

有限核の α 凝縮 (特に ^{16}O)

Yasuro Funaki (RIKEN)

Taiichi Yamada (Kanto Gakuin Univ.)

Peter Schuck (IPN, Orsay, Paris-Sud Univ.)

Hisashi Horiuchi (RCNP)

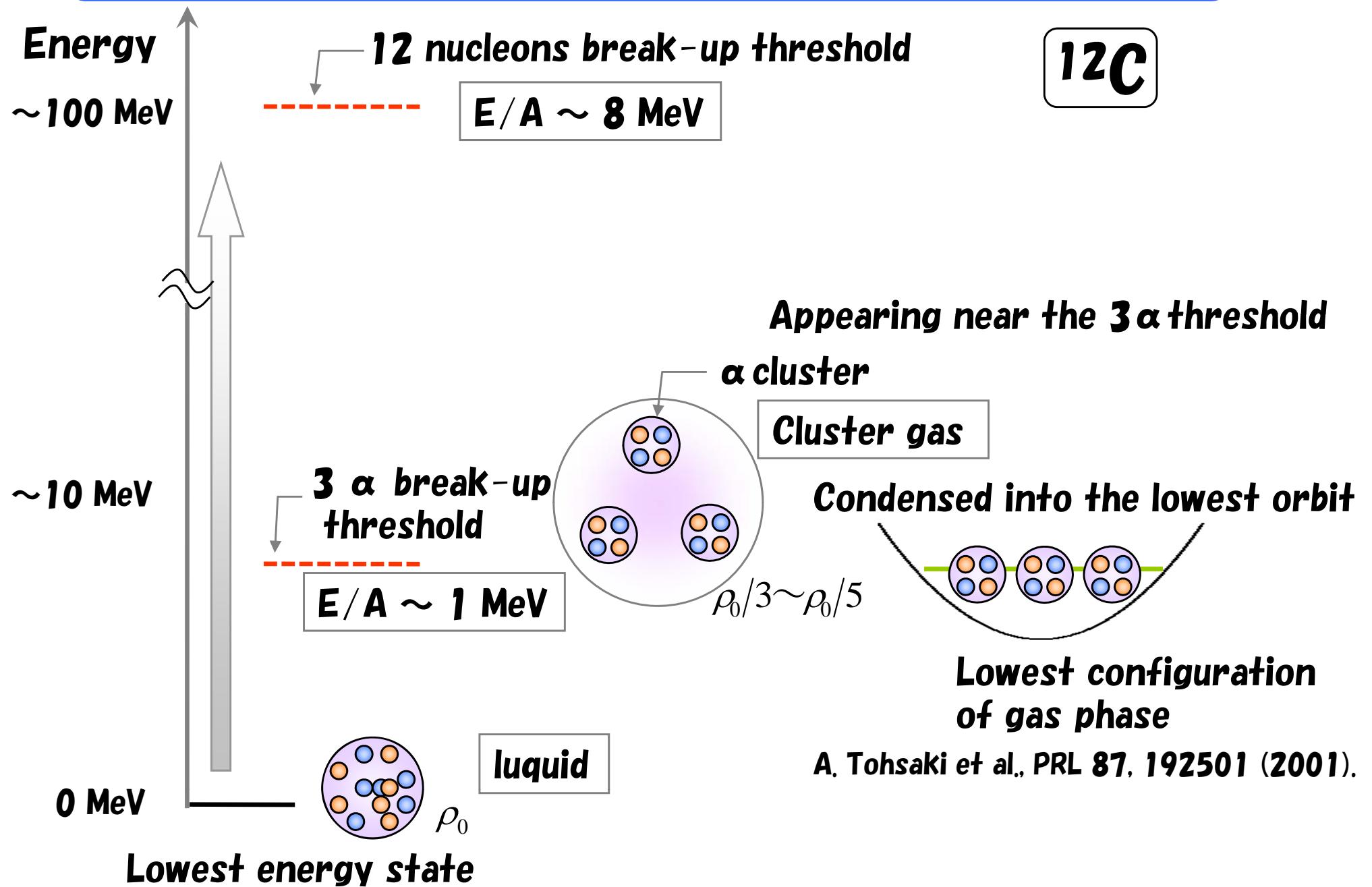
Akihiro Tohsaki (RCNP)

Gerd Röpke (Rostock Univ.)

and with T. Wakasa, M. Takashina, etc.

*RCNP ワークショップ、「核子と中間子の多体問題の統一的描像に向けて」
2007年、12月14日—15日*

Appearing of cluster gas state and ‘‘BEC’’ state in finite nuclei

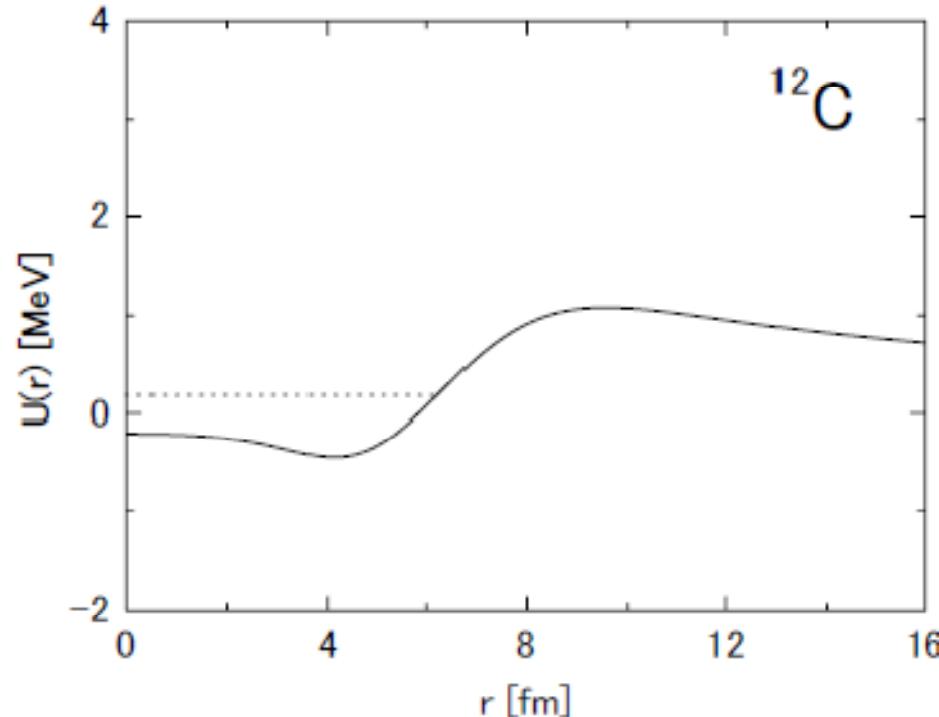


$\alpha - \alpha$ interaction range and potentials α particles feel in the gas states

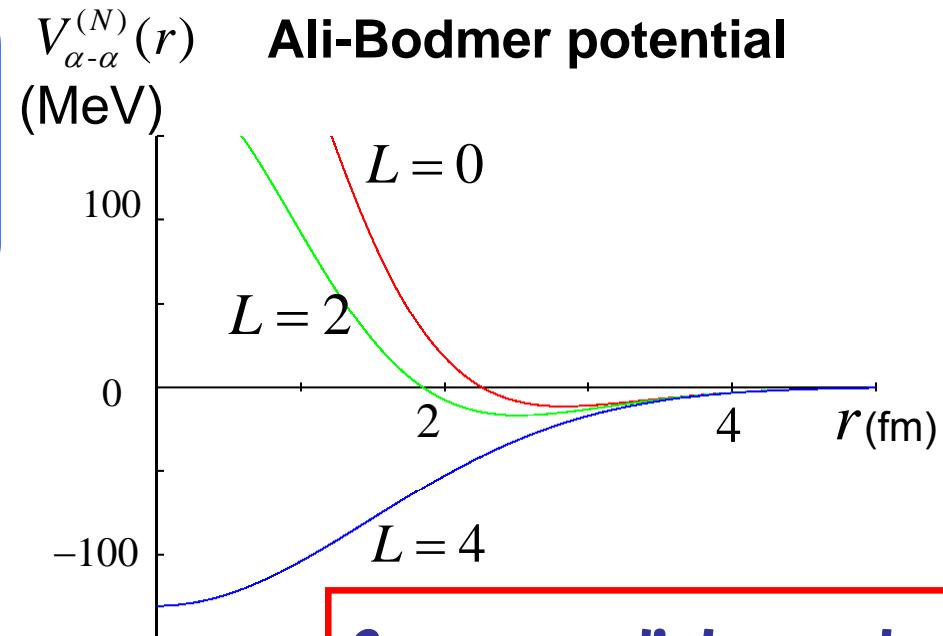
Coulomb barrier

Position is outside the $\alpha - \alpha$ interaction range (~ 4 fm)

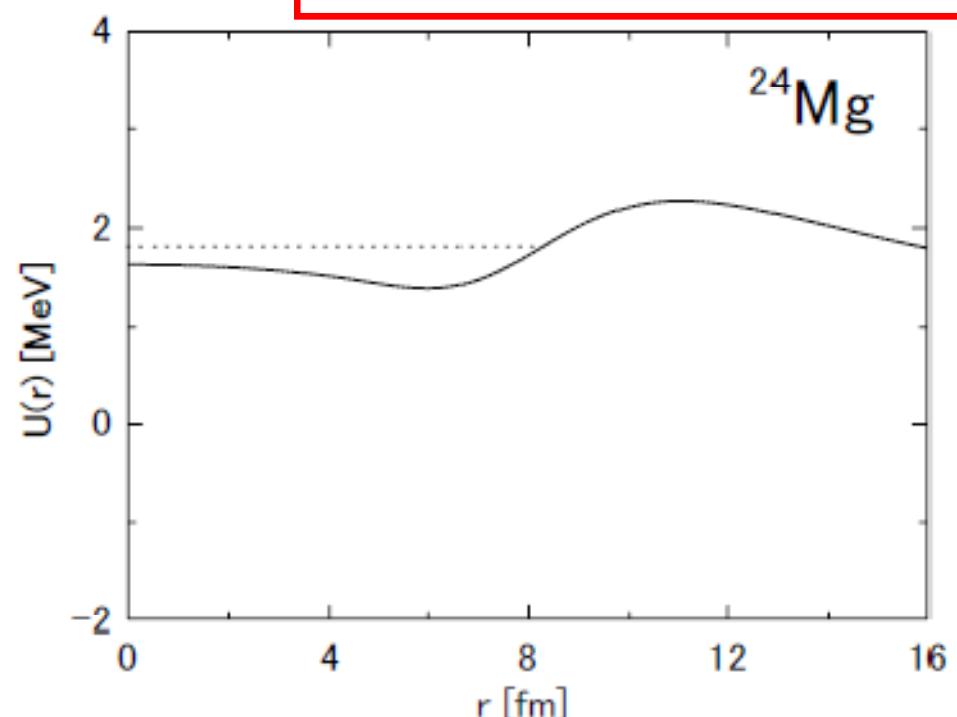
Stabilization of α condensate state



Gross-Pitaevsky eq.



Gas: $\alpha - \alpha$ distance should be more than 4 fm



T. Yamada and P. Schuck, PRC 69, 024309 (2004)

$n\alpha$ condensate wave function (THSR-w.f.)

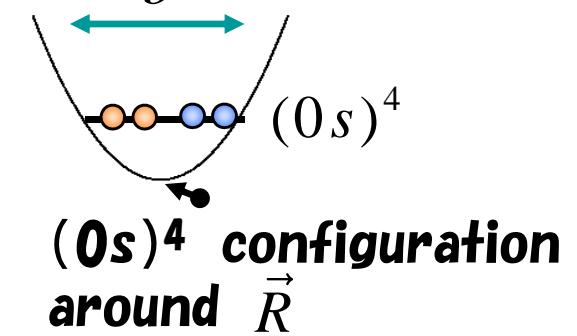
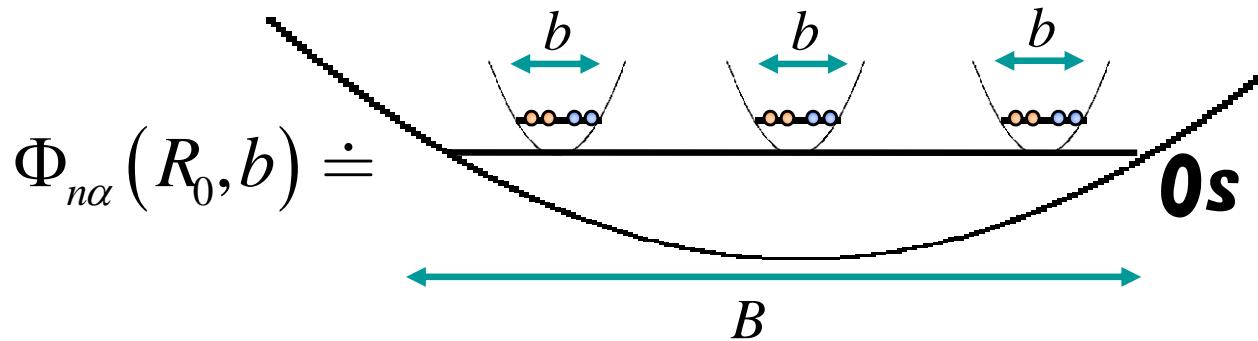
$$\Phi_{n\alpha}(\beta, b) = \mathcal{A} \left\{ \prod_{i=1}^n \left(\exp \left(-\frac{2}{B^2} \vec{X}_i^2 \right) \phi(\alpha_i) \right) \right\} \quad (B^2 = b^2 + 2R_0^2)$$

$$\propto \langle \vec{r}_1 i_1, \dots, \vec{r}_{4n} i_{4n} | (C_\alpha^\dagger)^n | \text{VAC} \rangle$$

$$C_\alpha^\dagger = \int d^3 \vec{R} \exp \left(-\frac{\vec{R}^2}{R_0^2} \right) B_\alpha^\dagger(\vec{R})$$

Brink's wave function

$$\phi(\alpha) \propto \langle \vec{r}_1 i_1, \dots, \vec{r}_4 i_4 | B_\alpha^\dagger(\vec{R}) | \text{VAC} \rangle \doteq$$



Calculation of matrix elements
is owing to Tohsaki's technique

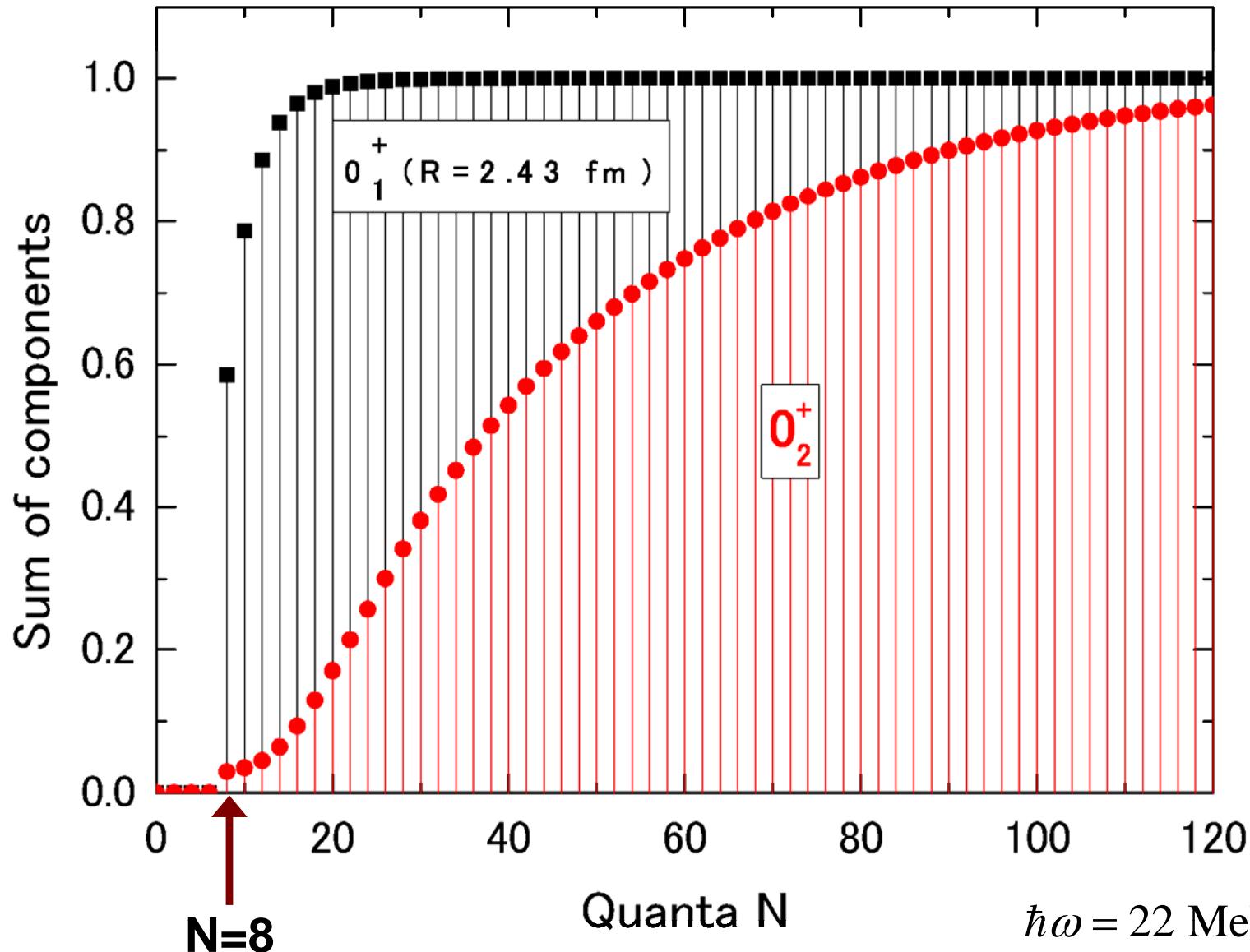
Hill-Wheeler equation (R_0 : generator coordinate, b : fixed)

$$\sum_{R_0} \langle \Phi_{n\alpha}(R_0, b) | H - E^\lambda | \Phi_{n\alpha}(R_0, b) \rangle f_{R_0}^\lambda = 0$$

$$\Psi_{n\alpha}^\lambda = \sum_{R_0} f_{R_0}^\lambda \Phi_{n\alpha}(R_0, b)$$

^{12}C

Expansion of 0^+_1 and 0^+_2 wfs with H.O. basis



The expansion was done wrt relative motions of α 's.

$\hbar\omega = 22 \text{ MeV}$

Calculated by T. Yamada

First example of α condensate state in finite nuclei

Hoyle state (0_2^+ state in ^{12}C (excitation energy : 7.65 MeV))

Indicating 3α condensate character

Microscopic approach (3 α cond. model w.f.)

The Solution of 3 α RGM eq. of motion, RGM

$$\langle \phi^3(\alpha) | H - E | \mathcal{A}[\chi(s, r) \phi^3(\alpha)] \rangle = 0$$

M. Kamimura, NPA 351, 456 (1981).

is almost equivalent to the 3 α cond. w.f.

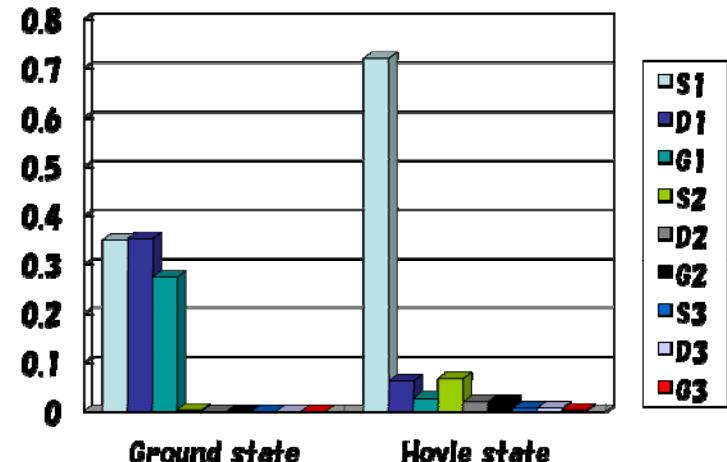
$$\chi(s, r) = \exp\left(-\frac{2}{B^2} \sum_{i=1}^3 (X_i - X_G)^2\right)$$

X_i : c.o.m of α -particle

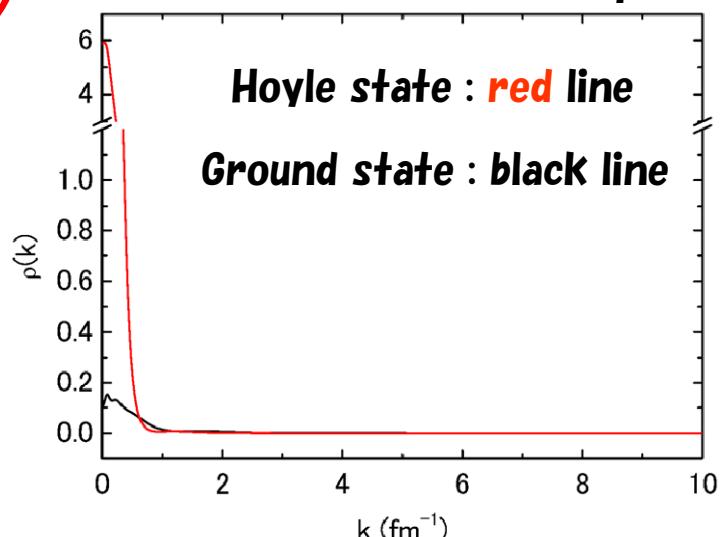
Y. F et al., PRC 67, 051306(R) (2003).

3 α break-up threshold : 7.27 MeV

Occupation probability of α - particle orbit



Momentum distribution of α - particle



- Occupation probability of α - particle orbit

Huge OS occupancy (> 70 %)

- Momentum distribution

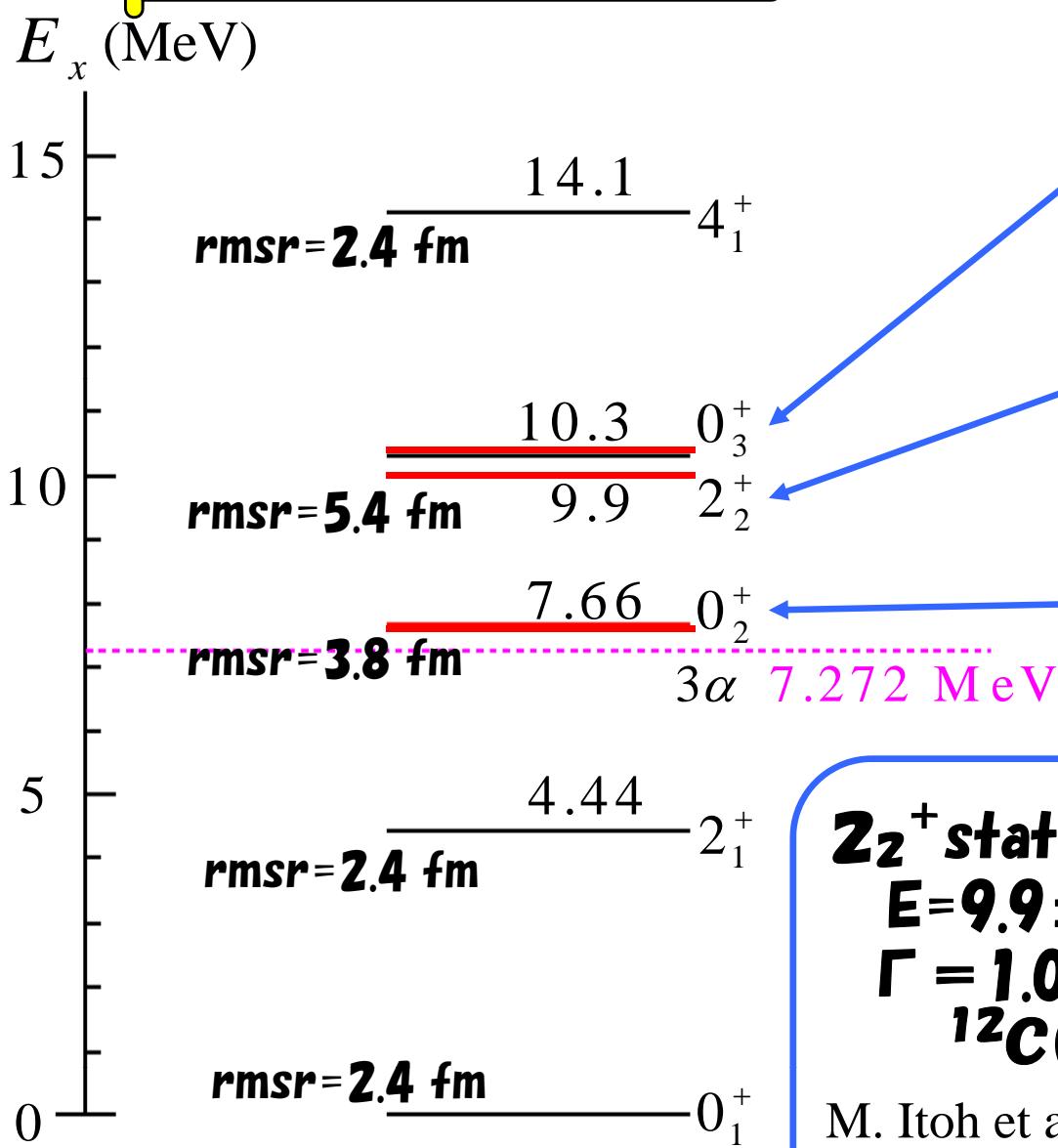
Delta-function-like behavior

H. Matsumura and Y. Suzuki, NPA 739, 238 (2004).

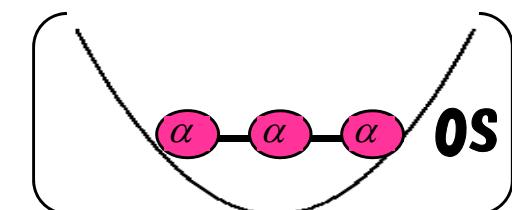
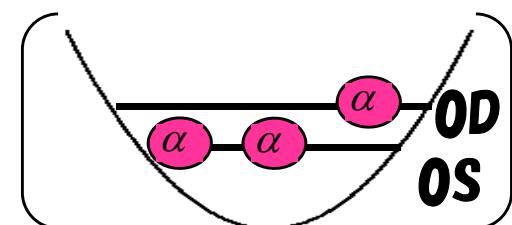
T. Yamada and P. Schuck EPJA 26, 185 (2005).

“BEC” in ^{12}C

Observed levels of ^{12}C



??
A
C. Kurokawa and K. Katō,
PRC 71, 021301 (2005).



2_2^+ state :
 $E = 9.9 \pm 0.3 \text{ MeV}$
 $\Gamma = 1.0 \pm 0.3 \text{ MeV}$
 $^{12}\text{C}(\alpha, \alpha')$

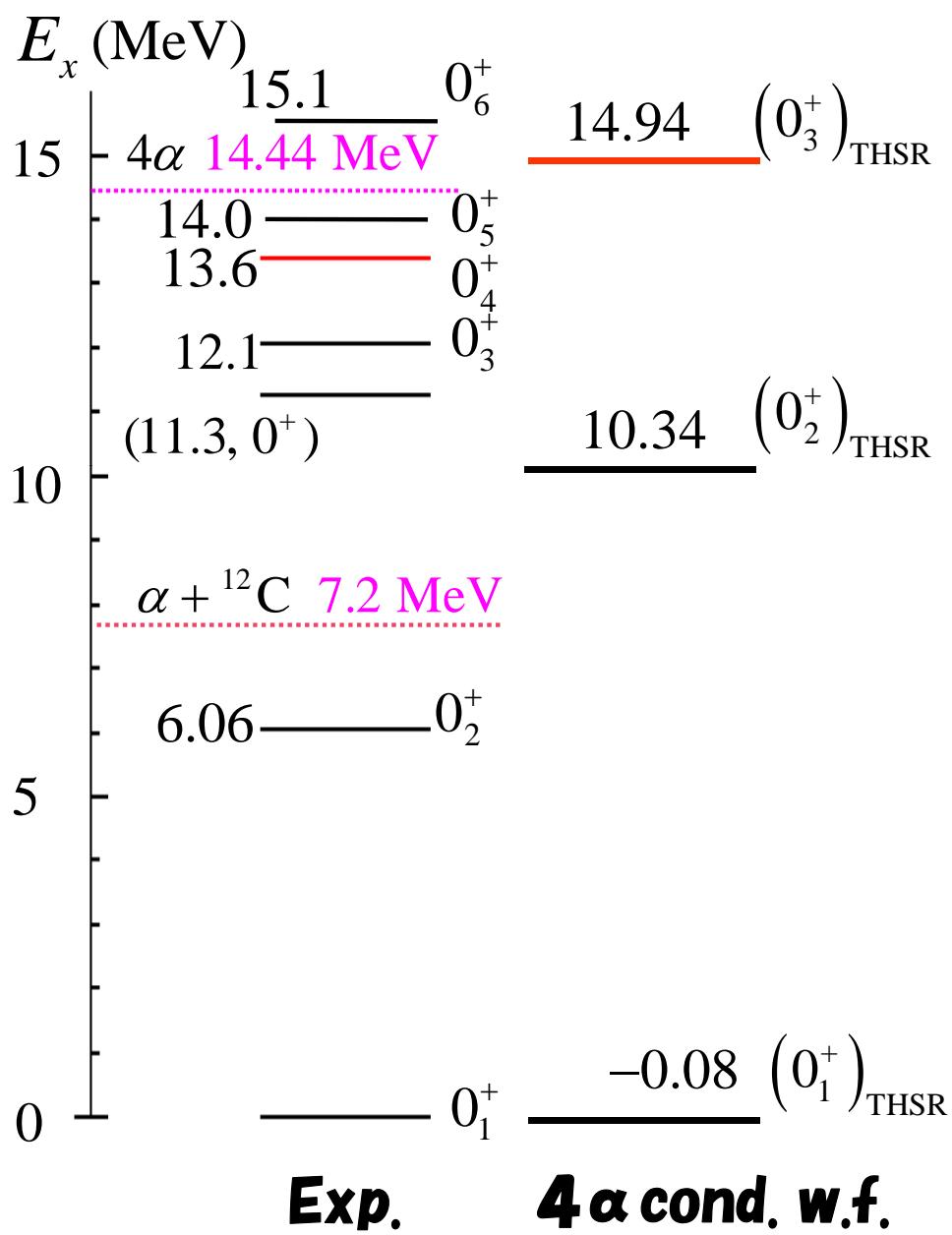
M. Itoh et al., Nucl. Phys. A 738 (2004) 268-272

α cond. + ACCC
 $E = 9.38 \text{ MeV}$
 $\Gamma = 0.64 \text{ MeV}$
Volkov No. 1 force
is adopted

Y. F. et al., EPJA 24, 321 (2005).

First attempt to explore 4α condensate state in $^{16}0$

Low lying 0^+ levels of $^{16}0$

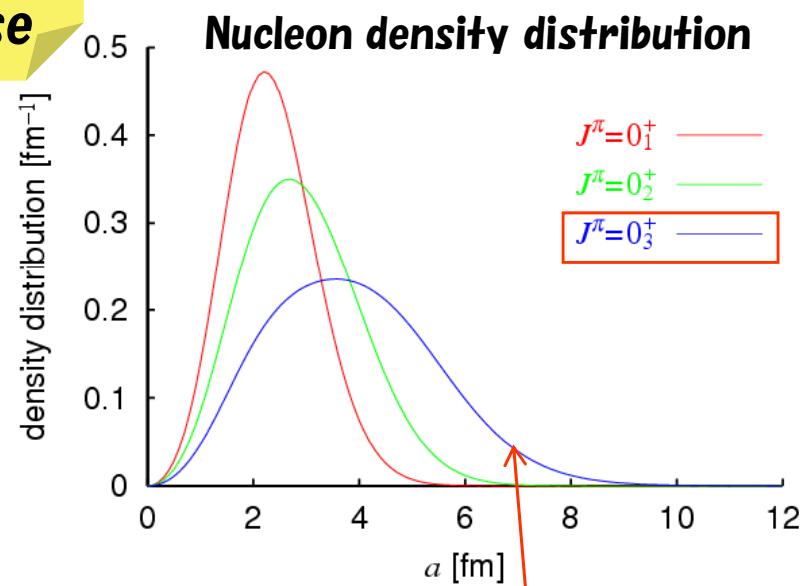


$n\alpha$ condensate model w.f. (THSR-w.f.)

$$\Phi_{n\alpha} = \mathcal{A} \left\{ \prod_{i=1}^n \left(\exp \left(-\frac{2}{B^2} X_i^2 \right) \phi(\alpha_i) \right) \right\}$$

A. Tohsaki, H. Horiuchi, P. Schuck and G. Röpke,
PRL 87, 192501 (2001).

$n=4$ case



$(0_3^+)_\text{THSR}$: Very dilute density.
~~4 α condensate state~~

$\alpha + {}^{16}\text{O}$ inelastic scattering

0_5^+ state:

A candidate of 4 α condensate

$E = 13.6 \text{ MeV}$

$\Gamma = 0.8 \text{ MeV}$

${}^{16}\text{O}(\alpha, \alpha')$ Wakasa et al.

The result of the calculation is consistent with the experimental data.

The 0^+ state at $E_x = 13.5 \text{ MeV}$ can be assigned to the four- α condensed state.

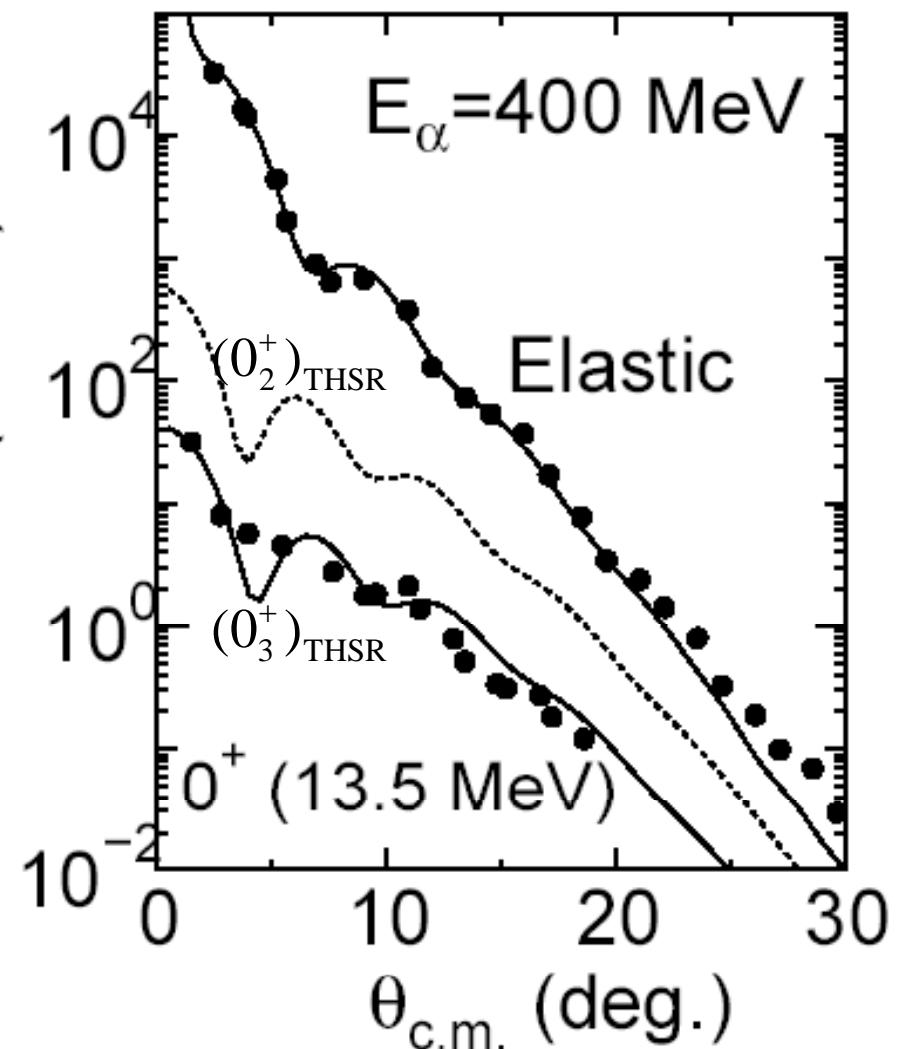
The 0^+ state wave function obtained at $E_x = 10.3 \text{ MeV}$ leads to a largely different absolute value.

$(0_3^+)_\text{THSR}$:

$E = 14.9 \text{ MeV}$

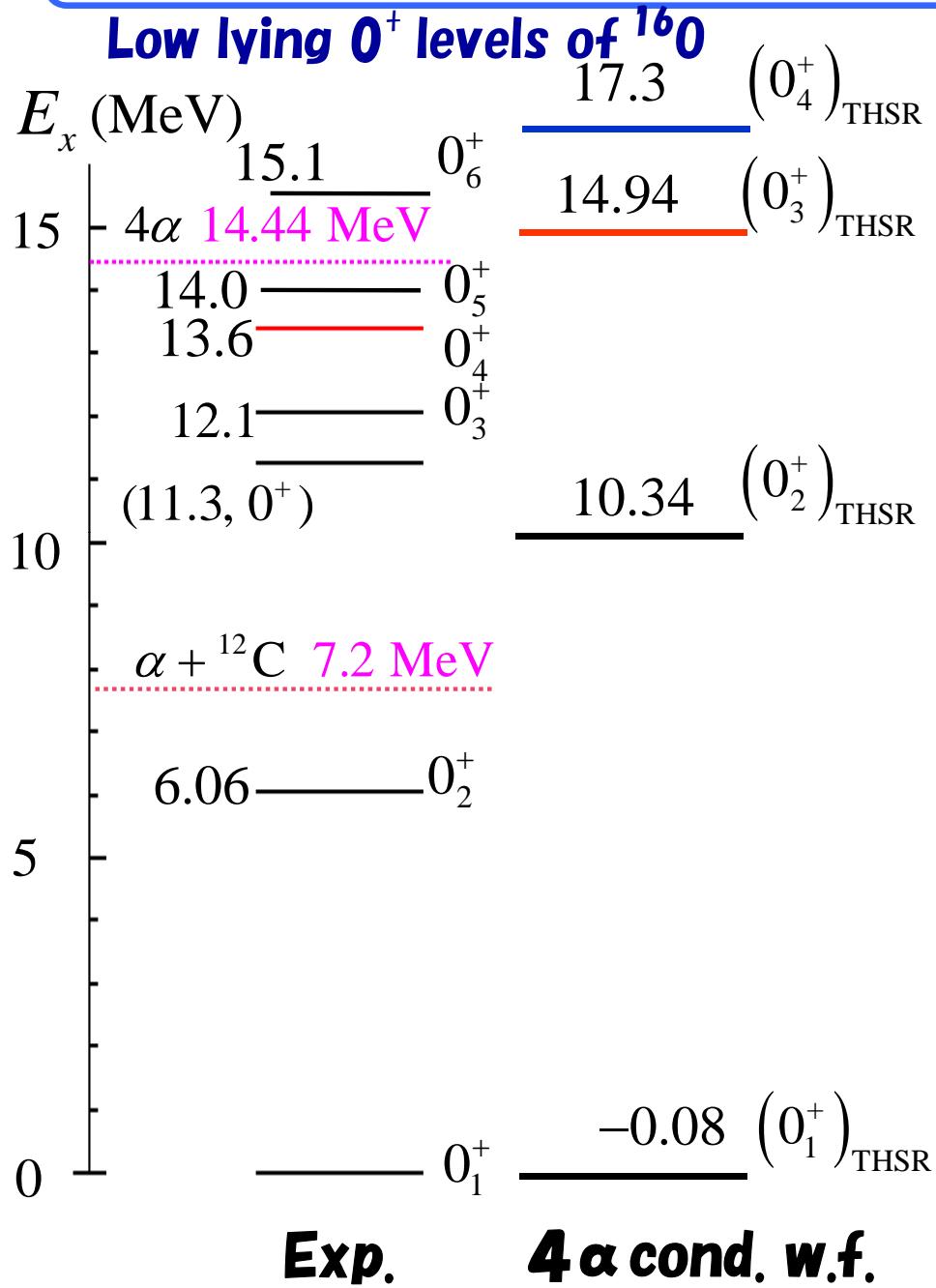
$\Gamma = 1.5 \text{ MeV}$

(based on R-matrix theory)



T. Wakasa, E. Ihara, M. Takashina and Y. F. et al,
PLB 653, 173 (2007).

First attempt to explore 4α condensate state in $^{16}0$

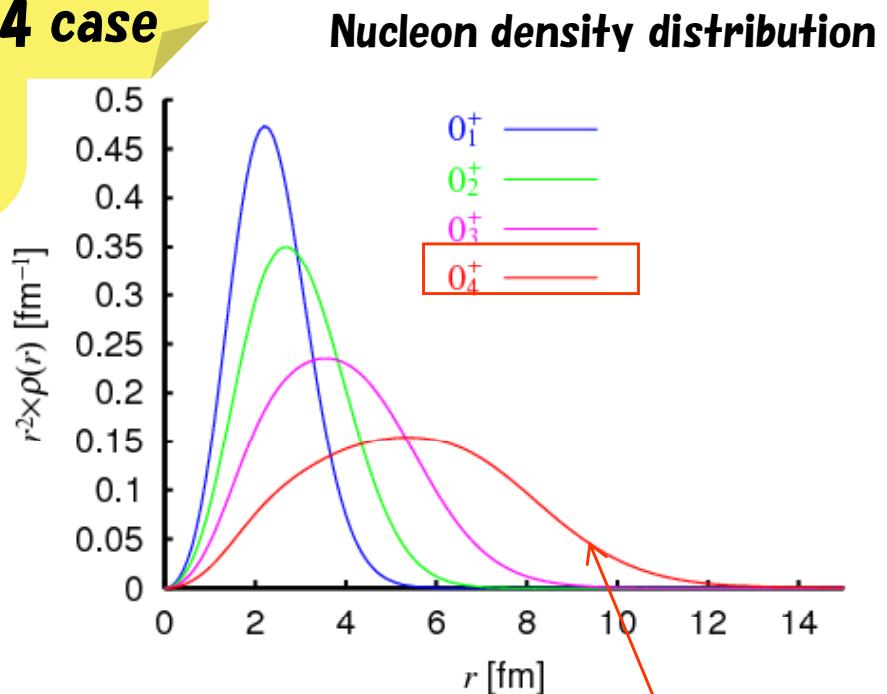


$n\alpha$ condensate model w.f. (THSR-w.f.)

$$\Phi_{n\alpha} = \mathcal{A} \left\{ \prod_{i=1}^n \left(\exp \left(-\frac{2}{B^2} X_i^2 \right) \phi(\alpha_i) \right) \right\}$$

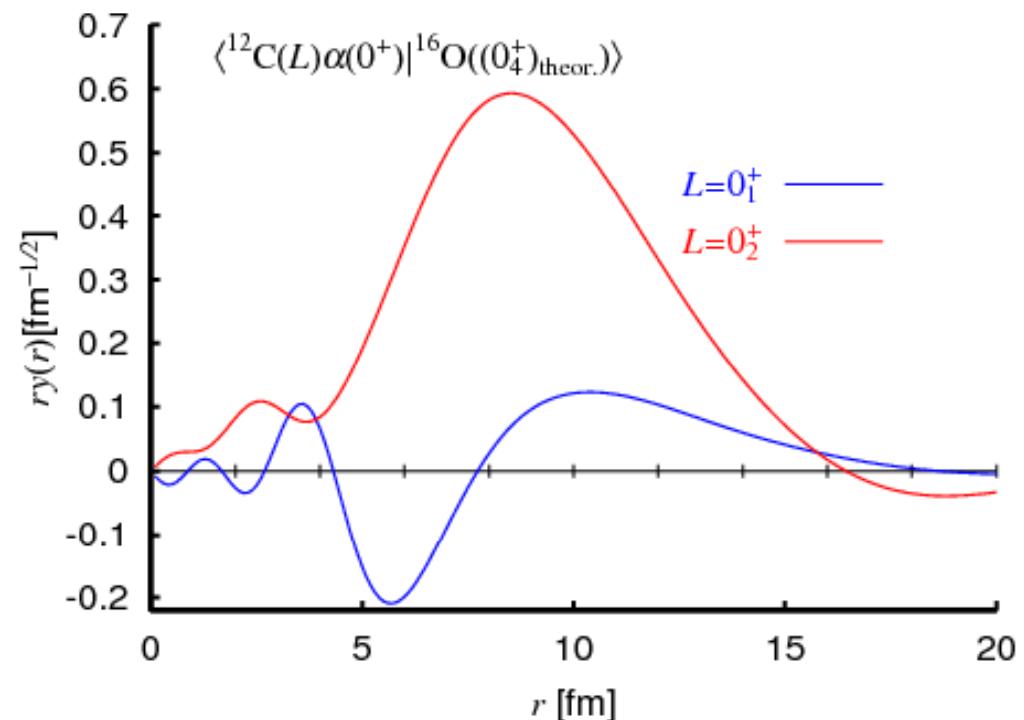
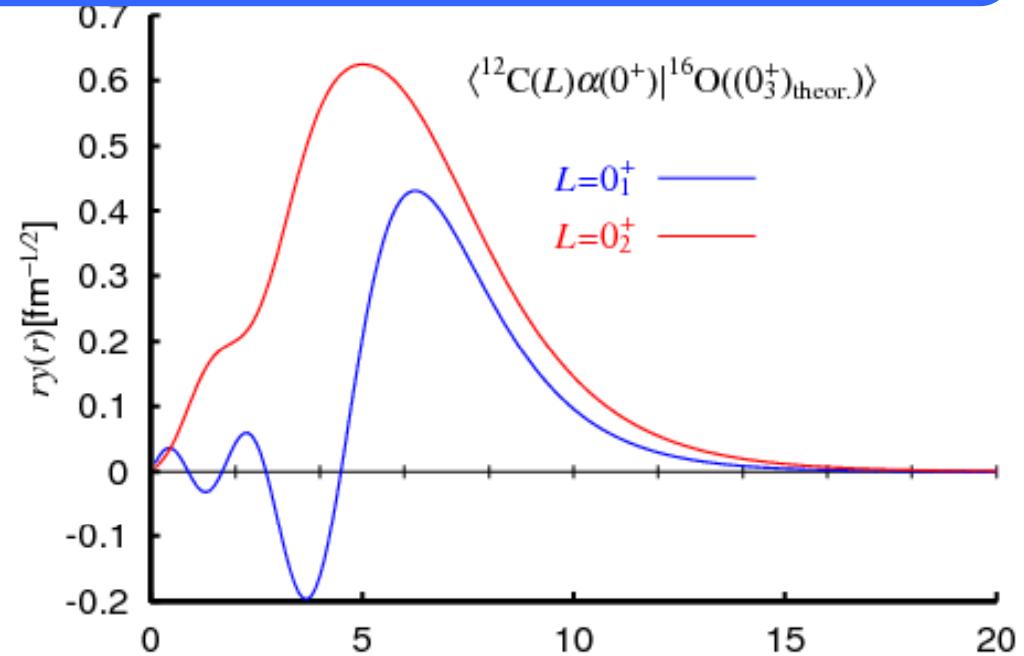
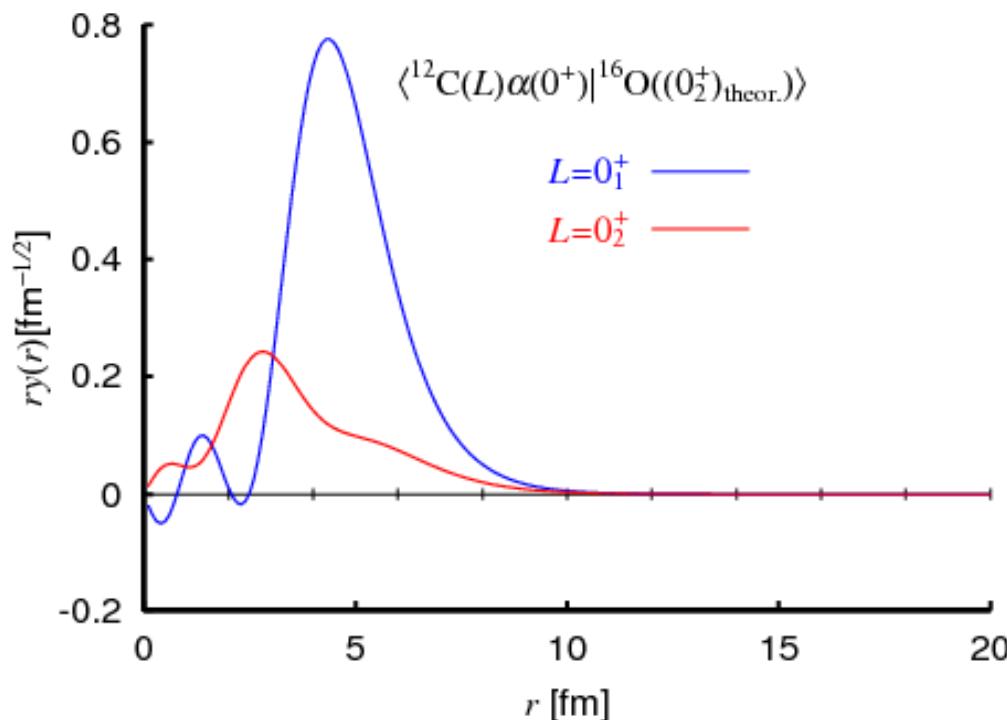
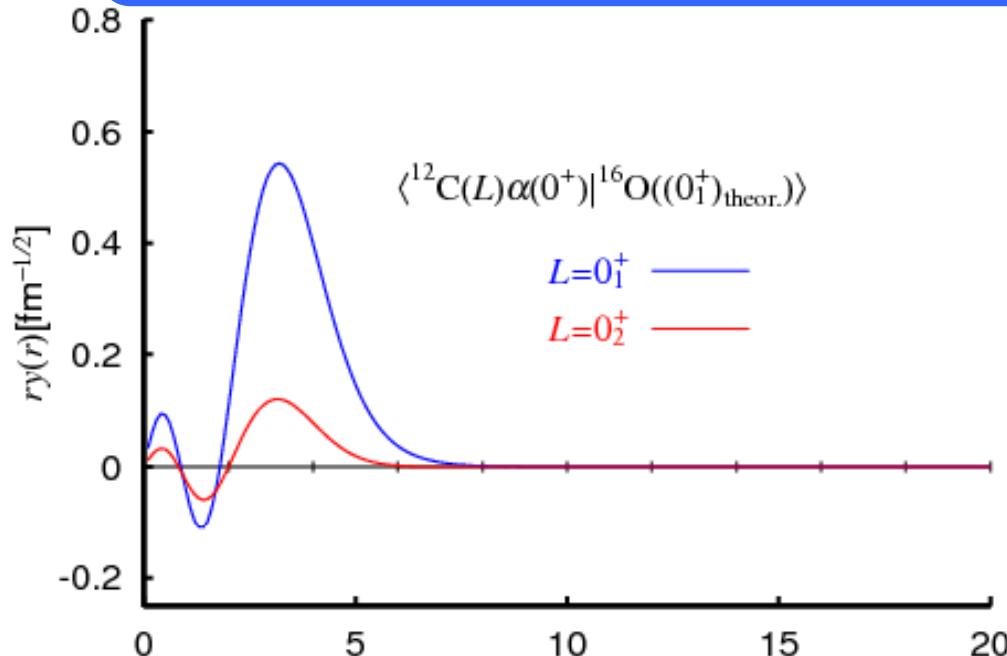
A. Tohsaki, H. Horiuchi, P. Schuck and G. Röpke,
PRL 87, 192501 (2001).

$n=4$ case

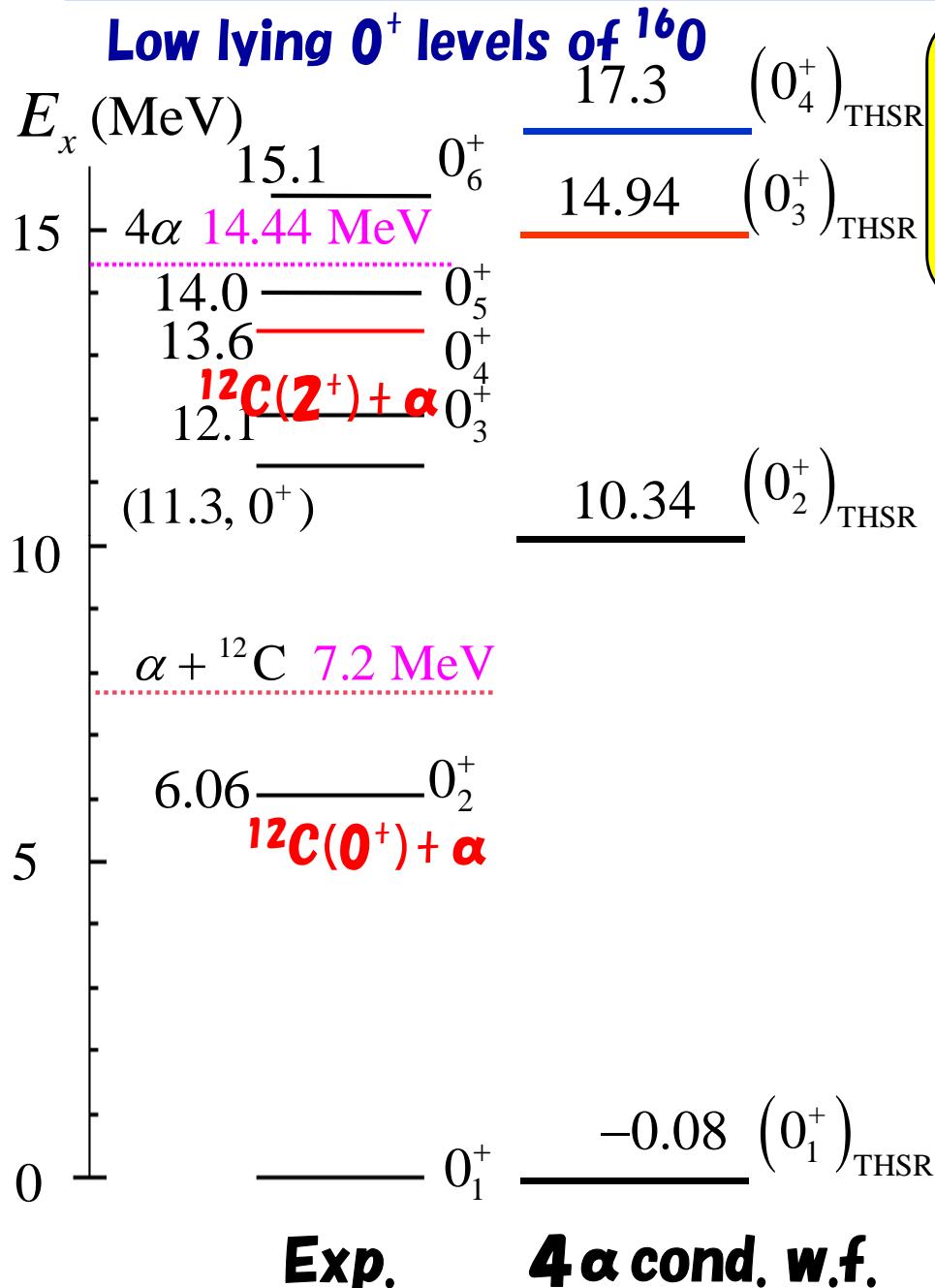


$(0_4^+)_\text{THSR}$: Very much dilute density.
 4α condensate state

**Reduced width amplitudes of 0_1^+ - 0_4^+ states obtained with THSR
w.f.(overlap amplitude between α plus $^{12}\text{C}(0_1^+ \text{ or } 0_2^+)$ and ^{16}O wfs)**



Motivation for 4α OCM (Orthogonality Condition Model)



$n \alpha$ cond. w. f. (microscopic)

$$\Phi_{n\alpha}(R_0, b) = \mathcal{A} \left\{ \prod_{i=1}^n \left(\exp \left(-\frac{2}{B^2} \vec{X}_i^2 \right) \phi(\alpha_i) \right) \right\}$$

- Need to check the existence of 4α cond. state in larger (4α) model space.
- 4α cond. w. f. may have a difficulty to represent the ${}^{12}\text{C} + \alpha$ structure.

4α cond. w.f. hardly describes these states

- $0^+_2(\alpha + {}^{12}\text{C}(0^+))$ Δ
- $(0^+_3) ?$
- $0^+_4(\alpha + {}^{12}\text{C}(2^+))$ \times

Need to solve full 4α problem,
 4α OCM (semi microscopic)

Fully solving 4 α -particles relative motions (4 α OCM)

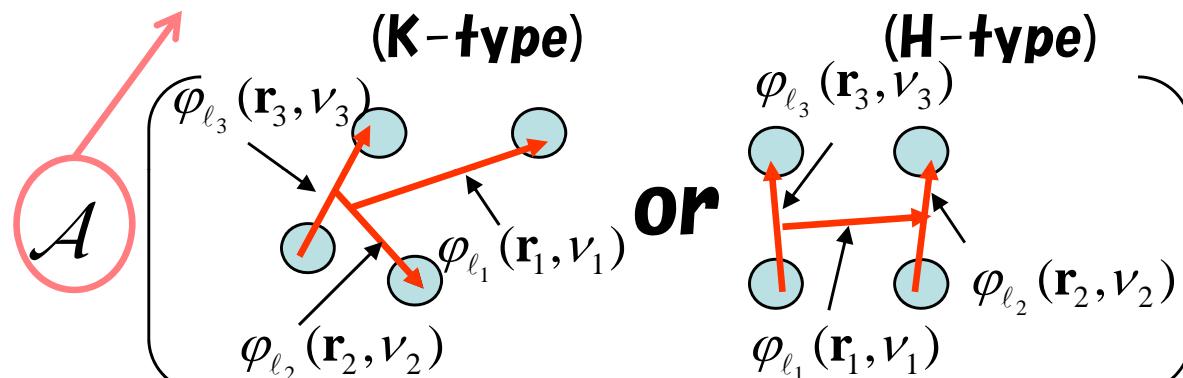
Present: Larger model space

$$\varphi_{\ell m}(\mathbf{r}, \nu) = N_\ell(\nu) r^\ell \exp(-\nu r^2) Y_{\ell m}(\mathbf{r})$$

Gaussian basis (GEM)

E. Hiyama et al. Prog. Part. Phys. 51, 223(2003).

Approximately taken into account



Adopted angular momentum channels, totally 10 channels

K-type and H-type, $[[l_1, l_2]_{l_2}, l_3]_L : [[0, 0]_0, 0]_0 [[2, 0]_2, 2]_0 [[0, 2]_2, 2]_0 [[2, 2]_0, 0]_0$

Hamiltonian

$$H = T + \sum_{i < j} \left[V_{2\alpha}(r_{ij}) + V_{2\alpha}^{Coul}(r_{ij}) \right] + V_{3\alpha} + V_{4\alpha} + V_{Pauli}$$

$V_{2\alpha}(r)$: 2-body force (folding MHN), $V_{2\alpha}^{Coul}$: Coulomb force

$V_{3\alpha}$ & $V_{4\alpha}$: phenomenological 3-body and 4-body forces (repulsive)

V_{pauli} : pauli operator (to eliminate 0s, 1s, 0d w.r.t. $\alpha - \bar{\alpha}$ rel. motion)

$^{12}\text{C} + \alpha$: succeeded
dilute 4α : not reproduced

4α OCM(H, O, basis)

K. Fukatsu and K. Kato, PTP 87, 151 (1992).

$^{12}\text{C} \pm \alpha$ coupled channel OCM

(H, O, basis)

**Y. Suzuki, PTP 55, 1751 (1976);
56, 111 (1976).**

$[[0,1]_1,1]_0$ $[[2,1]_1,1]_0$ **Added in only K-type**

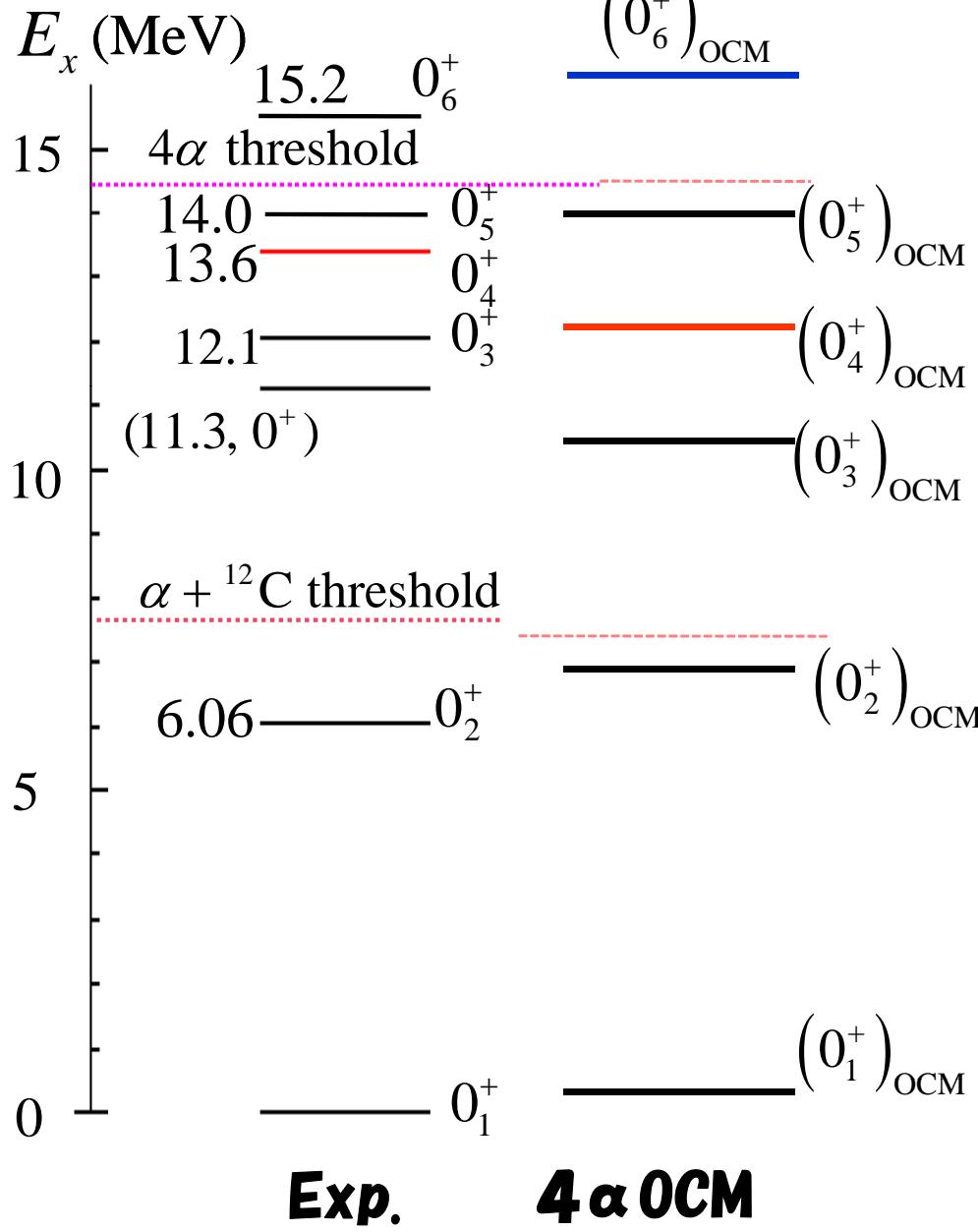
Energies from 4α threshold

	Cal. (MeV)	Exp. (MeV)
$^{12}\text{C(g.s.)}$	- 7.32	- 7.27
$^{16}\text{O(g.s.)}$	- 14.2	- 14.44

Correct $^{12}\text{C}(\text{g.s.}) + \alpha$ and 4α threshold energies

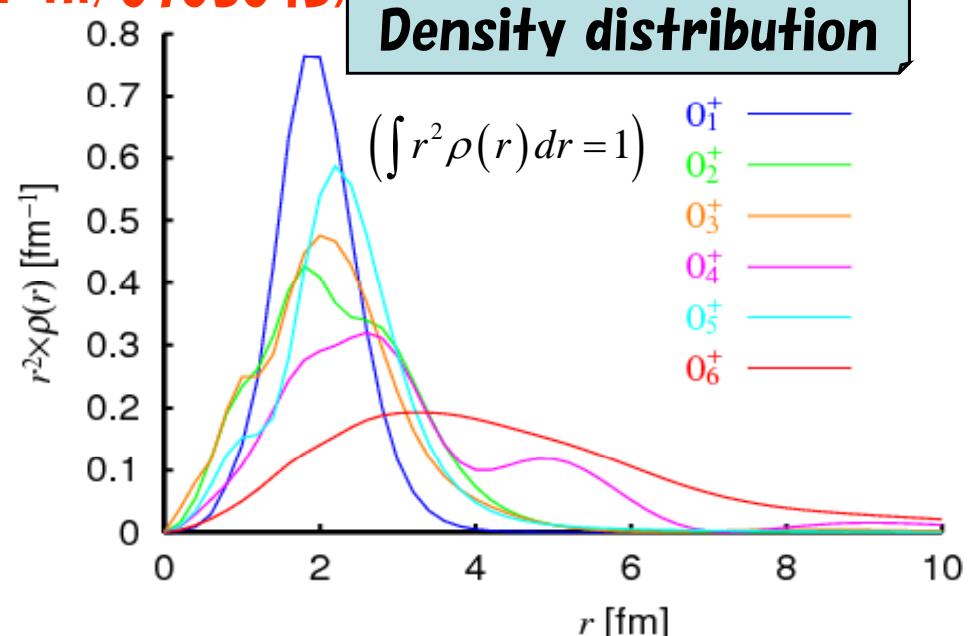
Energy levels, rms radii, monopole matrix elements and density distribution.

Low lying 0^+ levels of ^{16}O



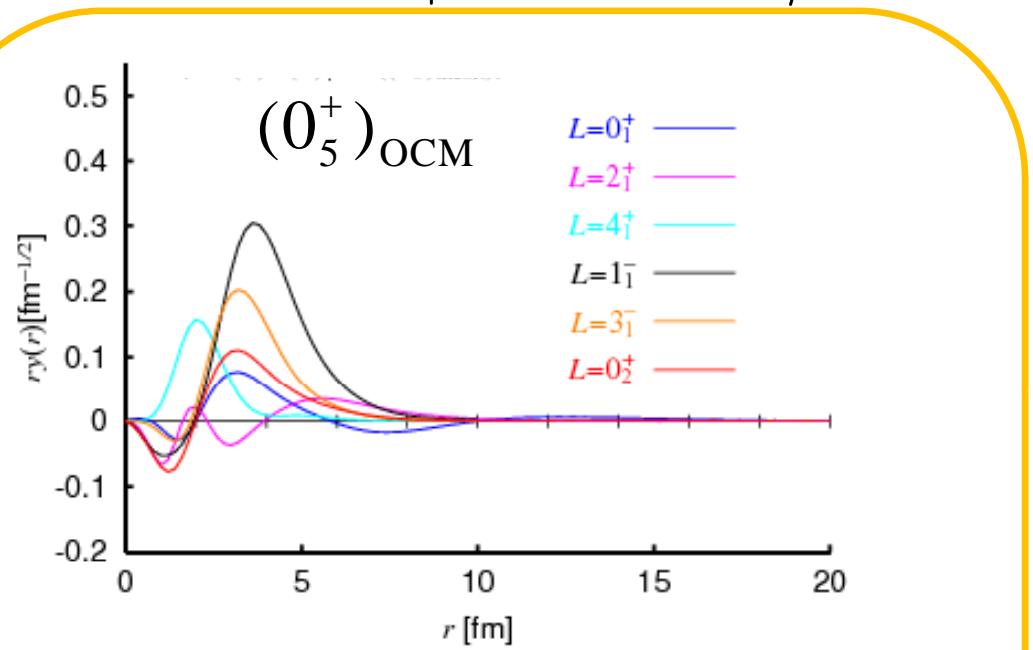
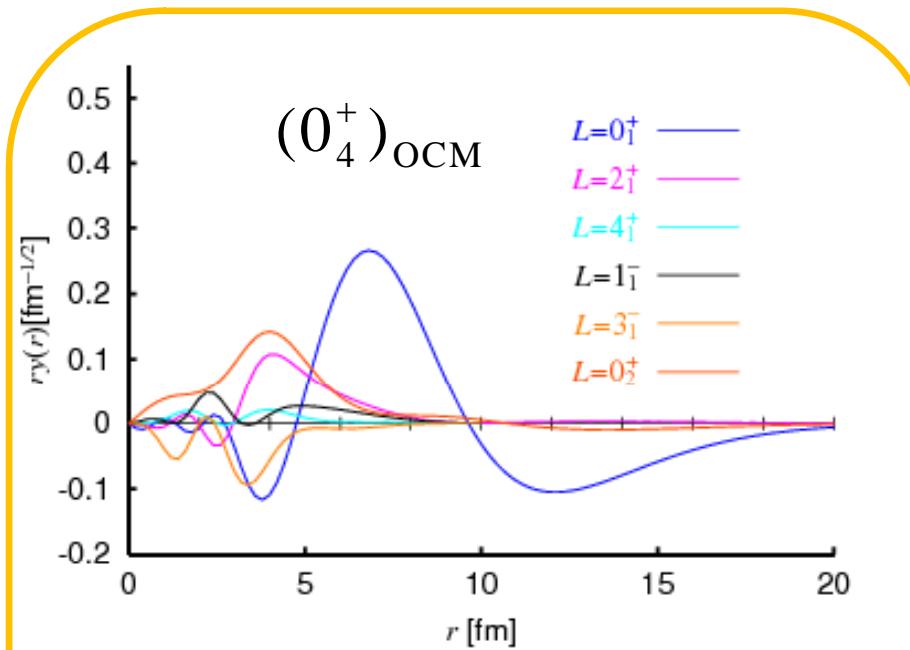
	R_{rms} (fm)	$M(E0)$ (fm ²)	$M(E0)$ (fm ²) Exp.
$(0_1^+)_{\text{OCM}}$	2.7		
$(0_2^+)_{\text{OCM}}$	3.1	4.2	$0_2^+:$ 3.55
$(0_3^+)_{\text{OCM}}$	2.9	4.1	$0_3^+:$ 4.03
$(0_4^+)_{\text{OCM}}$	3.9	2.4	$0_4^+:$ no data
$(0_5^+)_{\text{OCM}}$	3.1	2.0	$0_5^+:$ 3.3
$(0_6^+)_{\text{OCM}}$	5.4	1.4	

Large monopole matrix element can be the evidence of cluster states (Yamada, Y.F. et al., nucl-th/0703045)



Reduced width amplitudes of 0_4^+ and 0_5^+ states obtained with 4α OCM

Defined as $r \times \gamma_{L,J=0_k}(r) = r \times \left\langle \left[Y_L(\mathbf{r}) \Phi_L(^{12}\text{C}) \right]_0 \middle| \Phi_{J_k=0_k}(^{16}\text{O}) \right\rangle$



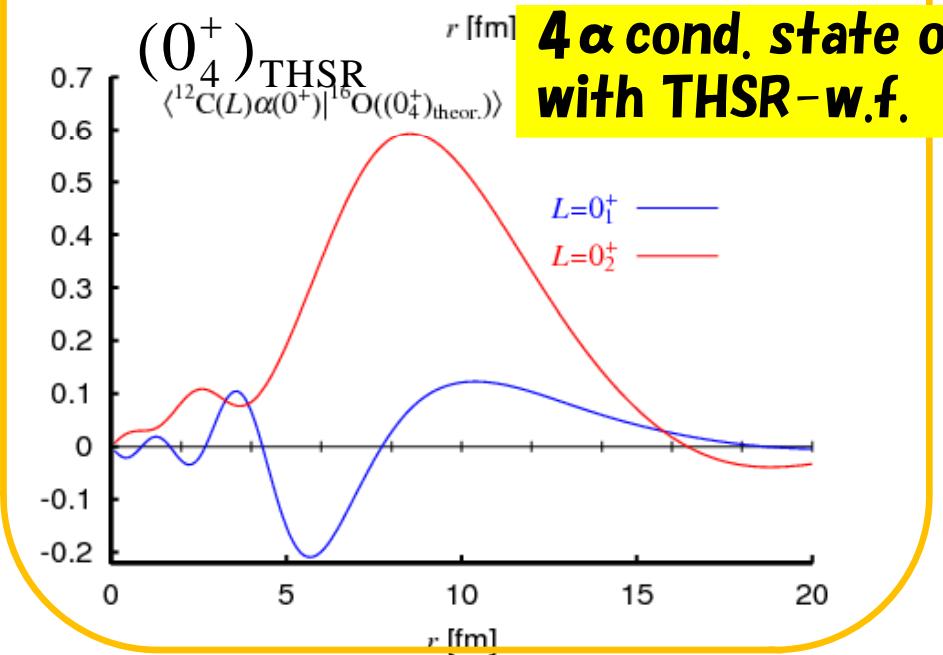
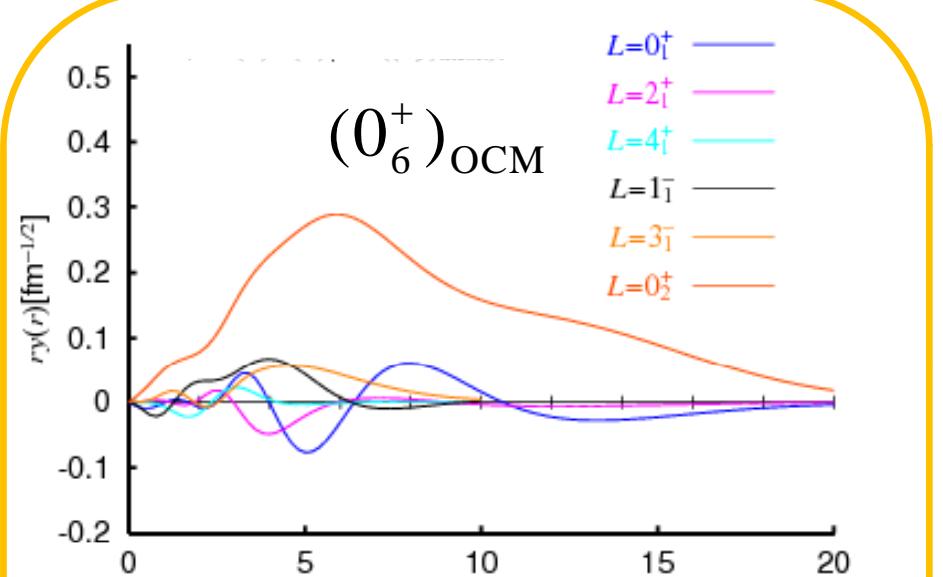
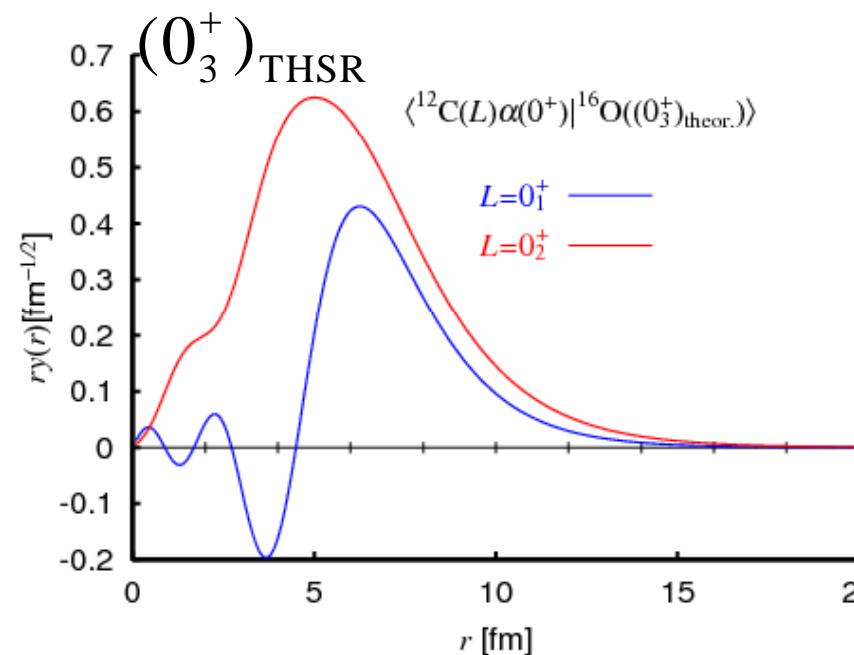
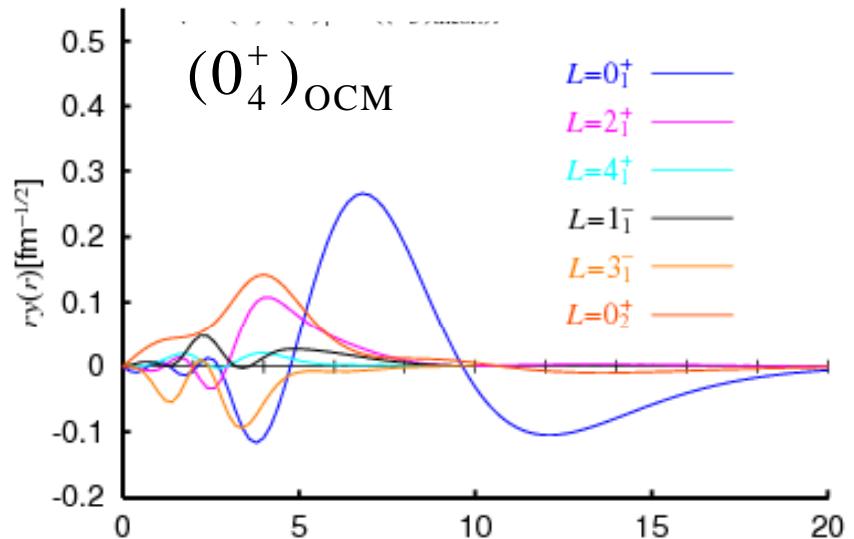
- Very well developed α cluster structure
- $^{12}\text{C}(\text{g.s.}) + \alpha$ component is dominant. higher nodal structure

- Large amplitude in surface region (~ 5 fm)
- developed α cluster structure
- mixing of $^{12}\text{C}(1^-) + \alpha$, $^{12}\text{C}(3^-) + \alpha$ and $^{12}\text{C}(0_2^+) + \alpha$ structures

- New (not discussed so far) $\alpha + ^{12}\text{C}$ cluster states.
- $\alpha + ^{12}\text{C}$ dynamics survives up to around the 4α threshold.

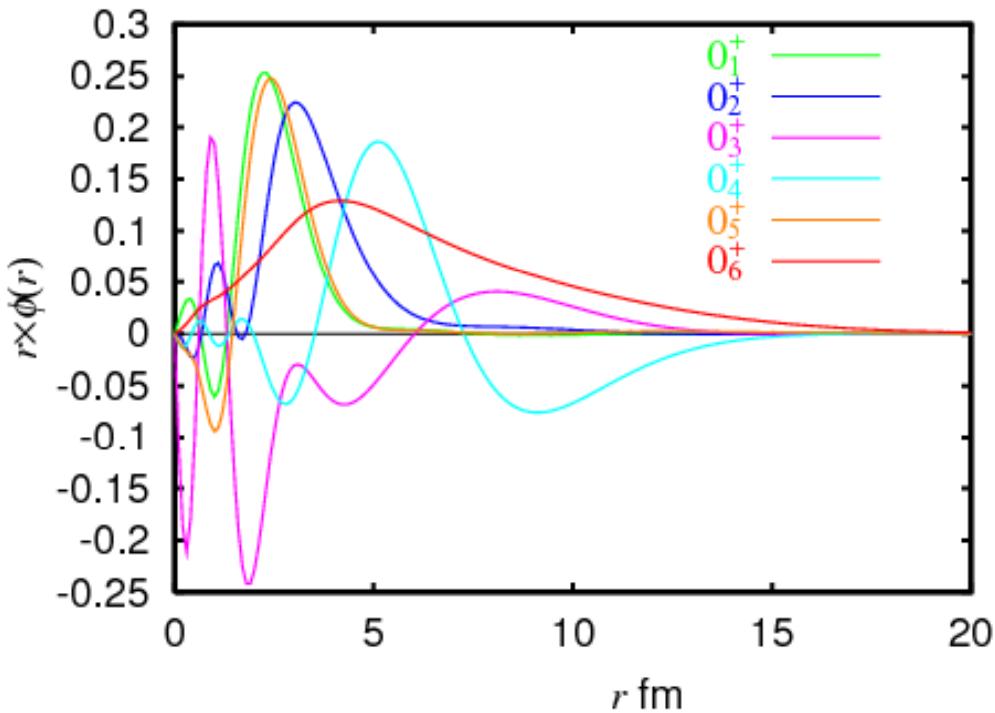
Reduced width amplitudes of 0_4^+ and 0_6^+ states obtained with 4α OCM

Defined as $r \times \gamma_{L,J=0_k}(r) = r \times \left\langle \left[Y_L(\mathbf{r}) \Phi_L(^{12}\text{C}) \right]_0 \middle| \Phi_{J_k=0_k}(^{16}\text{O}) \right\rangle$



4 α cond. state obtained with THSR-w.f.

Single particle occupancy and single particle orbit for the 0_1^+ - 0_6^+ states obtained with 4α OCM (Only the S orbit ($L=0$))



The largest (for $L = 0$) $\mu^\lambda = 2.33$

0_6^+ state : $2.33/4 = 58\%$

Large $0S$ occupancy !

Strongly evidence that the 0_6^+ state is the 4α condensate

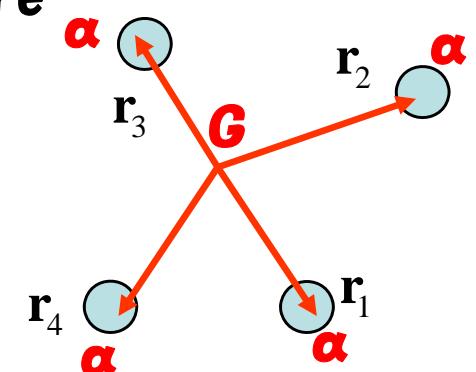
The largest occupancies for the other states are less than 20 %.

$$\int dr' \rho(r, r') f^\lambda(r') = \mu^\lambda f^\lambda(r)$$

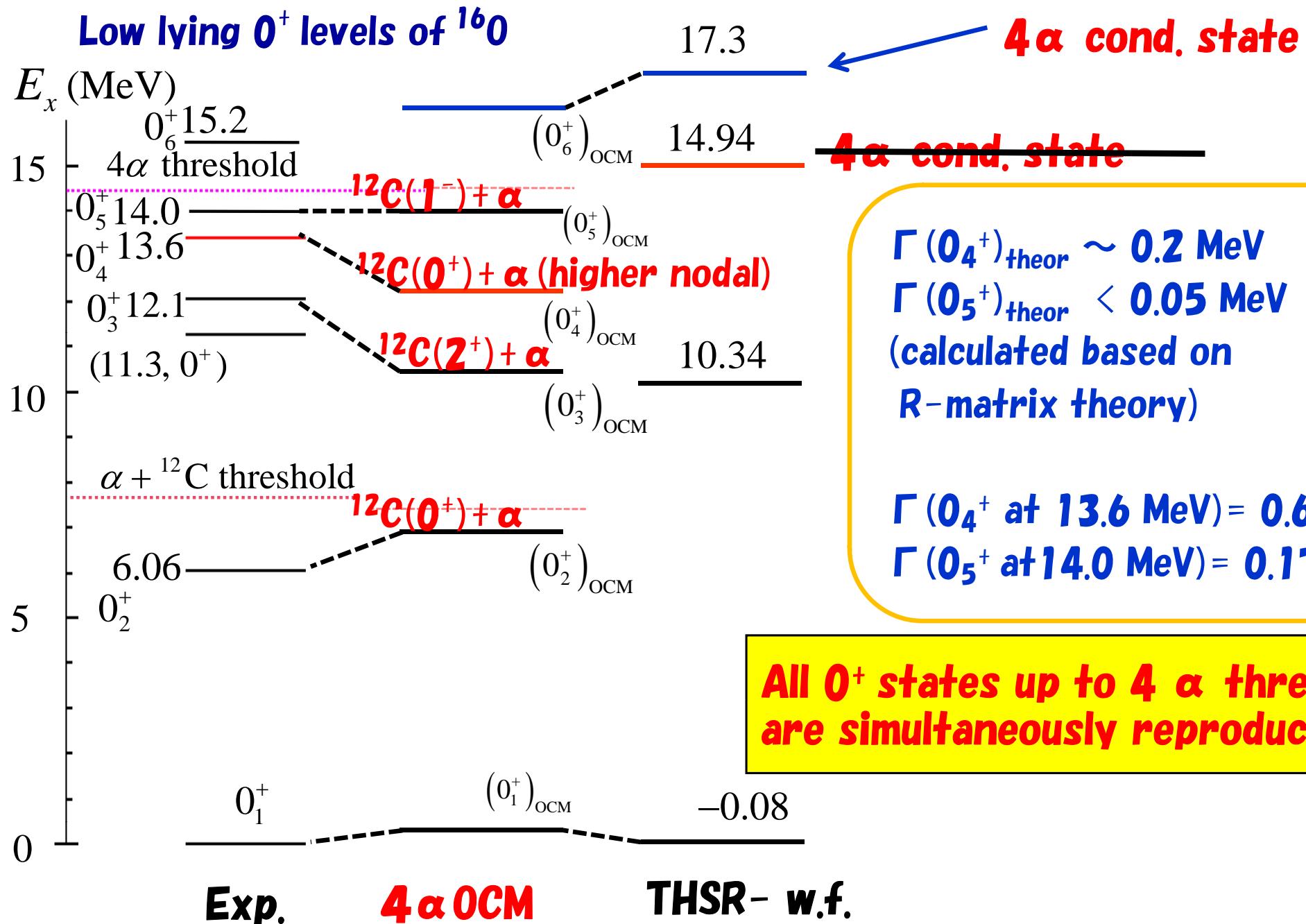
$$\rho(r, r') = 4 \int \Phi_{0_k^+}^*(r, r_2, r_3, r_4) \Phi_{0_k^+}(r', r_2, r_3, r_4) dr_2 dr_3 dr_4$$

$\lambda = (n, L)$, μ^λ : single particle occupancy, $f^\lambda(r)$: single particle orbit

$$\sum_\lambda \mu^\lambda = 4$$



Possible assignment of the two calculations and observations



Summary

- Beyond doubt the Hoyle state is the 3α condensate state. (THSR and OCM)
- 4α condensate w. f. (4α THSR-w.f.) predicts the existence of 4α condensate state. (not as the **third** 0^+ state but as the **fourth** 0^+ state)

Analysis by using 4α OCM(orthogonality condition model) in order to describe both $^{12}\text{C} + \alpha$, 4α gas states and others, if any, in larger model space.

4α condensate state (0_6^+) and other cluster states are simultaneously obtained. Large OS-occupancy, 60 %

- Two new resonance states are obtained near the 4α threshold. One has a developed α cluster structure ($R_{rms} \sim 3.0$ fm) in which $^{12}\text{C}(1^-) + \alpha$, $^{12}\text{C}(3^-) + \alpha$ and $^{12}\text{C}(0_2^+) + \alpha$ components are mixed. The other has a very well developed α cluster structure ($R_{rms} \sim 4.0$ fm). $^{12}\text{C}(0^+) + \alpha$ (higher nodal)
⇒ corresponding to the observed 0_4^+ and 0_5^+ states, respectively
- Successfully reproducing the well known $^{12}\text{C}(\text{g.s.}) + \alpha$ (6.05 MeV) (0_2^+) and $^{12}\text{C}(2^+) + \alpha$ (12.05 MeV) (0_3^+) structures,

For future work,

- Analyses of condensate fraction and 4α CSM are necessary for more reliable conclusion.
- 4α linear chain structure (the band head is estimated at 16.7 MeV)
Another type of 4α structure !
We are able to discuss it simultaneously with the lower structures !