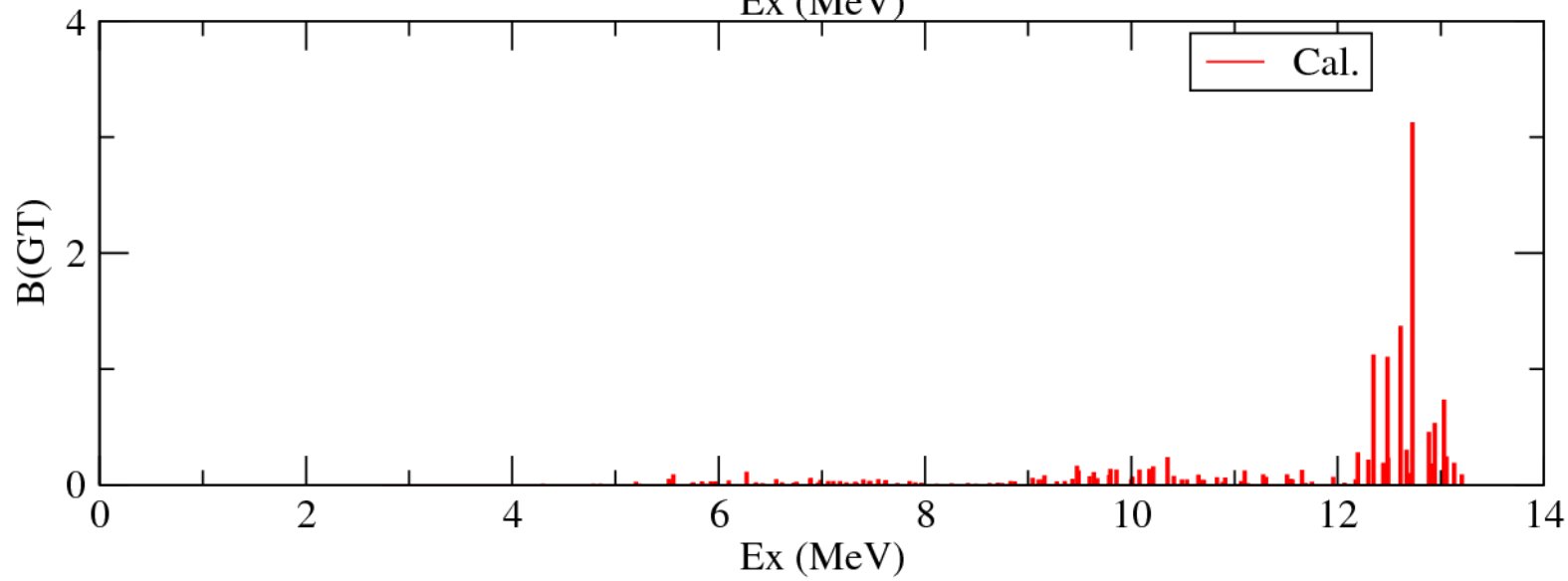
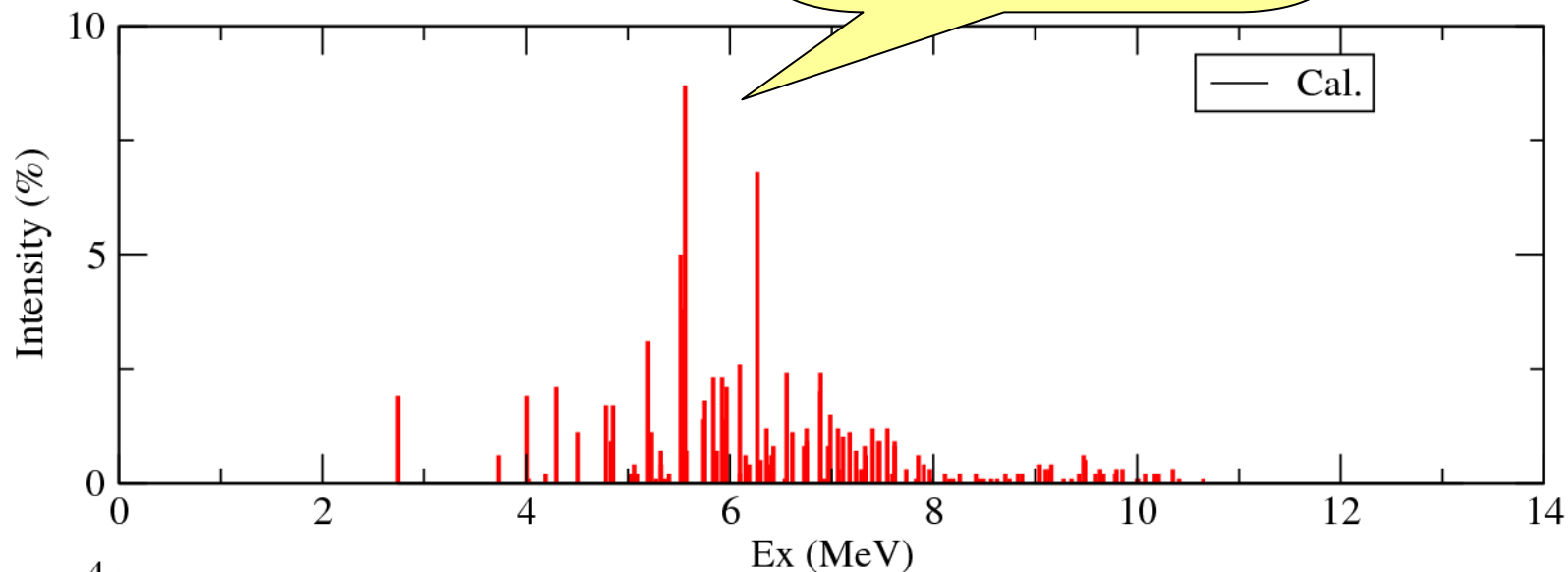
E_{level}	$T_{1/2}$	Comments
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0.0	342 ms 11	$\% \beta^- = 100$
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$T_{1/2}$: from [1987LuZX](#)

^{78}Cu

Beta strengths are
predicted at around
~6MeV



^{79}Zn

Shell model

9/2+ ground state

$Q(\text{cal}) = 8.49\text{MeV}$

$T_{1/2}(\text{cal.}) = 0.54\text{ s}$

(9/2+) 0.0 **0.995 s**

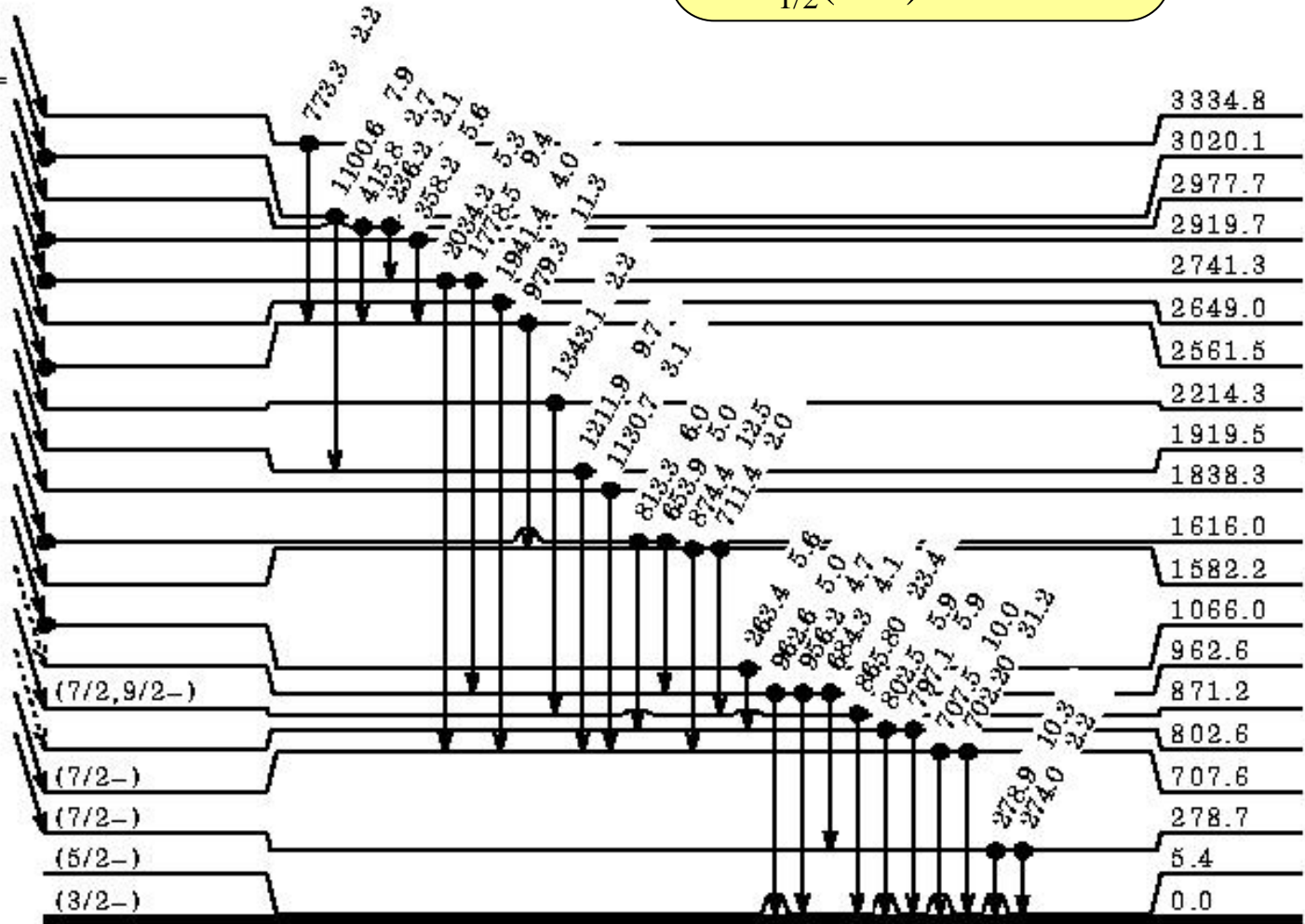
$^{79}\text{Zn}_{49}$

$Q^- = 9090\text{SY}$

%B=100

$I\pi^-$	Log ft
2.2	6.0
7.4	5.5
4.7	5.8
5.6	5.7
13	5.4
4.0	5.9
≤ 2	≥ 6.3
2.2	6.3
1.7	6.5
3.1	6.3
11	5.8
3.2	6.3
5.6	6.2
< 2	> 6.7
19	5.7
< 1.4	> 6.9
7	6.2
8.4	6.2

Forbidden ?

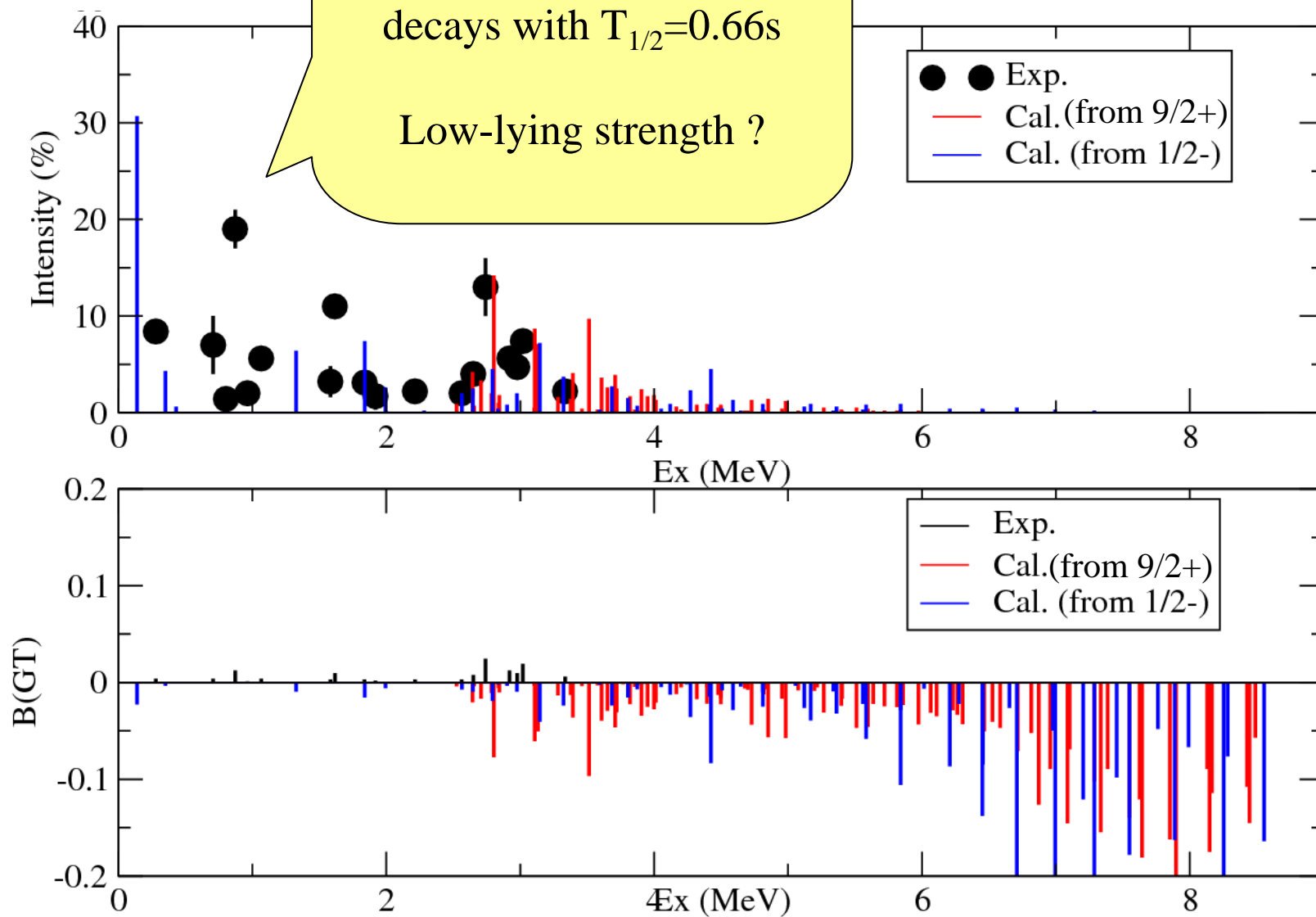


$^{79}\text{Ga}_{48}$

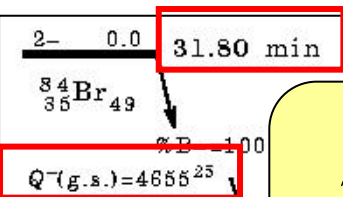
^{79}Zn

Shell model predicts 1/2-
isomer at 65keV which
decays with $T_{1/2}=0.66\text{s}$

Low-lying strength ?

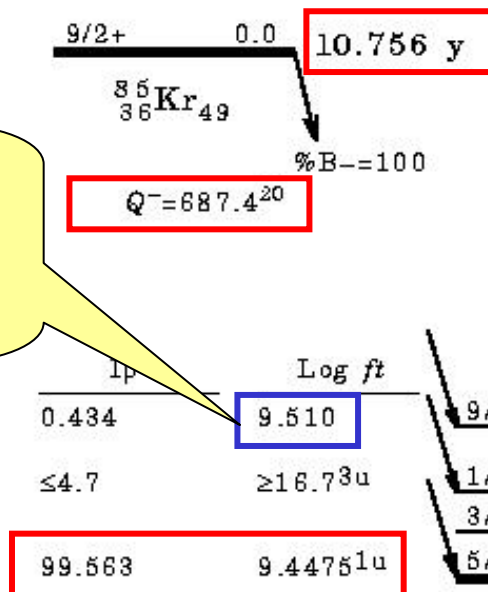
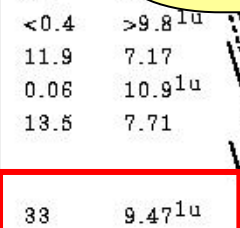


^{84}Br , ^{85}Kr , ^{86}Rb

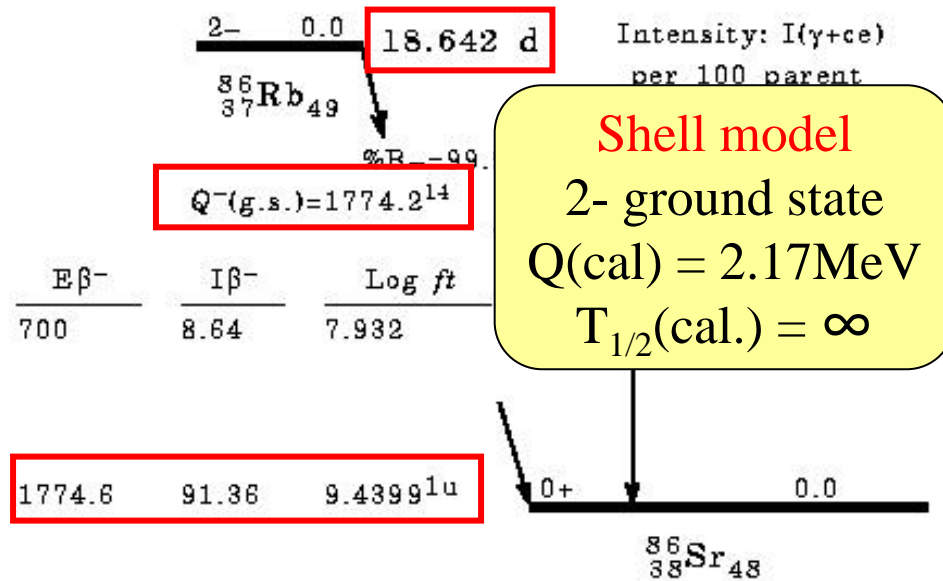


Shell model
 2- ground state
 $Q(\text{cal}) = 4.59 \text{ MeV}$
 $T_{1/2}(\text{cal.}) = 174 \text{ min}$

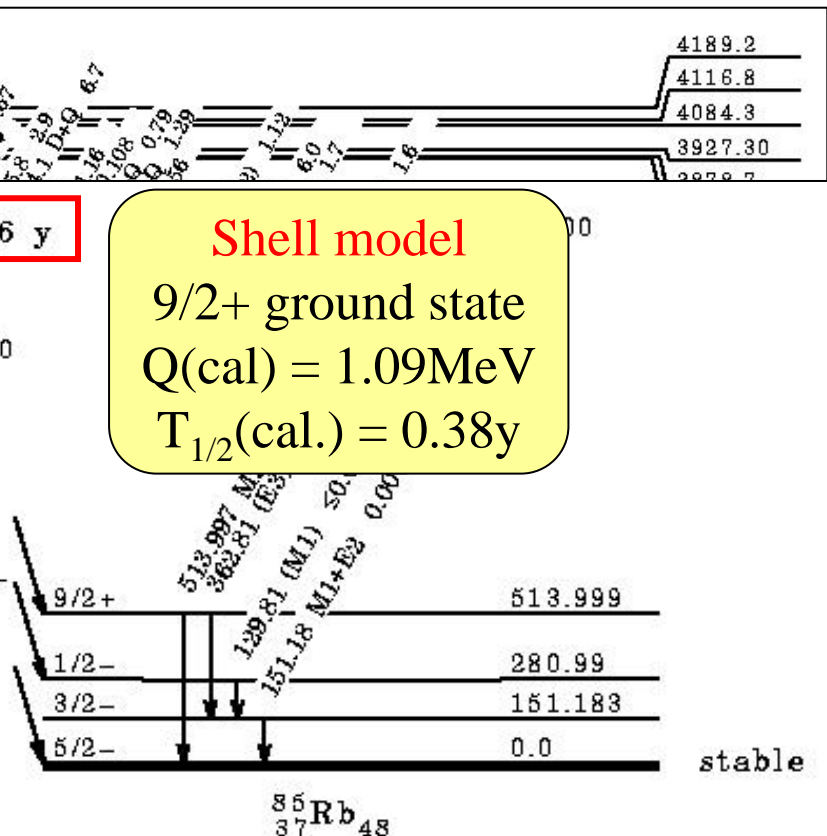
Significantly hindered
 $B(\text{GT})_{\text{exp}} = 0.0000019$
 $B(\text{GT})_{\text{cal}} = 0.00038$



Shell model
 9/2+ ground state
 $Q(\text{cal}) = 1.09 \text{ MeV}$
 $T_{1/2}(\text{cal.}) = 0.38 \text{ y}$



Shell model
 2- ground state
 $Q(\text{cal}) = 2.17 \text{ MeV}$
 $T_{1/2}(\text{cal.}) = \infty$



Unique forbidden decay

- ^{84}Br $(2-) \rightarrow 0+$ 33% $Q=4655(25)\text{keV}$
Exp. $\log f^{1\text{u}}t = 9.47$ $T_{1/2}=31.80(8)\text{min}$
Cal. 8.7 174min \rightarrow 14.2min
- ^{85}Kr $9/2+ \rightarrow 5/2-$ 99.563% $Q=687.4(20)\text{keV}$
Exp. $\log f^{1\text{u}}t = 9.4475$ $T_{1/2}=10.756(18)\text{y}$
Cal. 8.5 0.38y \rightarrow 0.28y
- ^{86}Rb $2- \rightarrow 0+$ 91.36% $Q=1774.2(14)\text{keV}$
Exp. $\log f^{1\text{u}}t = 9.4399$ $T_{1/2}=18.642(18)\text{d}$
Cal. 8.5 $\infty \rightarrow 2.0\text{d}$

Calculation underestimates the half-life

Summary

- Beta-decay properties for $N=50$ and 49 nuclei can be described reasonably well by the shell model
- f5pg9-shell with JUN45 interaction
- Quenching factor $q=0.6$ for Gamow-Teller operator
- Underestimate Q -values by ~ 1 MeV for neutron-rich nuclei
- Predicted half-life agrees with experimental data within a factor of 2 \sim 3 near the stable nuclei, while it systematically underestimates for neutron-rich nuclei probably because of insufficient correlations
- Allowed transitions already give reasonable predictions for $N=50$ cases, but forbidden decays sometimes play a crucial role for $N=49$ cases near the stable nuclei.
- Enlarged model space is desired for better description