

温度効果を考慮した
重い中性子過剰核の核分裂障壁
Thermal effects on the Fission Barrier
of
neutron-rich nuclei

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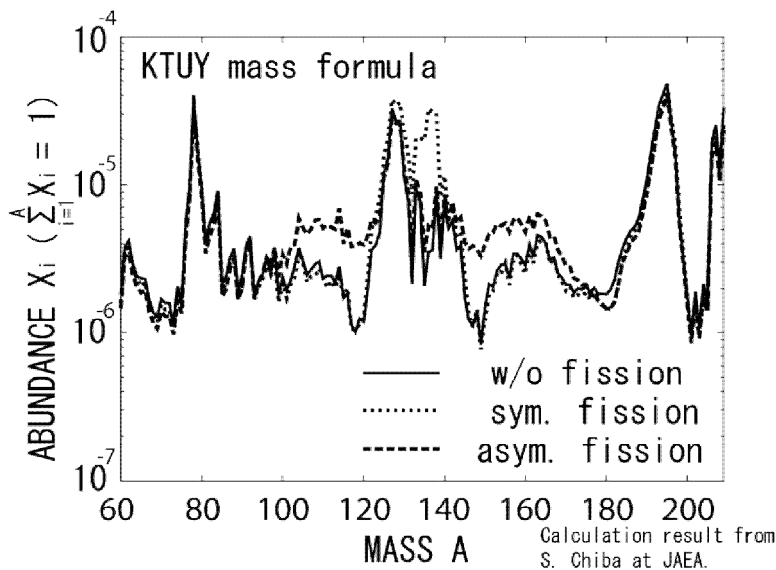
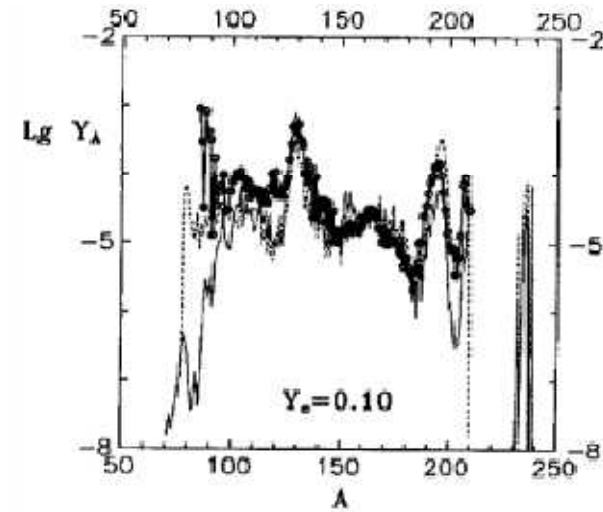
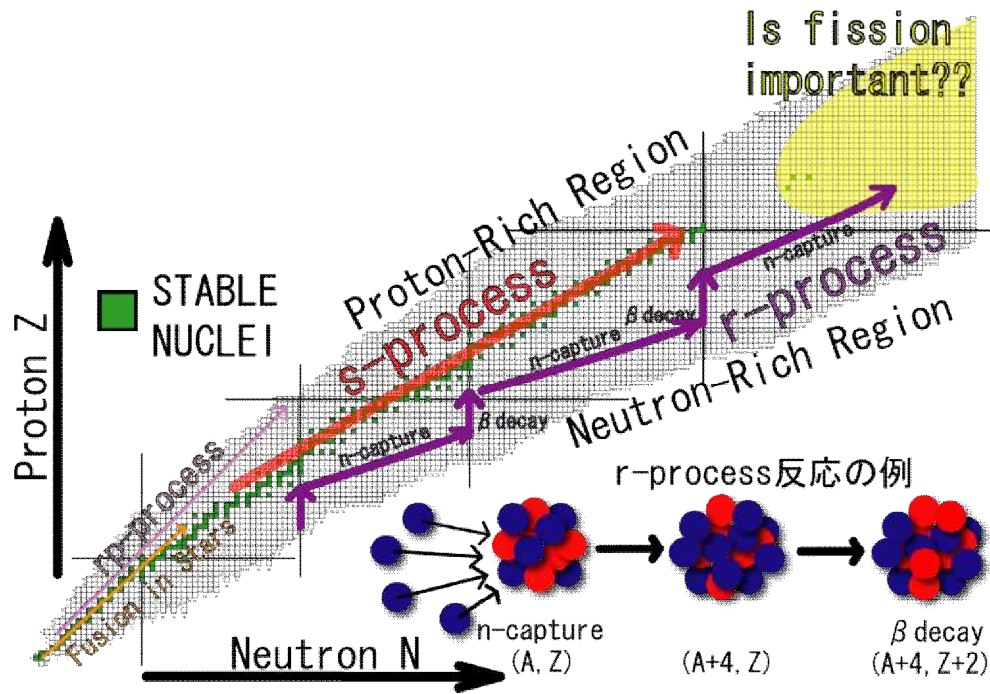
宇宙核物理連絡協議会主催 第1回研究戦略ワークショップ
[「核物理から見た宇宙-rプロセス研究の新時代に向かって」](#)

"New Era of Nuclear Physics in the Cosmos -the r-process nucleosynthesis"

日時: 2008年9月25,26日

場所: 理化学研究所 RIBF棟2階大会議室(201号室)

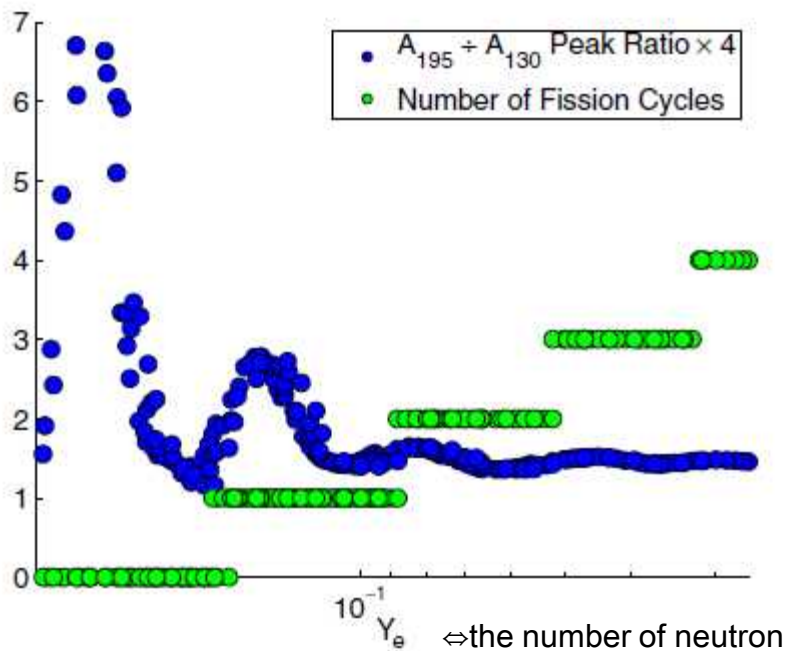
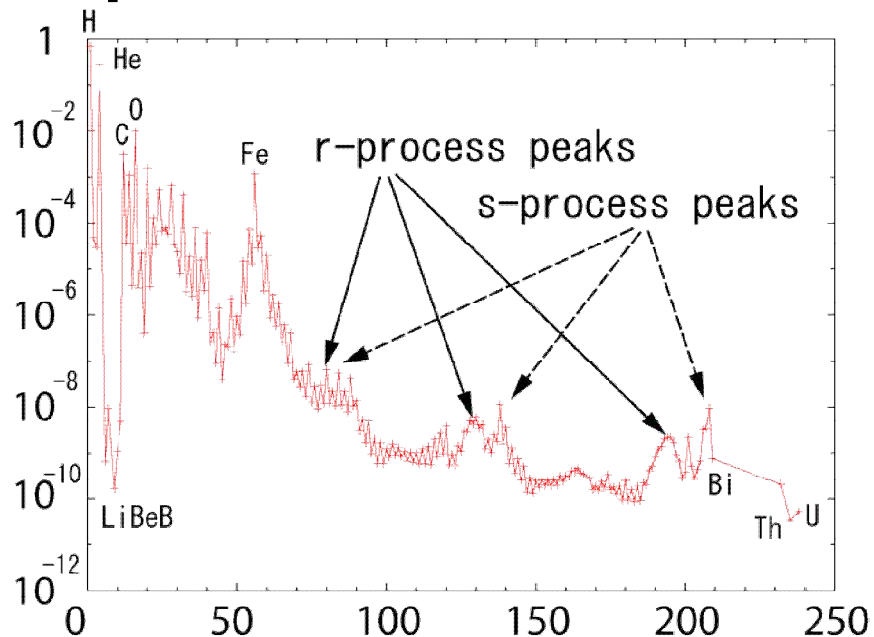
INTRODUCTION



Nucl. Phys. A688, 587, (2001),
I.V. Panov, C. Freiburghaus, F.-K. Thielemann

Circle : solar abundance.
Dashed and solid lines : r-abundance with different mass formulae.

r-process with fission



1. Cosmochronometer

^{238}U : half life = 4.468×10^9 year,
 ^{232}Th : half life = 1.405×10^{10} year.

2. Termination of r-process

3. Final dist. of the abundance

4. Fission cycling

(Peak ratio A_{195} / A_{130})

cf. J. Phys. G : Nucl. Part. Phys. 35 (2008) 014059,
 R. Surman, J. Beun, G. C. McLaughlin, S Kane, and W. R. Hix

Ambiguities in Fission!!

1. Fission Fragments
2. Fission Barrier height
3. Curvature

MOTIVATION

clarify the role of fission in the r-process

- **Fission barriers in r-process have been derived with**
 1. ETFSI (Extended Thomas Fermi Strutinsky Integral)
 2. Macroscopic-microscopic model
- **Heavy Neutron-rich nuclei (exotic nuclei)**
- **Assumption of Nuclei at Zero temperature
(Ground-state)**

cf. "Effect of beta-decay of excited-state nuclei"
J. Phys. G : Nucl. Part. Phys. **35** (2008) 025203
M. A. Famiano, R. N. Boyd, T. Kajino, K. Otsuki, M. Terasawa, and G. J. Mathews

Influence on abundance of heavy nuclei $A > 160$

Our approach

◆ Microscopic Model

Skyrme-Hartree-Fock-Bogoliubov Method (HFB) (symmetric fission)

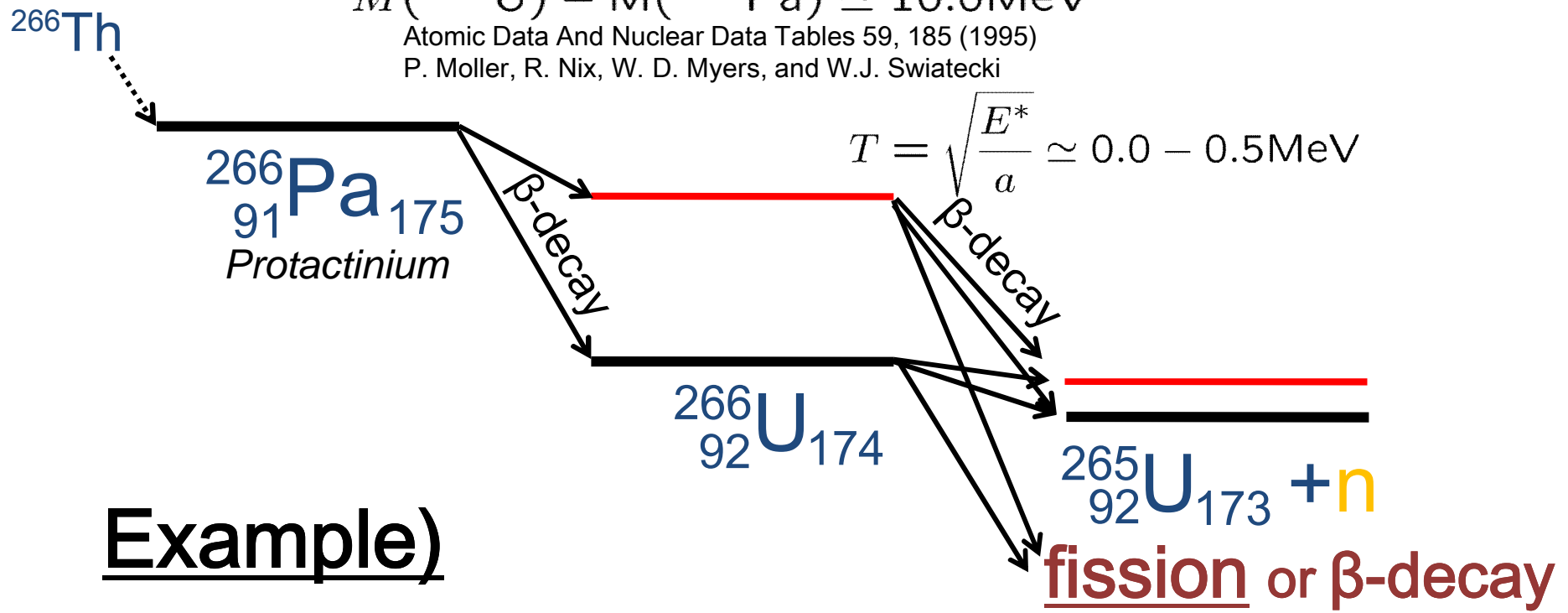
◆ Thermal effect (Highly Excited State)

SCENARIO

$$M(^{266}\text{U}) - M(^{266}\text{Pa}) \simeq 10.0\text{MeV}$$

Atomic Data And Nuclear Data Tables 59, 185 (1995)

P. Moller, R. Nix, W. D. Myers, and W.J. Swiatecki



Example)

The effect of highly excited nucleus

- Pairing phase-transition ($T_{\text{crit}} \simeq \frac{\Delta_0}{1.76}$) "Quantum Theory of Many Particle Systems" Fetter and Walecka, p. 448

$$\Delta_0 \simeq 0.6 \text{ MeV} \quad (T_{\text{crit}} = 0.34 \text{ MeV}) \text{ for } ^{236}\text{U}$$

Influence → Fission Barrier Height,

Beta-decay, and Neutron Emission

Finite Temperature HFB (THFB)

Nucl. Phys. A352 (1981) 30,
 Alan L. Goodman
 Phys. Rev. C70, 025801(2004)
 Nicolae Sandulescu

Density operator

$$D_{\text{HFB}} = \frac{\exp(-\beta \sum_i E_i \hat{n}_i)}{Z_{\text{HFB}}} = \prod_i (f_i \hat{n}_i + (1 - f_i)(1 - \hat{n}_i))$$

■ Grand Partition Function Z : $Z_{\text{HFB}} = \text{Tr} \left(\exp(-\beta \sum_i E_i \hat{n}_i) \right) = \prod_i (1 + e^{-\beta E_i})$

■ Expectation value of an operator O : $O = \langle \hat{O} \rangle = \text{Tr}(\hat{O} D)$

$$\rho = U^T f U^* + V^\dagger (1 - f) V \quad \kappa = U^T f V^* + V^\dagger (1 - f) U$$

■ Quasi Particle Occupation probability : $f_k = \frac{1}{1 + \exp(\beta E_k)}$

THFB equation from variation $\Omega = E - TS$ E_k : Quasi Particle Energy

$$\begin{pmatrix} h(T) - \lambda & \Delta(T) \\ -\Delta^*(T) & -h^*(T) + \lambda \end{pmatrix} \begin{pmatrix} U_k \\ V_k \end{pmatrix} = E_k \begin{pmatrix} U_k \\ V_k \end{pmatrix}$$

$$h_{ij}(T) = K_{ij}(T) + \sum_{kl} v_{ikjl} \rho_{lk}(T) \quad \Delta_{ij}(T) = \frac{1}{2} \sum_{kl} v_{ijkl} \kappa_{kl}(T)$$

CONSTRAINT THFB

■ Constraining operator

$$Q = \sum_i r_i^2 Y_{20}(\hat{r}_i) \quad \beta = \sqrt{\frac{5}{16\pi} \frac{4\pi}{3AR_0^2}} \langle Q \rangle$$

■ Total Energy and Average Gap Parameter

Average gap del ta: $\langle \Delta \rangle = \int dr \kappa_T(r) \Delta_T(r) / \int dr \kappa_T(r)$

Total Energy: $E = \langle H \rangle$

Barri er hei ght: $V_f = E_{max} - E_{min}$

Parameter set : SLy4

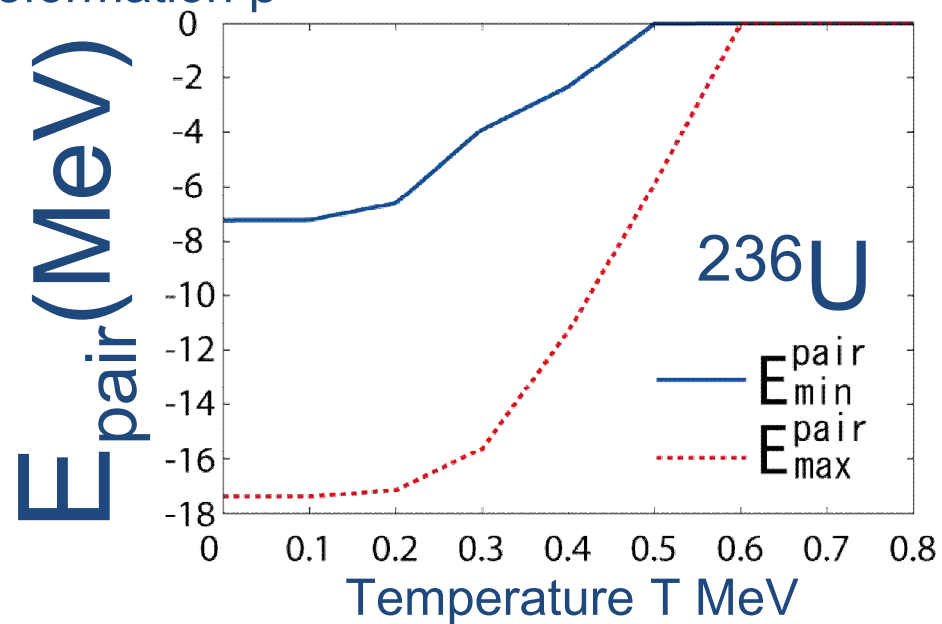
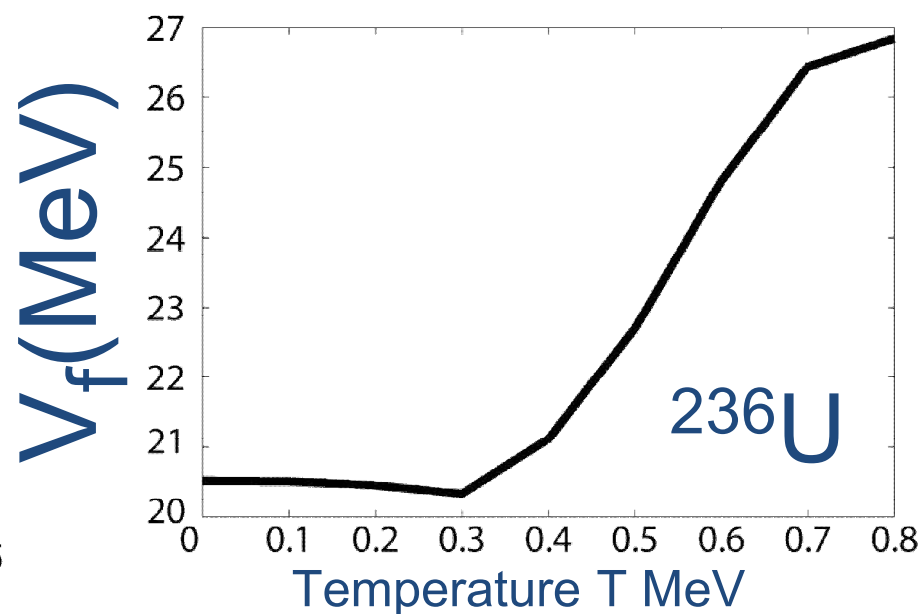
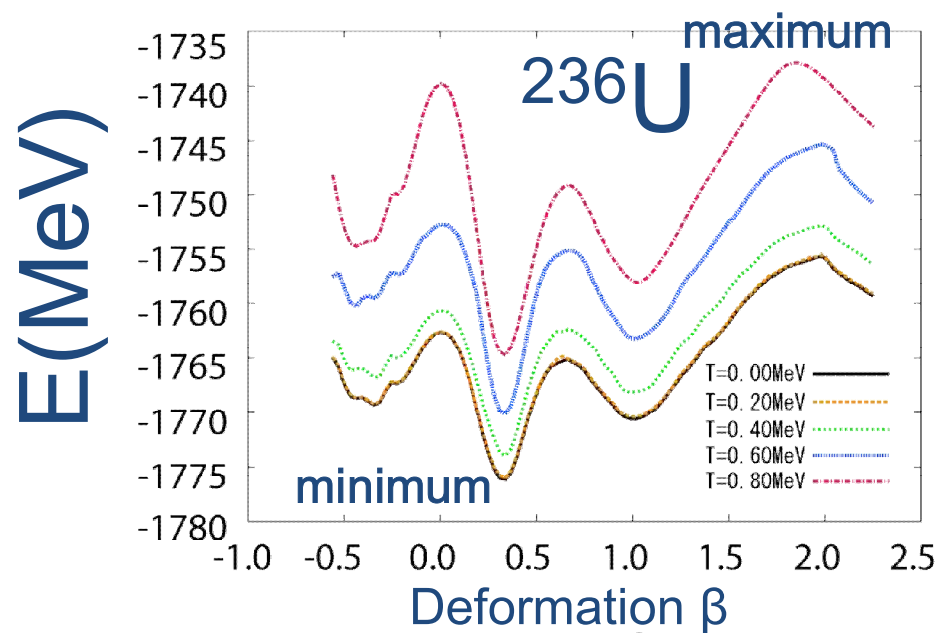
Volume pairing : $V_{vol}^p = 295.369$,

$V_{vol}^n = 286.669$ MeV,

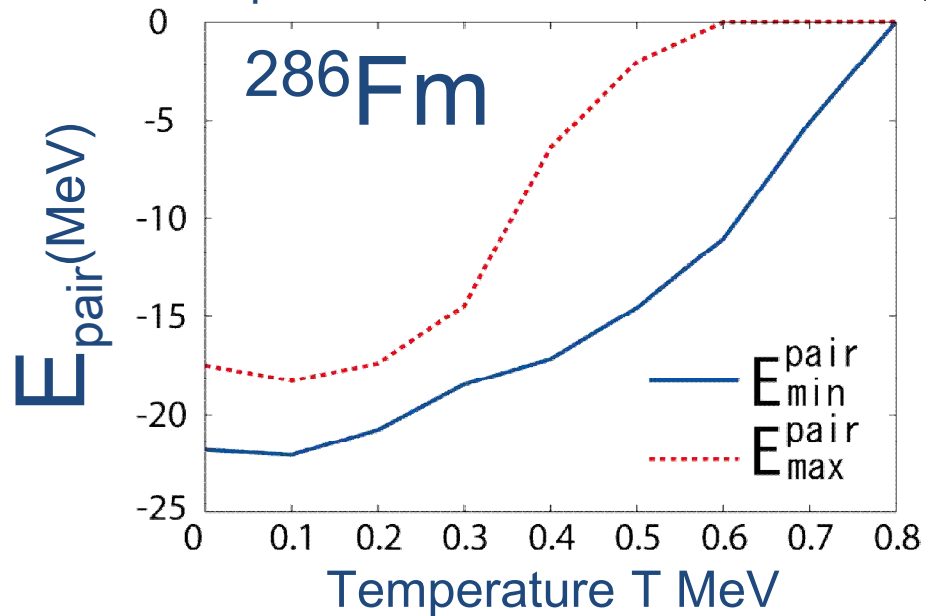
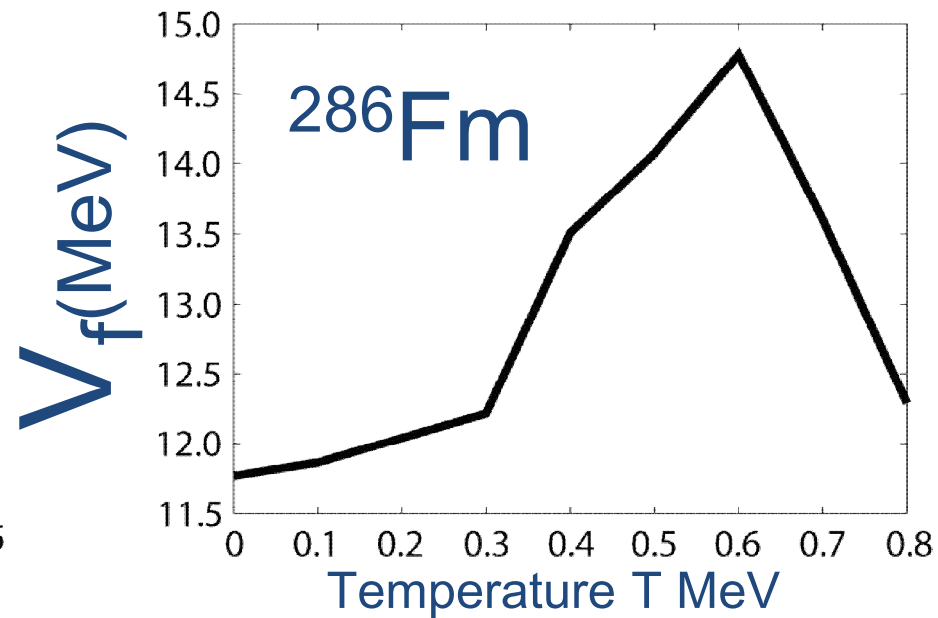
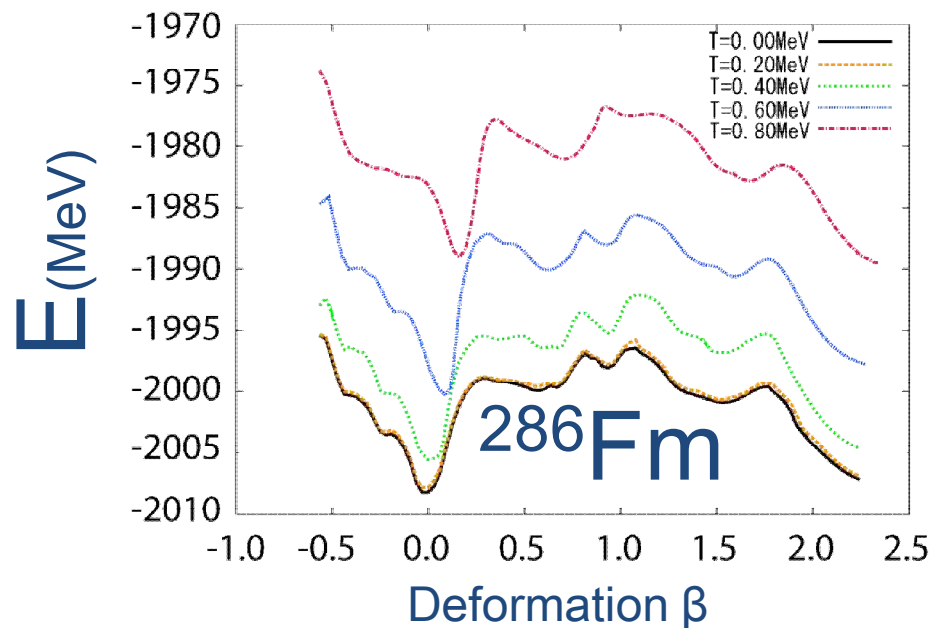
determined so as to reproduce

the gap parameter $\Delta(^{236}\text{U})$

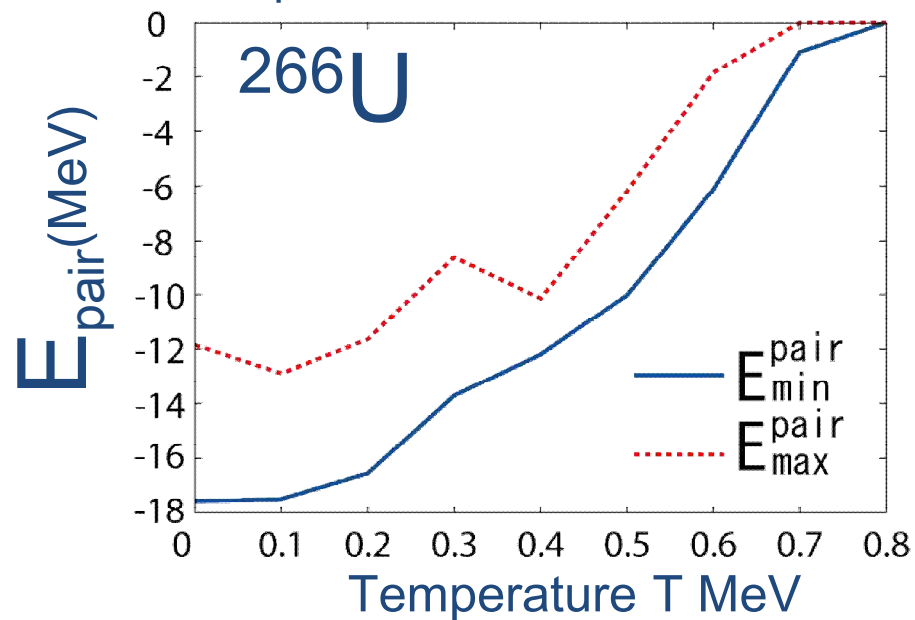
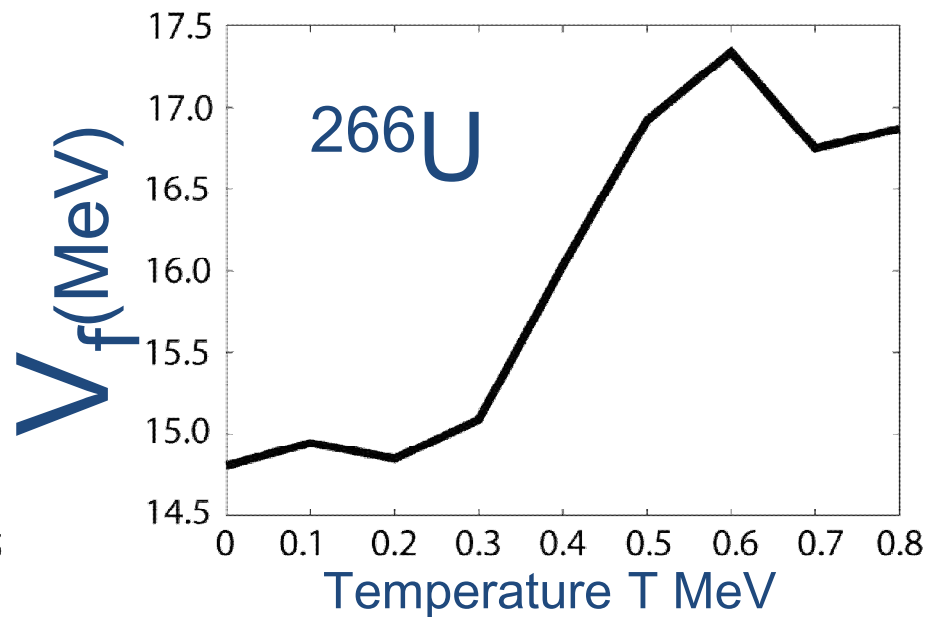
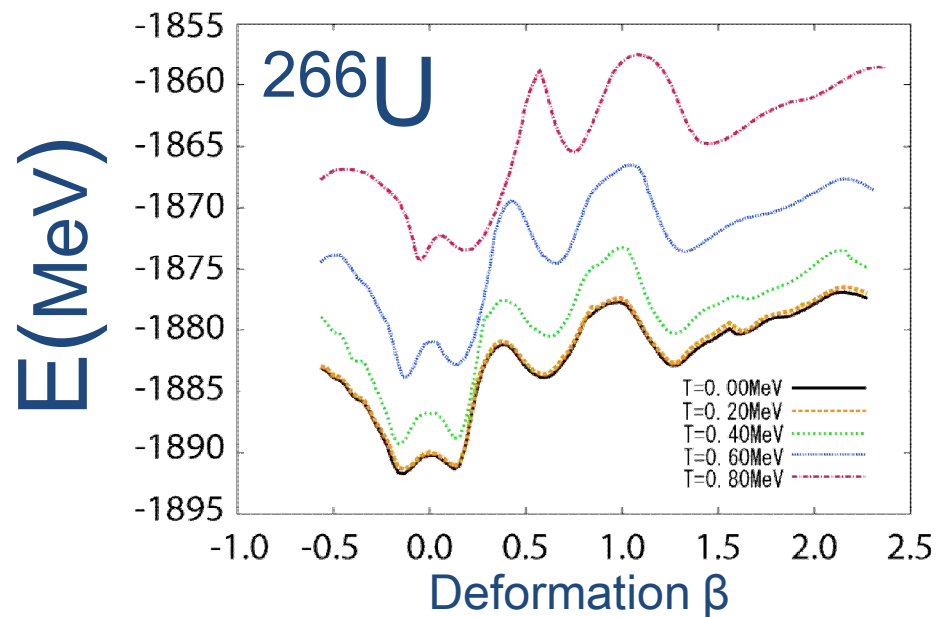
RESULT $\sim 236\text{U}$



RESULT $\sim 286\text{Fm}$



RESULT $\sim 266\text{U}$



RESULT ~ Decay Rate

- Decay width derived from Bohr-Wheeler formula

$$\Gamma_f = \frac{\omega_0}{2\pi} e^{-V_b/T}$$

E^* is approximated with

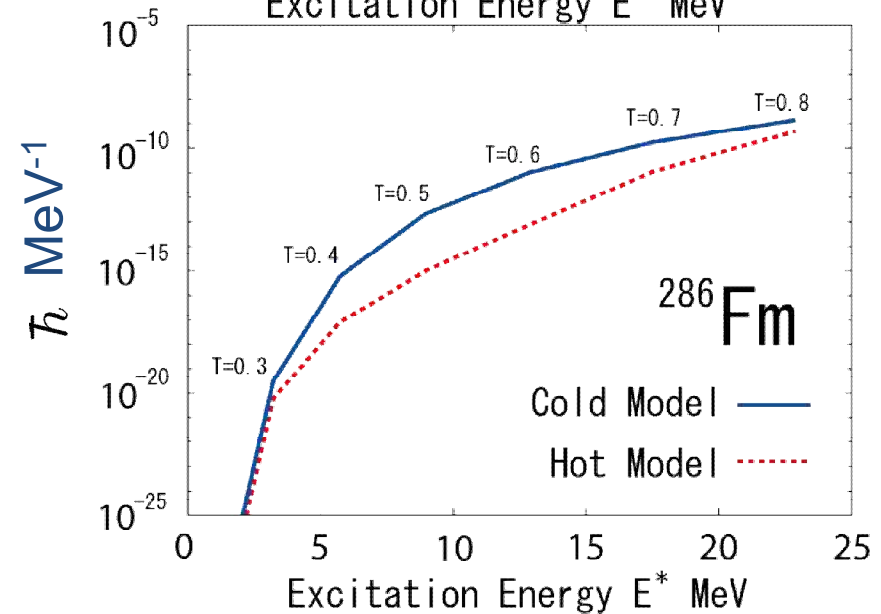
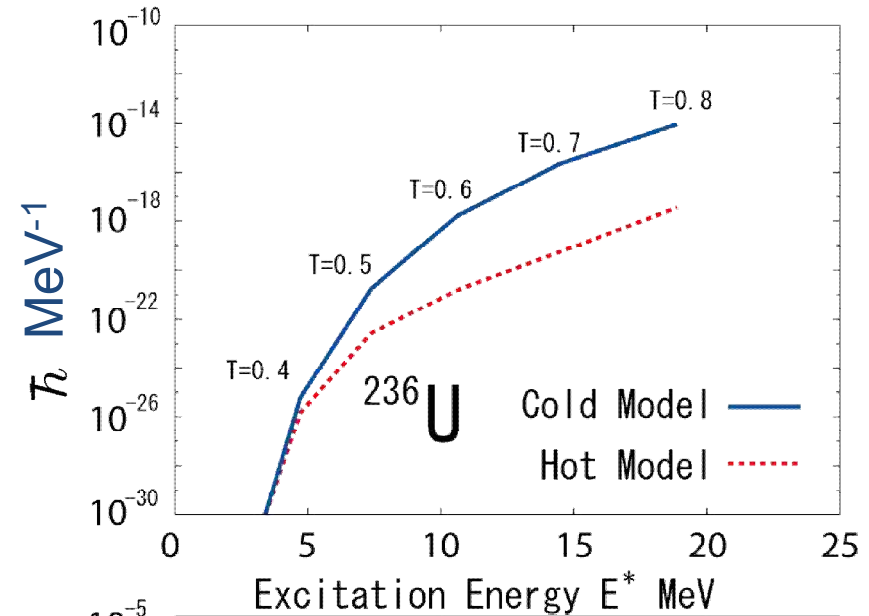
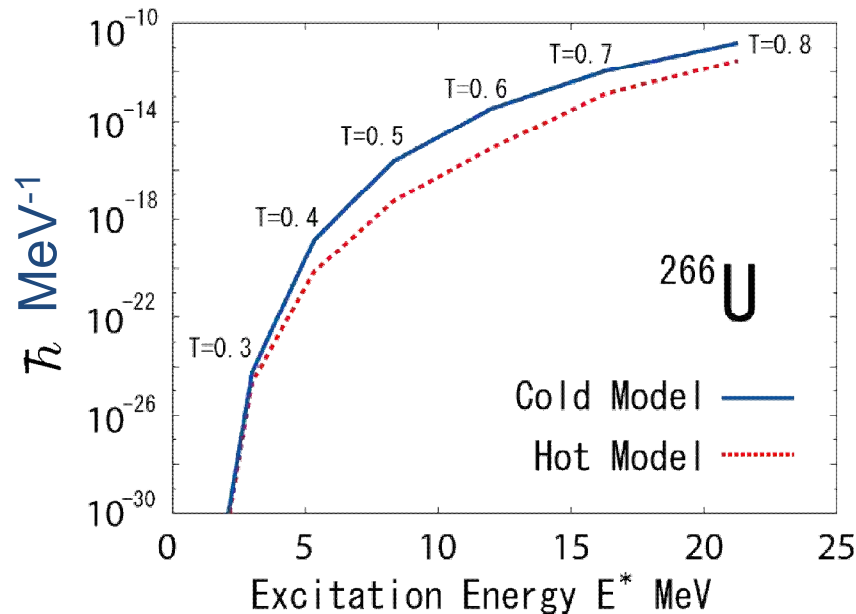
$$E^* = aT^2 \simeq \frac{A}{8}T^2$$

$$M(^{236}\text{U}) - M(^{236}\text{Pa}) \simeq 3.0\text{MeV}$$

$$M(^{266}\text{U}) - M(^{266}\text{Pa}) \simeq 10.0\text{MeV}$$

$$M(^{286}\text{Fm}) - M(^{286}\text{Es}) \simeq 11.1\text{MeV}$$

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SUMMARY

Calculate the Fission Barrier Height for neutron-rich nuclei with THFB method

■ Pairing Phase transit. leads to important effects on V_f .
depend on pairing energy at g.s and saddle point.

■ Fission rates of hot neutron-rich nuclei are suppressed.

■ Fission of hot nuclei is important ?

excitation energy of r-process nuclei(neutron-rich nuclei),
mass, pairing strength, level density parameter

---- *Future work* ----

- ◆ Asymmetric fission for the practical r-process application
- ◆ Correction of Temperature at Saddle point