

Complex Scaling Method : Principles and Applications

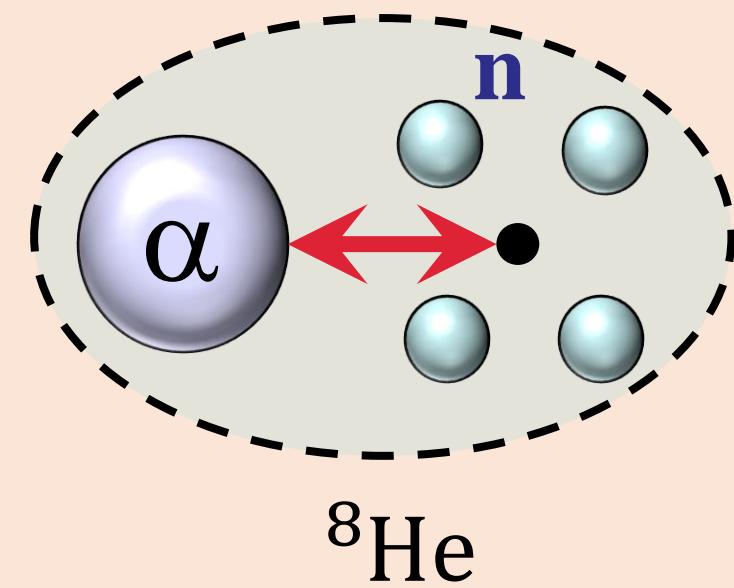
Myo, Takayuki
明 孝之



大阪工業大学

Outline

- Physics of light unstable nuclei
- Resonances and non-resonant continuum states
- Complex scaling Method (CSM)
複素（座標）スケーリング法
- Nuclear model : Cluster model
- Basic properties of CSM
- Many-body resonances
- Level density, Green's function
- Strength functions, breakup reactions



1. Resonance spectroscopy with CSM
2. Scattering states with CSM
 - Complex-scaled Green's function

Continuum Level Density (CLD) in complex scaling

$$\Delta(E) = \rho(E) - \rho_0(E)$$

$$= -\frac{1}{\pi} \text{Im} \left[\text{Tr} [G(E) - G_0(E)] \right]$$

$$\Delta(E) = \frac{1}{2\pi i} \text{Tr} \left[S(E)^\dagger \frac{d}{dE} S(E) \right]$$

S. Shlomo, NPA539('92)17

A. Kruppa, PLB431('98) 237, K. Arai and A. Kruppa, PRC60('99)064315

R. Suzuki, T. Myo and K. Kato, PTP113('05)1273.

[R.D. Levine](#), Quantum Mechanics of Molecular Rate Processes, Oxford, 1969.

Complex Scaling

$$\Delta(E) = -\frac{1}{\pi} \text{Im} \left[\text{Tr} [G^\theta(E) - G_0^\theta(E)] \right]$$

Level density

$$\rho(E) = \sum_n \delta(E - E_n), \quad G_{(0)}(E) = \frac{1}{E - H_{(0)}}$$

Single channel case

$$\Delta(E) = \frac{1}{\pi} \frac{d\delta_\ell}{dE}$$

$$\frac{1}{x - x_0} = \text{P} \left(\frac{1}{x - x_0} \right) - i\pi \cdot \delta(x - x_0)$$

$$G^\theta(E) = \frac{1}{E - H^\theta} \text{ with interaction}$$

$$G_0^\theta(E) = \frac{1}{E - H_0^\theta} \text{ Asymptotic}$$

Continuum Level Density (CLD) in CSM using discretized continuum states

$$\Delta(E) = \rho(E) - \rho_0(E)$$

$$\rho_{(0)}^\theta = U(\theta)\rho_{(0)}U^{-1}(\theta)$$

$$\begin{aligned}\rho^\theta &\cong \rho_N^\theta = -\frac{1}{\pi} \text{Im} \left[\sum_n^N \frac{1}{E - E_n(\theta)} \right] \\ &= \frac{1}{\pi} \sum_n^N \frac{E_I}{(E - E_R)^2 + E_I^2}\end{aligned}$$

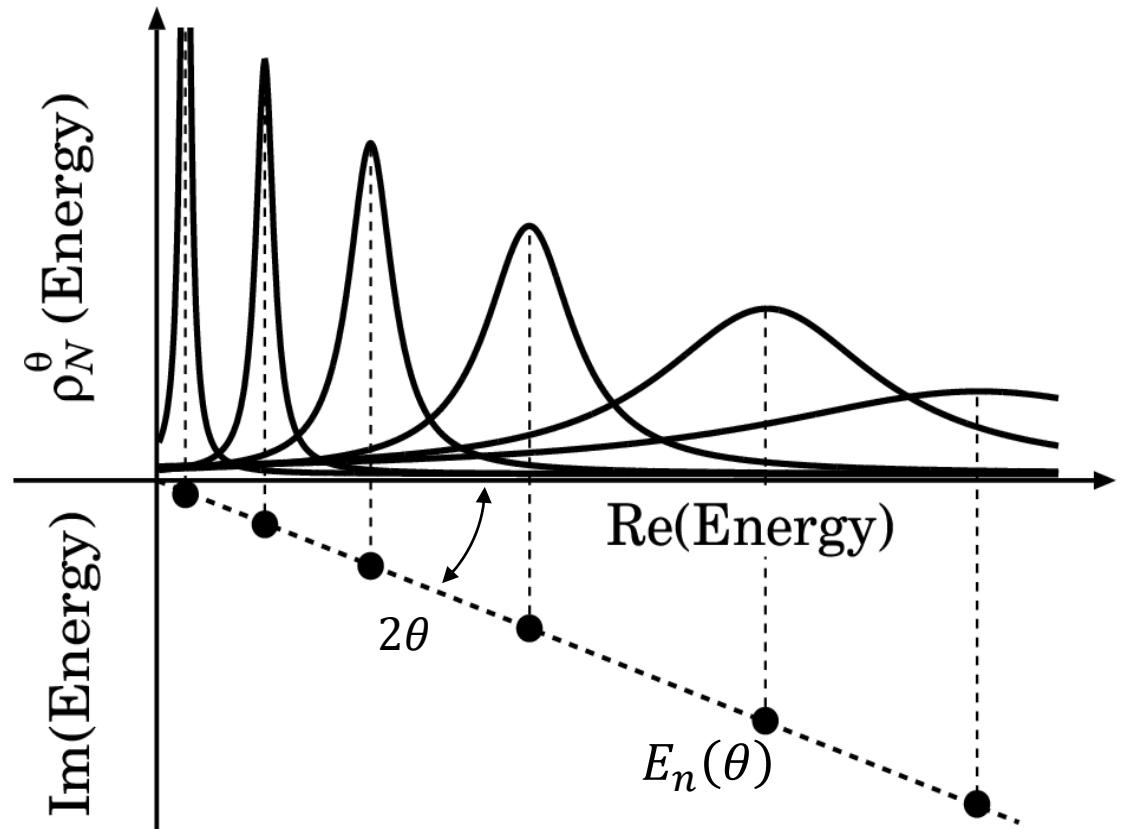
Number of
discretized states

Level density

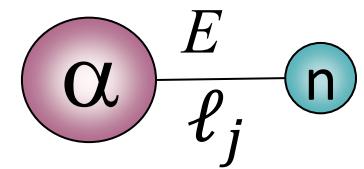
$$\rho(E) = \sum_n \delta(E - E_n)$$

Complex-scaled
energy eigenvalue

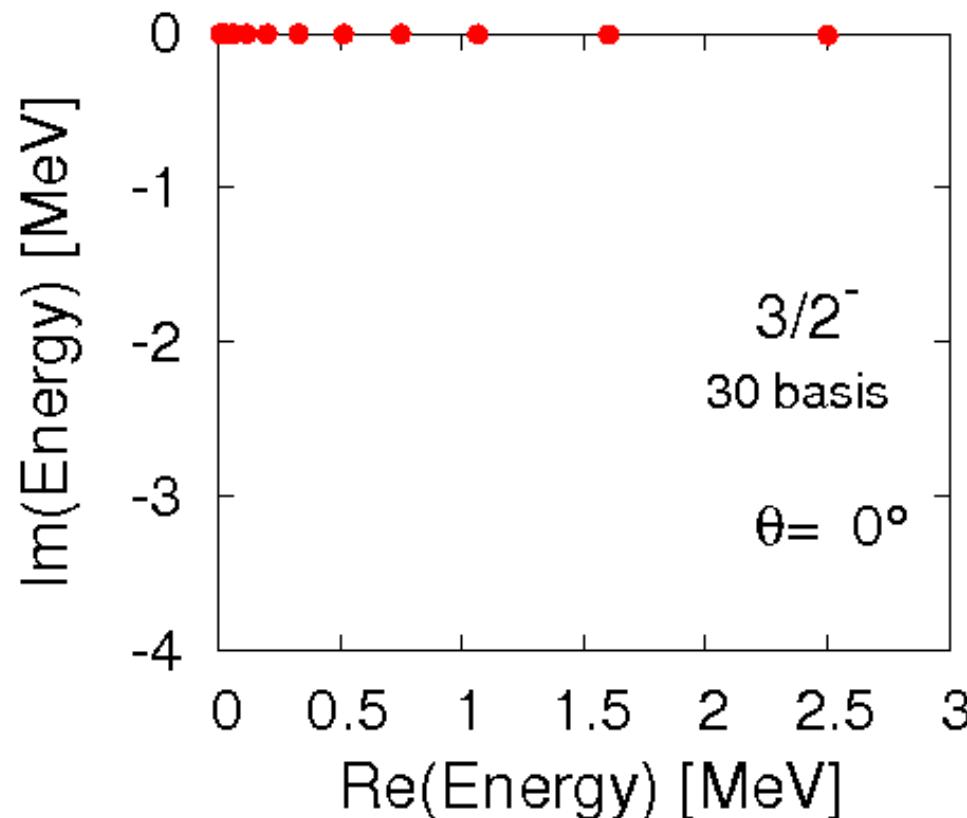
$$E_n(\theta) = E_R - iE_I$$



$\alpha+n$ scattering using discretized states

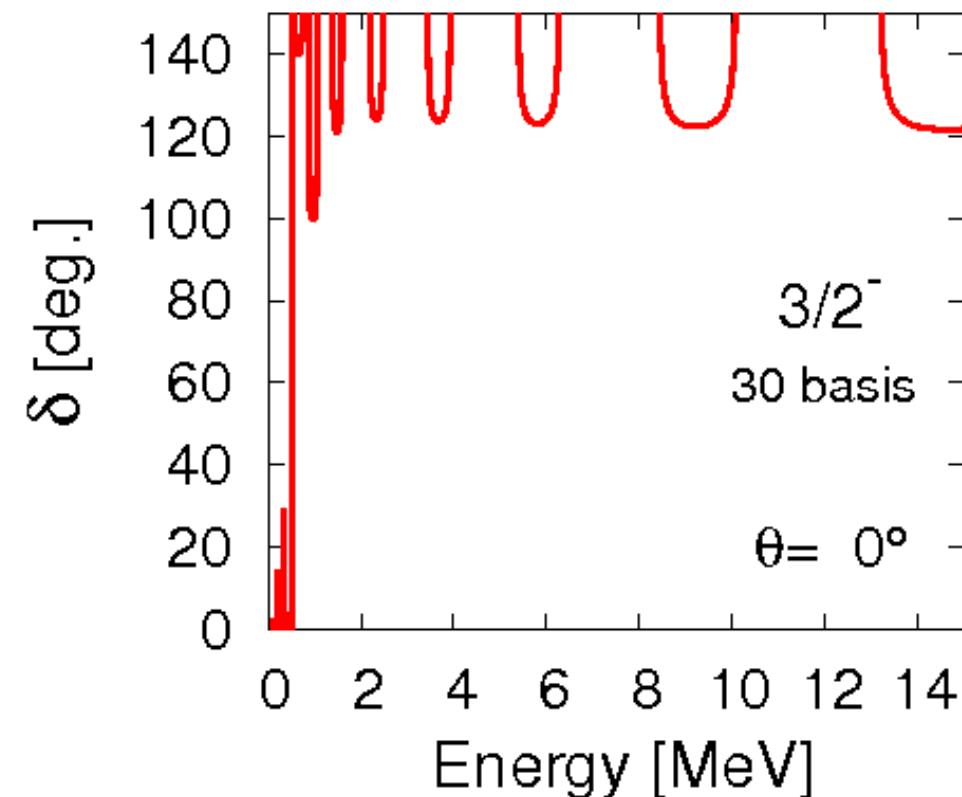


energy eigenvalues



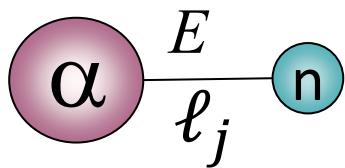
30 Gaussian basis functions

P-wave scattering phase shift



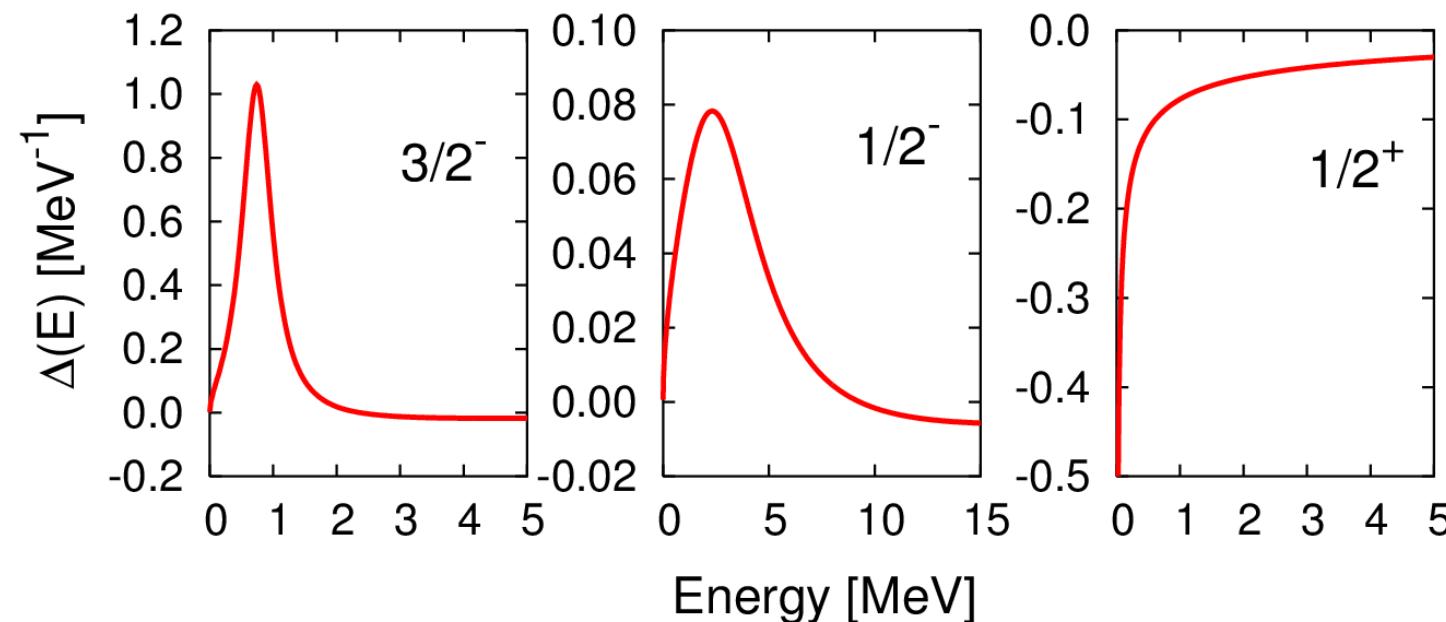
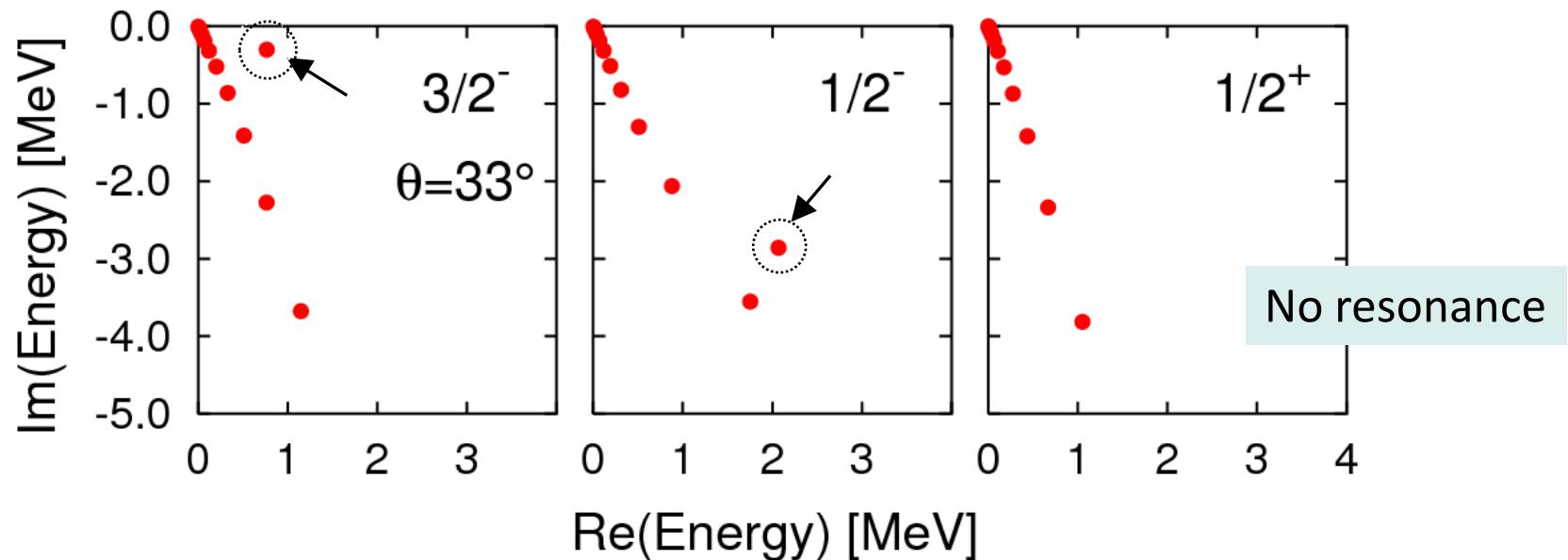
$\alpha+n$ scattering with discretized continuum states

Energy eigenvalues



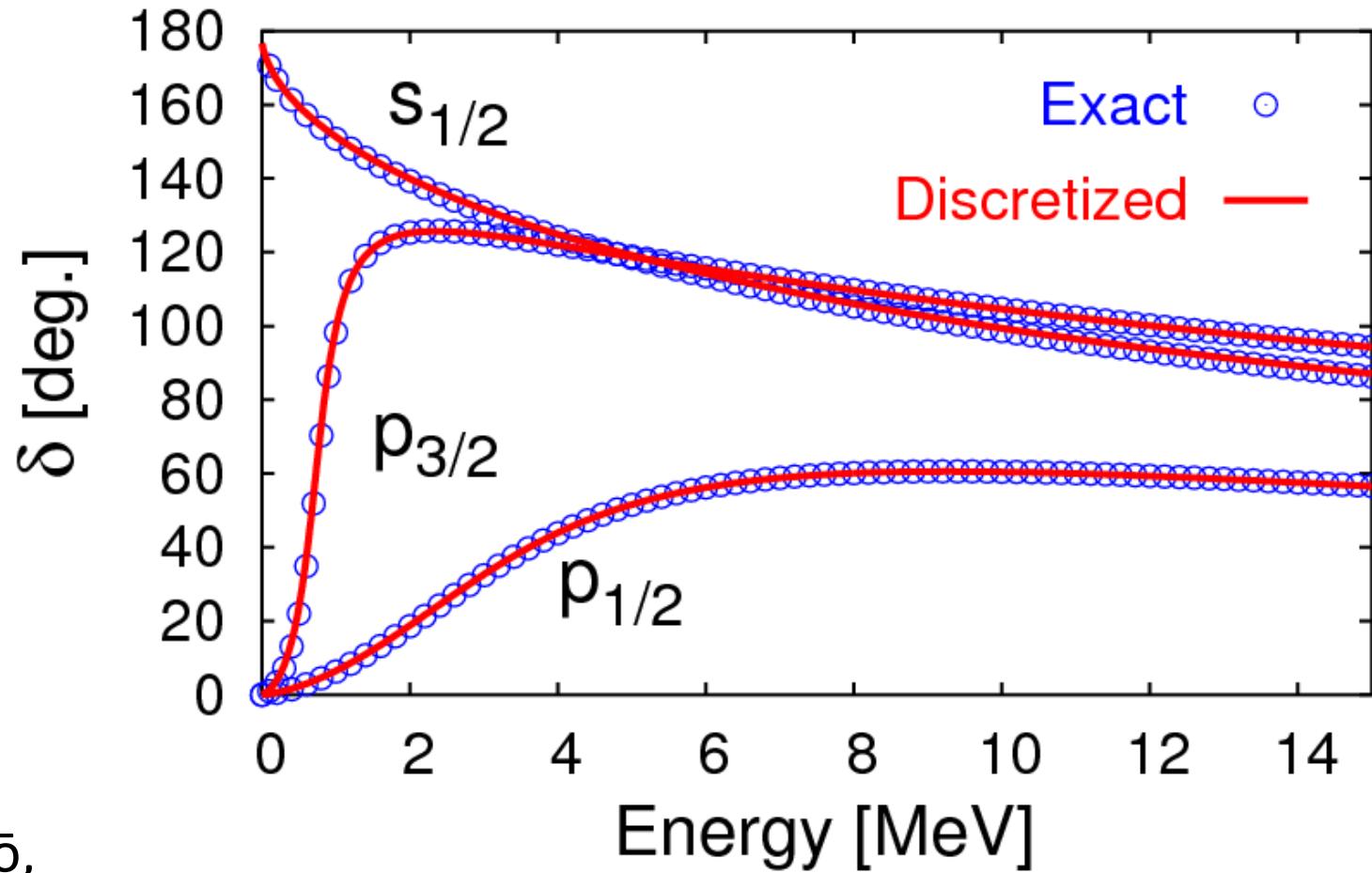
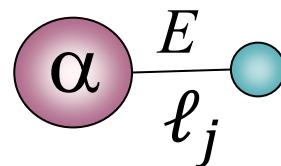
CLD $\Delta(E)$
(s,p-waves)

θ -independent



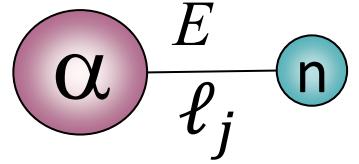
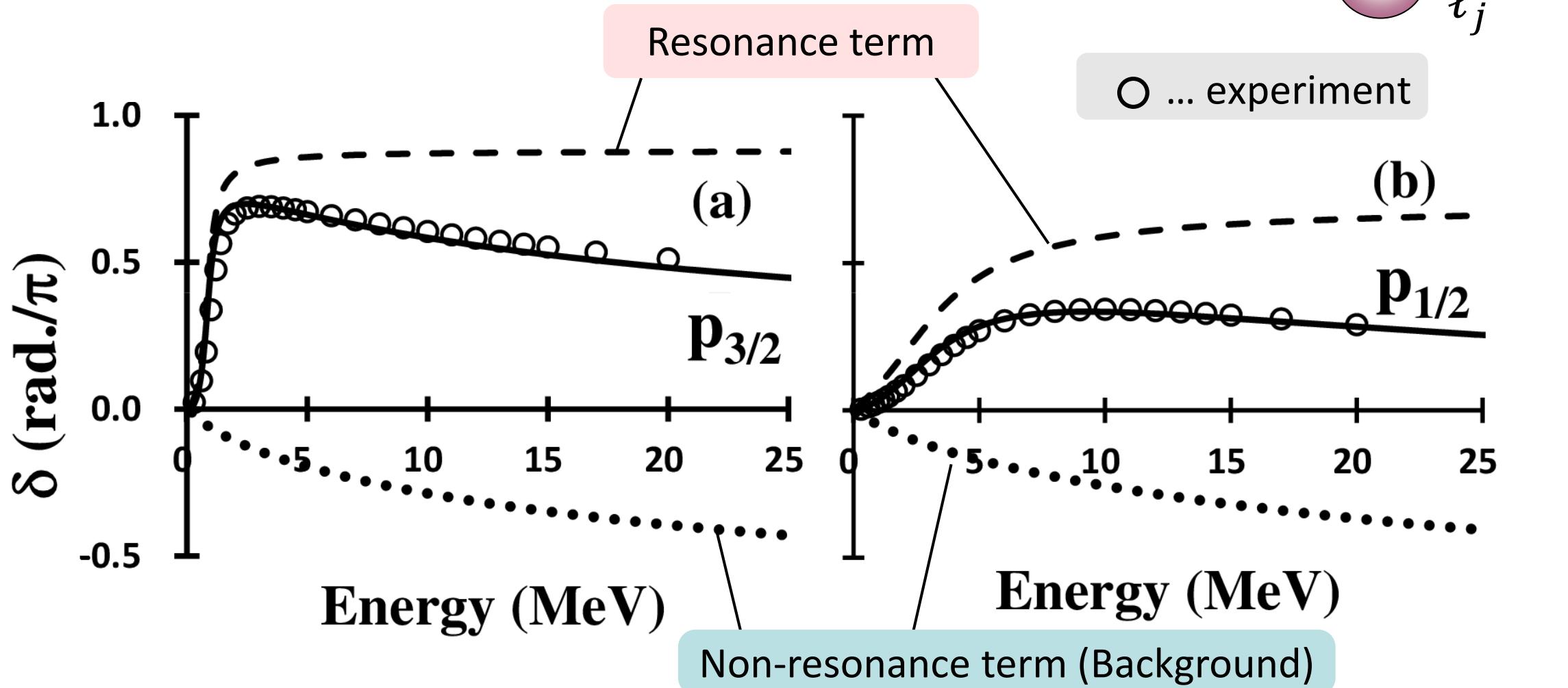
$\alpha+n$ scattering with discretized continuum states

Phase shifts (s,p)



R. Suzuki, T. Myo and K. Katō,
PTP113('05)1273.

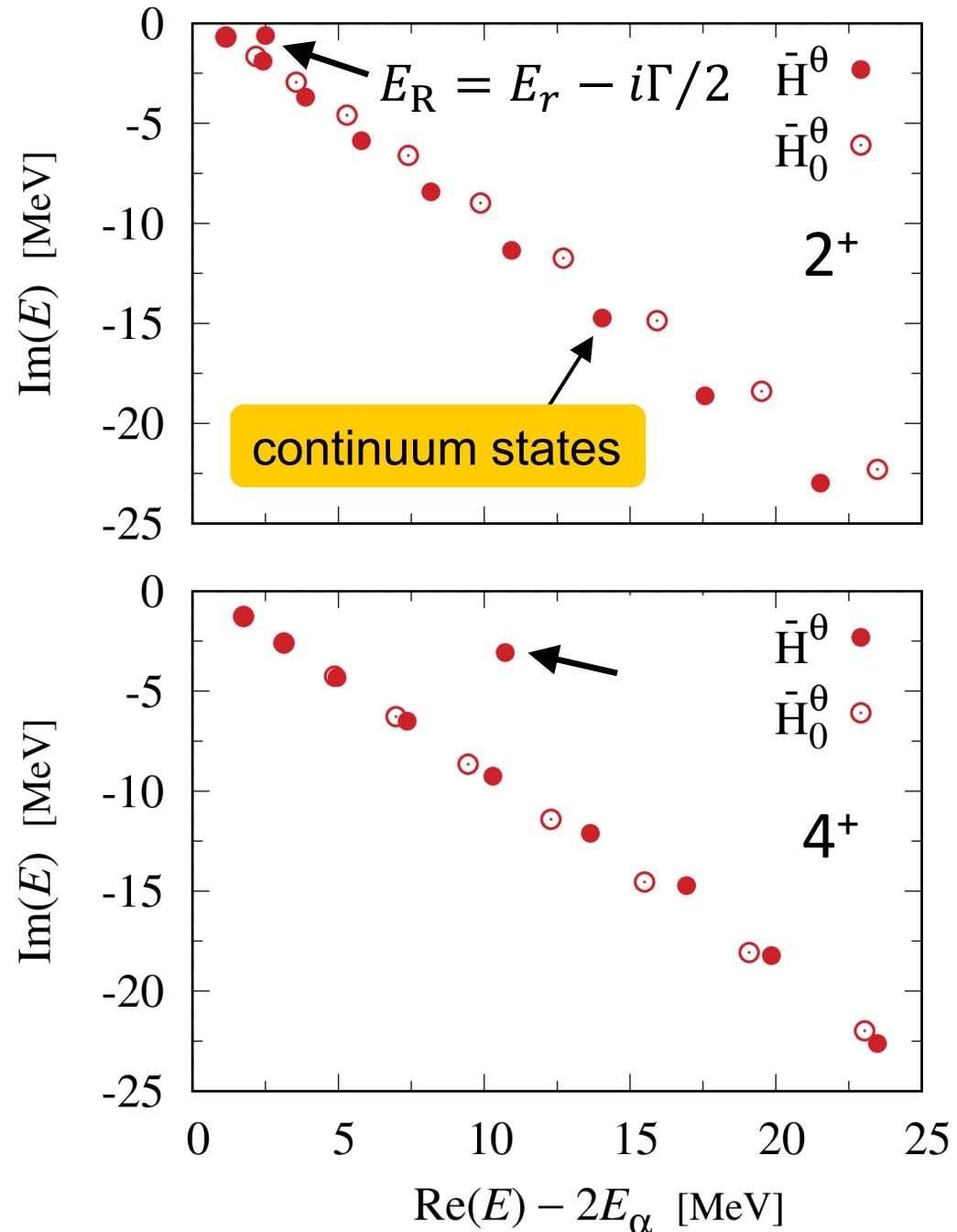
Decomposition of phase shift



Odsuren, Katō, Aikawa, Myo, Phys. Rev. C, 89, 034322 (2014).

Odsuren, Kikuchi, Myo, Khuukhenkhuu, Masui, K. Katō, Phys. Rev. C 95, 064305(2017).

Odsuren, Myo, Kikuchi, Teshigawara, Katō, Phys. Rev. C104, 014325 (2021).



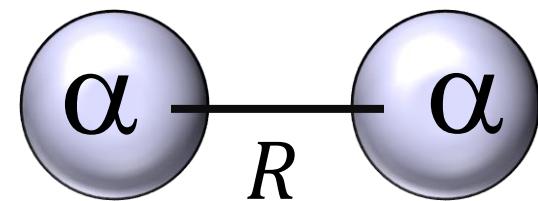
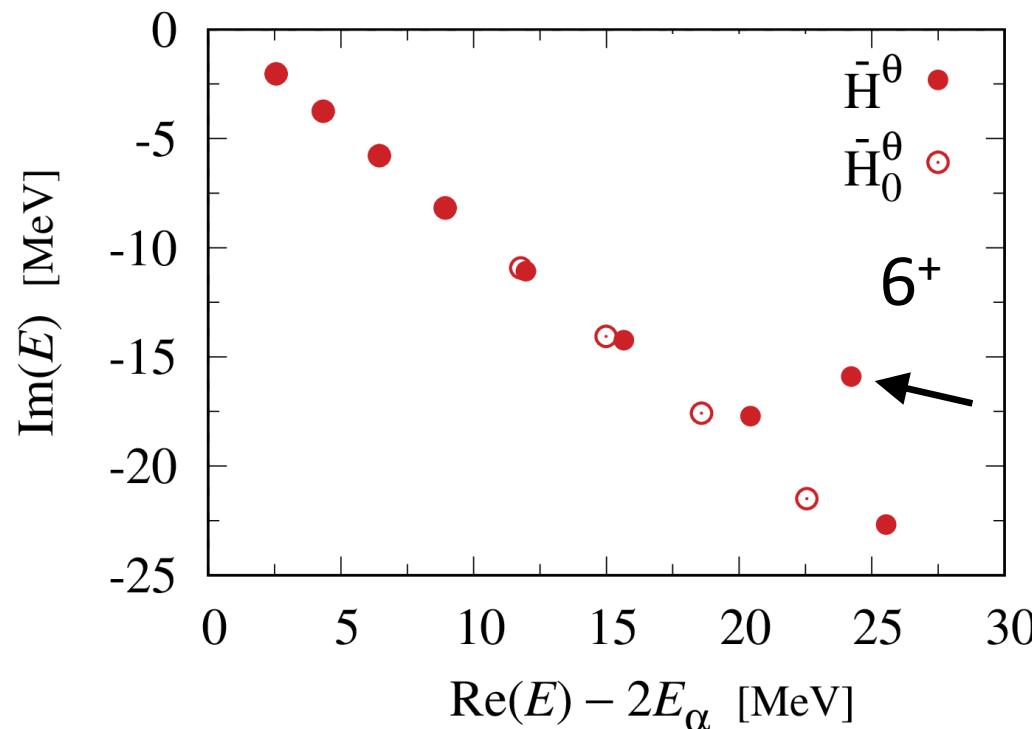
${}^8\text{Be}$ energy

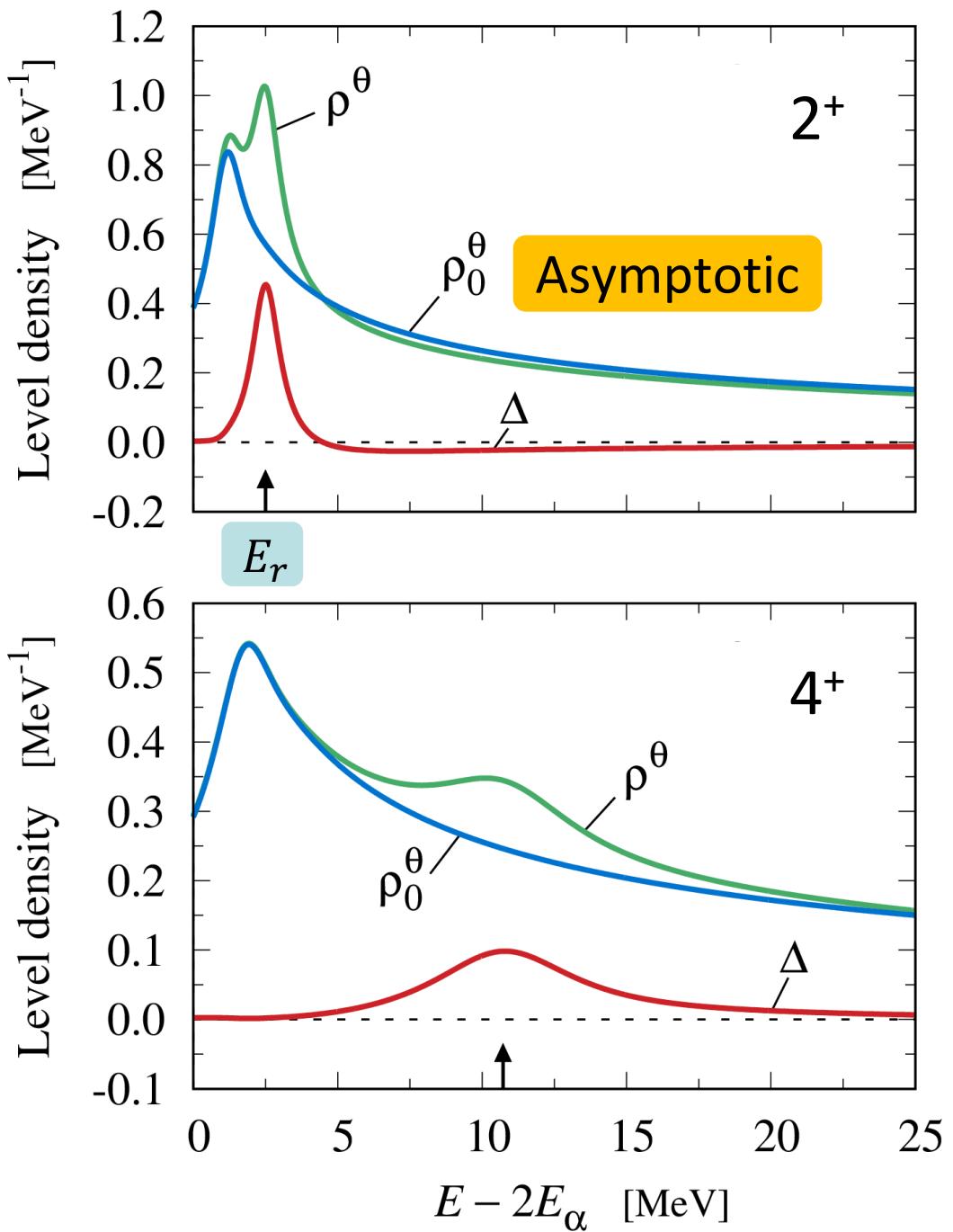
$$\theta = 27^\circ$$

E_ϑ measured from $2E_\alpha$

\bar{H}^θ : Complex-scaled Hamiltonian

\bar{H}_0^θ : Asymptotic Hamiltonian w/o nuclear force





${}^8\text{Be}$, level density

Asymptotic

$$\Delta(E) = \rho(E) - \rho_0(E)$$

$$= \sum_i^N \delta(E - E_i) - \sum_i^N \delta(E - E_{0,i})$$

$$= -\frac{1}{\pi} \text{Im} \left(\sum_i^N \frac{1}{E - E_{\theta,i}} \right) + \frac{1}{\pi} \text{Im} \left(\sum_i^N \frac{1}{E - E_{0,\theta,i}} \right)$$

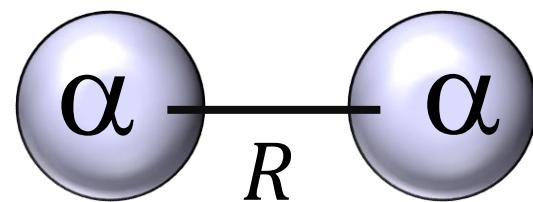
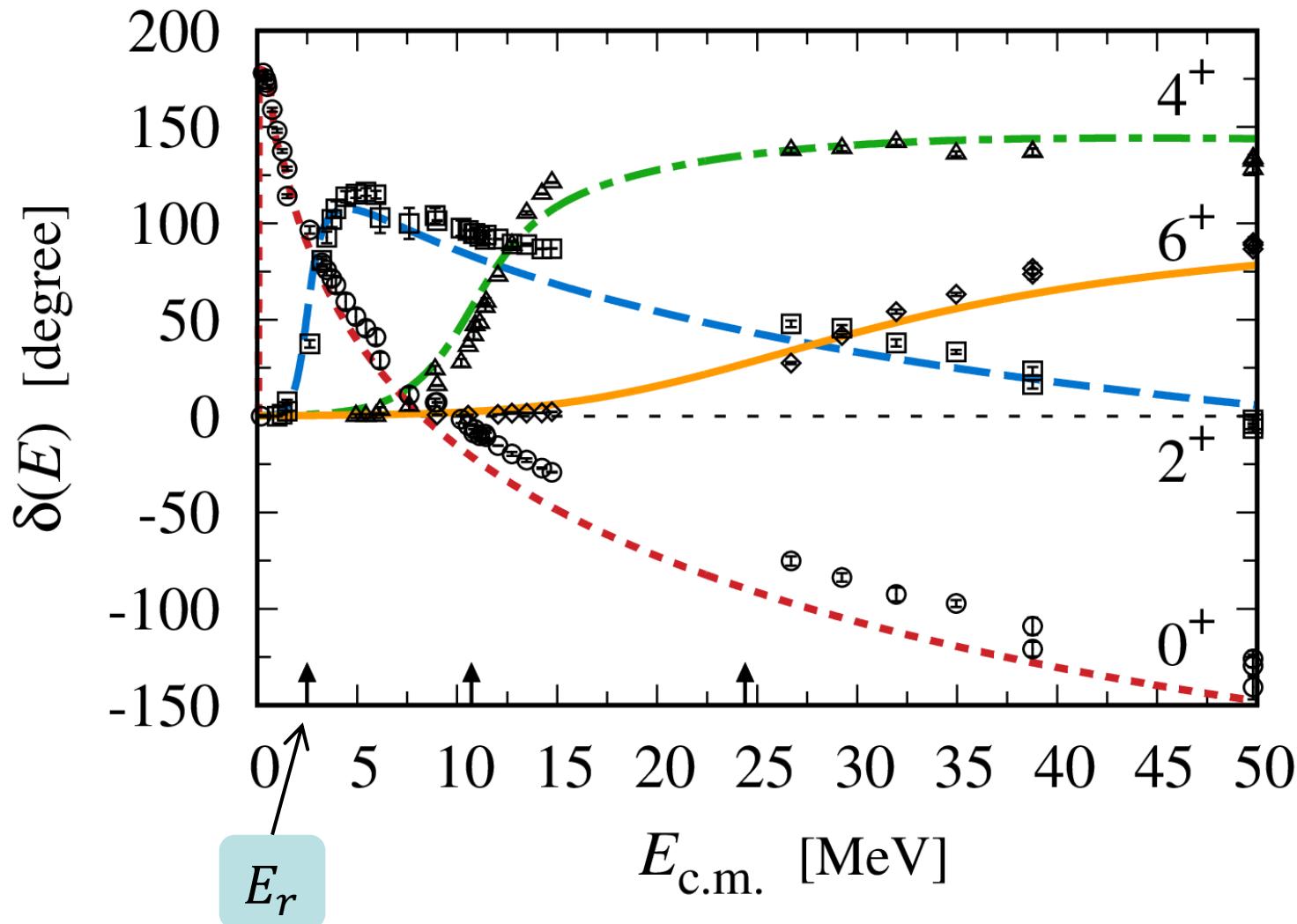
$$= \frac{1}{\pi} \frac{d\delta(E)}{dE}$$

complex scaling

- Continuum level density reveals resonances
- Measured from $\alpha+\alpha$ threshold energy

Myo, Kikuchi, Masui, Katō, Prog. Part. Nucl. Phys. **79**, 1 (2014)

$\alpha+\alpha$ phase shifts



- $\alpha+\alpha$ in Brink model with complex scaling $\theta = 27^\circ$
- NO channel radius

$$\Delta(E) = \rho(E) - \rho_0(E)$$

$$= \frac{1}{\pi} \frac{d\delta(E)}{dE}$$

Strength function $S(E)$ in complex scaling

$$S(E) = \sum_n \langle \tilde{\Phi}_0 | \hat{O}^\dagger | \varphi_n \rangle \langle \tilde{\varphi}_n | \hat{O} | \Phi_0 \rangle \cdot \delta(E - E_n)$$

\hat{O} : Transition operator

$$= -\frac{1}{\pi} \text{Im} [R(E)]$$

initial state
final state

$$R(E) = \sum_n \frac{\langle \tilde{\Phi}_0 | \hat{O}^\dagger | \varphi_n \rangle \langle \tilde{\varphi}_n | \hat{O} | \Phi_0 \rangle}{E - E_n}$$

Response function

- Complex-scaled Green's function

$$G^\theta(E) = \frac{1}{E - H_\theta} = \sum_n \frac{|\varphi_n^\theta\rangle\langle\tilde{\varphi}_n^\theta|}{E - E_n^\theta}$$

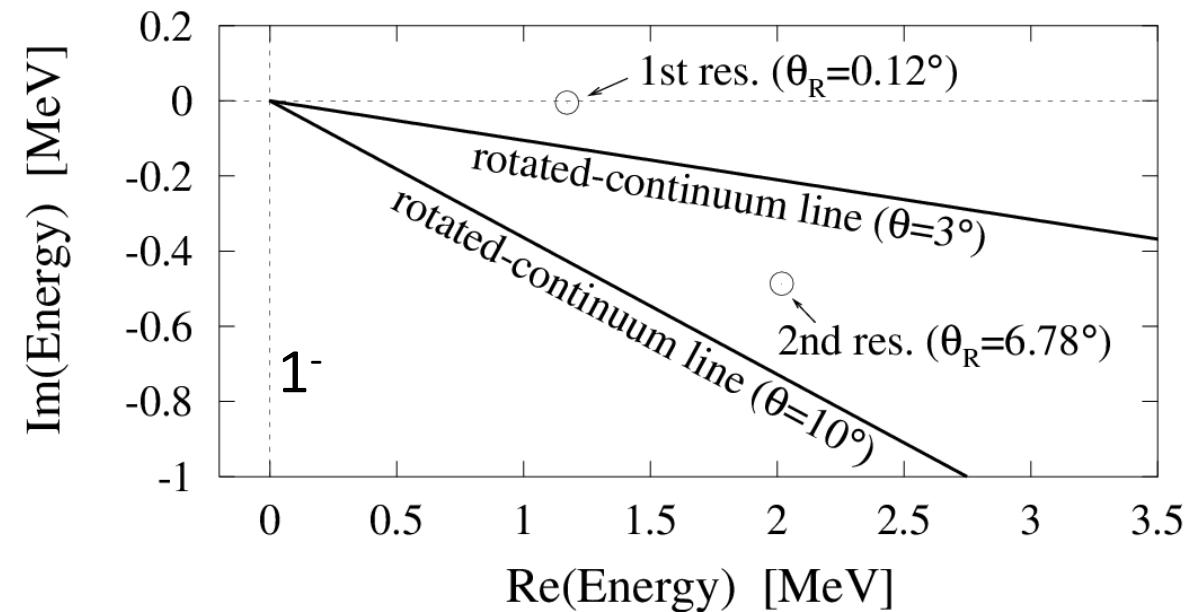
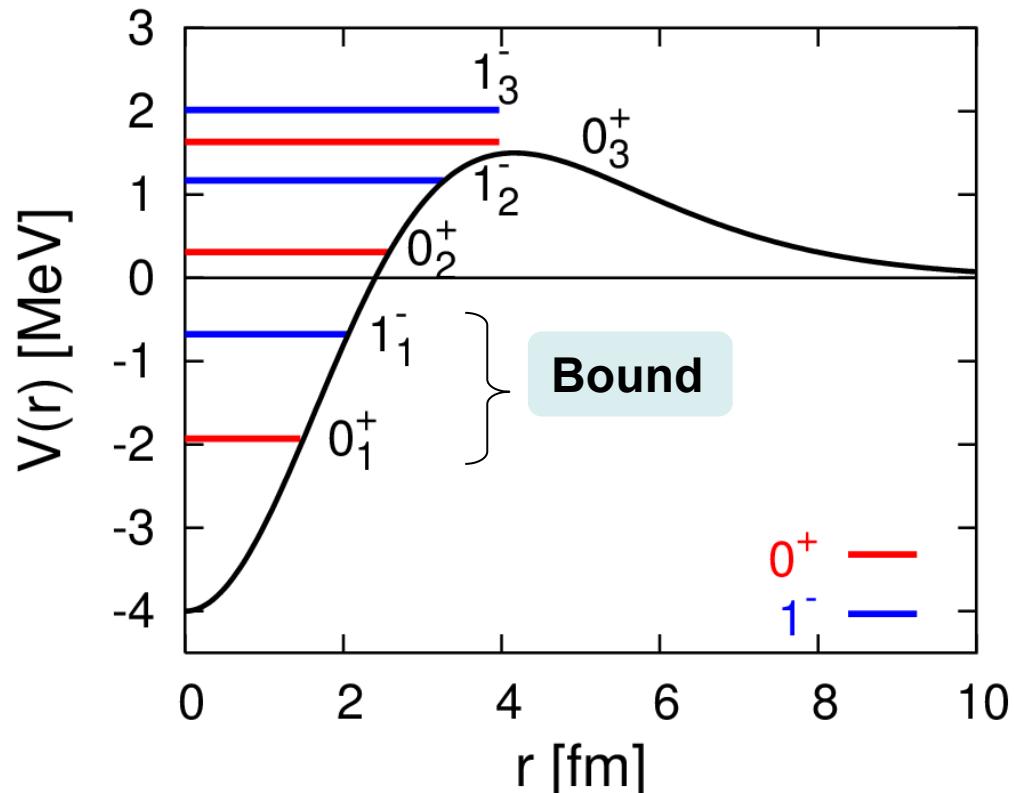
Bound+Resonance+Continuum

Reaction theory

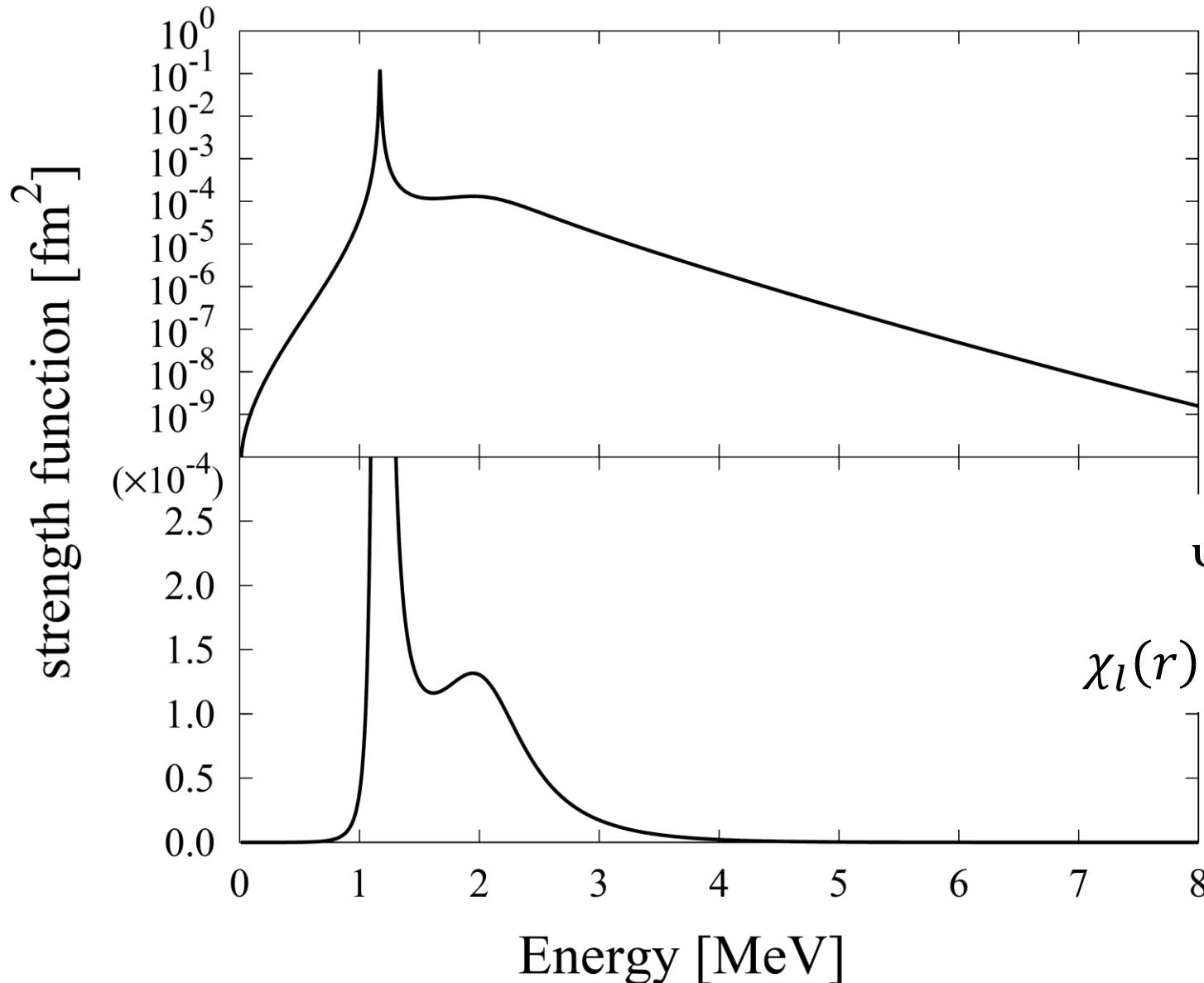
- Lippmann–Schwinger-eq. (Kikuchi)
- CDCC (Matsumoto)
- Scatt. Amp. (Kruppa, Dote(K^{bar}N))

Dipole transition : Schematic potential

$$H = -\frac{\hbar^2 \nabla^2}{2\mu} + V, \quad \frac{\hbar^2}{\mu} = 1, \quad V(r) = -8e^{-0.16r^2} + 4e^{-0.04r^2}$$



Dipole transition strength, $0^+ \rightarrow 1^-$



$$\hat{O}_m(E1) = r Y_{1m}(\hat{r})$$

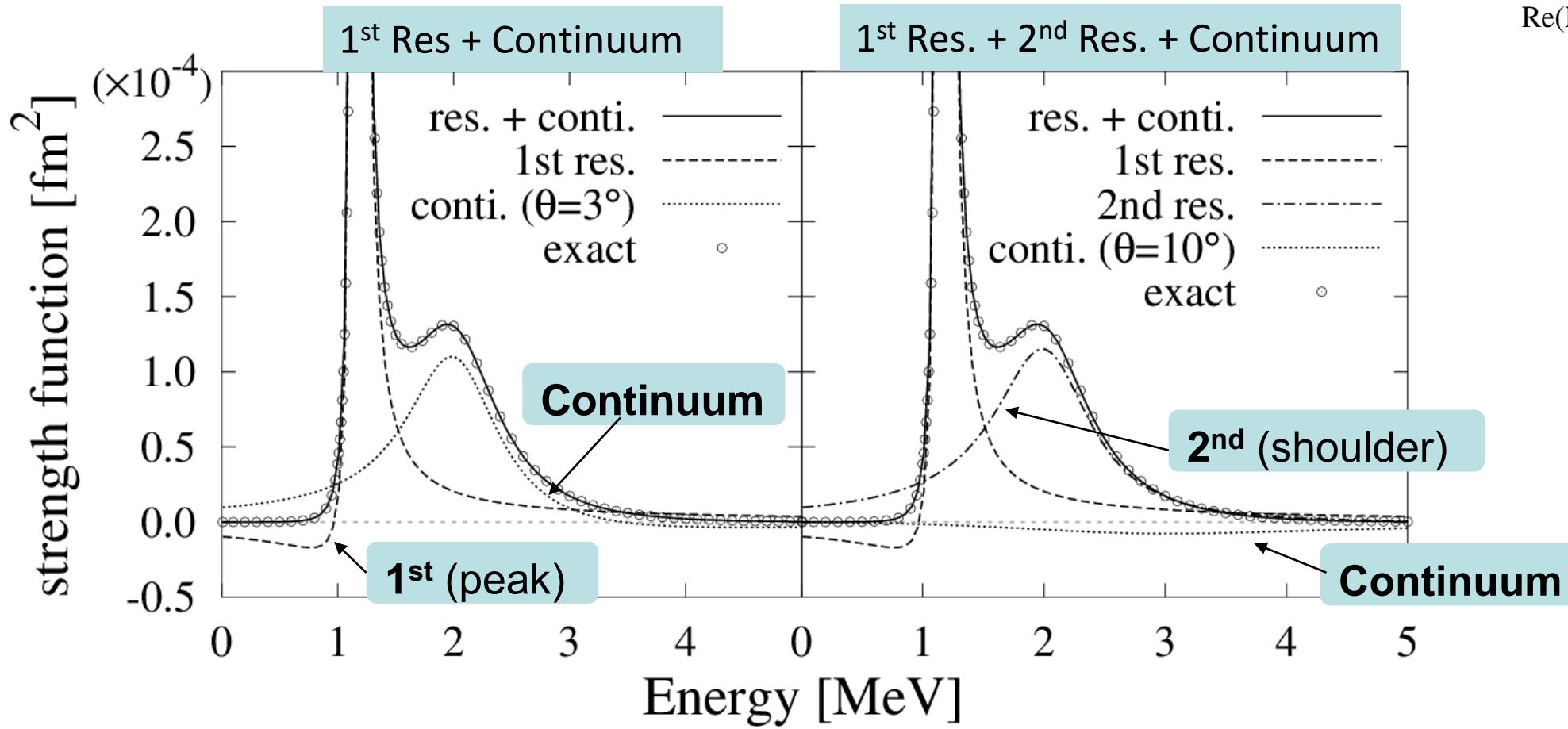
TM, A. Ohnishi and K. Katō,
PTP99(1998)801

$$\Psi_{lm}(\mathbf{r}) = C_l(k) \frac{\chi_l(r)}{kr} Y_{lm}(\hat{\mathbf{r}})$$

$$\chi_l(r) \xrightarrow[r \rightarrow \infty]{} u_l^{(-)}(kr) - S_l(k) u_l^{(+)}(kr)$$

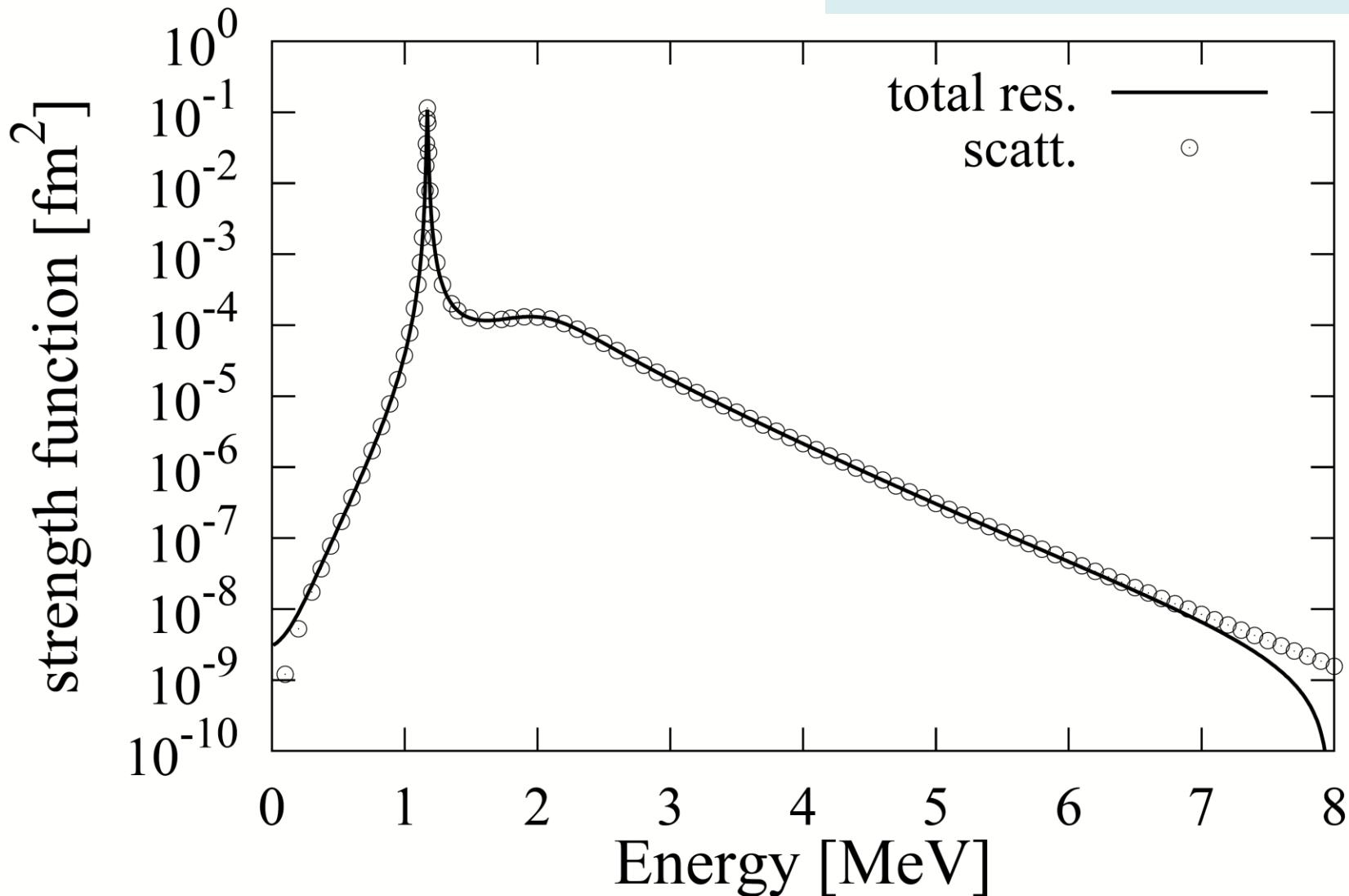
Solve scattering problem

Dipole transition with extended completeness relation



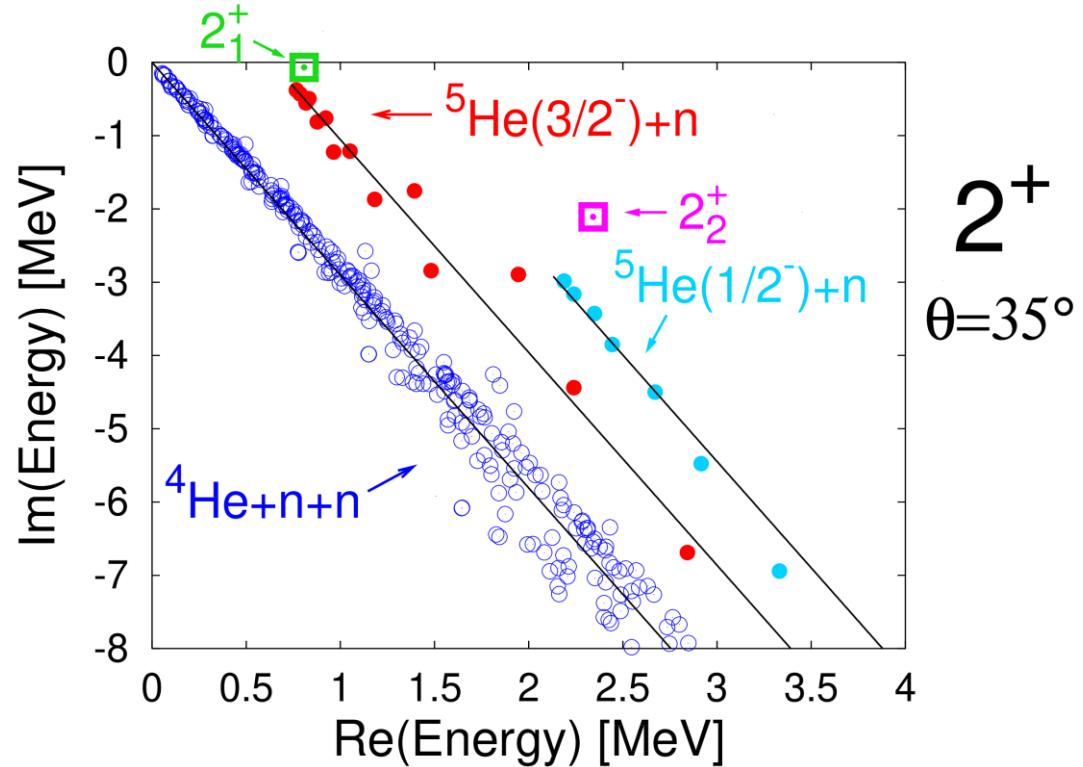
E1 strength from only resonances

9 resonances in 1^- state

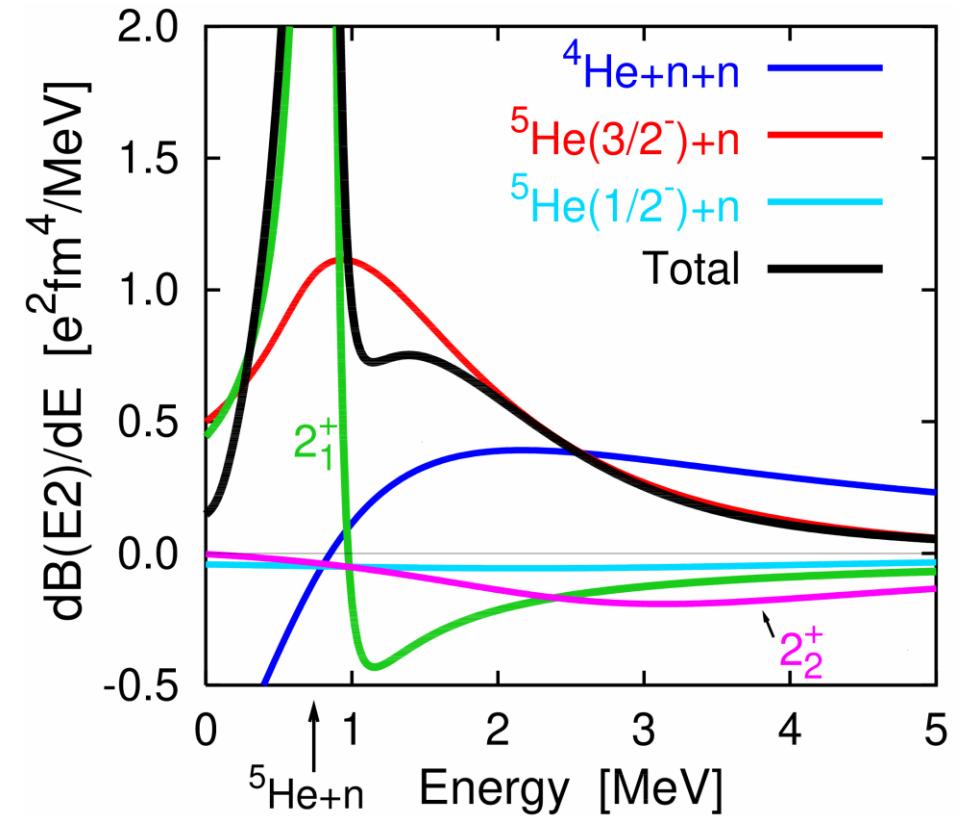


Quadrupole transition of ${}^6\text{He}$ into $\alpha + n + n$

- Zero energy : $\alpha + n + n$ threshold



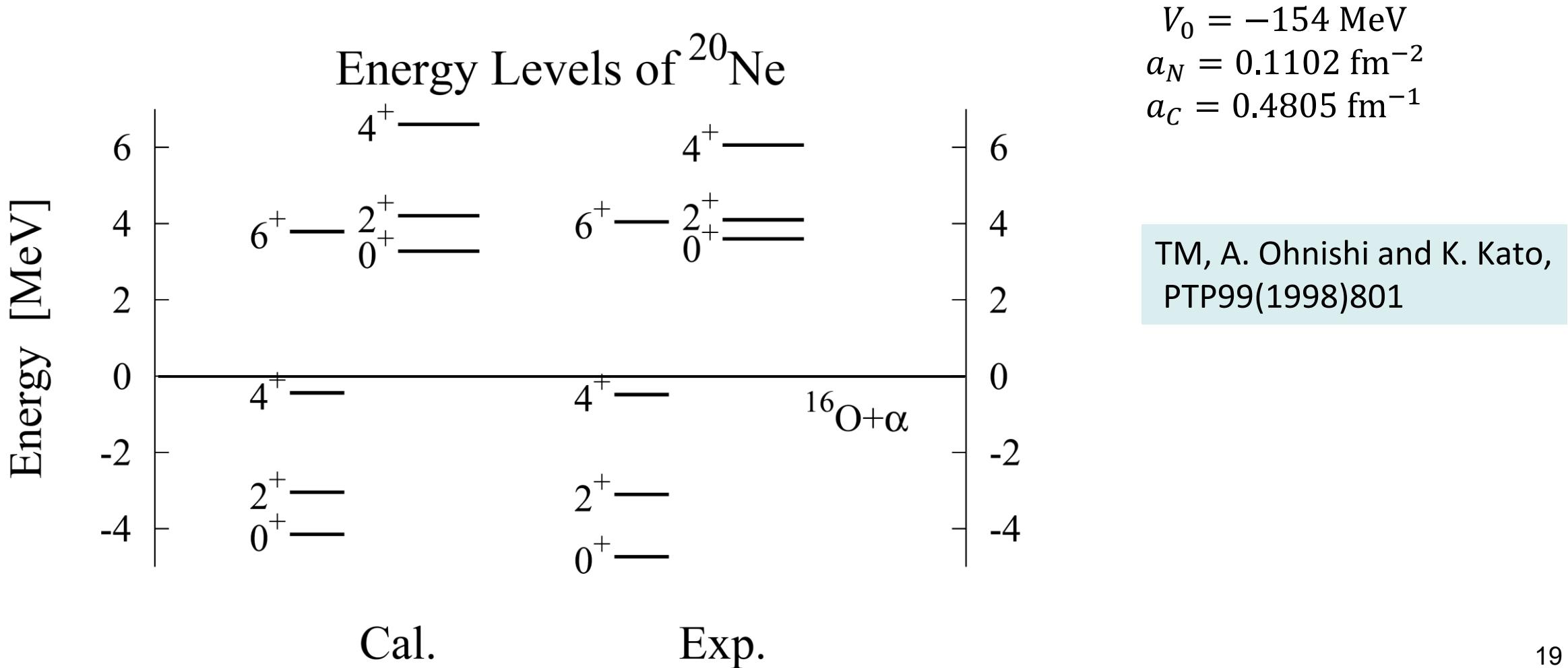
Transition by $er^2 Y_2(\hat{r})$



- Sharp peak : **1st resonance**
- Shoulder structure : **${}^5\text{He} + n$ continuum state**
- Minor effect : **2nd resonance**

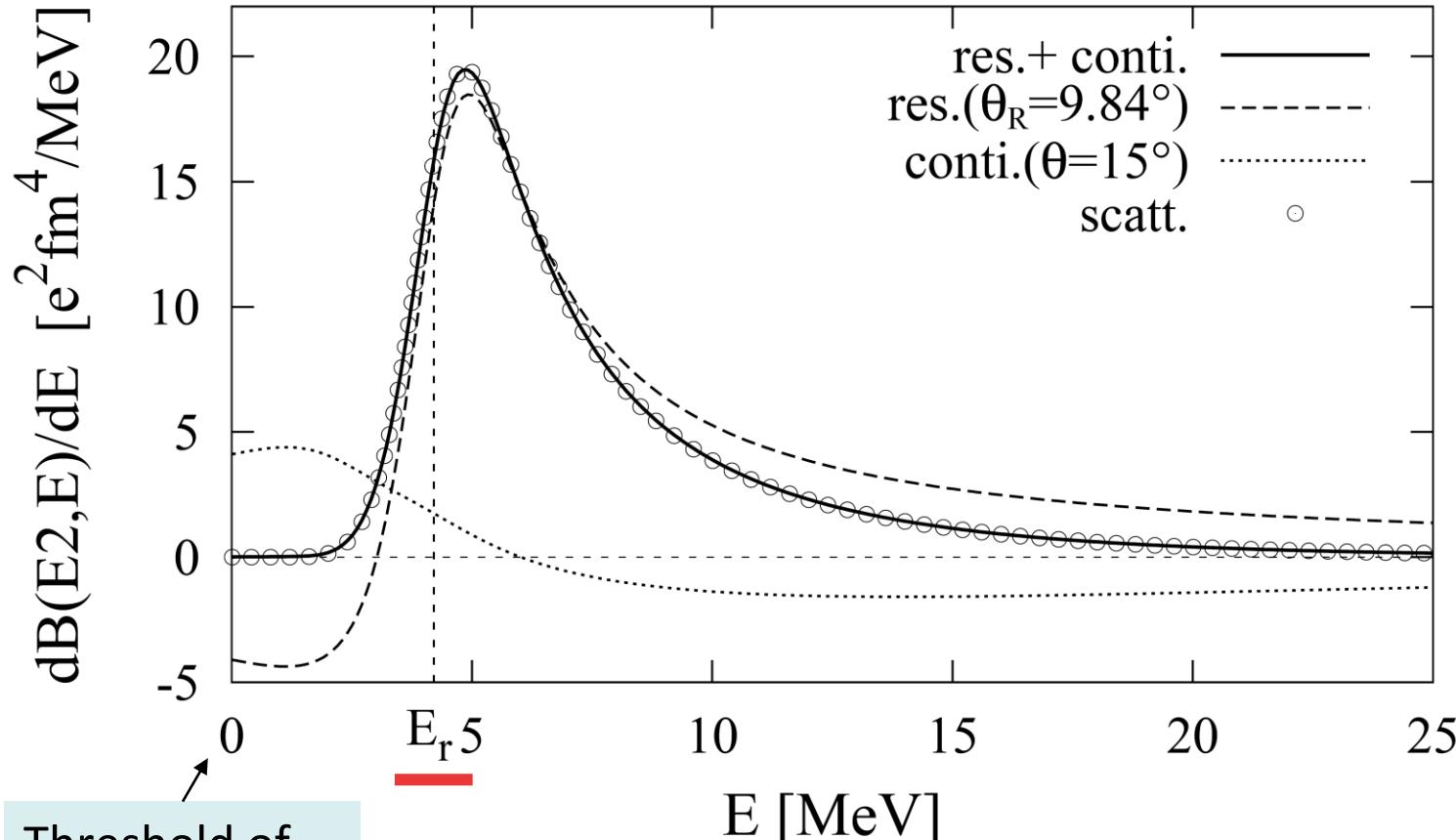
^{20}Ne as $^{16}\text{O} + ^4\text{He}$, E2 transition

$$H = -\frac{\hbar^2}{2\mu} \nabla^2 + V_N + V_C, \quad V_N = V_0 e^{-a_N r^2}, \quad V_C = 8 \cdot 2 \frac{e^2}{r} \operatorname{erf}(a_C r)$$



$^{20}\text{Ne} = ^{16}\text{O} + ^4\text{He}$, E2 transition

$$H = -\frac{\hbar^2}{2\mu} \nabla^2 + V_N + V_C, \quad V_N = V_0 e^{-a_N r^2}, \quad V_C = 8 \cdot 2 \frac{e^2}{r} \operatorname{erf}(a_C r)$$



Peak position is **shifted** from resonance energy E_r due to :

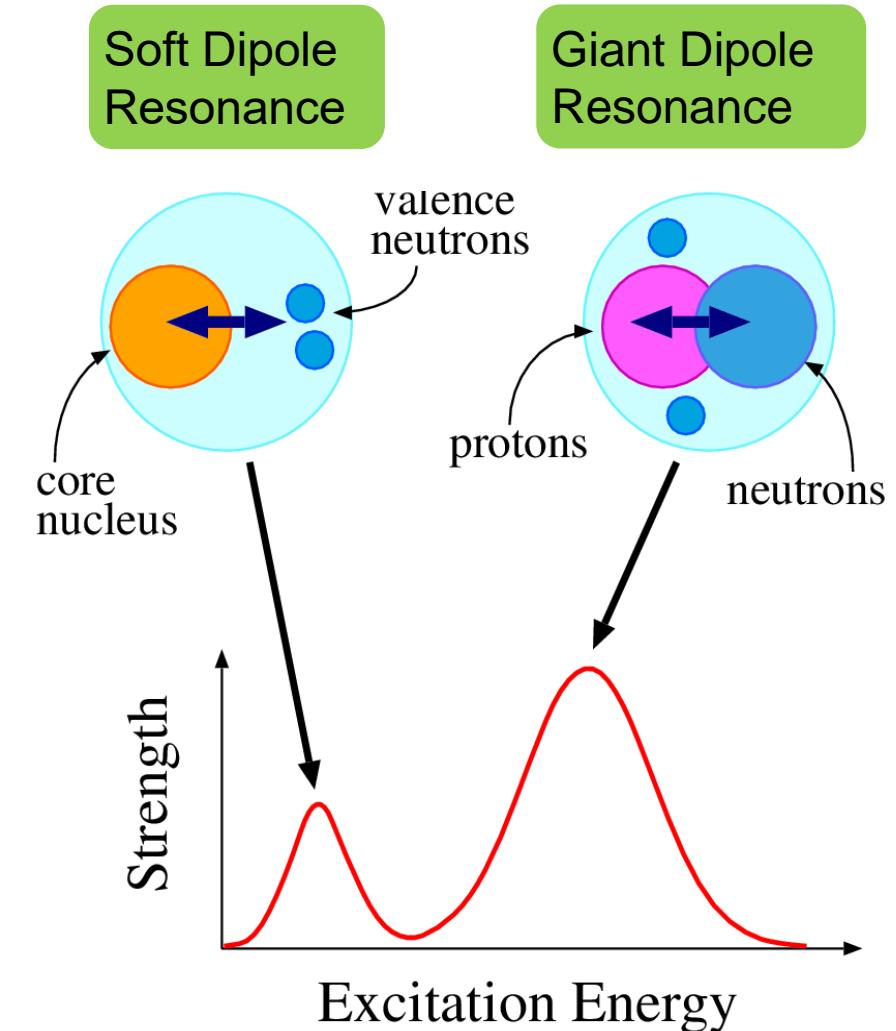
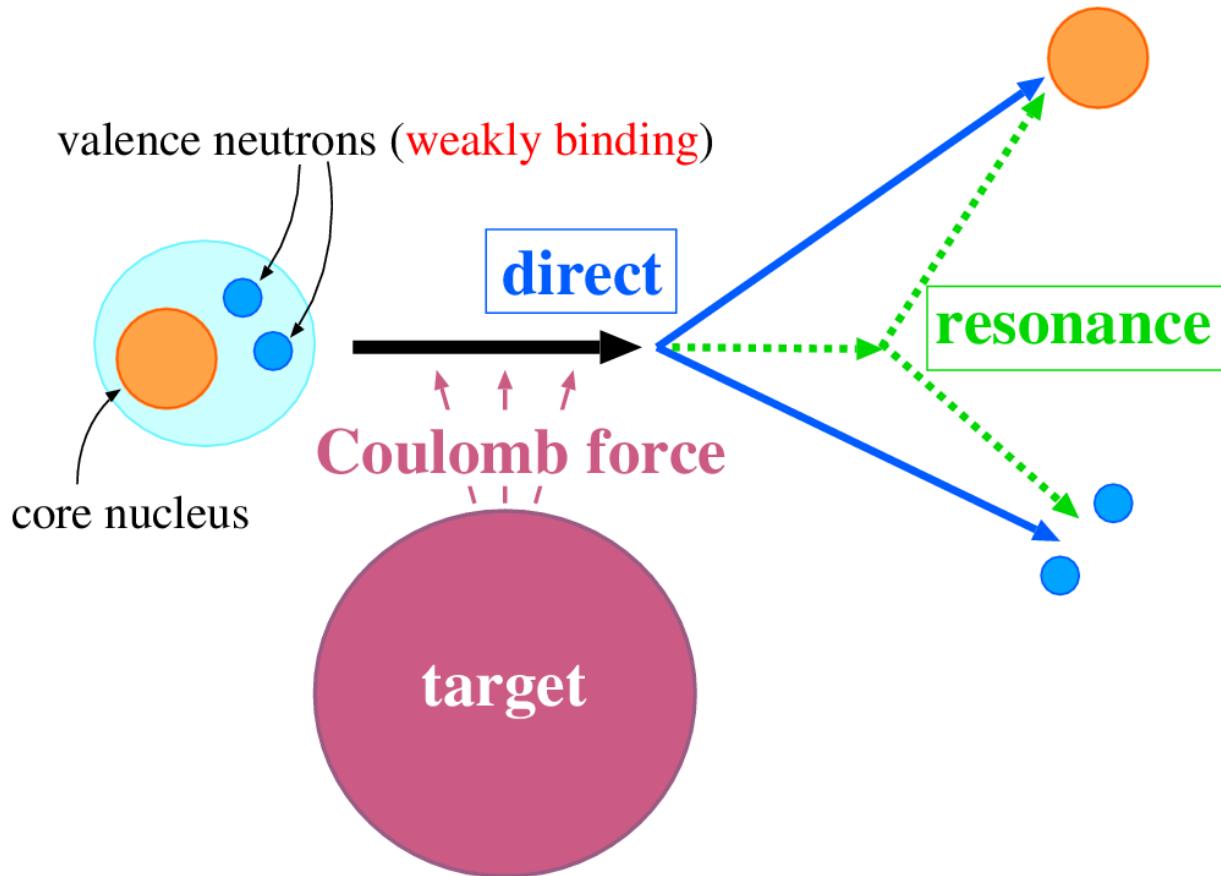
- 1) **Imaginary part** of the transition matrix element of resonance
- 2) **Continuum component** as background

Strength function $S(E)$: properties

- $E_\alpha = E_R - iE_I$: complex energy of state α , $E_R, E_I > 0$
- $M_\alpha = \langle \Phi_\alpha | O | \Phi_0 \rangle^2 = M_R + iM_I$: transition matrix element to the state α
- $S_\alpha(E) = -\frac{1}{\pi} \text{Im} \left(\frac{M_\alpha}{E - E_\alpha} \right) = -\frac{1}{\pi} \text{Im} \left(\frac{M_R + iM_I}{E - E_R + iE_I} \right) = \frac{1}{\pi} \frac{M_R E_I - M_I (E - E_R)}{(E - E_R)^2 + E_I^2}$
- $\int_{-\infty}^{\infty} S_\alpha(E) dE = M_R$: Strength of the distribution
- M_I : **Asymmetry** of the distribution measured from E_R
- $S_\alpha(E_R) = \frac{1}{\pi} \frac{M_R}{E_I}$ at central energy
- $S_\alpha(E_R \pm E_I) = \frac{1}{\pi} \frac{M_R \pm M_I}{2E_I}$: width of the distribution
- $E_{\max/\min} = E_R + \frac{E_I}{M_I} (M_R \mp |M_\alpha|)$: Max/Min energies of $S_\alpha(E)$

Myo, Katō, PRC 107 (2023)

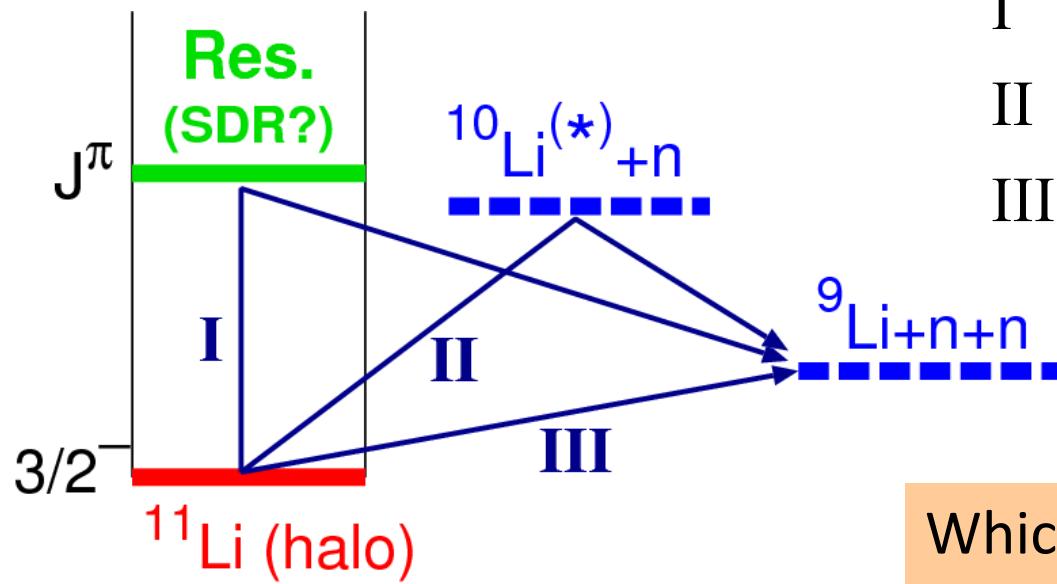
Coulomb breakup reactions & soft dipole resonance



P.G. Hansen and B.Jonson, Europhys. Lett. 4(1987)409.
K. Ikeda, NPA538(1992)355c.
R. Kanungo, TM et al., PRL114 (2015) 192502

Coulomb breakup reaction

- Structure and Responses of **unbound states**
 - Resonance (Gamow states) and Non-resonant continua
- Interesting excitation : soft dipole resonances?

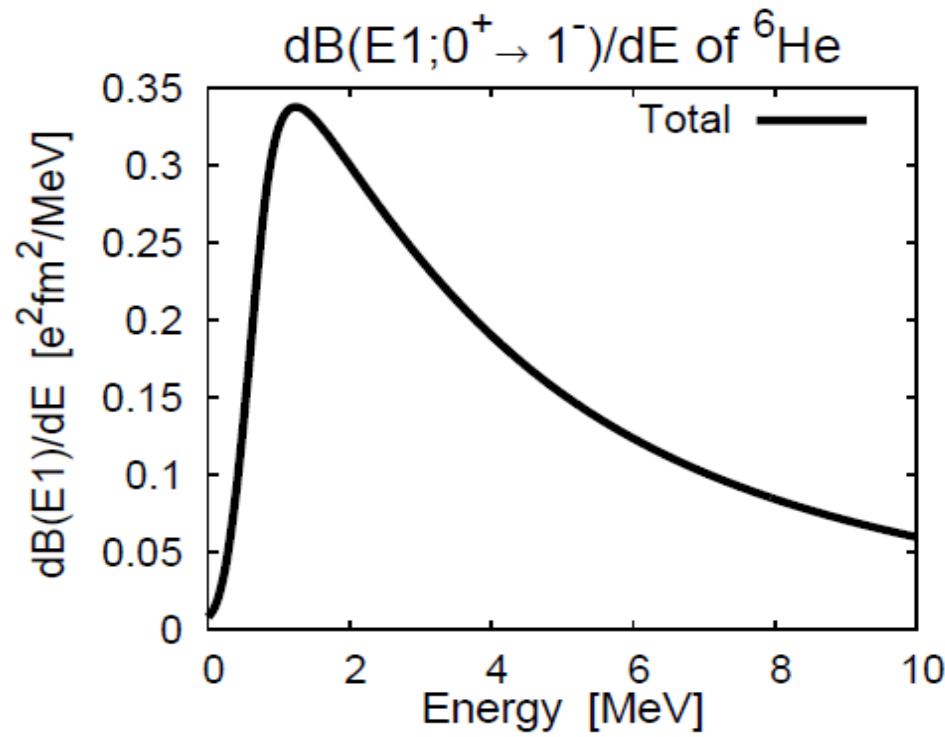
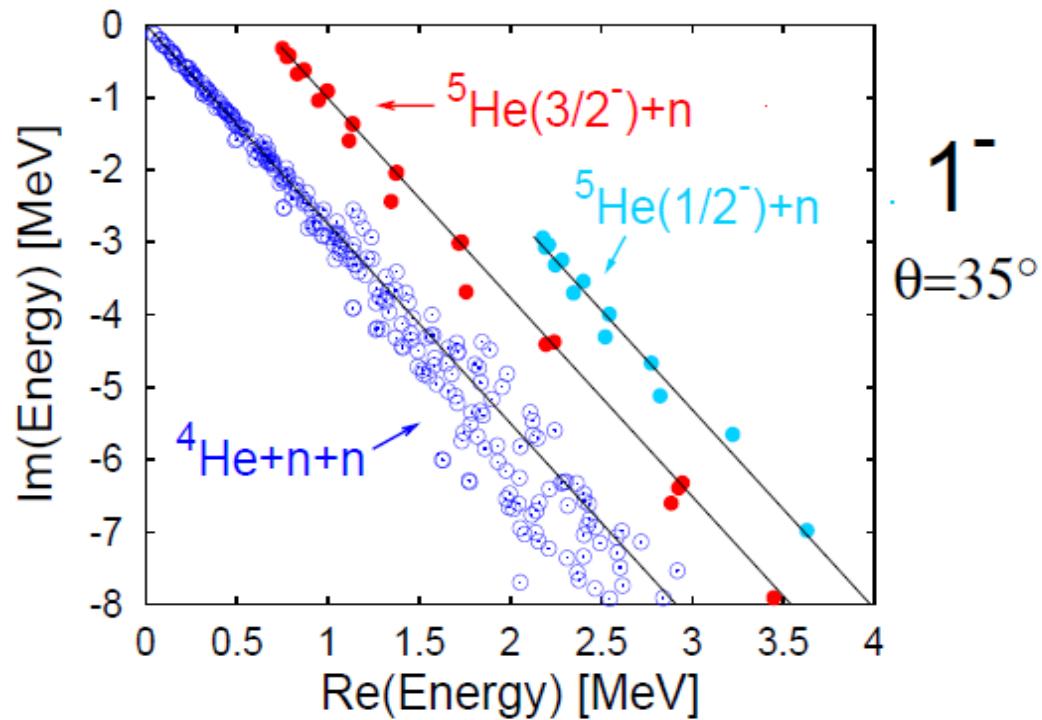


I : resonance of ^{11}Li
II : sequential via $^{10}\text{Li} + n$
III : 3-body direct breakup

Which breakup process is preferred and constructs the structures of strength?

(1) Case of ${}^6\text{He}$: E1 transition with CSM

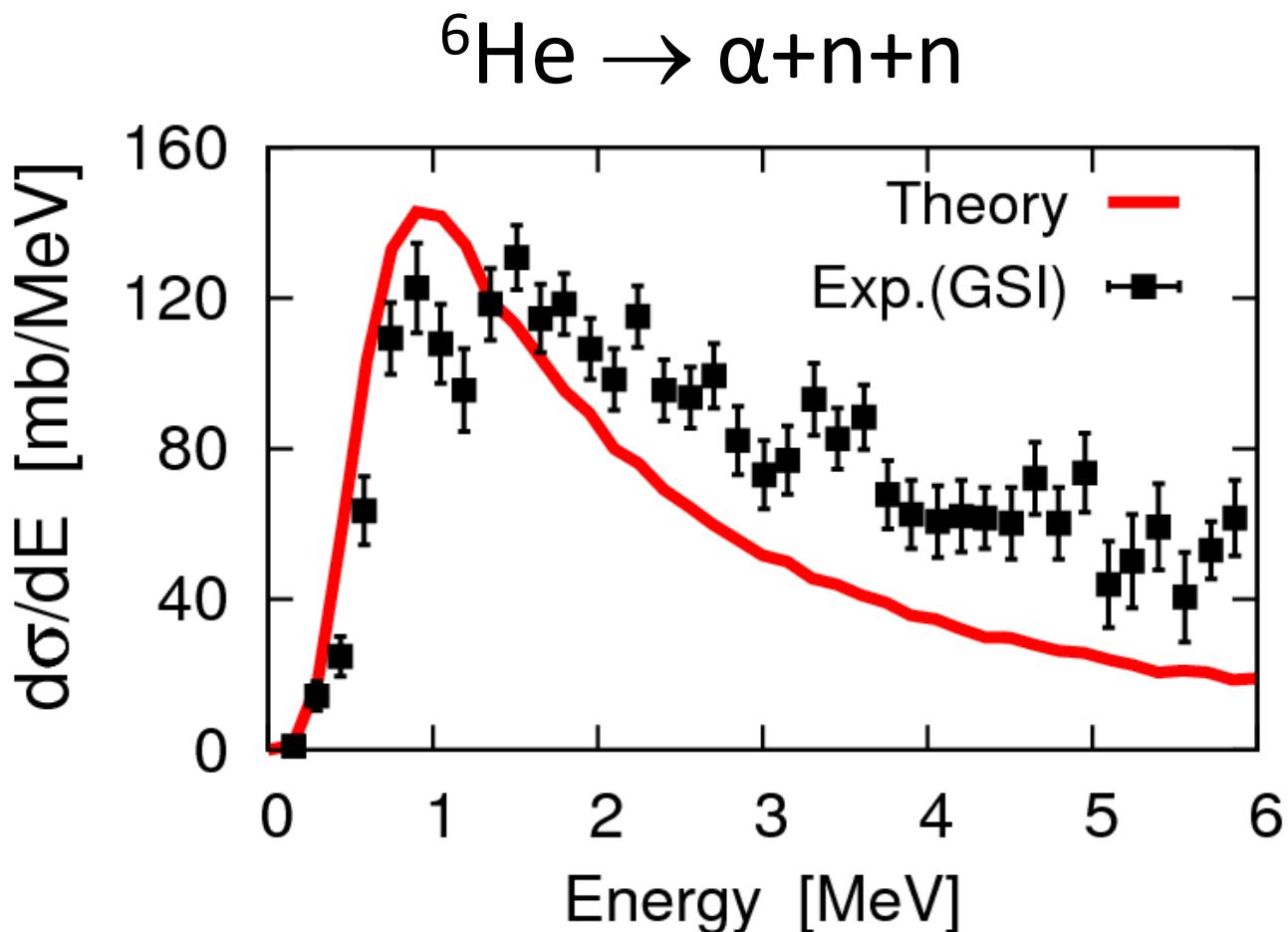
- $\Psi({}^6\text{He}) = \mathcal{A}\{ \Phi({}^4\text{He}) \Phi(\text{nn}) \}$ with OCM, $\Phi(\text{nn})$: Gaussian expansion (discretized continuum)
- ${}^4\text{He}-\text{n}$: KKNN pot. / n-n : Minnesota.



No three-body resonance

- Low energy enhancement

Coulomb breakup strength of ${}^6\text{He}$



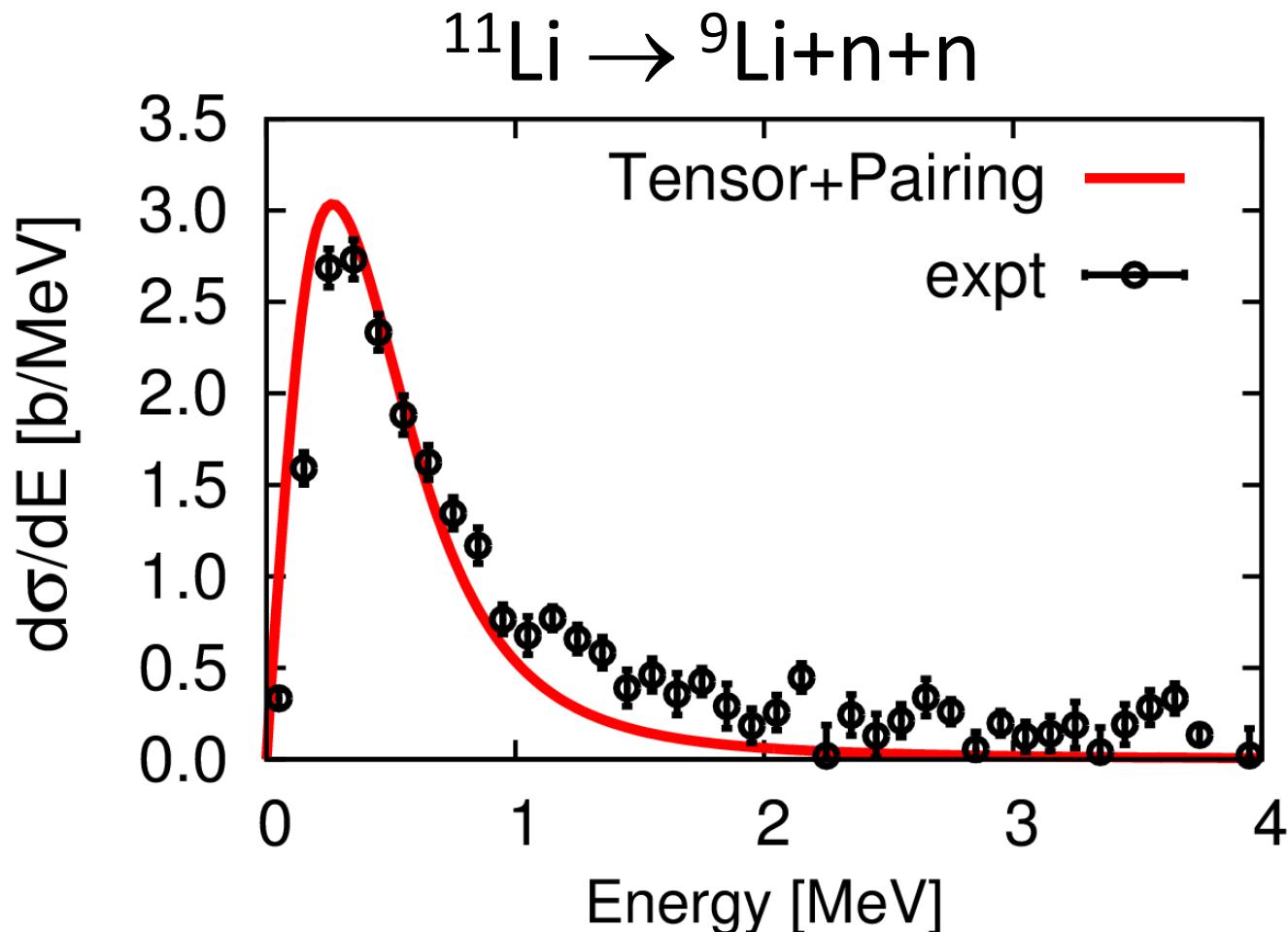
E1+E2 (complex scaling)
Equivalent photon method

TM, K. Katō, S. Aoyama and K. Ikeda
PRC63(2001)054313.

Kikuchi, TM, Takashina, Katō, Ikeda
PTP122(2009)499
PRC81(2010)044308.
(invariant mass of α - n & n - n)

${}^6\text{He}$: 240MeV/A, Pb Target (T. Aumann et.al, PRC59(1999)1252)

Coulomb breakup strength of halo nucleus ^{11}Li



No three-body resonance

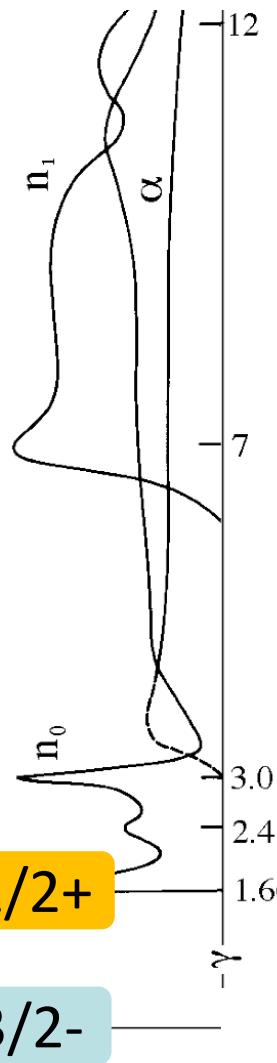
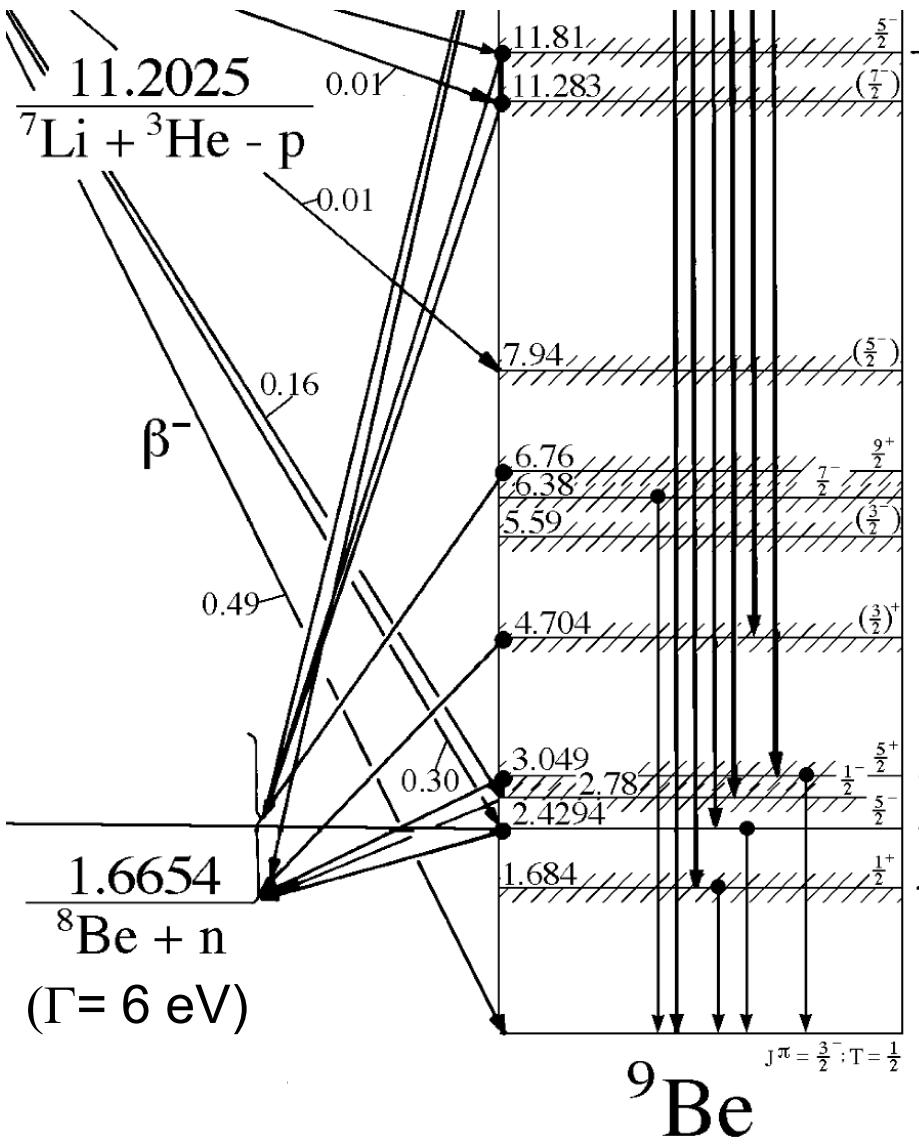
E1 strength

+ Complex scaling

+ Equivalent photon method

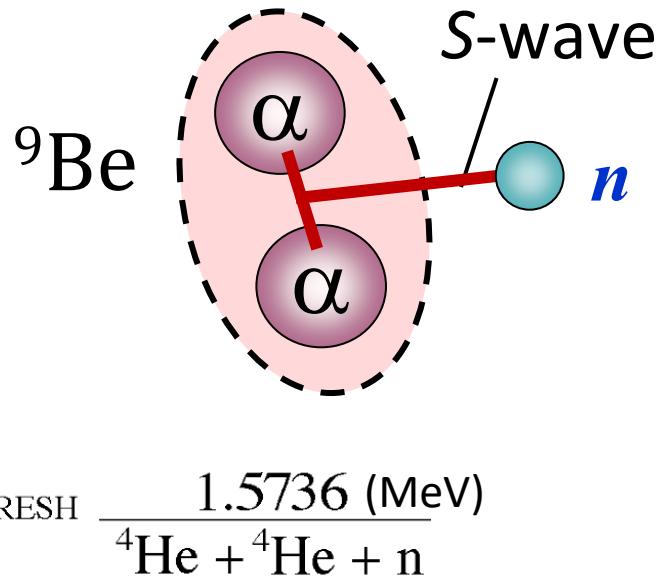
T.Myo, K.Katō, H.Toki, K.Ikeda
PRC76(2007)024305

- Expt: T. Nakamura et al. , PRL96,252502(2006)
- Energy resolution with $\sqrt{E} = 0.17$ MeV.



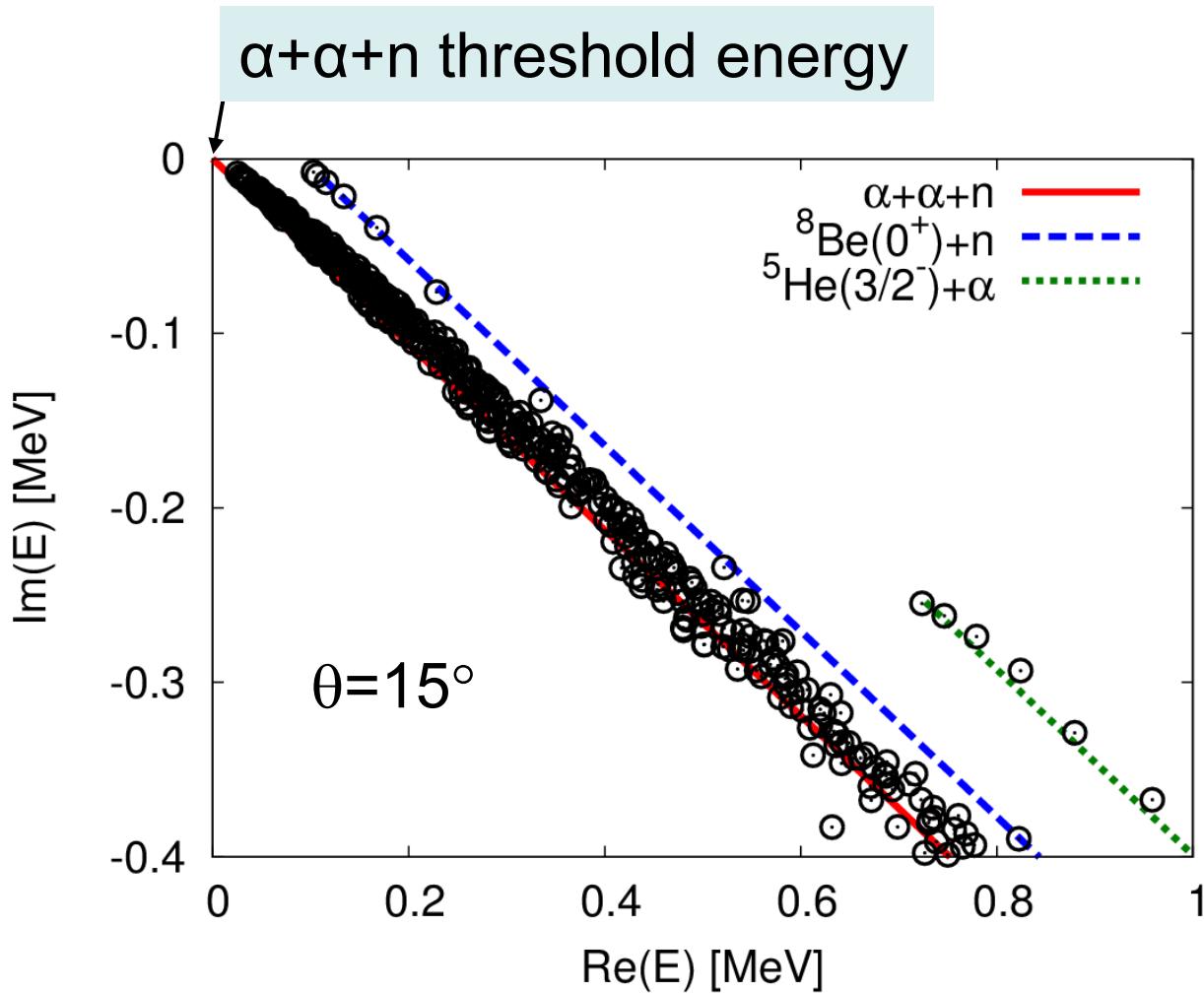
^9Be photodisintegration

TUNL Nuclear Data evaluation



- Property of $1/2^+$ S-state - resonance ? / virtual state ?
- Photodisintegration with $\alpha+\alpha+n$ 3-body approach

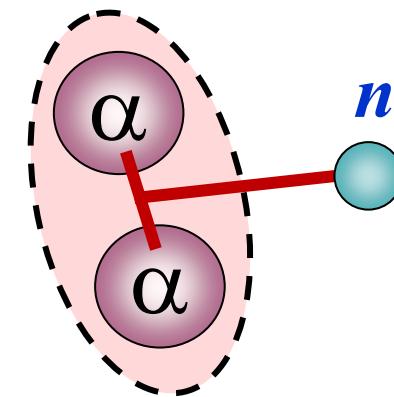
${}^9\text{Be}$: Property of $1/2^+$ state



$$H({}^9\text{Be}) = T_R + T_r + V_{\alpha_1 n} + V_{\alpha_2 n} + V_{\alpha \alpha} + V_{\text{PF}} + V_3$$

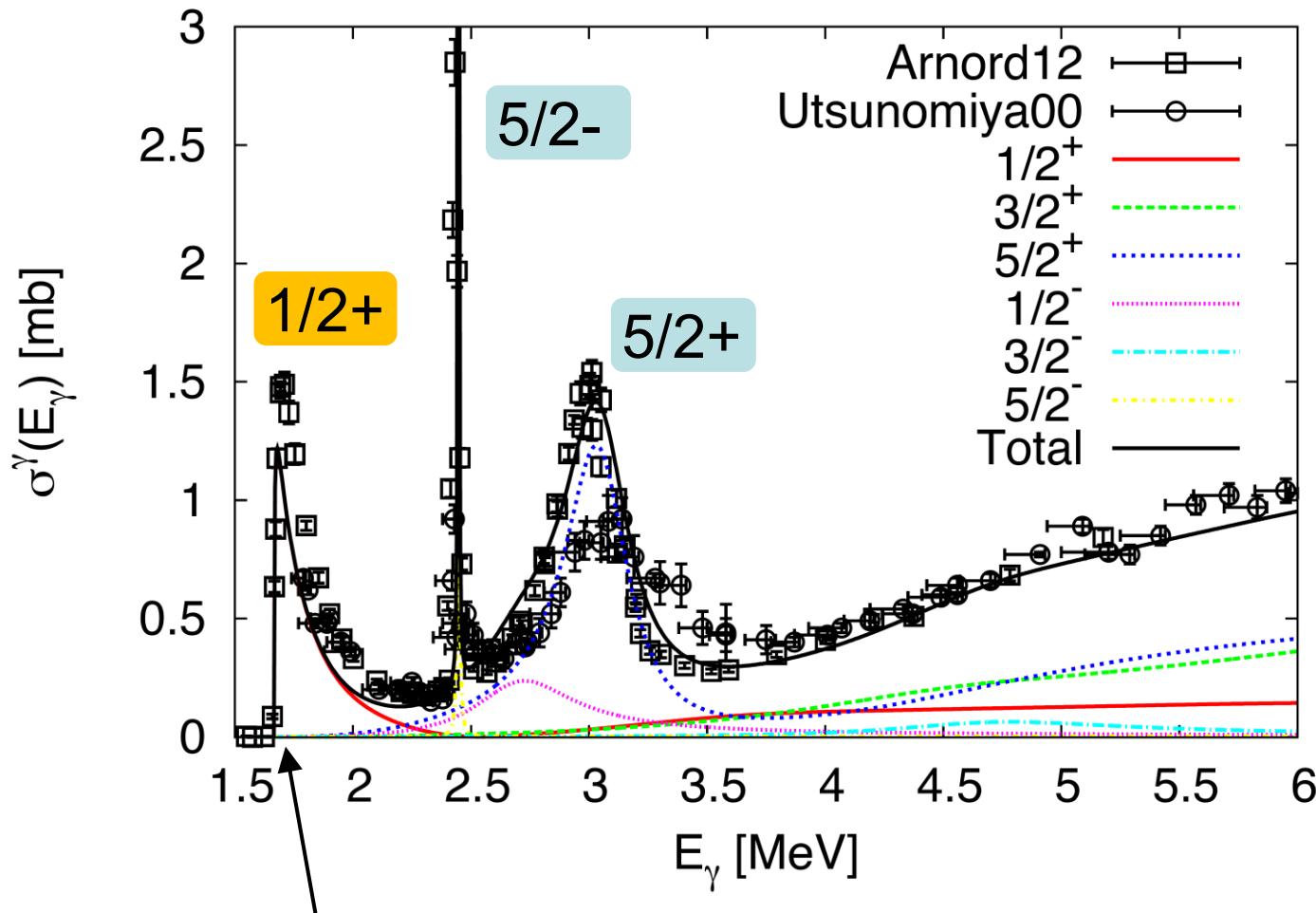
$\alpha+\alpha+n$ 3-body model (OCM+CSM)

No 3-body resonance



3-body potential
(1 range Gaussian
with hyper-radius)

${}^9\text{Be}$: Photodisintegration to $\alpha+\alpha+n$

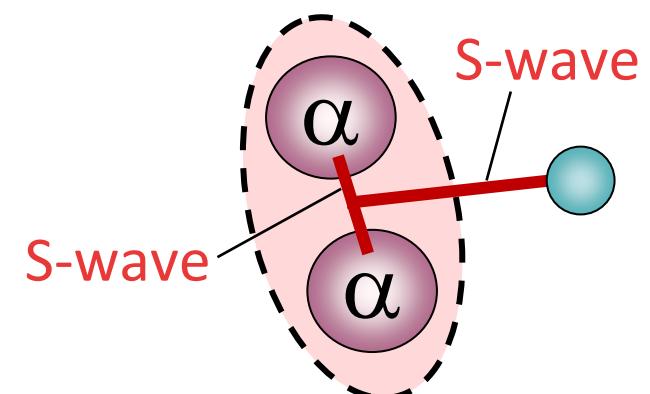


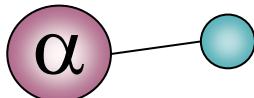
Low-lying $1/2^+$ enhancement suggests the ${}^8\text{Be}+n$ **S-wave virtual state** above the $\alpha+\alpha+n$ threshold.

Reaction rate of $\alpha(\alpha n, \gamma){}^9\text{Be}$ in supernova

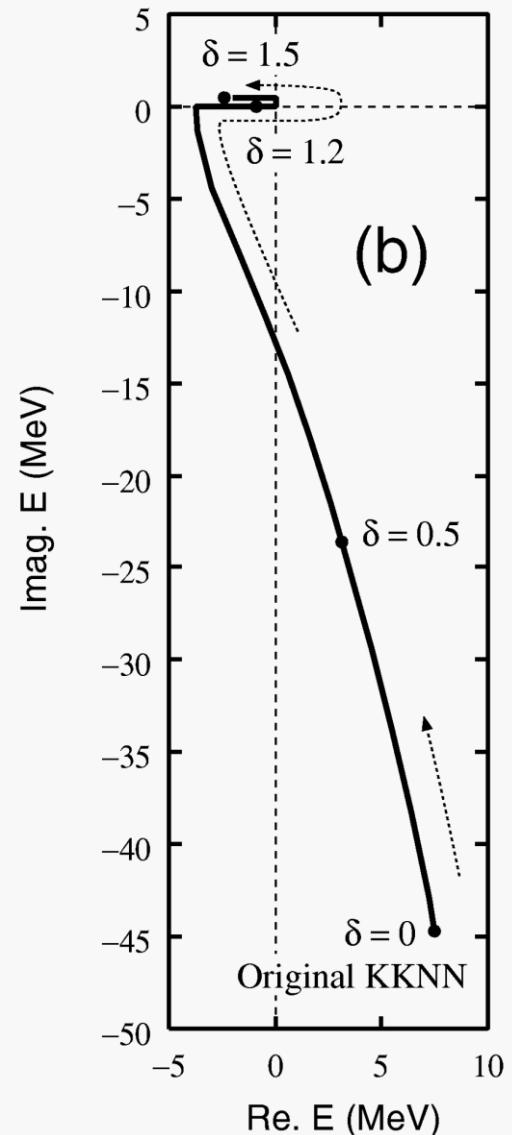
$\alpha+\alpha+n$ OCM

Kikuchi, Katō, Odsuren, T.Myo, Aikawa
PRC 93 (2016) 054605

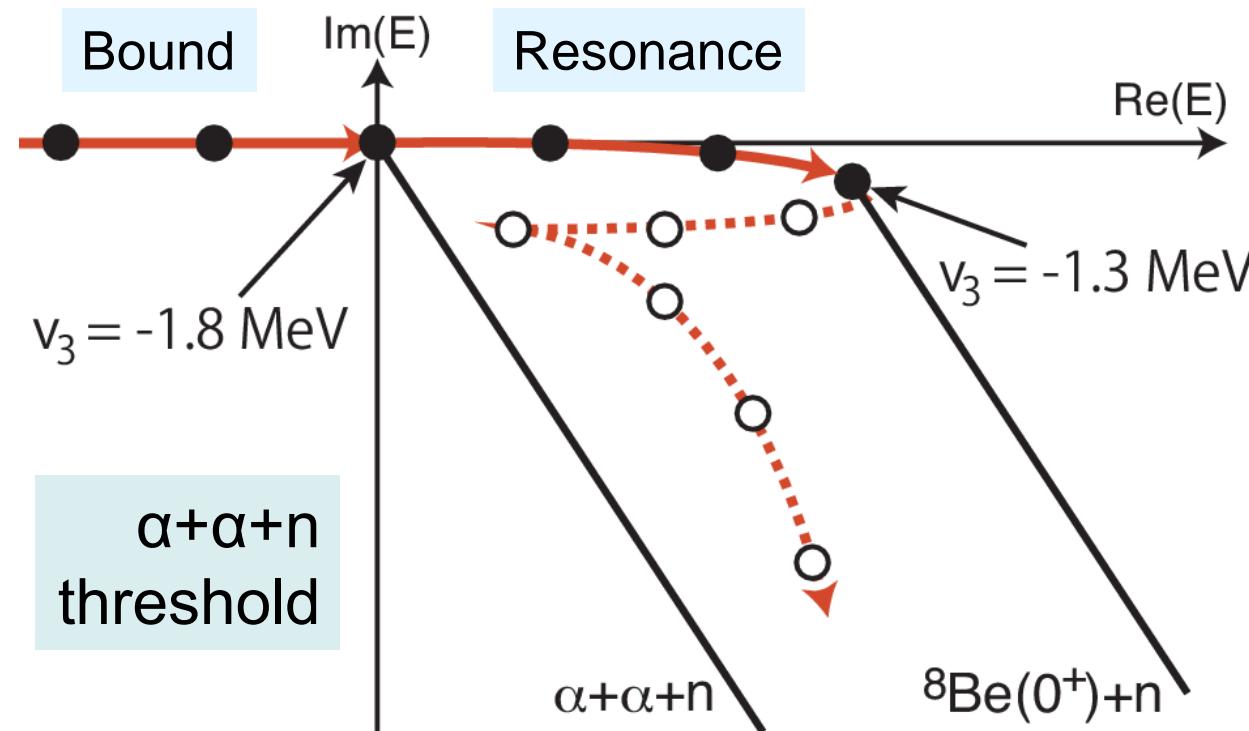




$\alpha+n$ (s-wave)



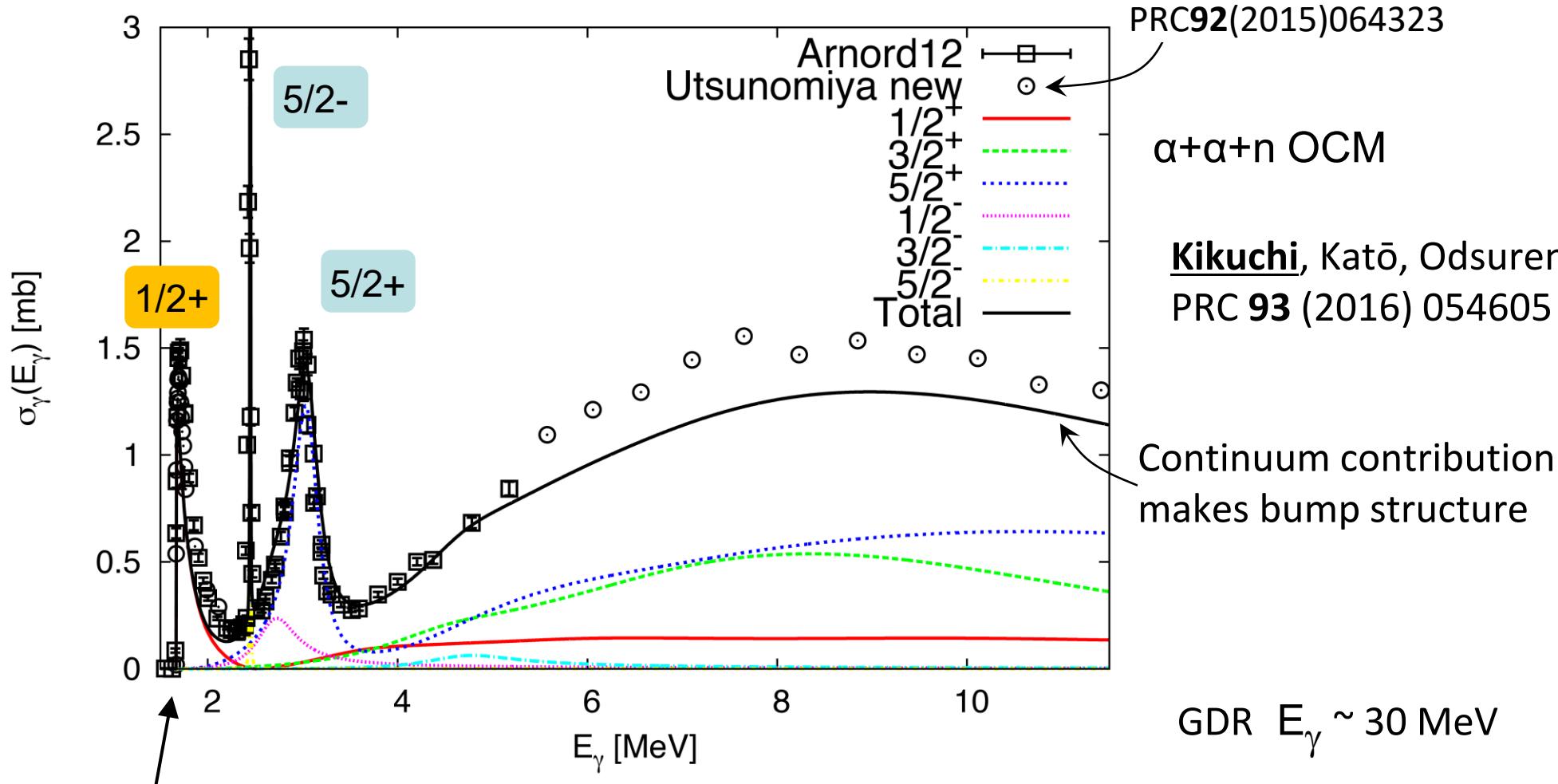
Pole trajectory of ${}^9\text{Be}$ $1/2^+$ by analytical continuation of interaction



S-wave pole in $\alpha+n$, ${}^9\text{Li}+n$ with Jost function method

H. Masui, S. Aoyama, T. Myo, K. Kato, K. Ikeda, Nucl. Phys. A 673 (2000) 207.

${}^9\text{Be}$: Photodisintegration to $\alpha+\alpha+n$



Low-lying $1/2^+$ enhancement suggests the ${}^8\text{Be}+n$
S-wave virtual state above the $\alpha+\alpha+n$ threshold.

Soft dipole resonance in ${}^8\text{He}$ and ${}^8\text{C}$

- Physical Review C104 (2021) 044306 Energy spectra of ${}^8\text{He}$ and ${}^8\text{C}$
- Physical Review C106 (2022) L021302 Soft dipole mode of ${}^8\text{He}$ (letter)
- Prog. Theor. Exp. Phys. 2022 (2022) 103D01
Soft dipole resonance in ${}^8\text{He}$ (full paper)
- Physical Review C107 (2023) 034305 Soft dipole resonance in ${}^8\text{C}$

Ikeda's idea on Soft Dipole Resonance

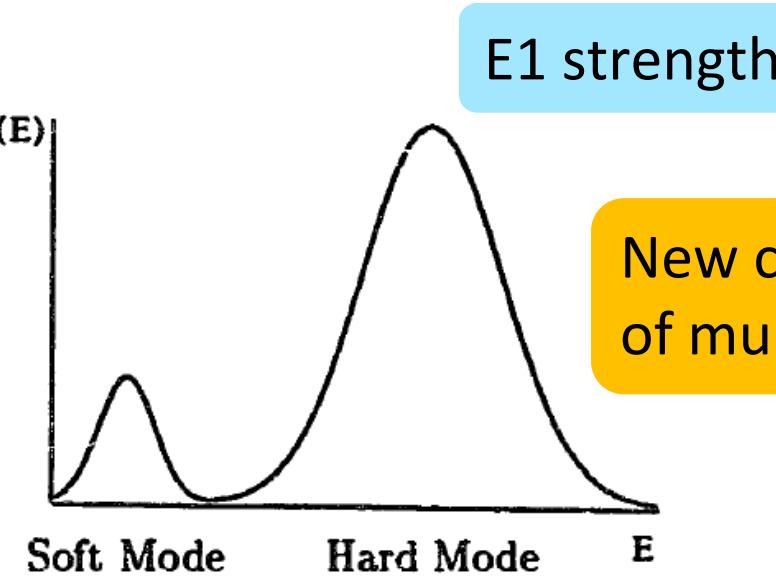
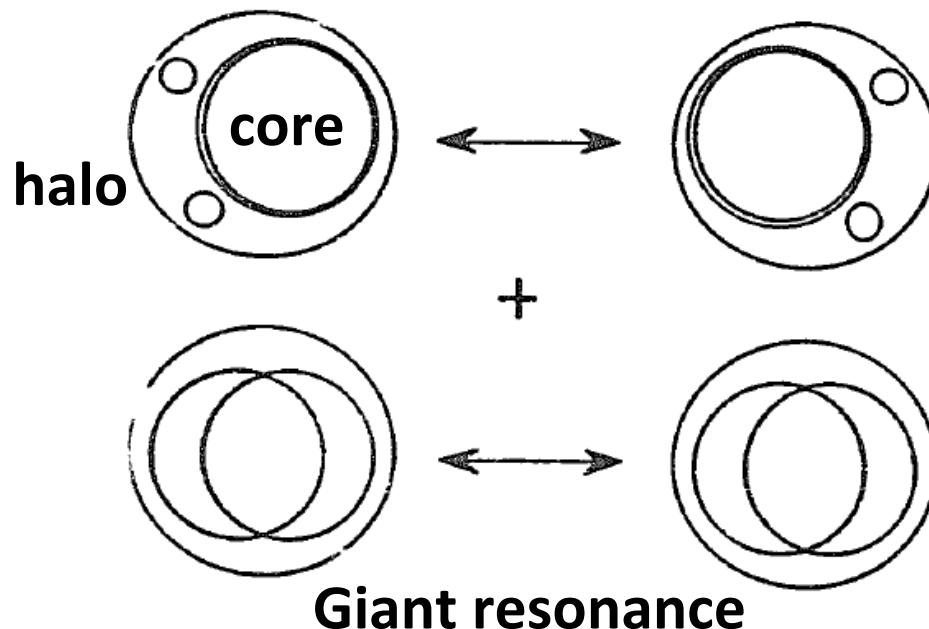
STRUCTURE OF NEUTRON RICH NUCLEI —A Typical Example of the Nucleus ^{11}Li —

NPA538 (1992) 355

Kiyomi IKEDA 池田 清美

Dipole oscillation of valence neutrons against core

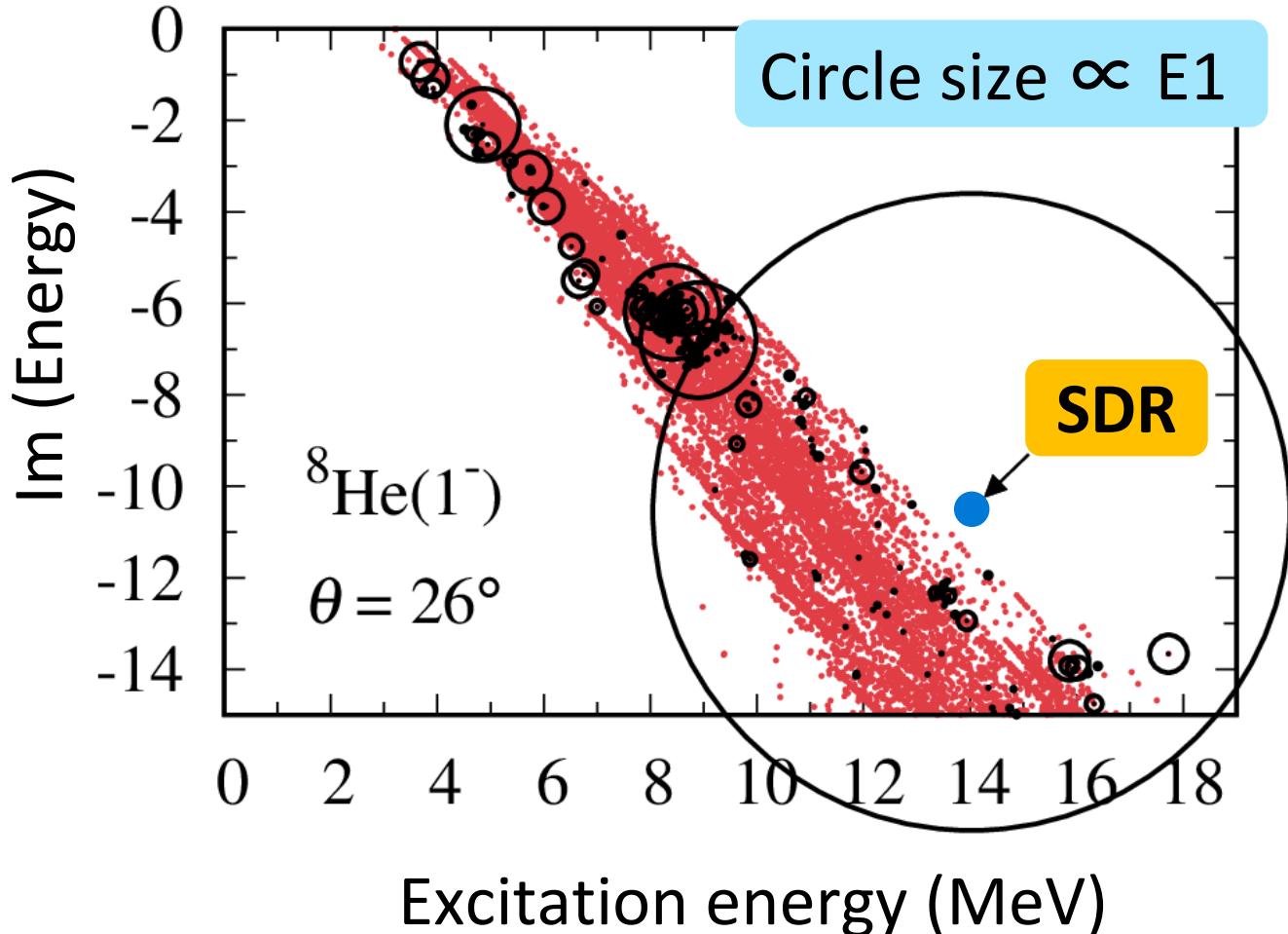
P.G. Hansen and B. Jonson,
Europhys. Lett. **4** (1987) 409.
 $^{11}\text{Li} = {}^9\text{Li} + \text{dineutron}$



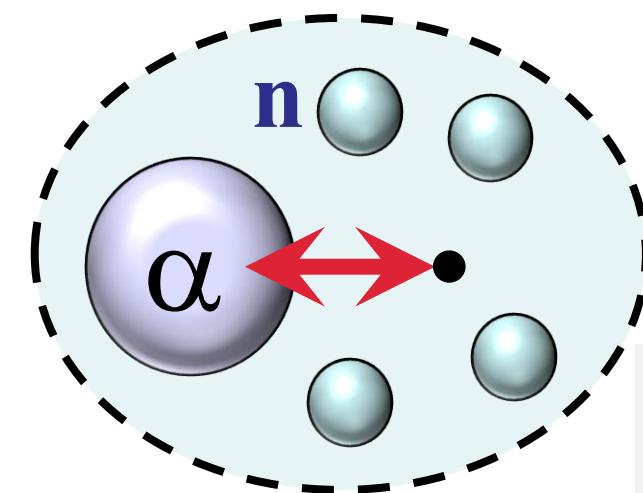
${}^8\text{He}$: Candidate of Soft Dipole Resonance

complex scaling

E1 transition (pole) = 0.55 ($e^2\text{fm}^2$)



- Cluster Sum-Rule Value
 $B_c(E1) = 1.01 e^2\text{fm}^2$
- $(E_x, \Gamma) = (14, 21) \text{ MeV}$
- α -4n distance : extended to 2.7 fm
- Possibility of SDR of ${}^5\Lambda\text{He} + 4\text{n}$



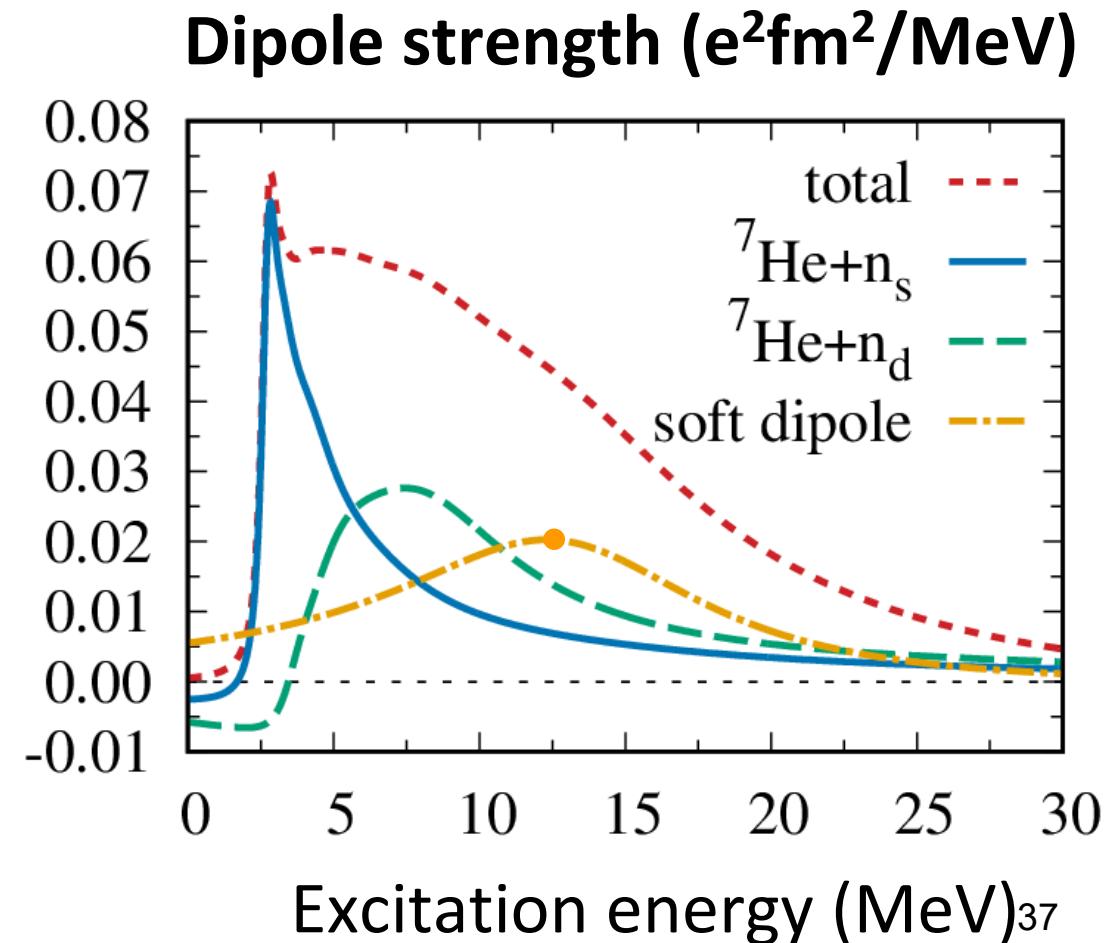
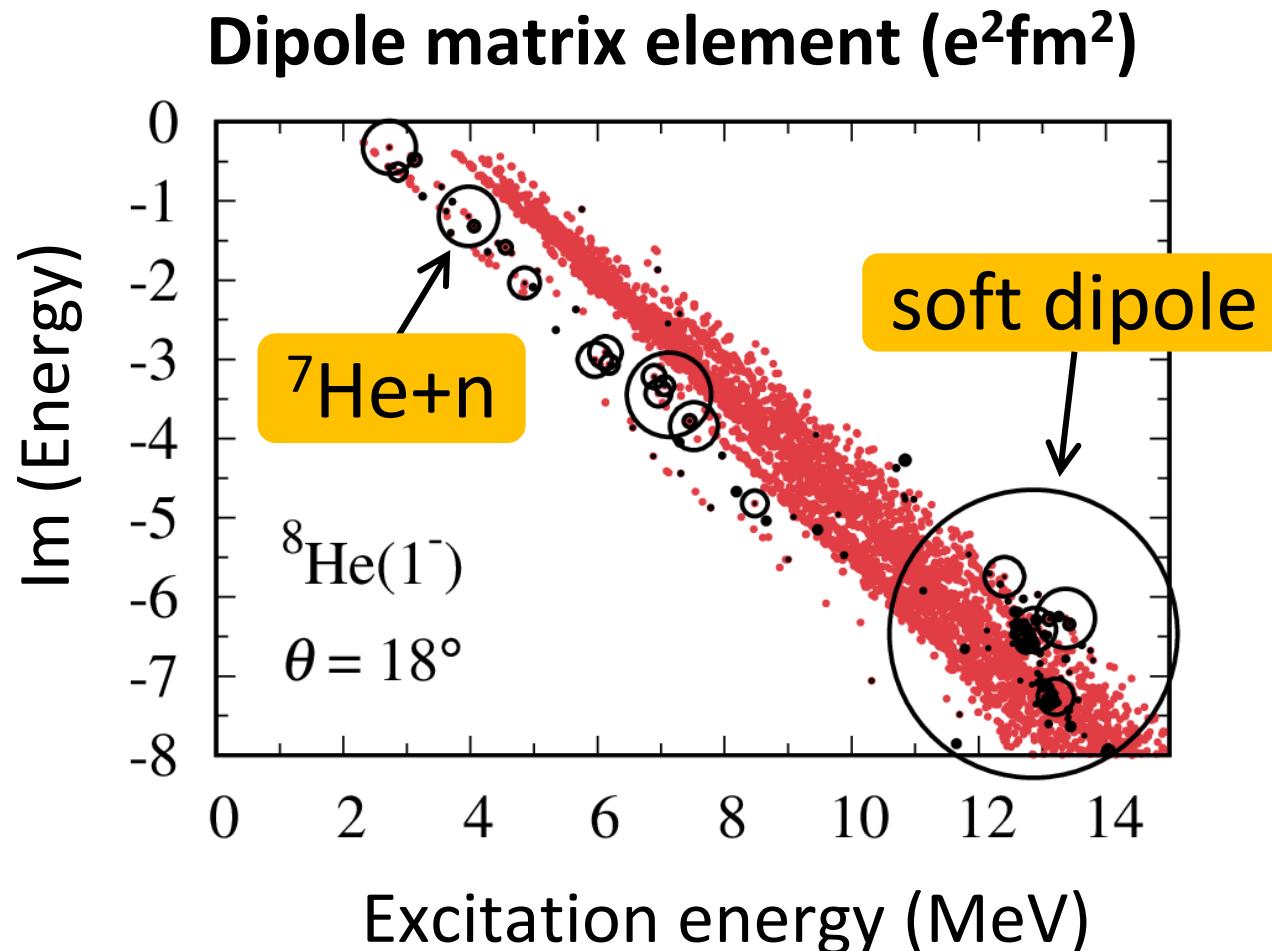
Myo, Odsuren,
Katō, PTEP2023

Dipole strength function of ${}^8\text{He}$

- Zero energy : ${}^8\text{He}_{\text{g.s.}}$

Circle size $\propto E1$

complex scaling



PhD Thesis by Christopher Lehr (TU Darmstadt)

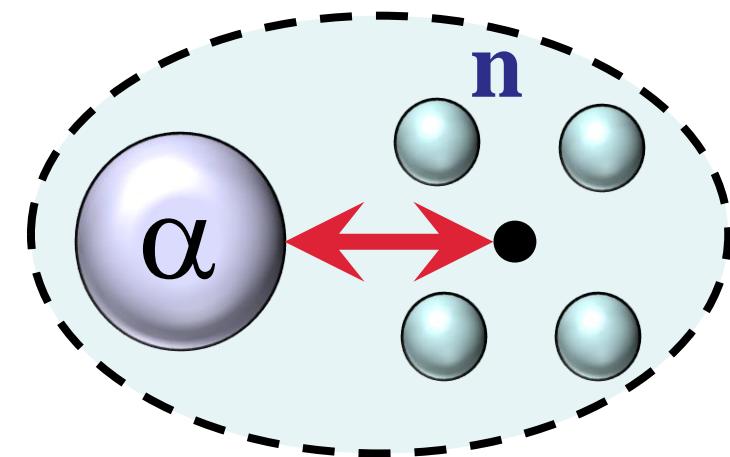
- Title : Low-energy dipole response of the halo nuclei $^{6,8}\text{He}$
- Coulomb breakup experiment of ^8He with Pb target with the help of C target (SAMURAI 37)

Radius of SDR in ${}^8\text{He}$ (fm)

	G.S.	SDR
Matter	2.53	$3.11 + i0.86$
Proton	1.81	$1.97 + i0.26$
Neutron	2.73	$3.41 + i0.99$
α -4n	2.05	$2.67 + i0.84$
4n	2.91	$3.71 + i1.14$

Burgers, Rost, J. Phys. B 29 (1996)
 Dote, Inoue, Myo, PLB784 (2018)
 Myo, Katō, PTEP2020 (2020)
 Myo, Katō, PRC107 (2023)

complex scaling



- α -4n : extend by **0.6 fm**
- 4n : expand by **0.8 fm**

Interpretation of imaginary part : Myo, Katō, PRC 107 (2023)

Summary

- Complex scaling not only for resonance, but also for level density
- Natural extension of completeness relation
- Asymptotically damping condition for resonances
- Apply complex scaling to few-body approach of nuclei
- Many-body unbound states are obtained in unstable nuclei.
- Remaining Problems
 - Many-body virtual state
 - Partial decay width